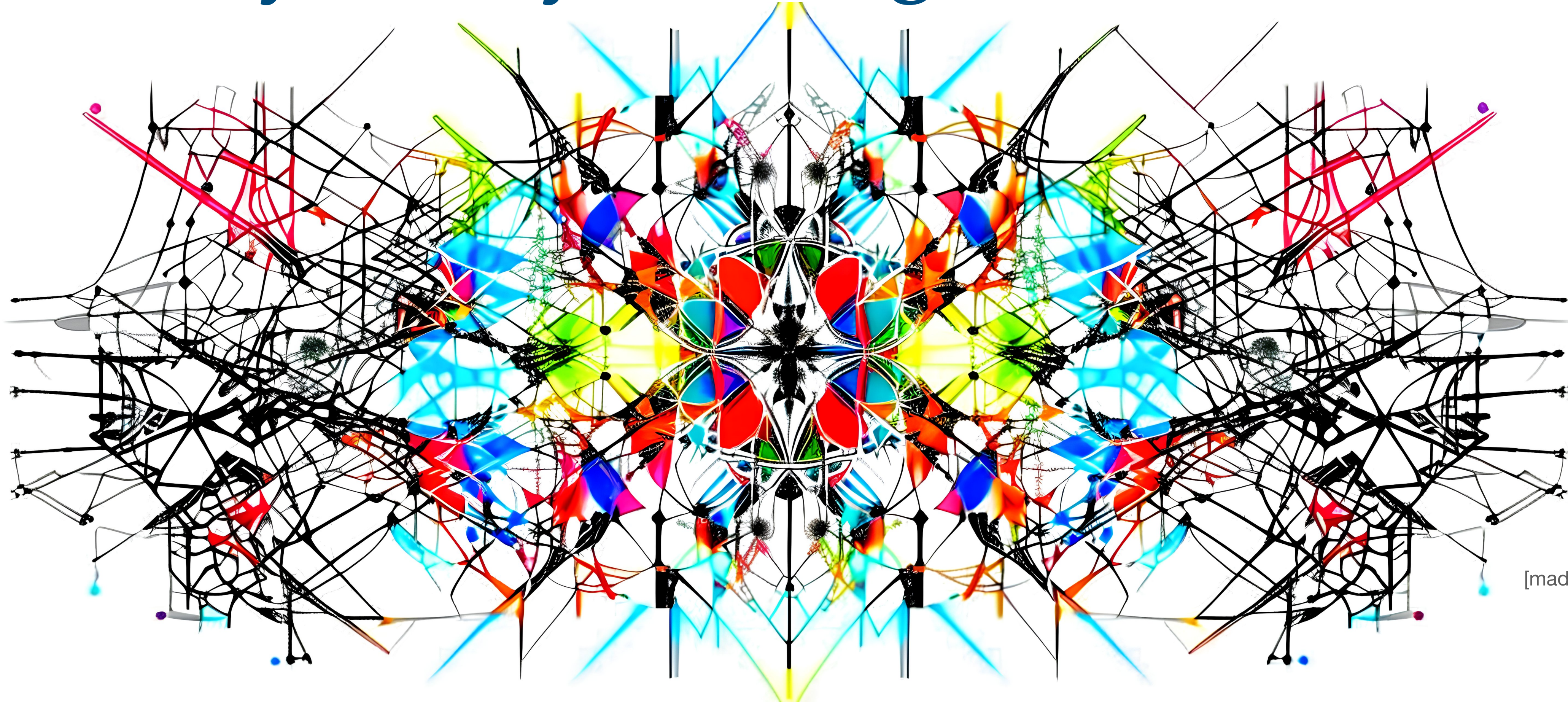


Imprints of Electro-Weak Symmetry Breaking at Colliders



[made with stablecog]

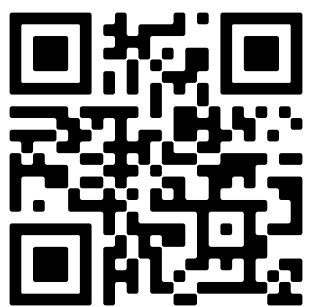
Andreas Papaefstathiou

Kennesaw State University, GA, USA

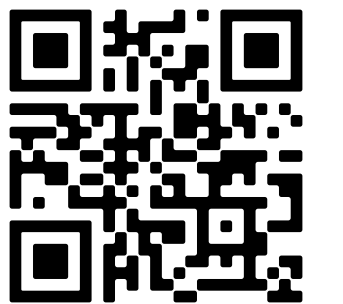
@ Georgia State University [September 26th 2023]



**KENNESAW STATE
UNIVERSITY**

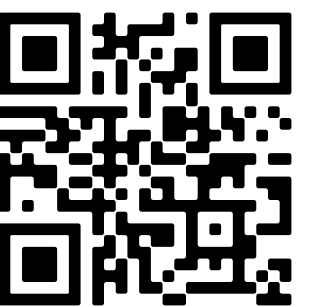


Today's Plan:



Today's Plan:

- 1 High Energy Physics
- 2 Open Questions,
- 3 → the Breaking of Symmetry,
- 4 Extended Scalar Sectors.



1

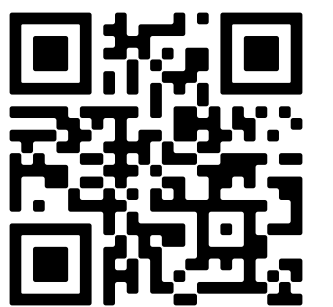
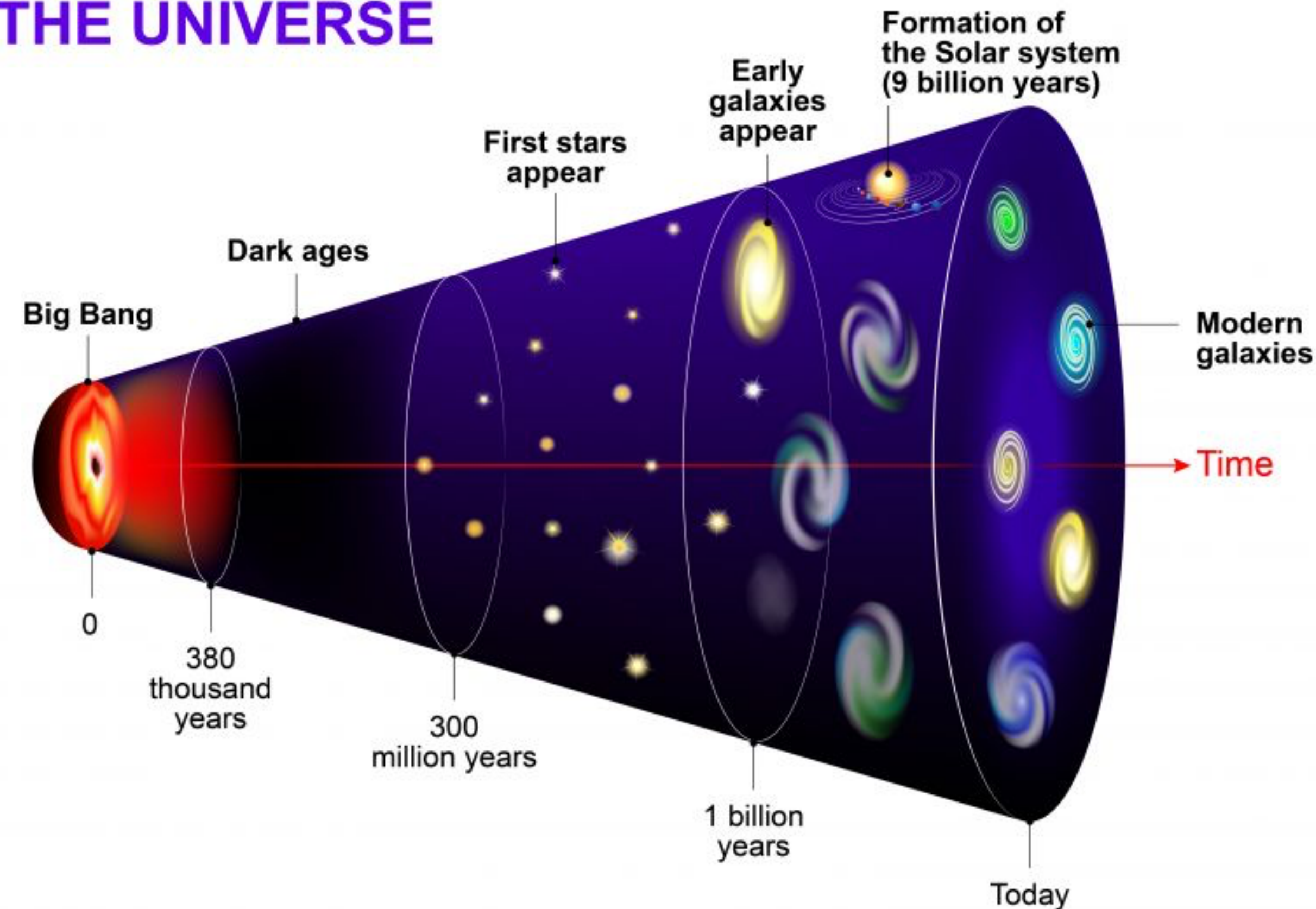
What is High Energy Physics?



What is High Energy Physics?

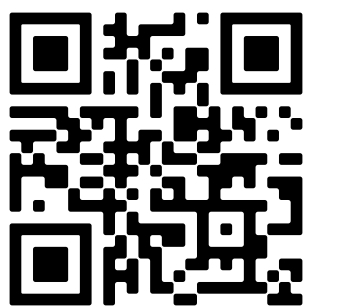
Aim: Smash particles together to “emulate” conditions closer to the Big Bang! (“Particle Archaeology”)

EVOLUTION OF THE UNIVERSE



What is High Energy Physics?

Aim: Smash particles together to understand the structure of matter today! (“Particle Sociology”)



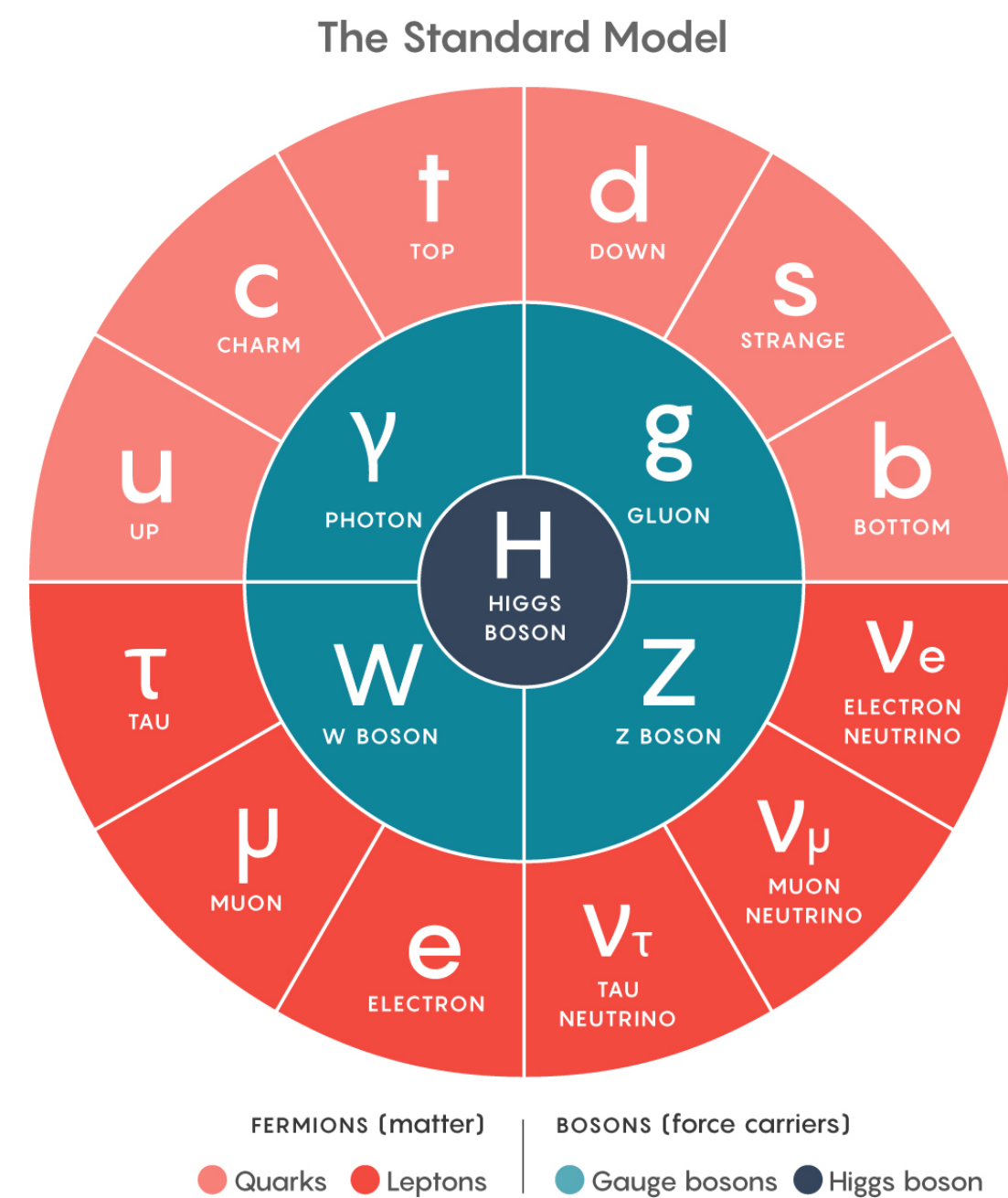
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e.g. **The “Standard Model” (SM)**



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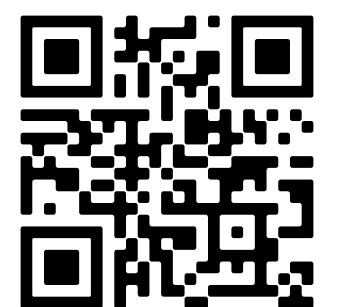
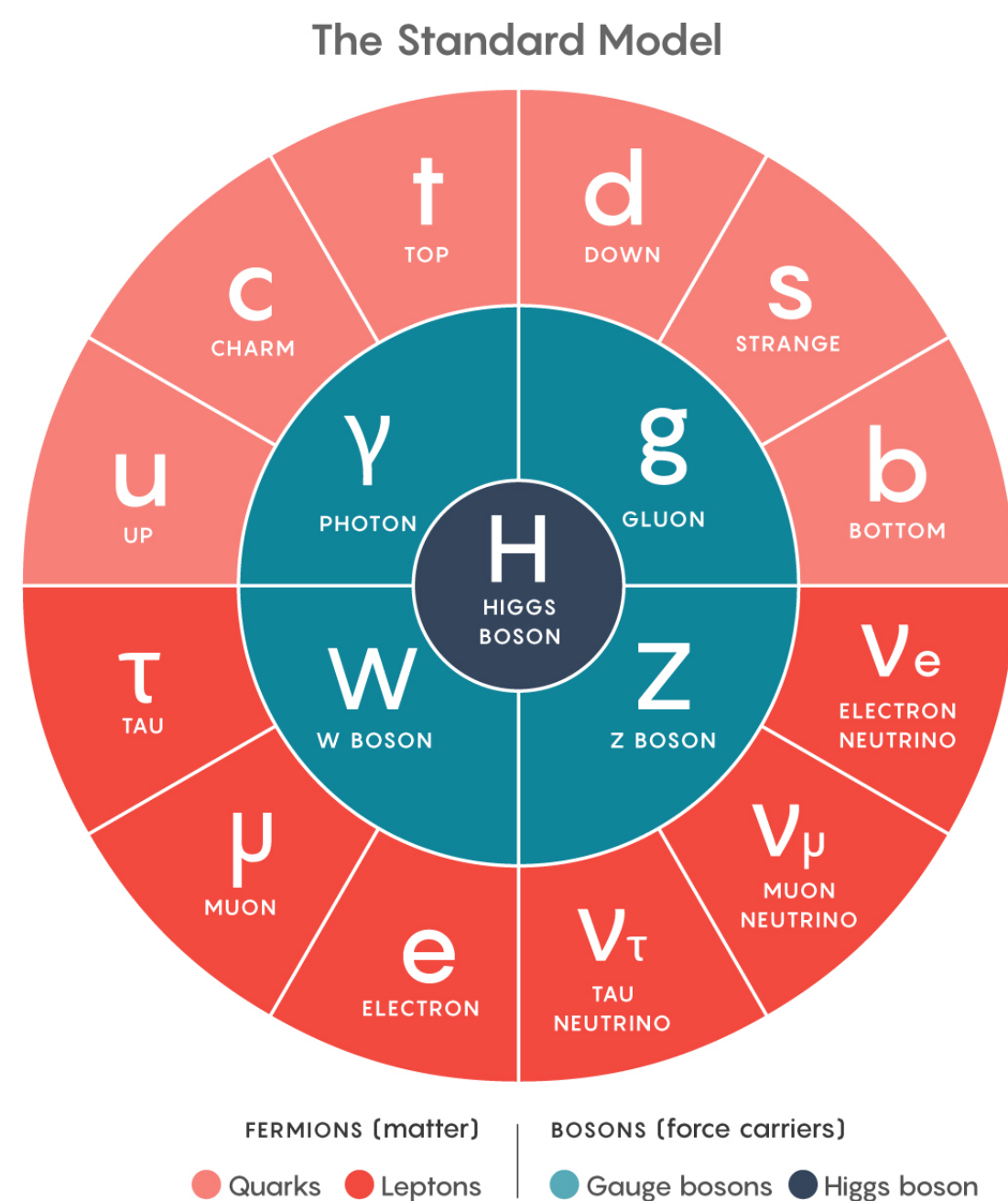
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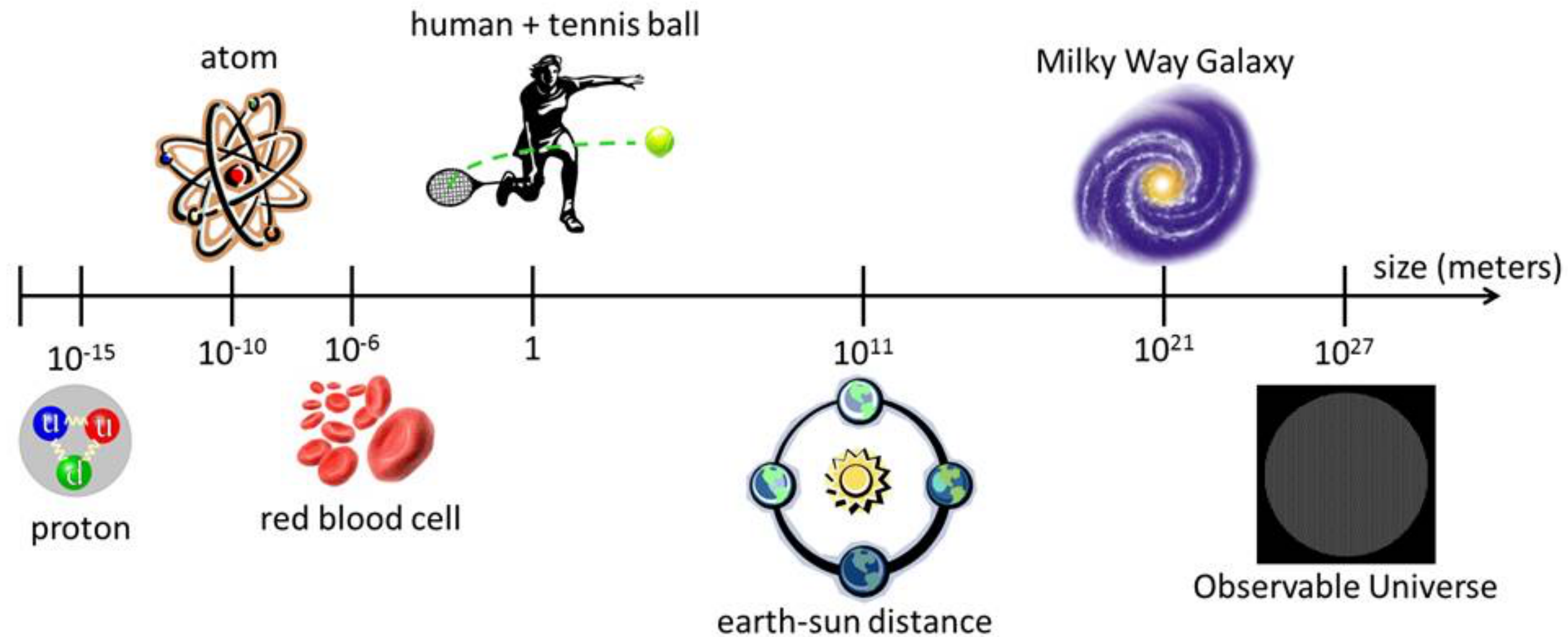
e.g. **The “Standard Model” (SM)**

and discover **exotic** phenomena!



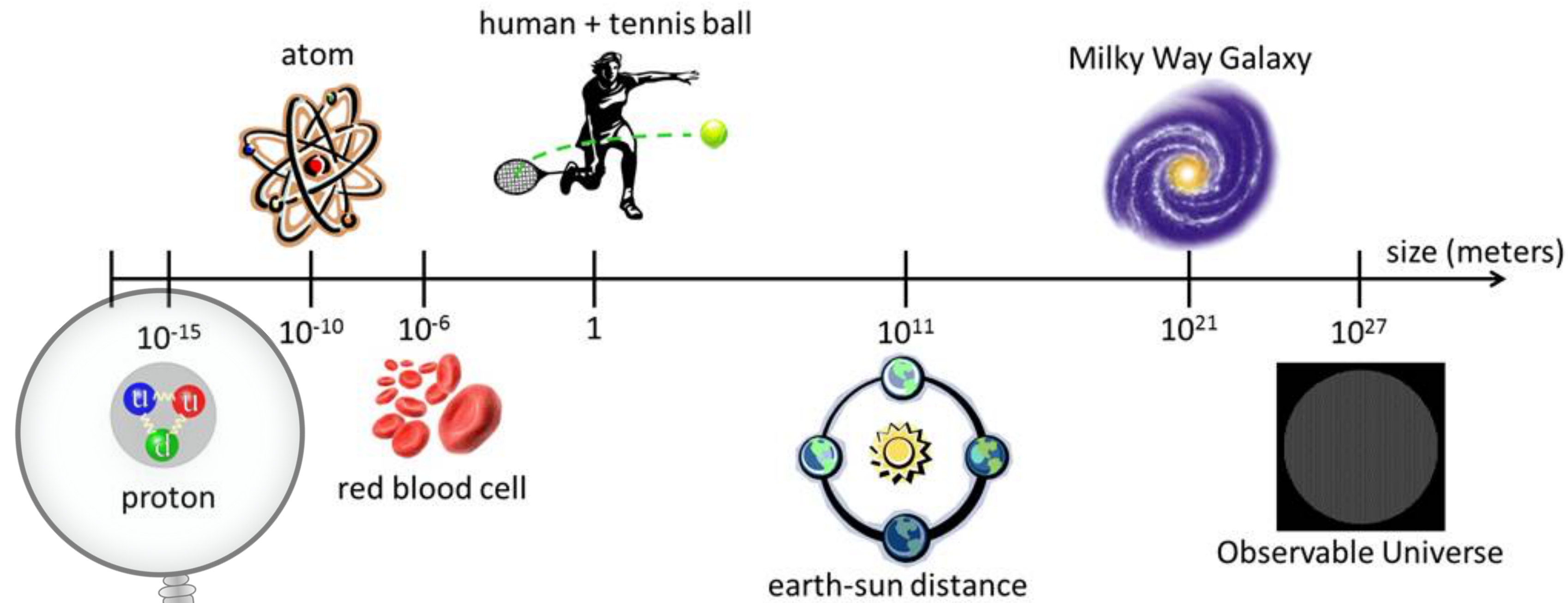
What is High Energy Physics?

Rule: Higher Energy \iff Smaller Scales!



What is High Energy Physics?

Rule: Higher Energy \iff Smaller Scales!



$m_p \sim 1 \text{ GeV} \Rightarrow$ To probe proton structure, go to higher energies than this!



High Energy Particle Colliders

Large Hadron Collider
@ CERN
in Geneva, Switzerland



Center-of-mass energy:
13,600 GeV = 13.6 TeV

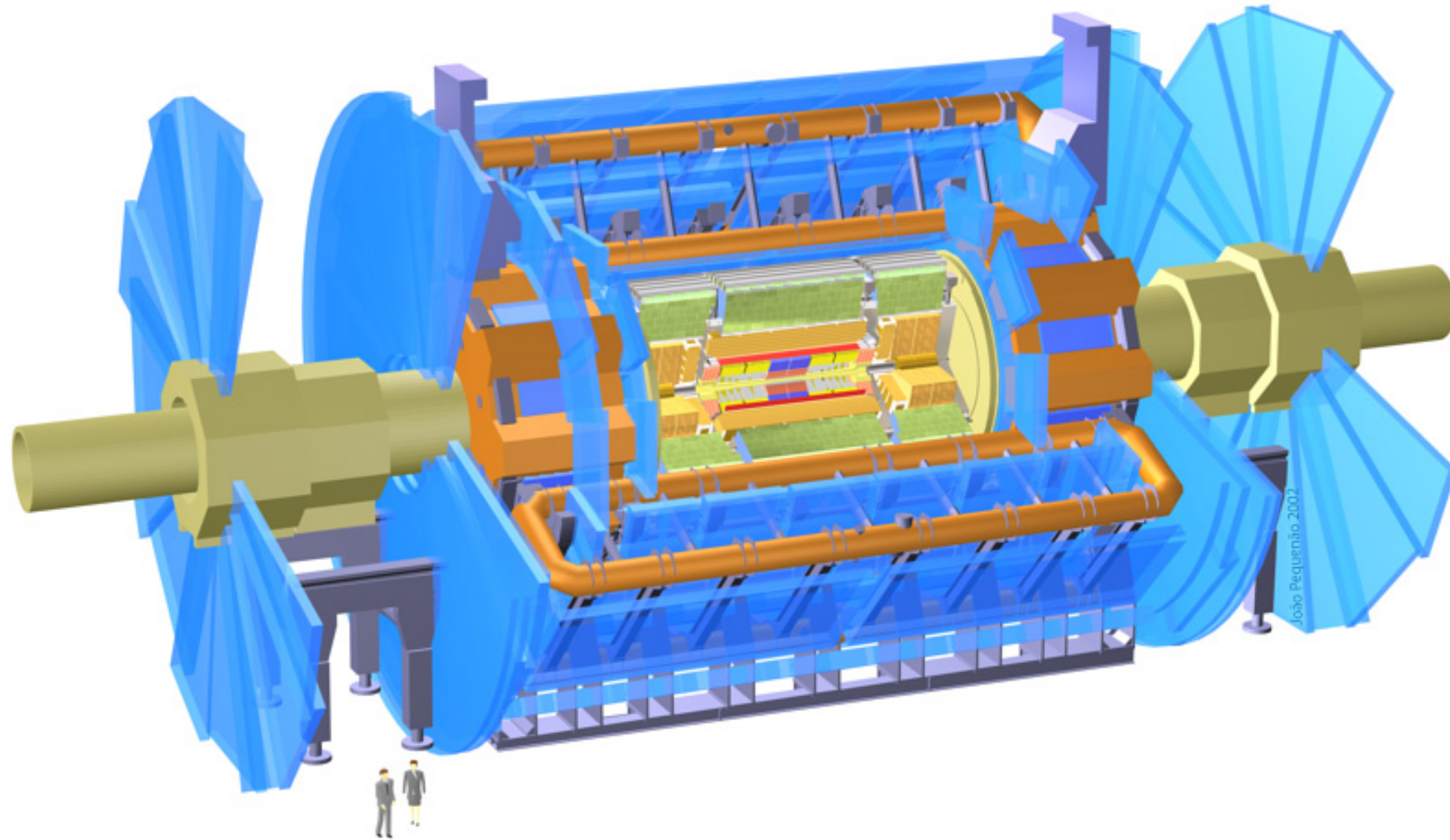


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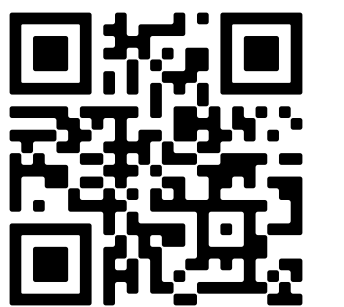
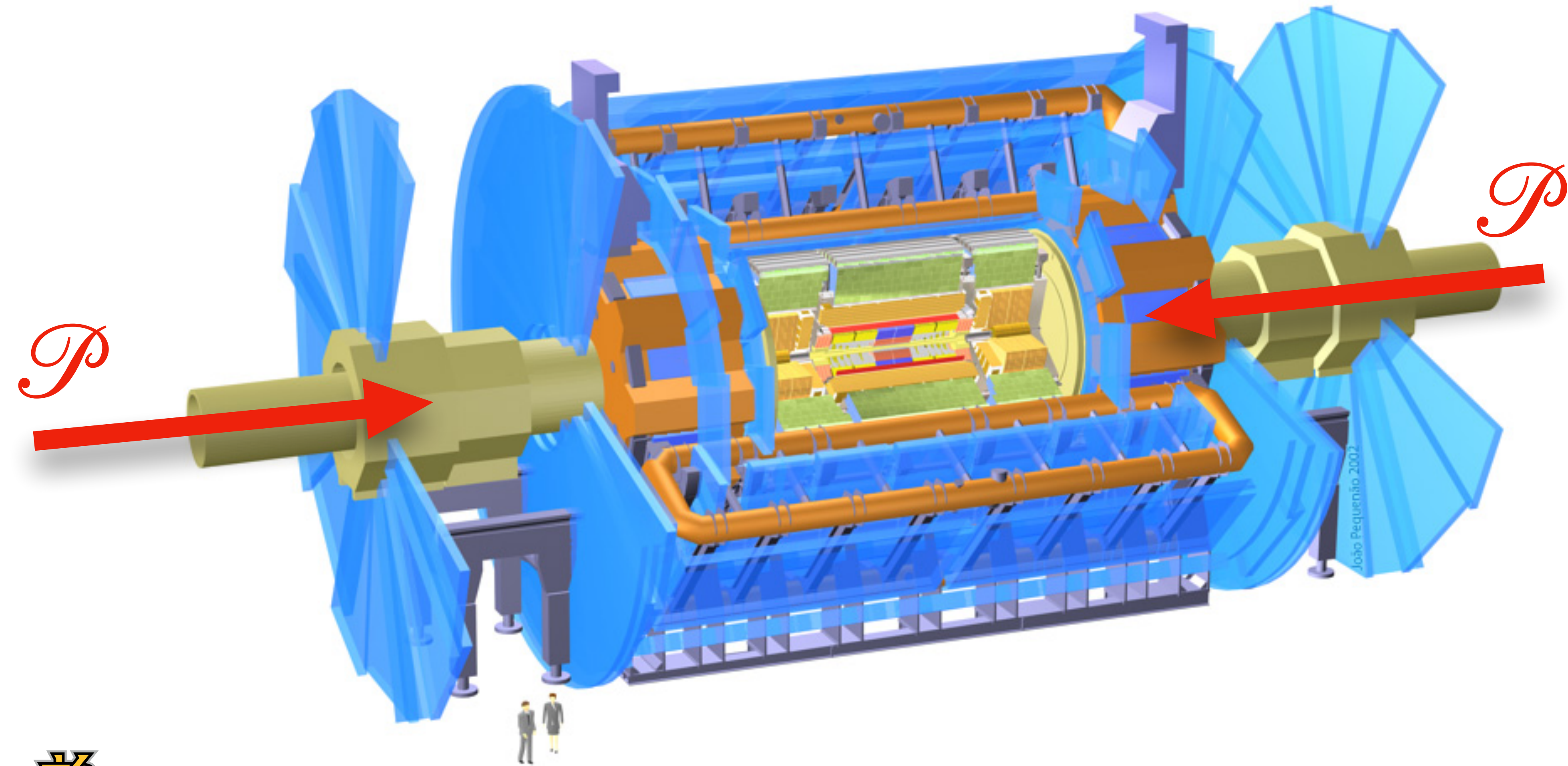
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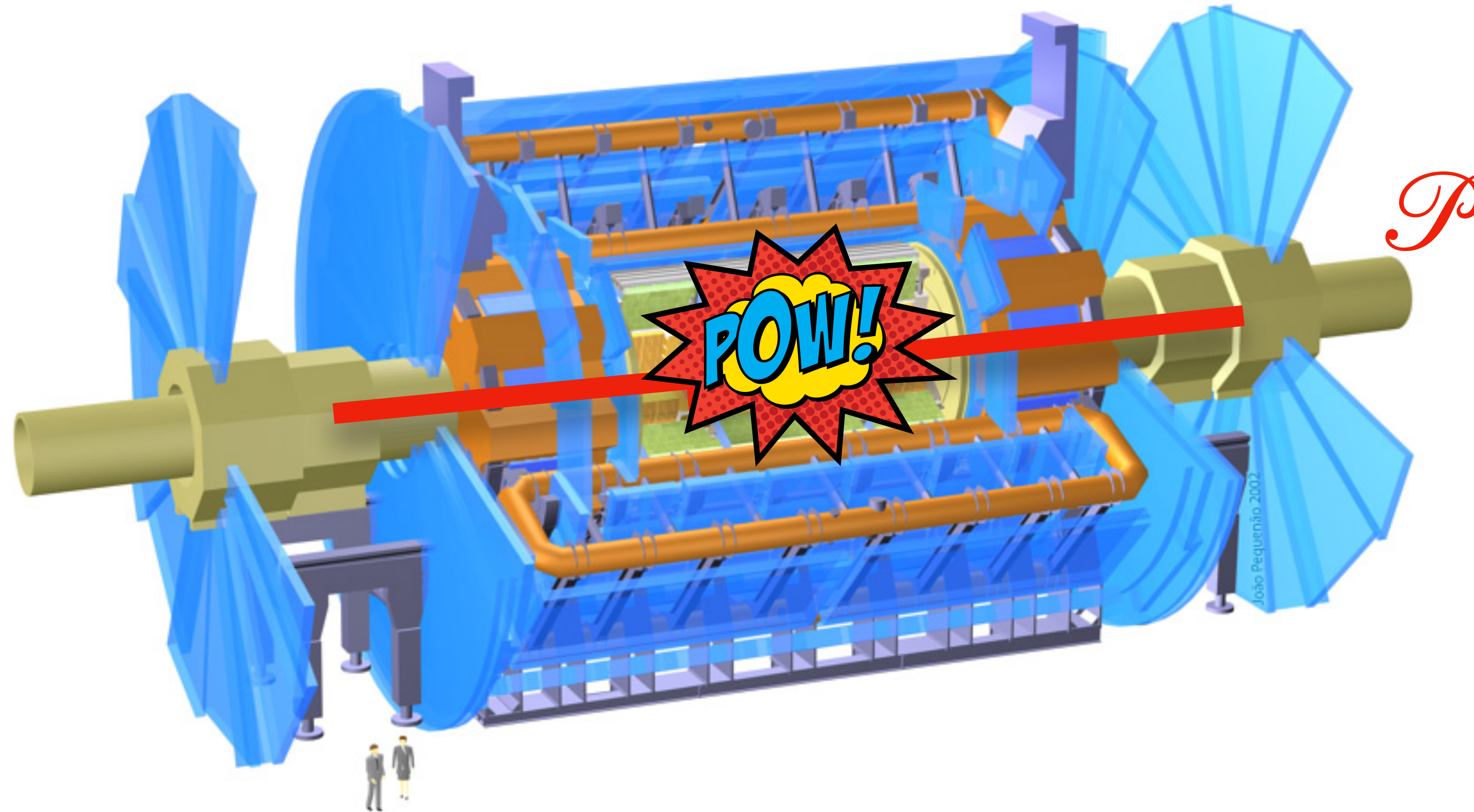
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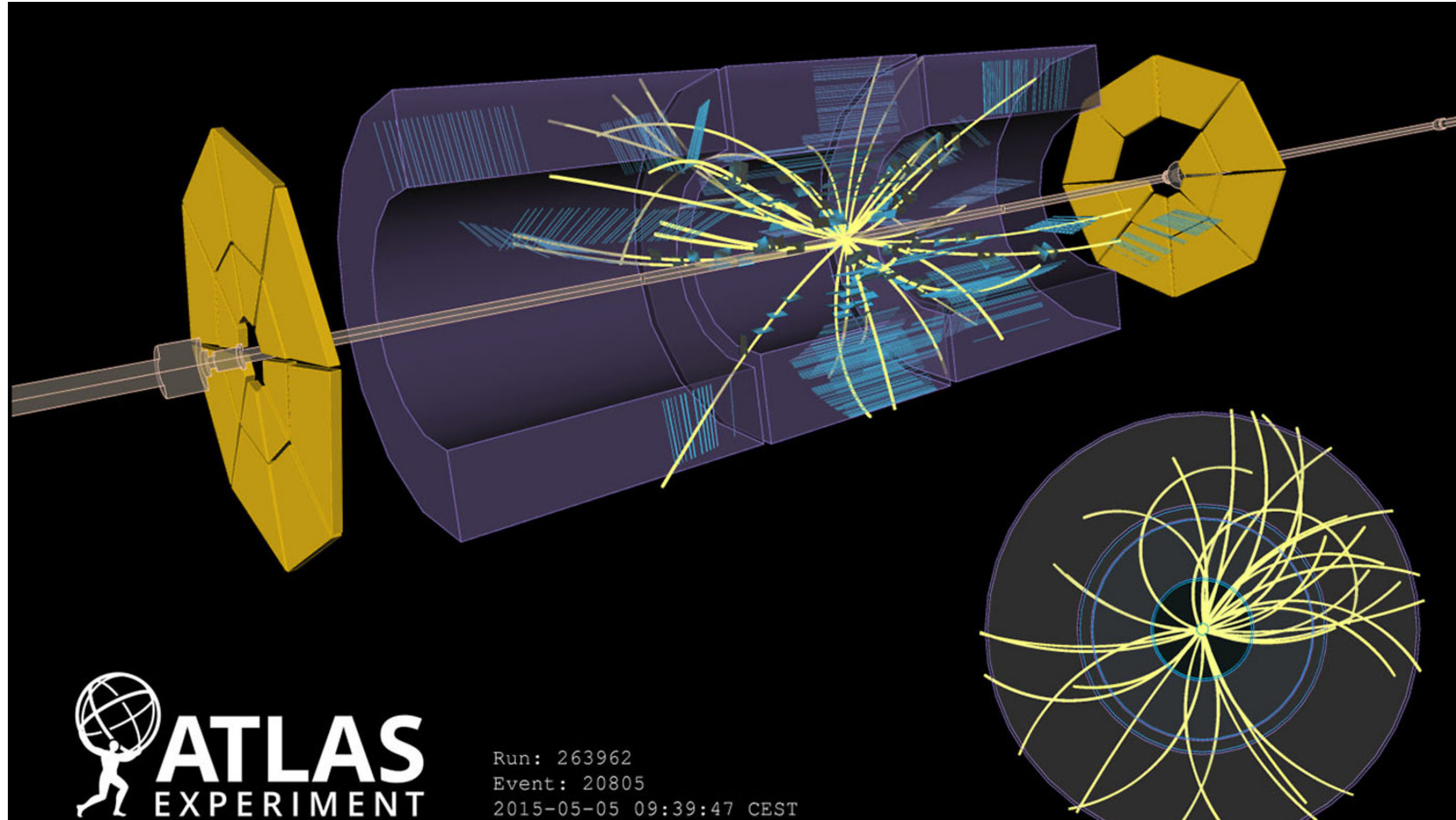


LHC: Fun Facts

- Circumference of **27 km** (17 miles).
- Lies **100 m** (330 feet) underground.
- Produces **hundreds of millions of particle collisions/second** (=“events”).
- **Energy stored** in particle beam = Energy of **1000 kg car @ 3000 km/h** (~1900 mph)!
- The data recorded can fill around **100,000 DVDs each year!**



This is a real “event”!



How do we make sense of it all?



Pictured: Richard Feynman playing the bongos.

How do we make sense of it all?

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + i\bar{\Psi}\not{D}\Psi + \Psi_i y_{ij} \Psi_j + \text{h.c.} + |D_\mu\Phi|^2 - V(\Phi)$$

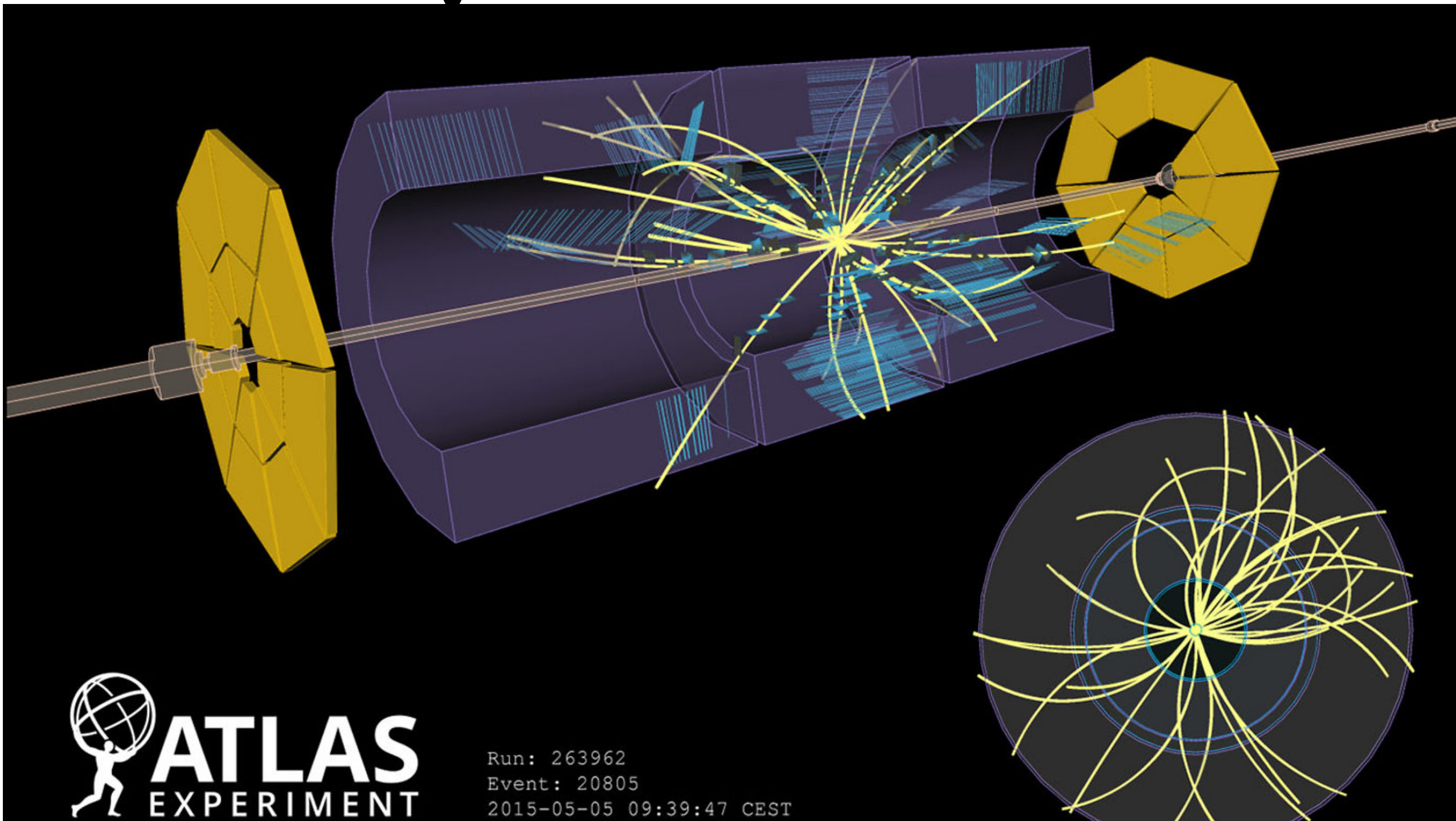


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Experiment



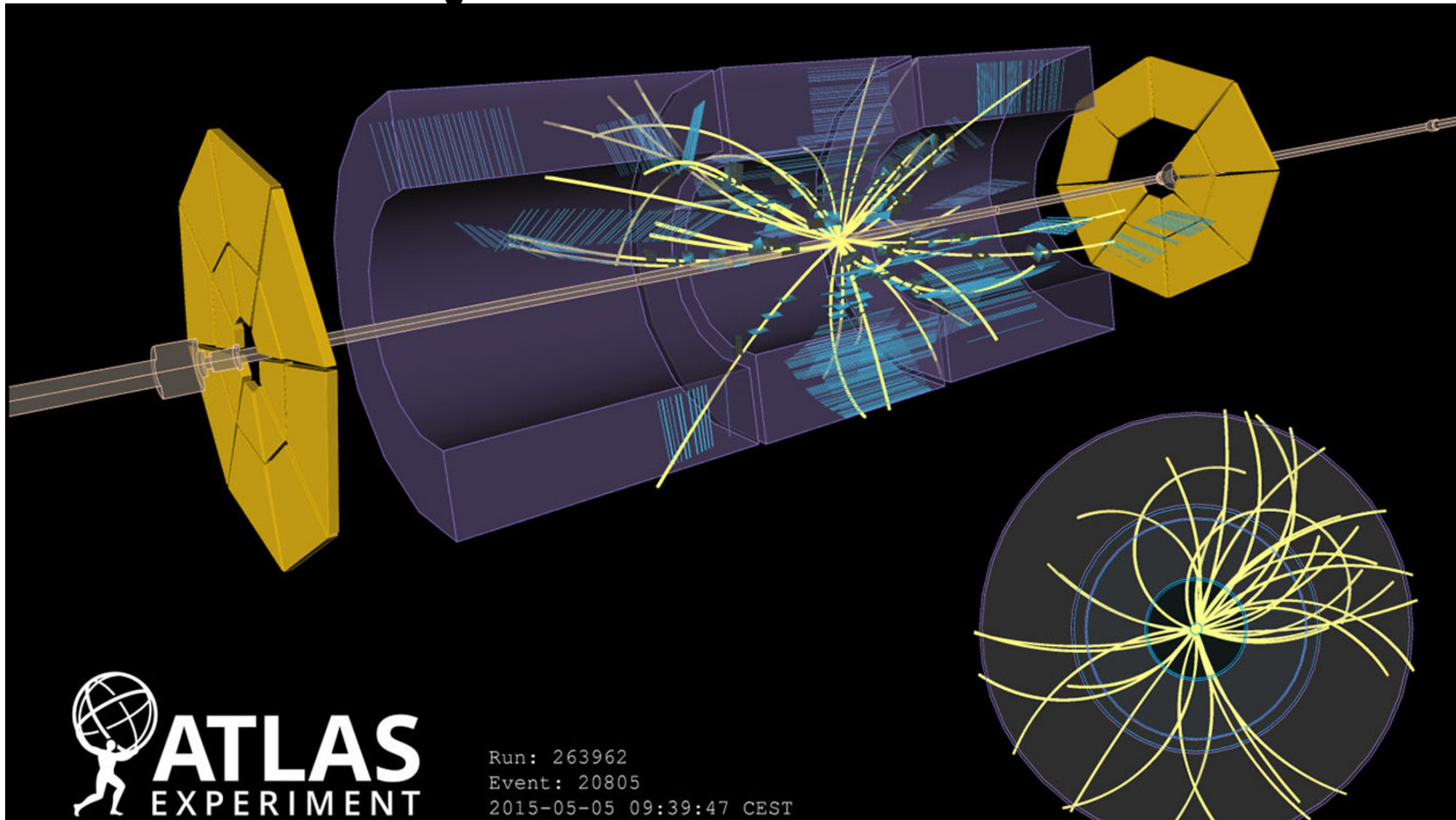
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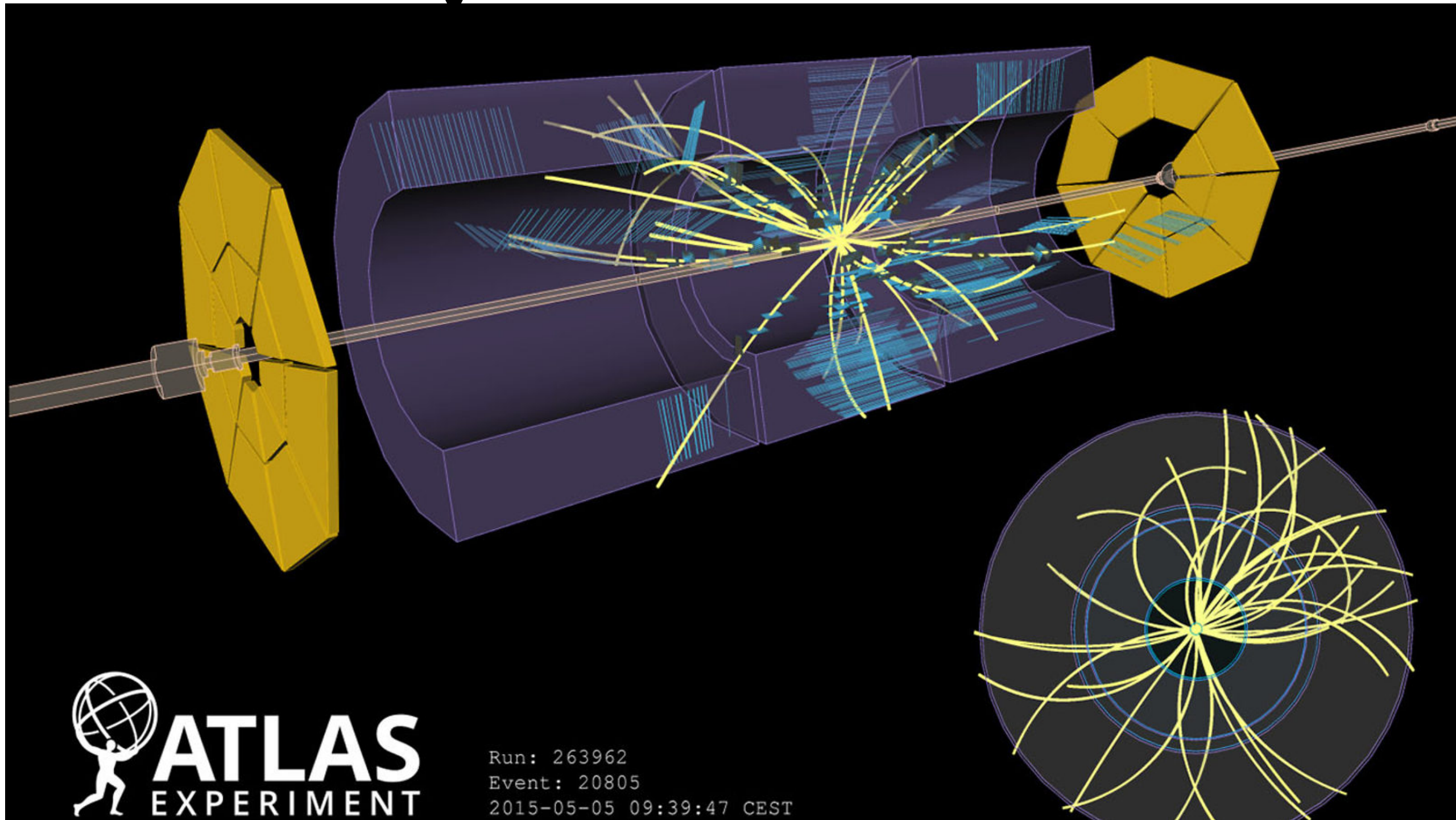
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How do we make sense of it all?

Shut up and Calculate!

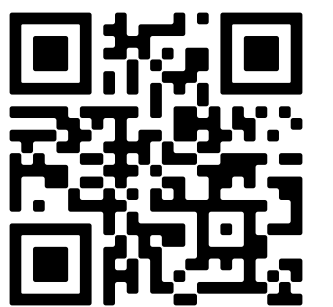
???

Experiment



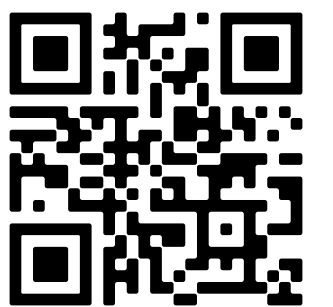
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(Shut up and Calculate) Using Quantum Field Theory



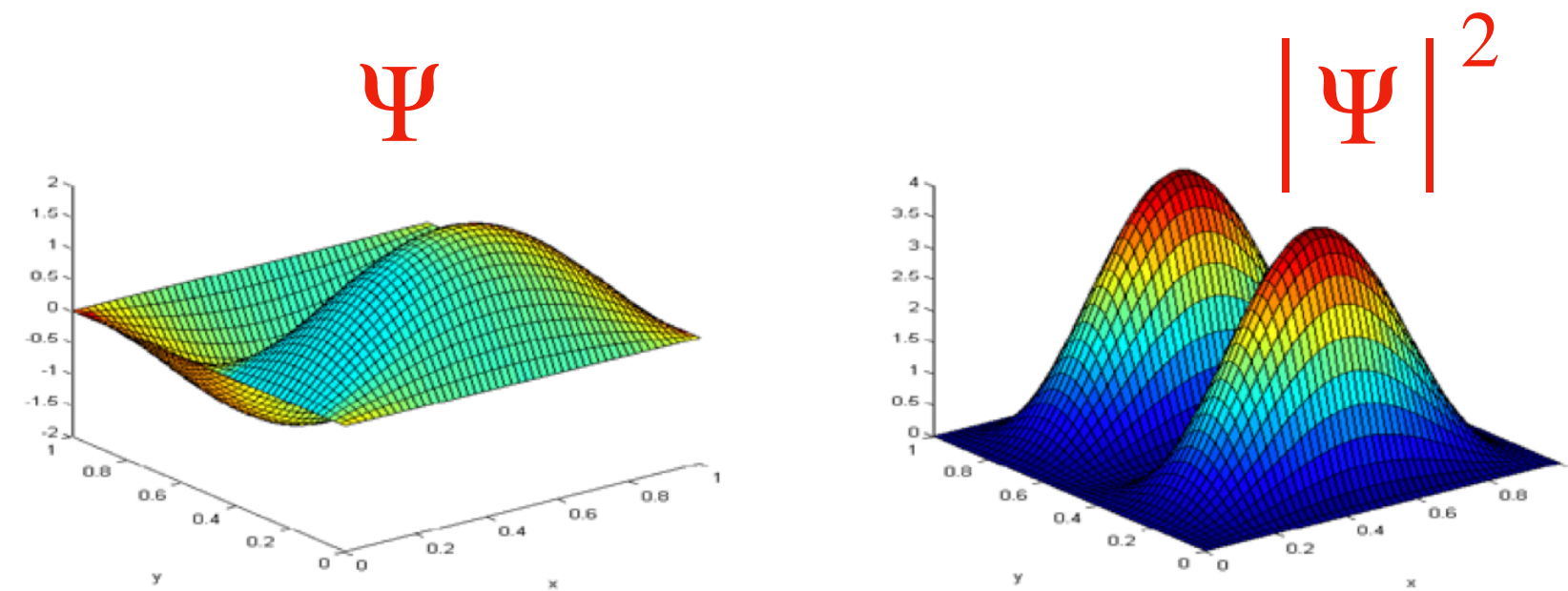
(Shut up and Calculate) Using Quantum Field Theory

- **Quantum Mechanics** → weird & wonderful world of the very small objects.



(Shut up and Calculate) Using Quantum Field Theory

- Quantum Mechanics → weird & wonderful world of the very small objects.



→ Calculate Ψ ,

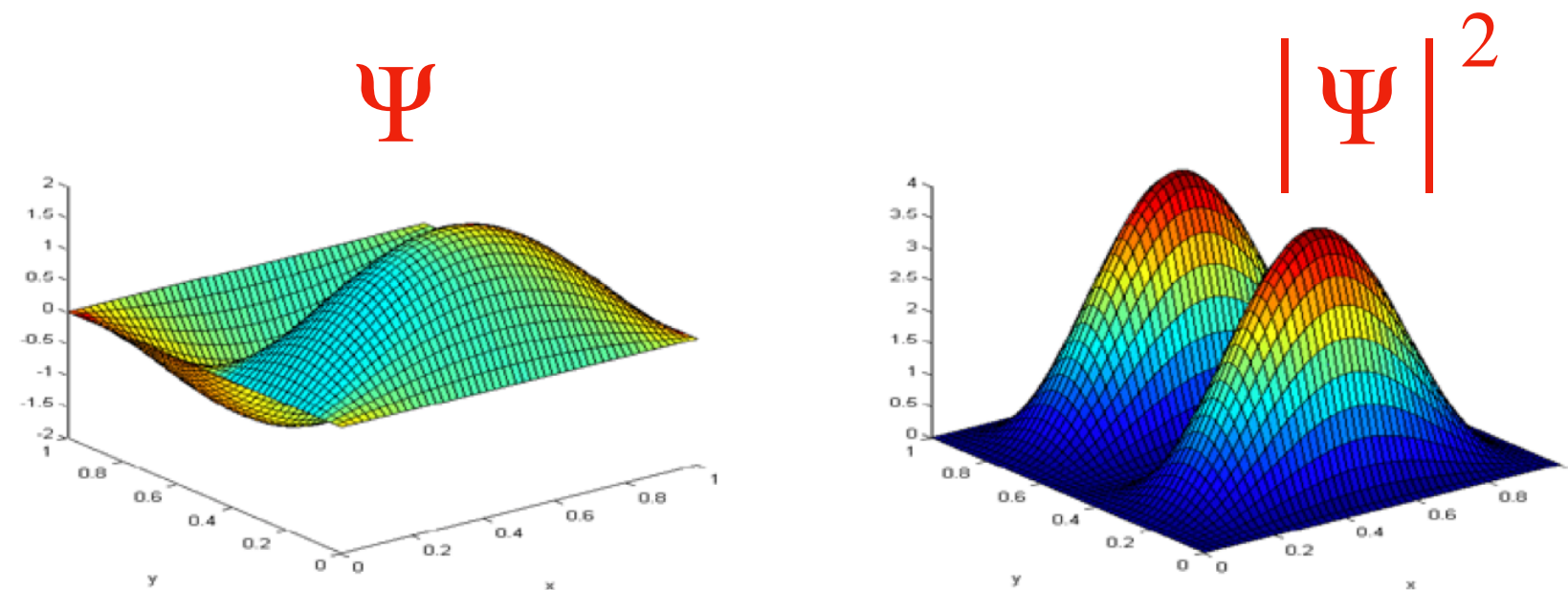
⇒ Probability to find particle somewhere $\propto |\Psi|^2$

[Ψ is called the “wave function”]



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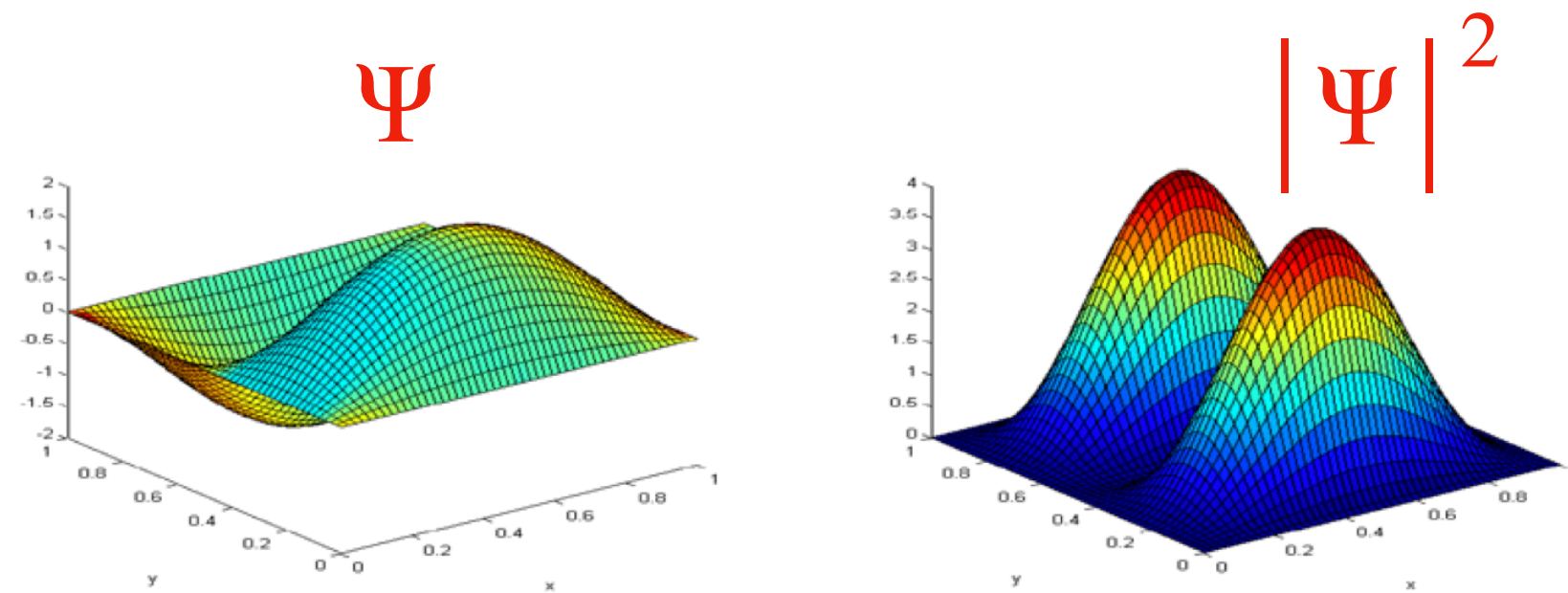
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- **Special Relativity** describes the world of the very **fast-moving objects**.



(Shut up and Calculate) Using Quantum Field Theory

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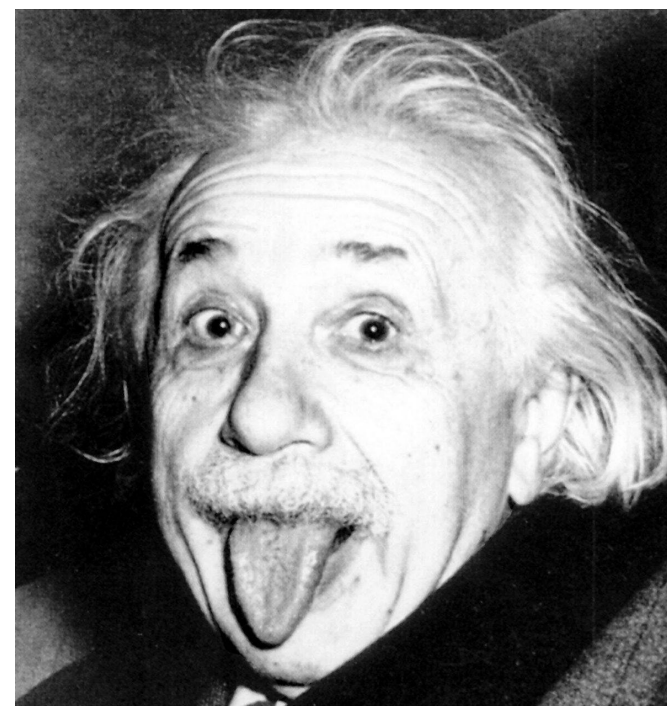


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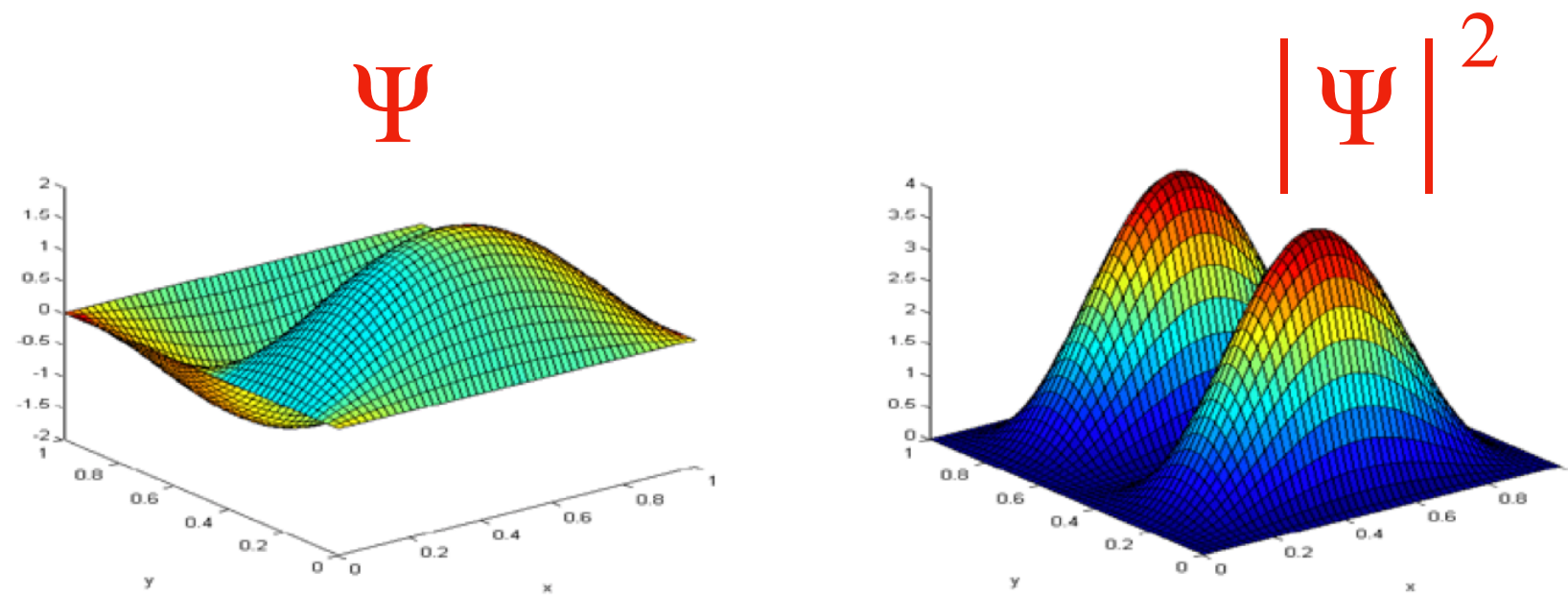


$$E = mc^2$$



(Shut up and Calculate) Using Quantum Field Theory

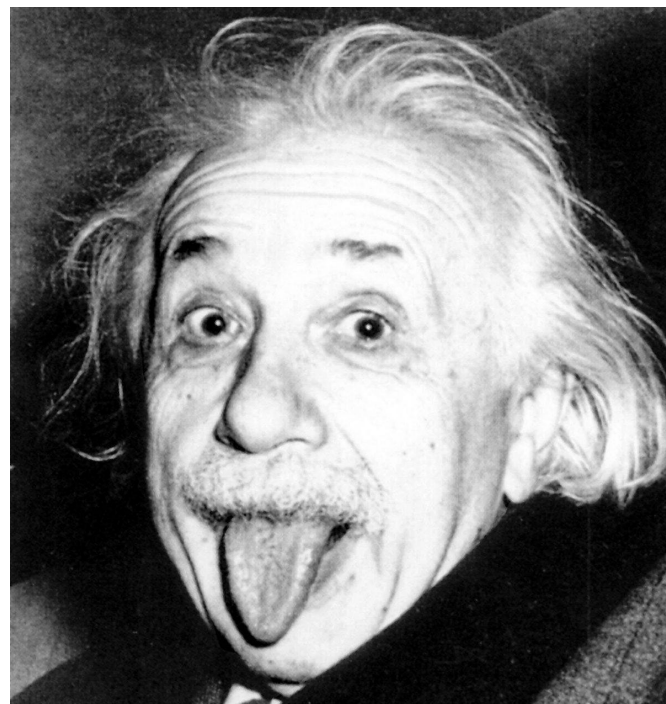
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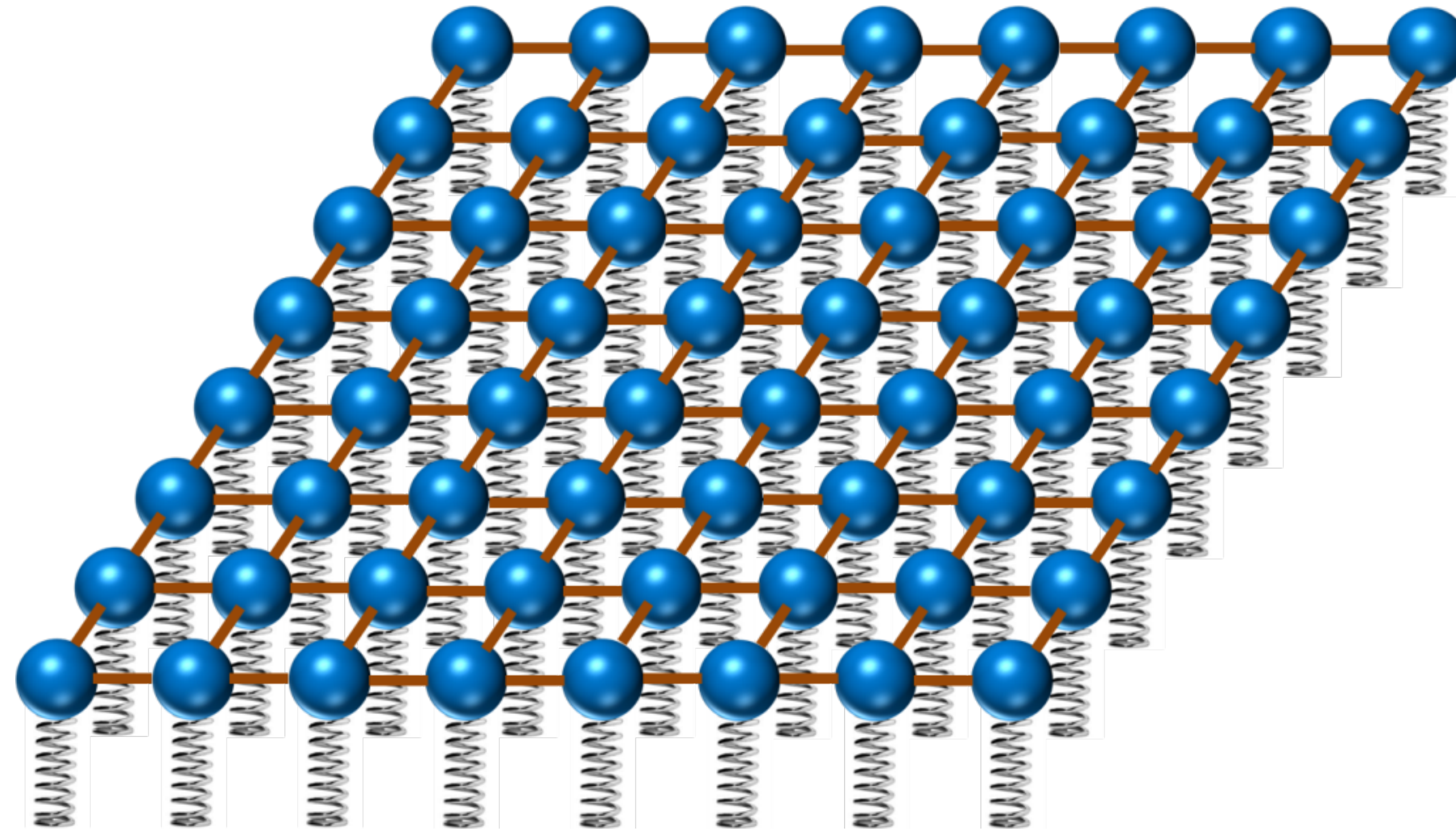


$$E = mc^2$$

- Mass and Energy are equivalent!
- ⇒ You can exchange one for the other!
- ⇒ Particle creation from Energy!



Quantum Mechanics + Special Relativity \approx **Quantum Field Theory**



From Theory to Experiment and Back Again



THEORY

EXPERIMENT



From Theory to Experiment and Back Again

$e^- \bullet$

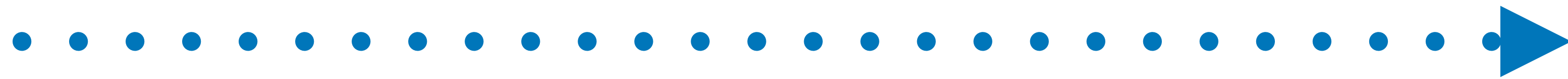
$e^- \bullet$



From Theory to Experiment and Back Again



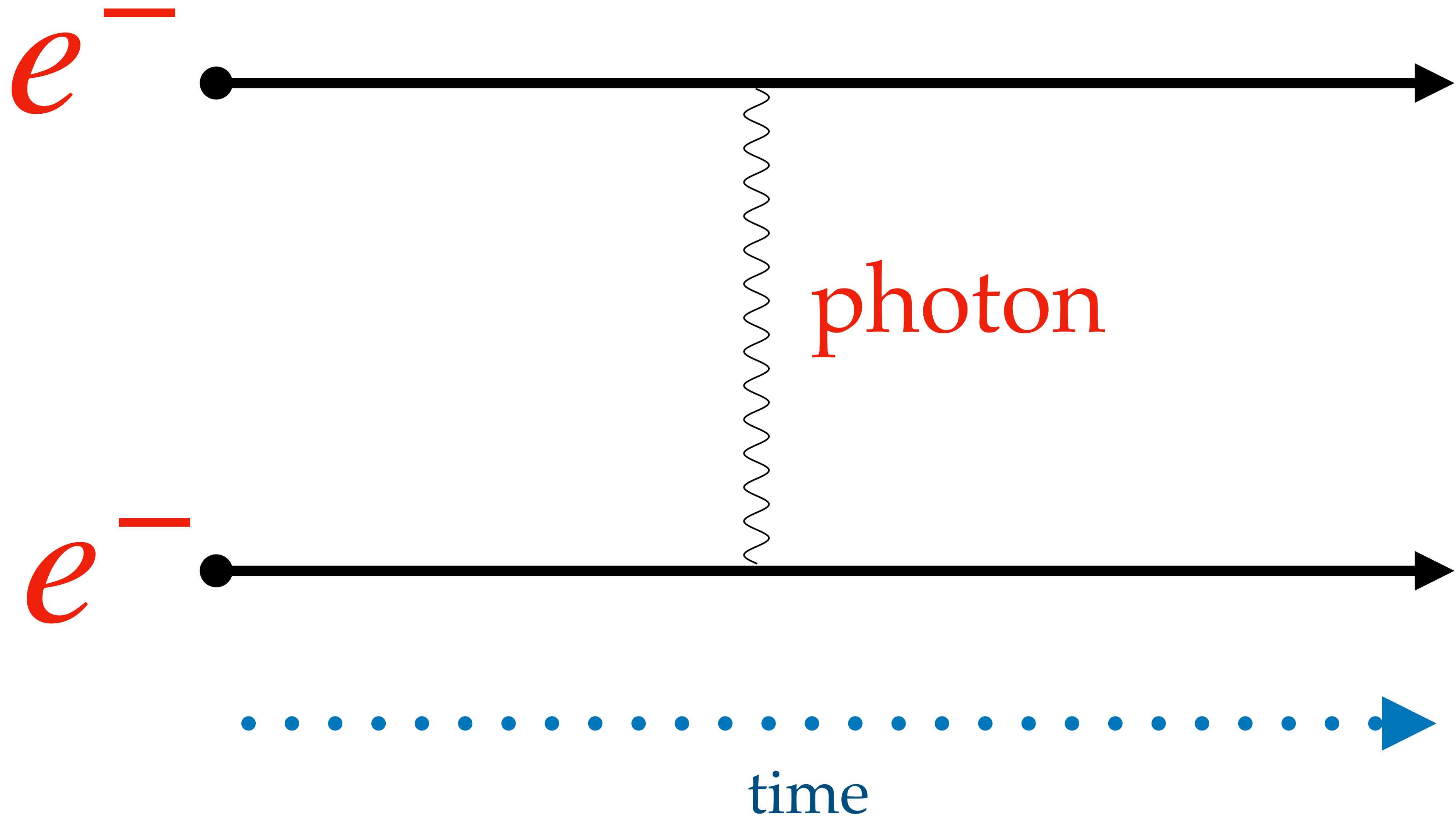
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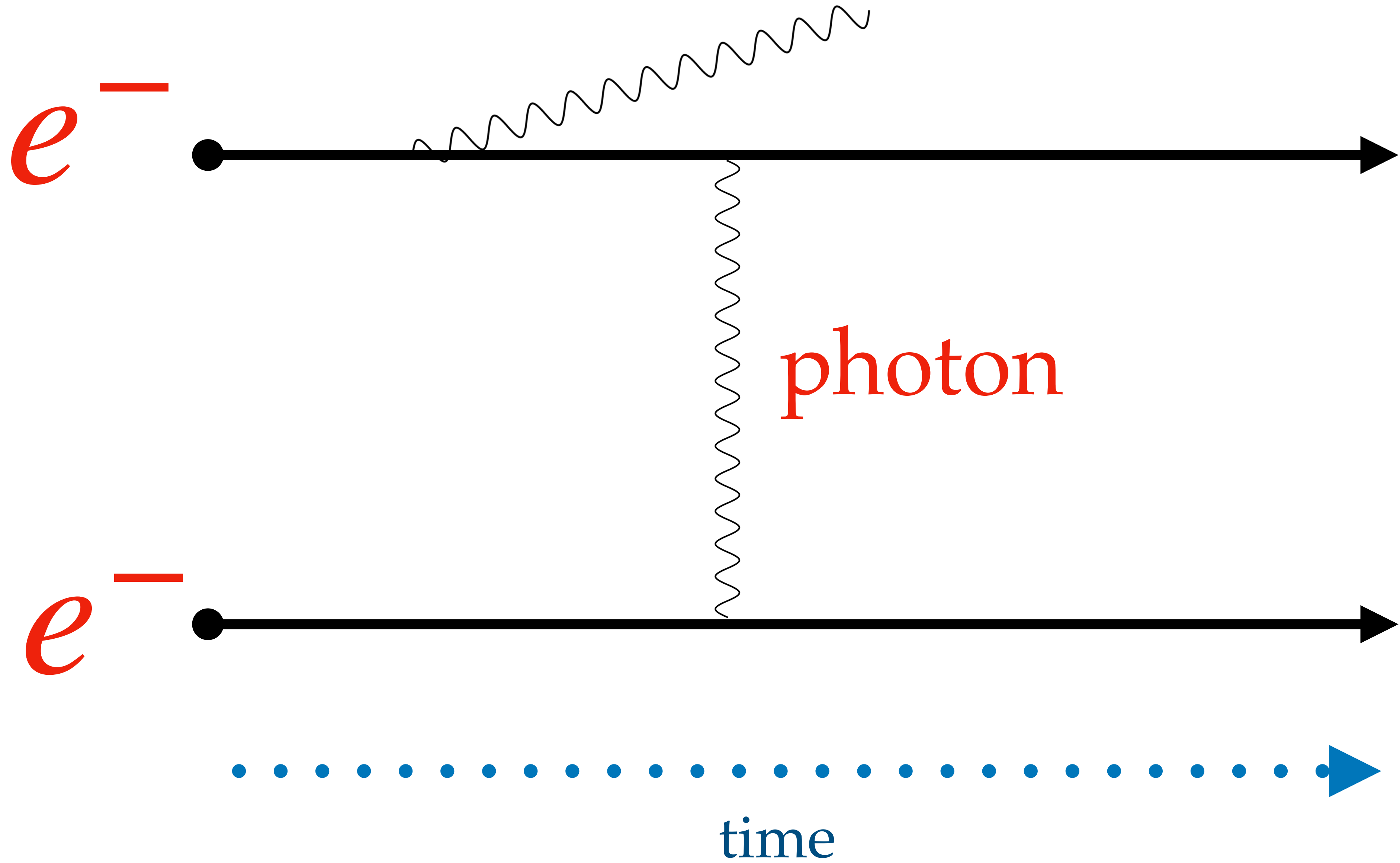
time



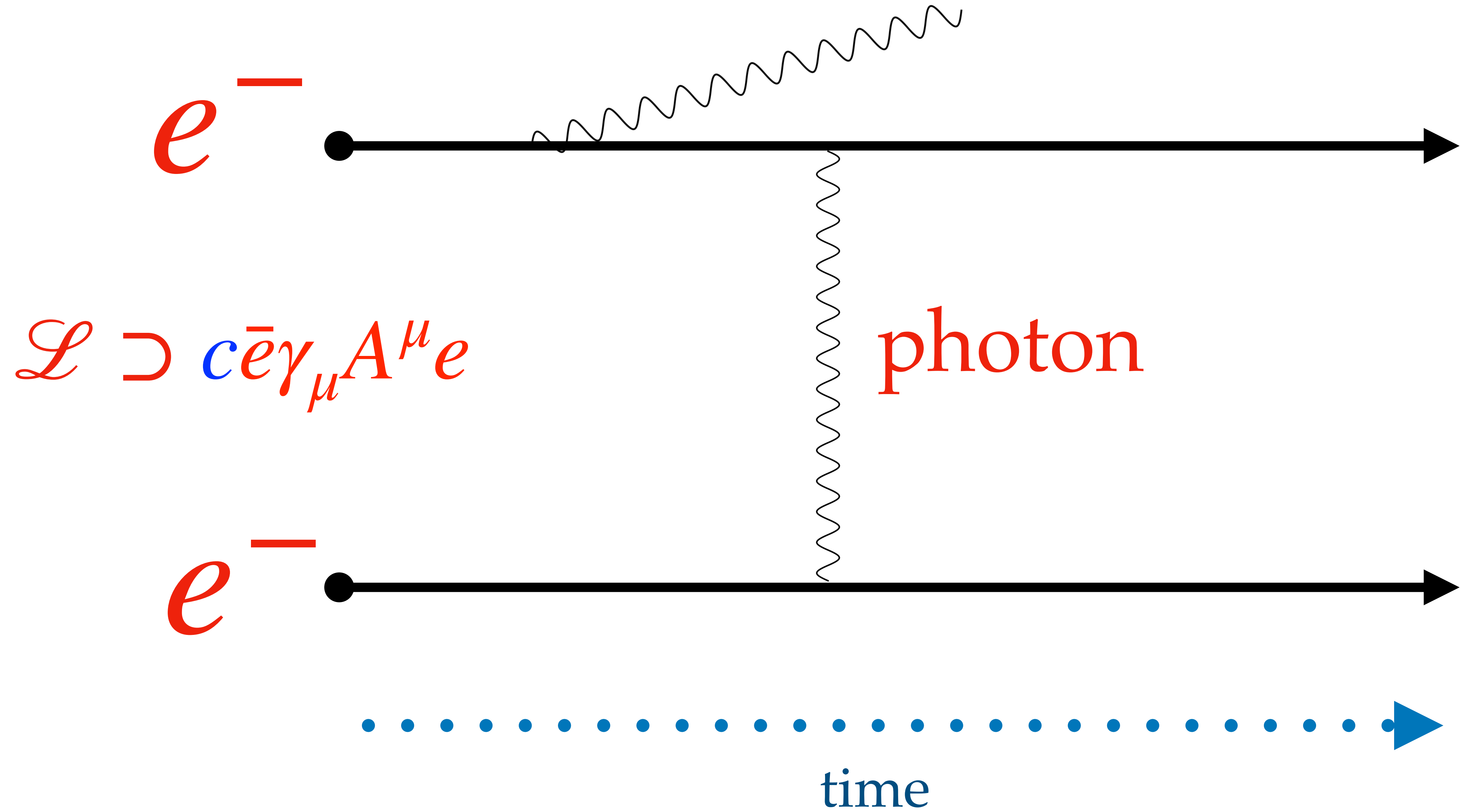
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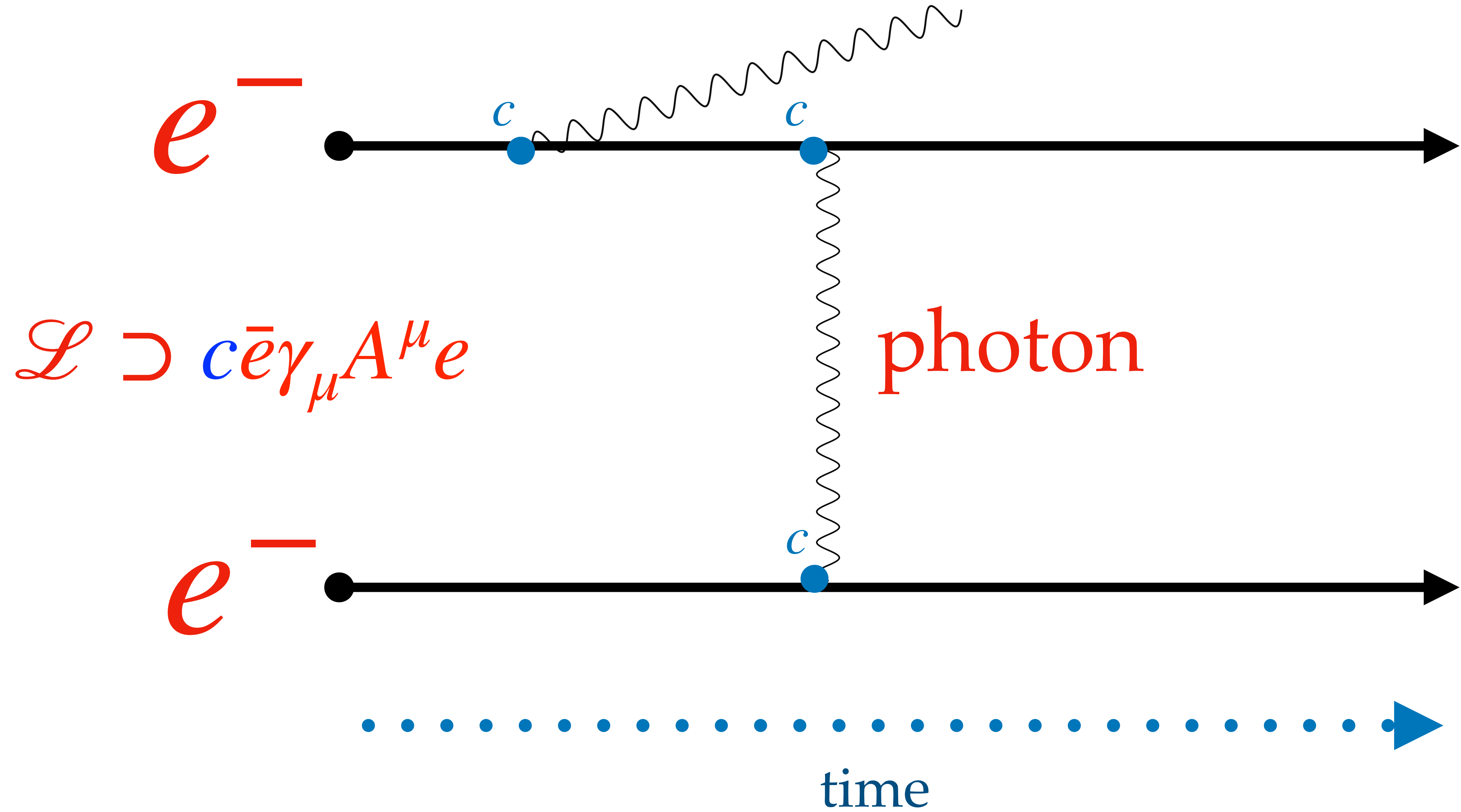
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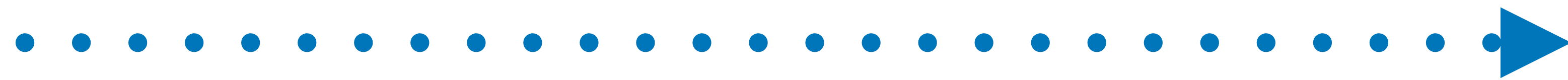
$$\mathcal{L} \supset c \bar{e} \gamma_{\mu} A^{\mu} e$$



From Theory to Experiment and Back Again



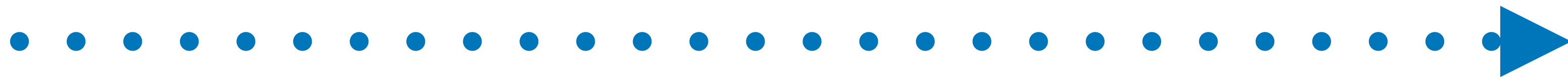
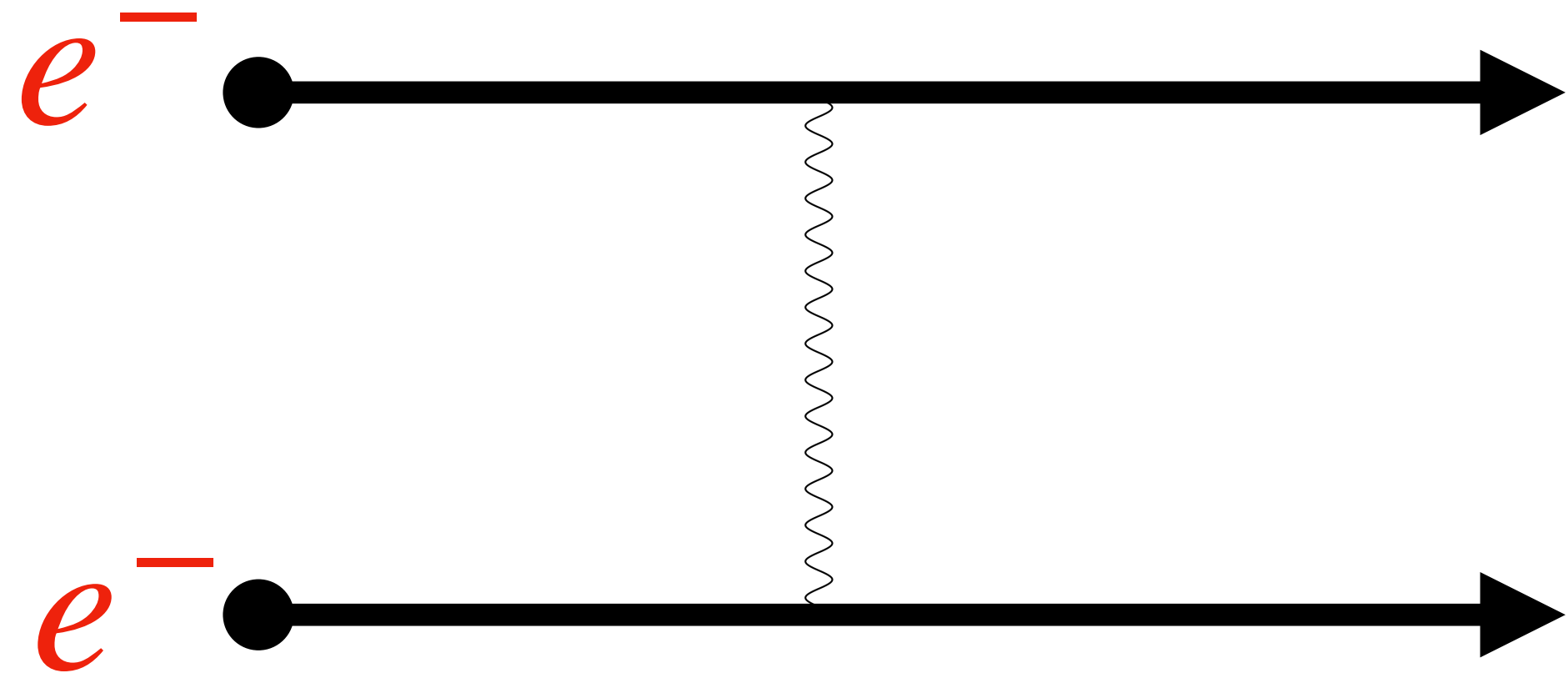
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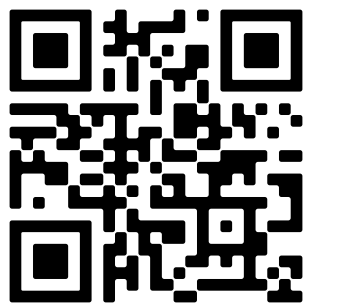
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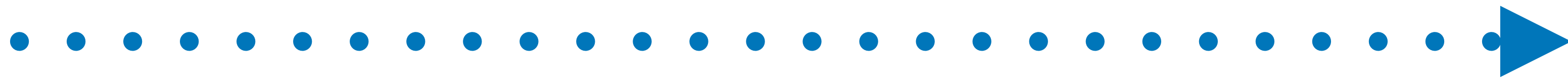
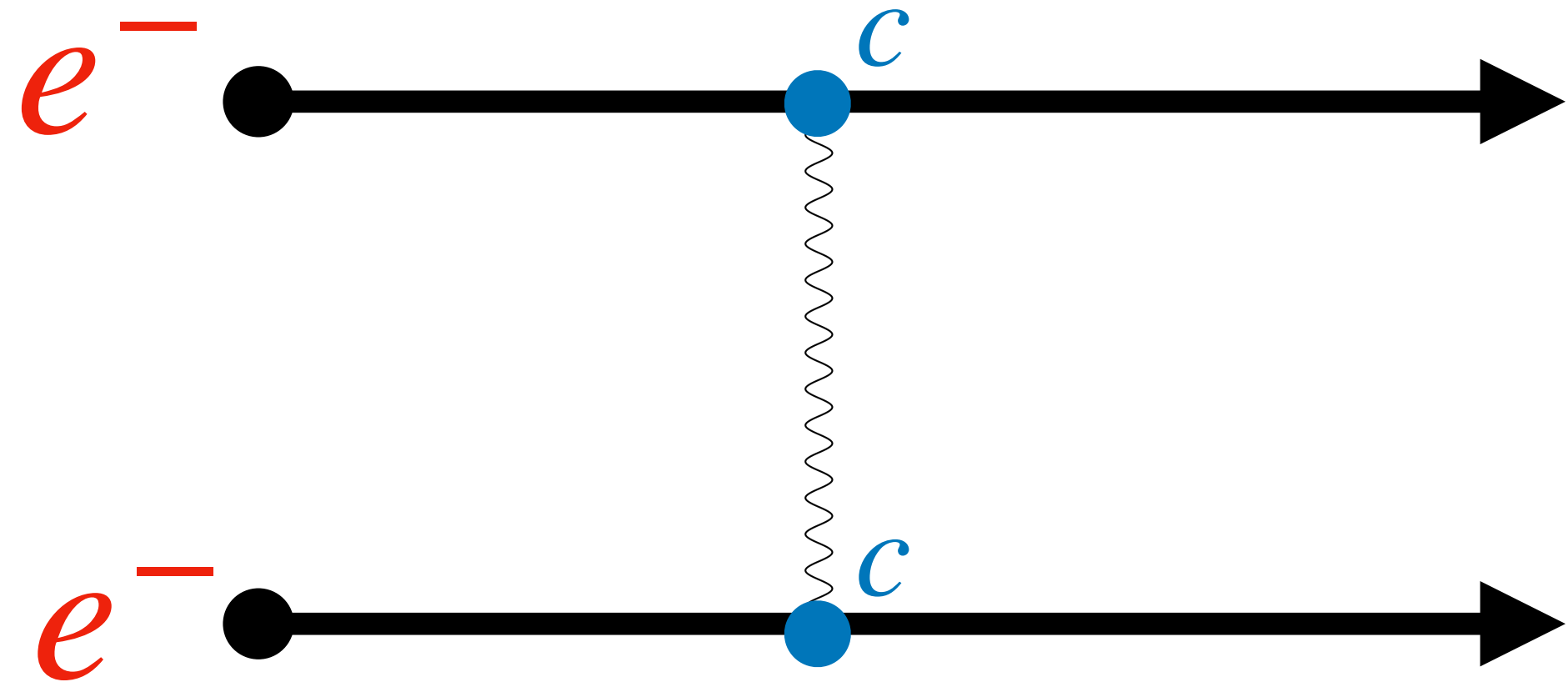
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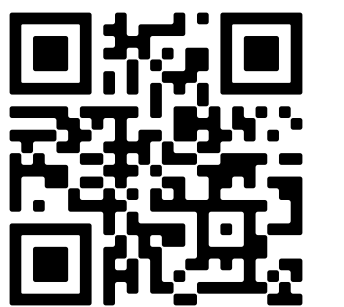
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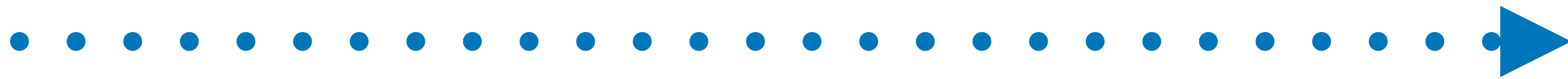
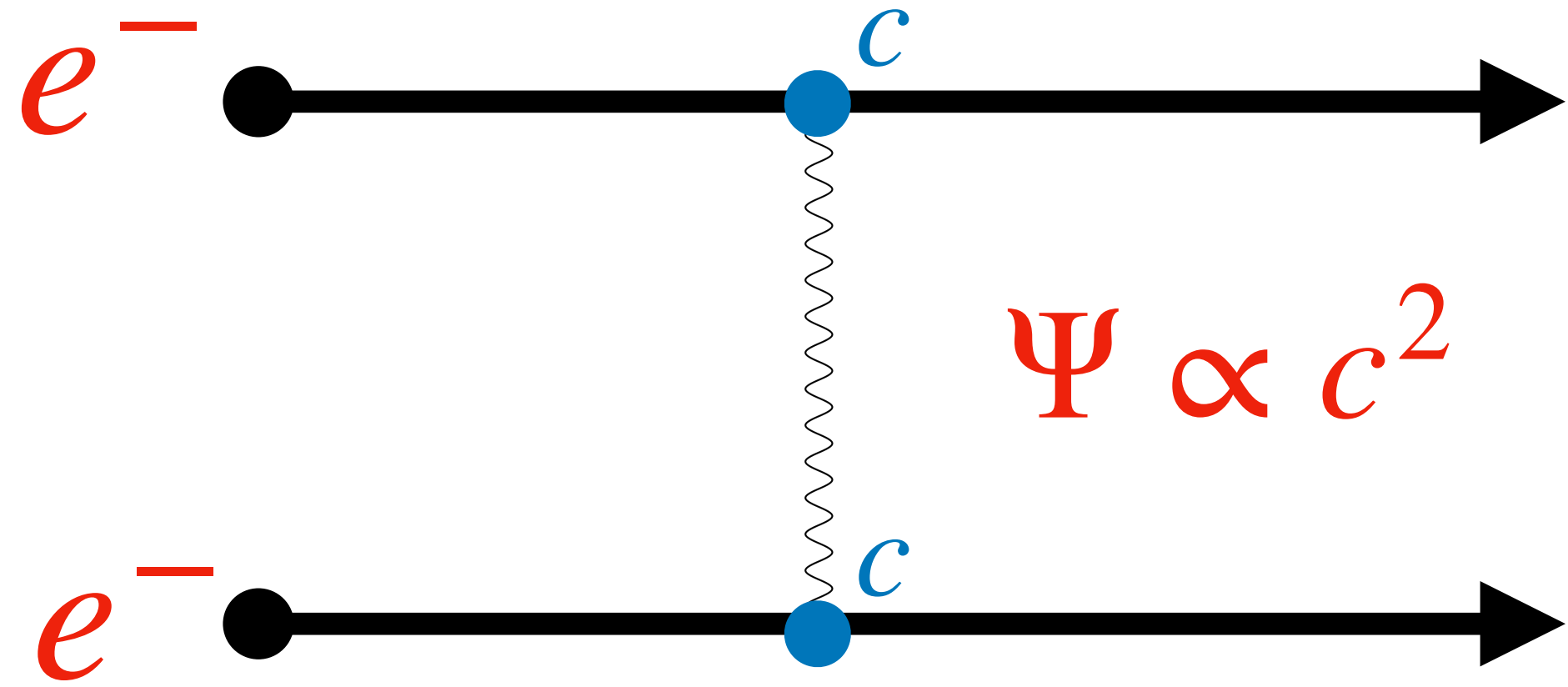
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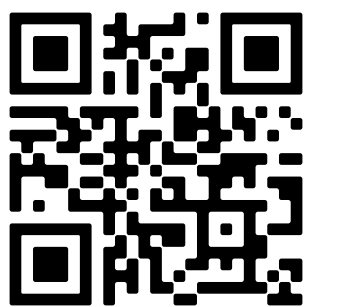
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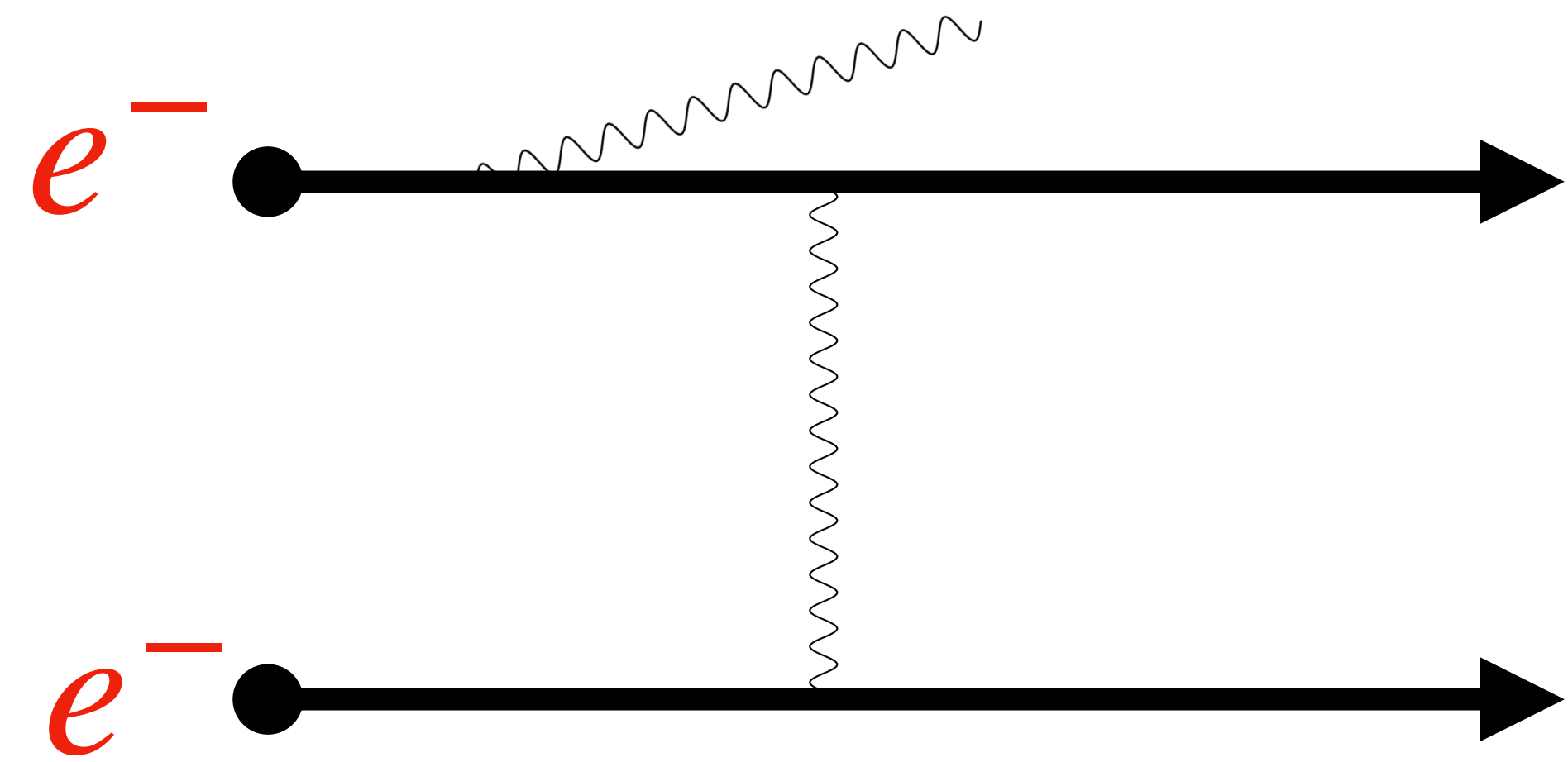
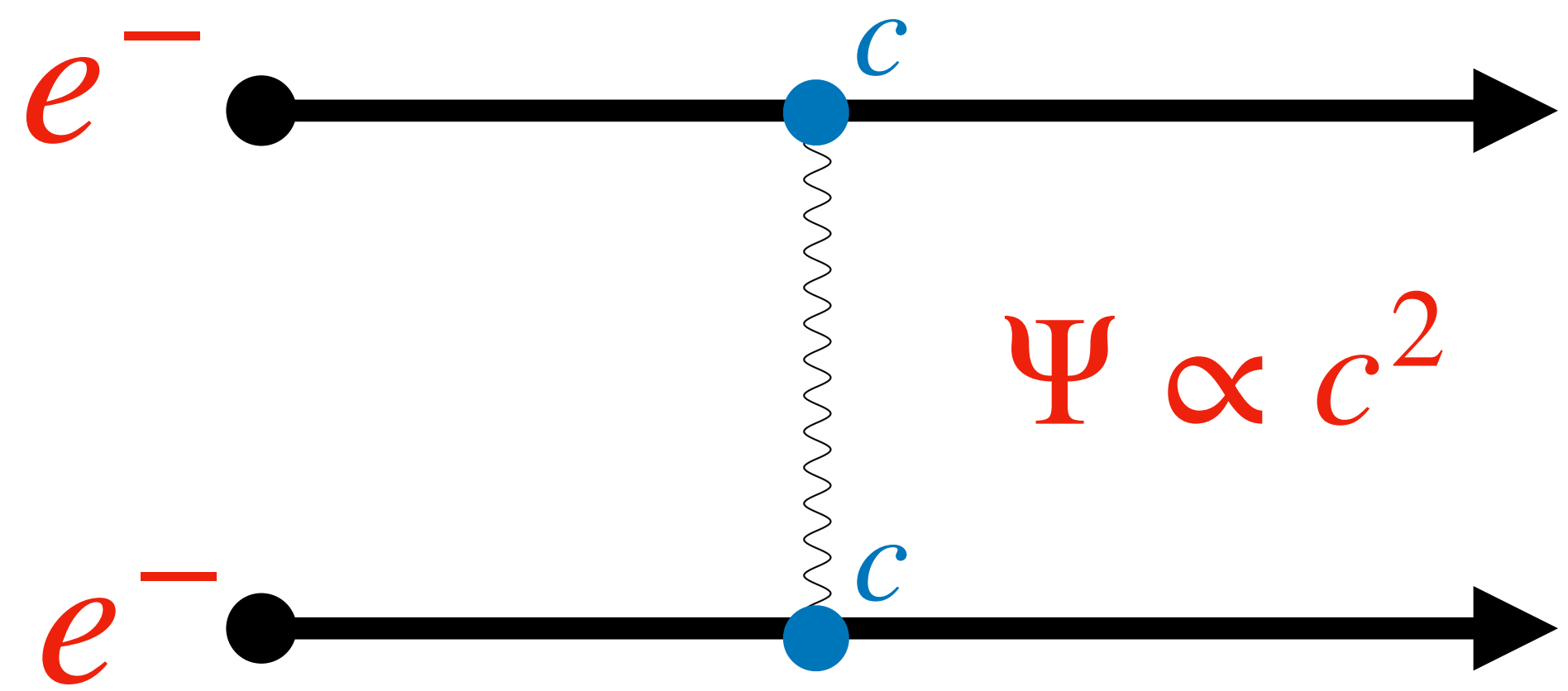
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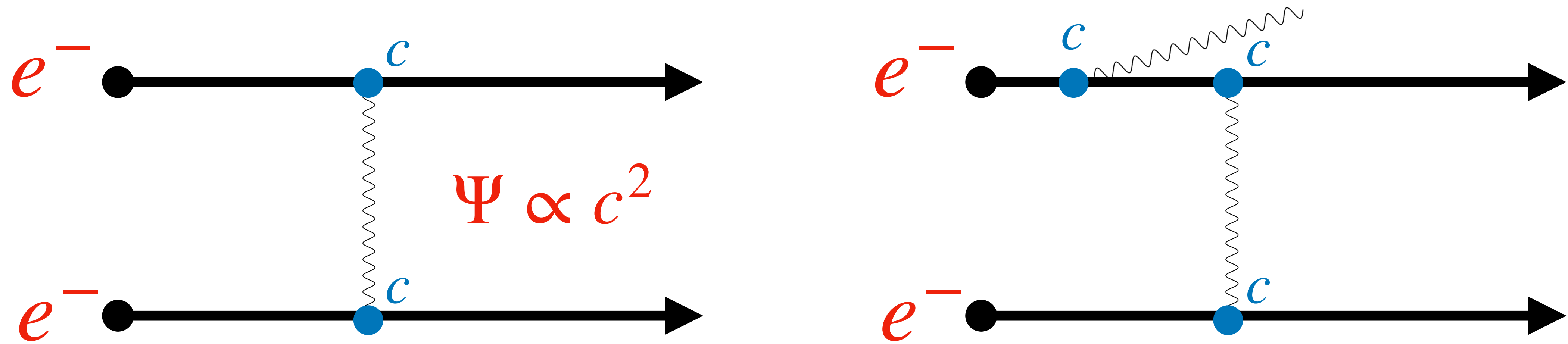
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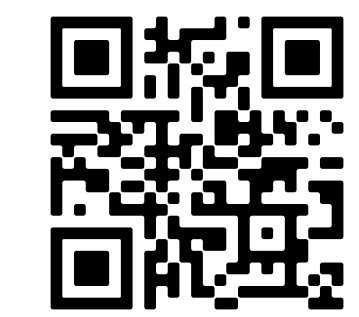
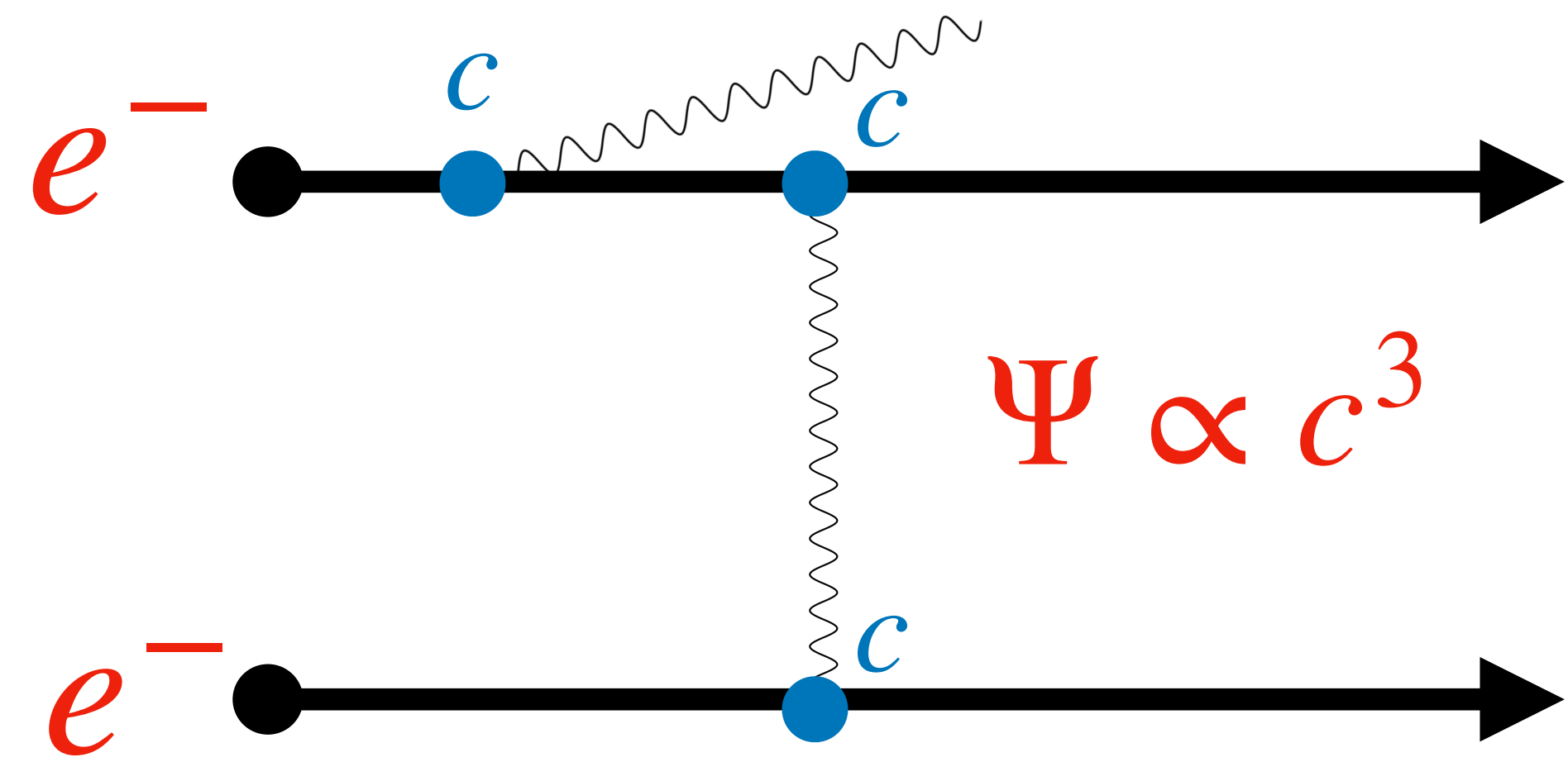
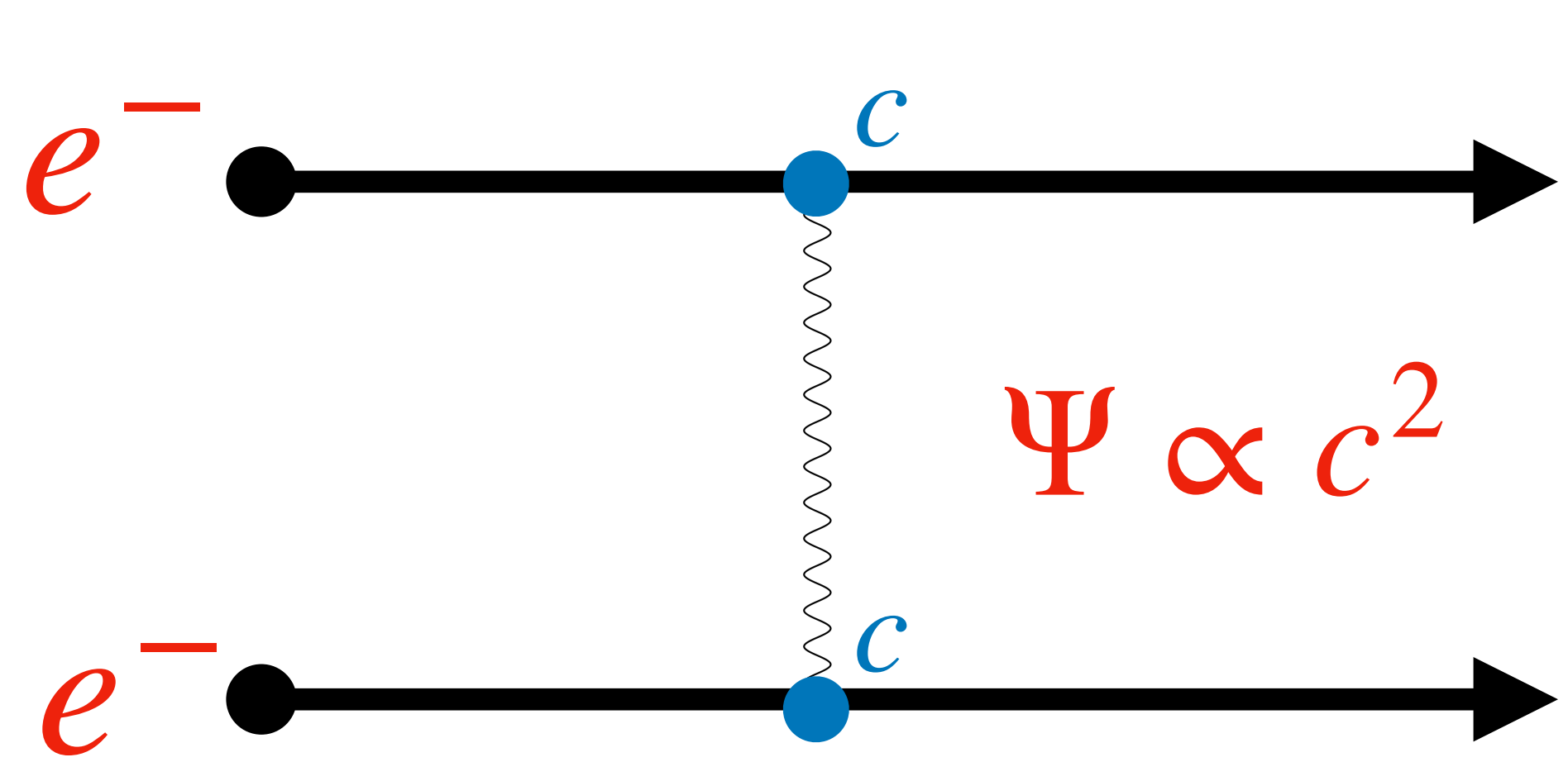
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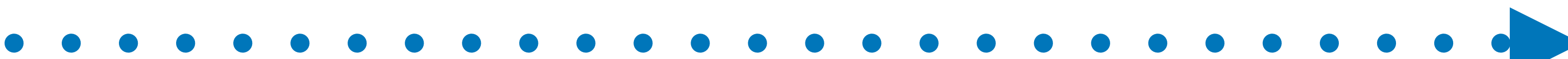
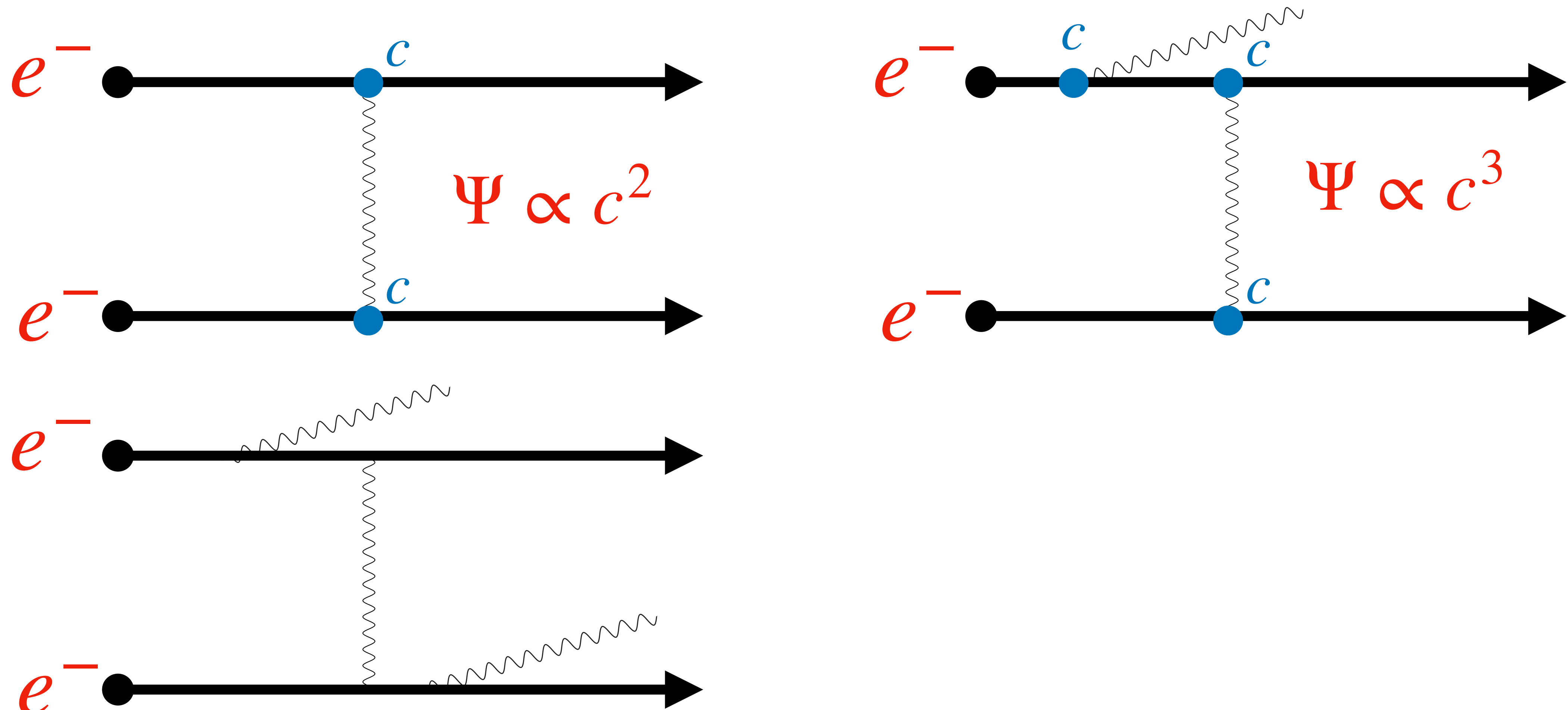
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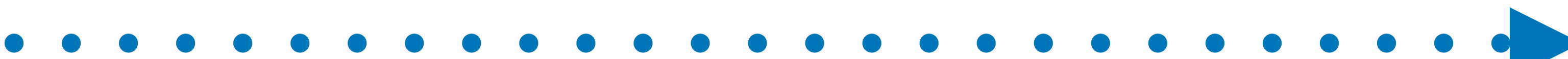
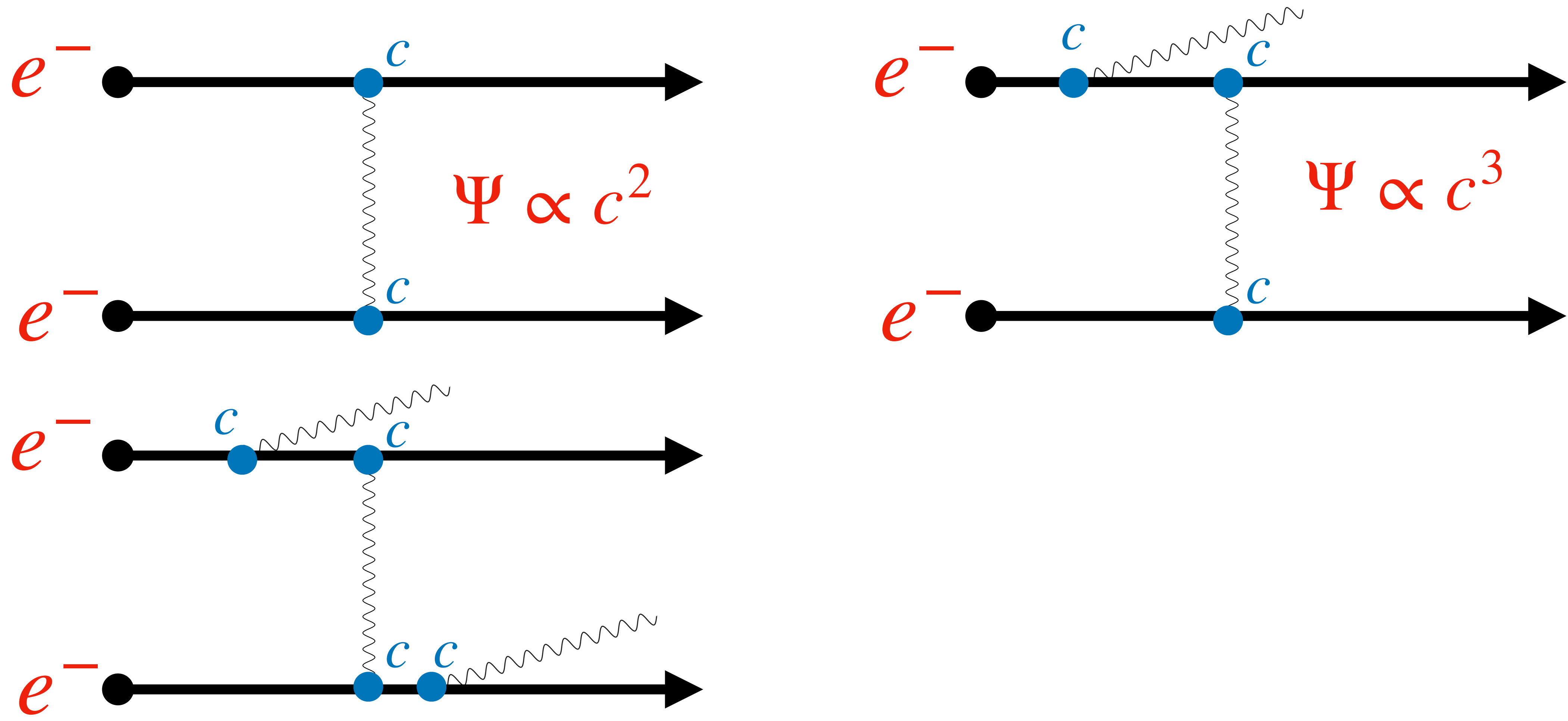
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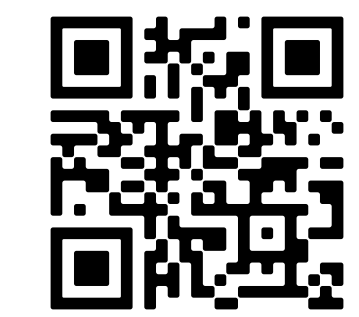
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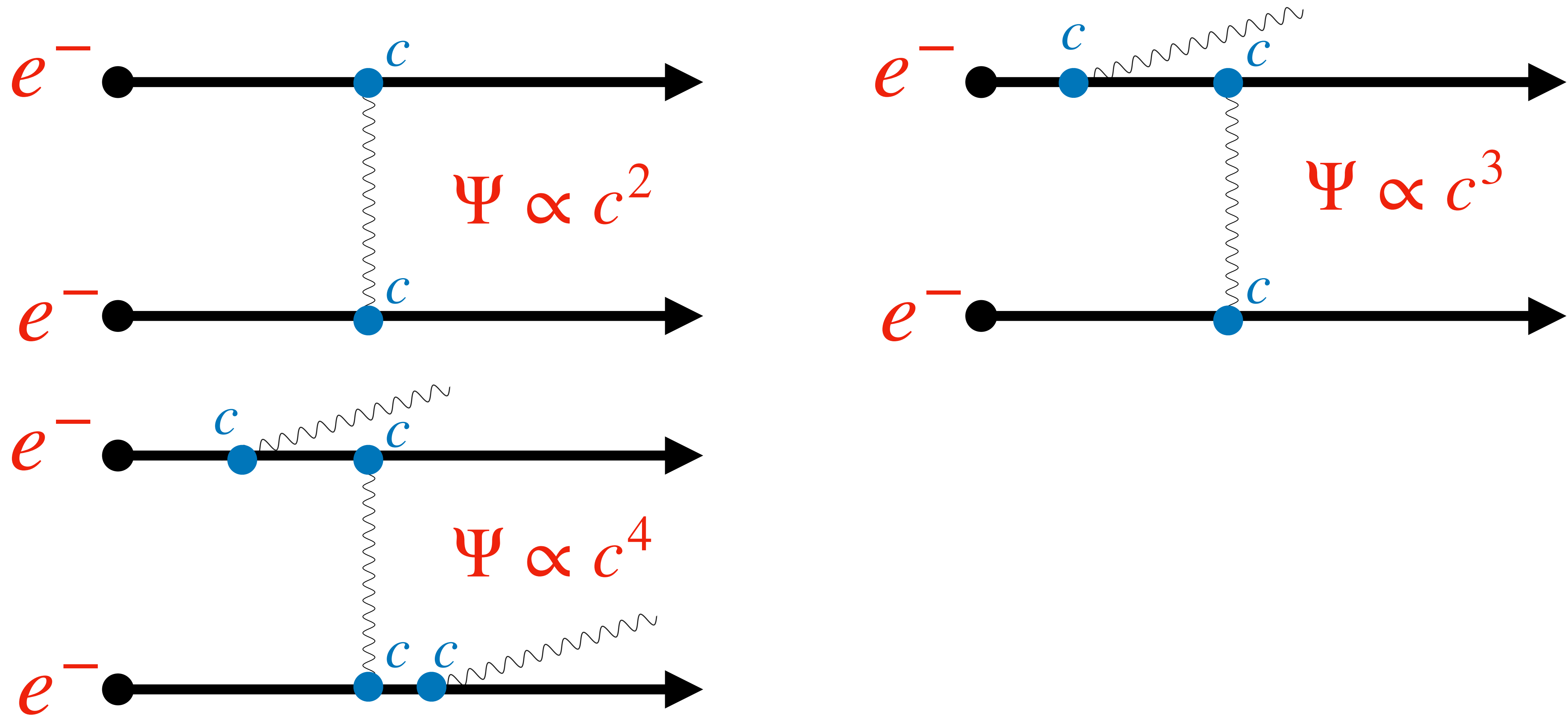
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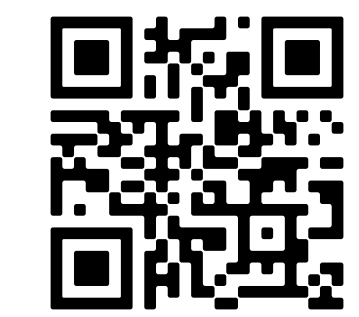
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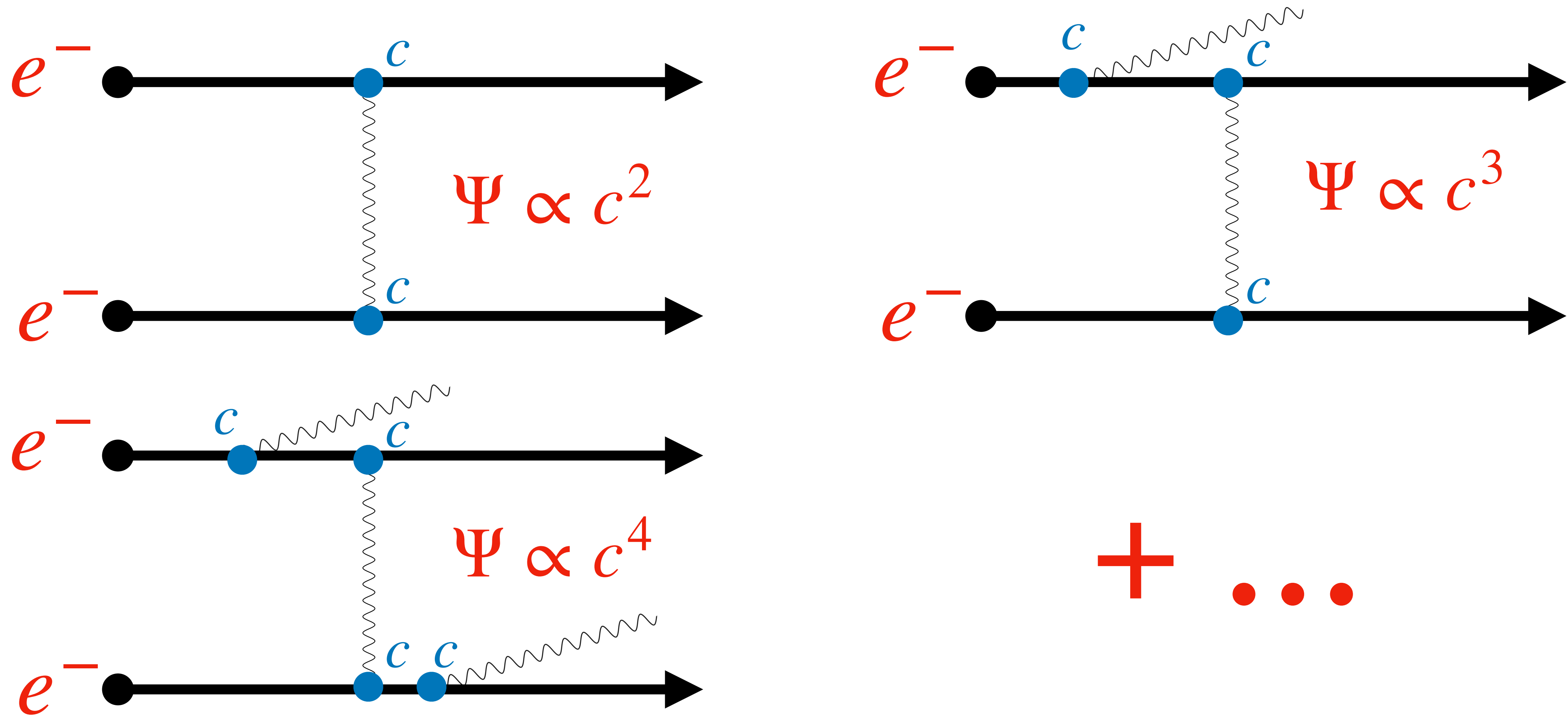
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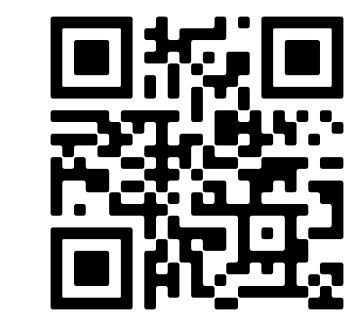
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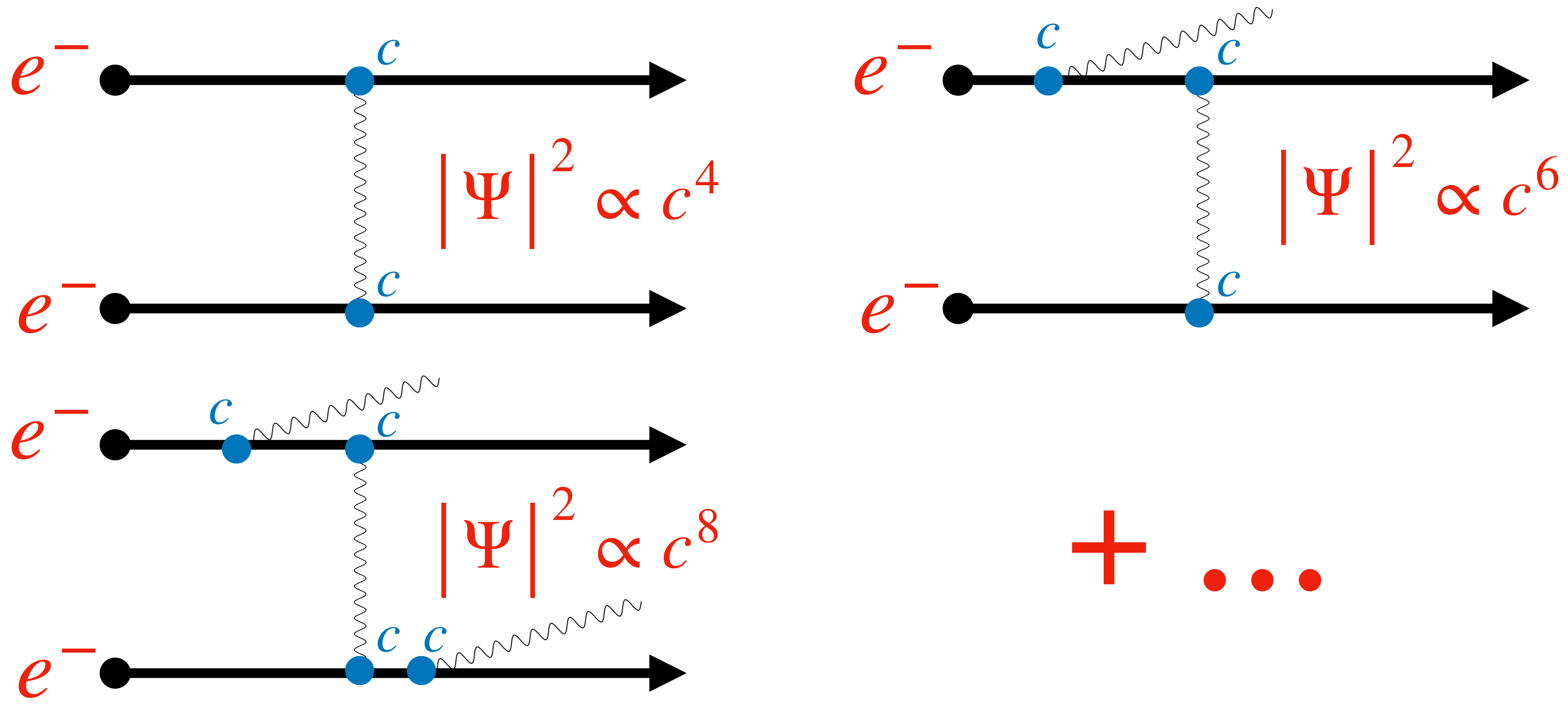
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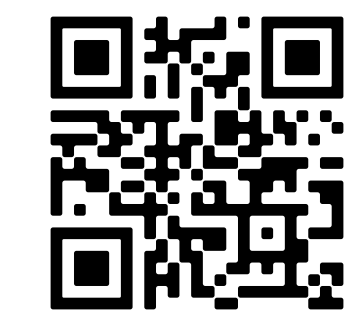
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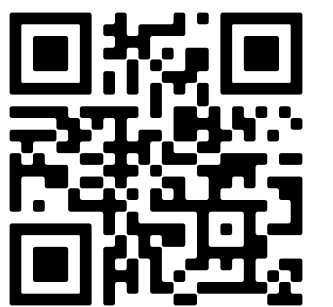


From Theory to Experiment and Back Again

We now know how to calculate **probabilities** in Quantum Field Theory!

$$|\Psi|^2 \propto c^4 + c^6 + c^8 + \dots$$

If $c^2 < 1$ then we can simply add terms for increased precision!



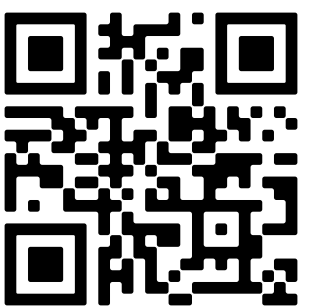
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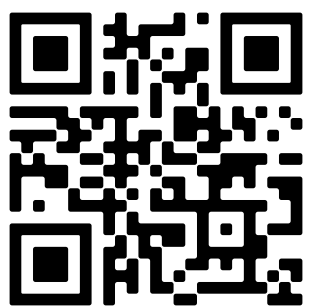
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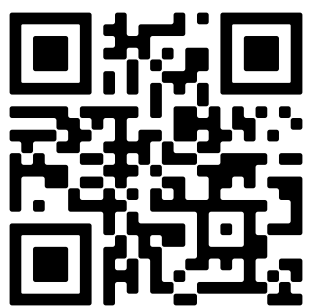
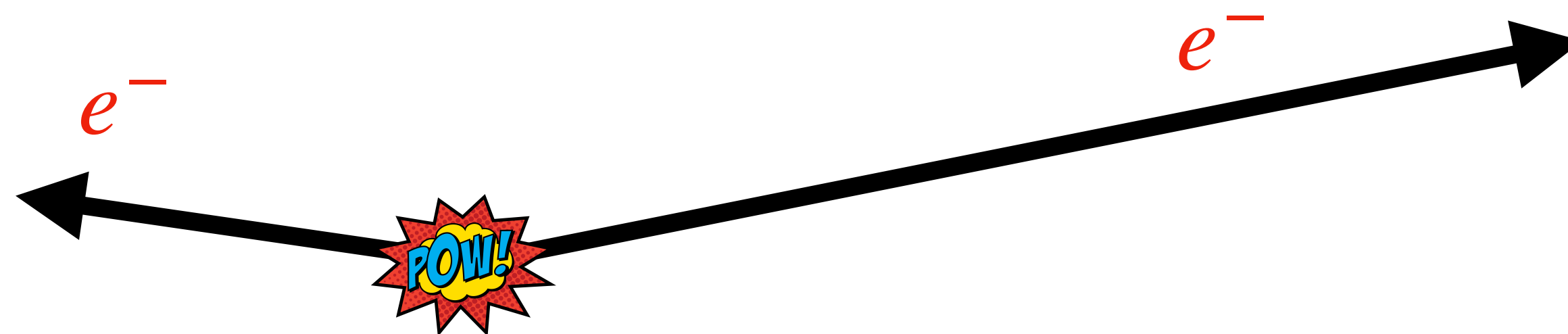
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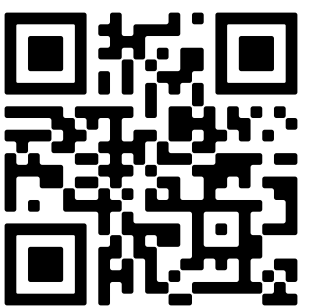
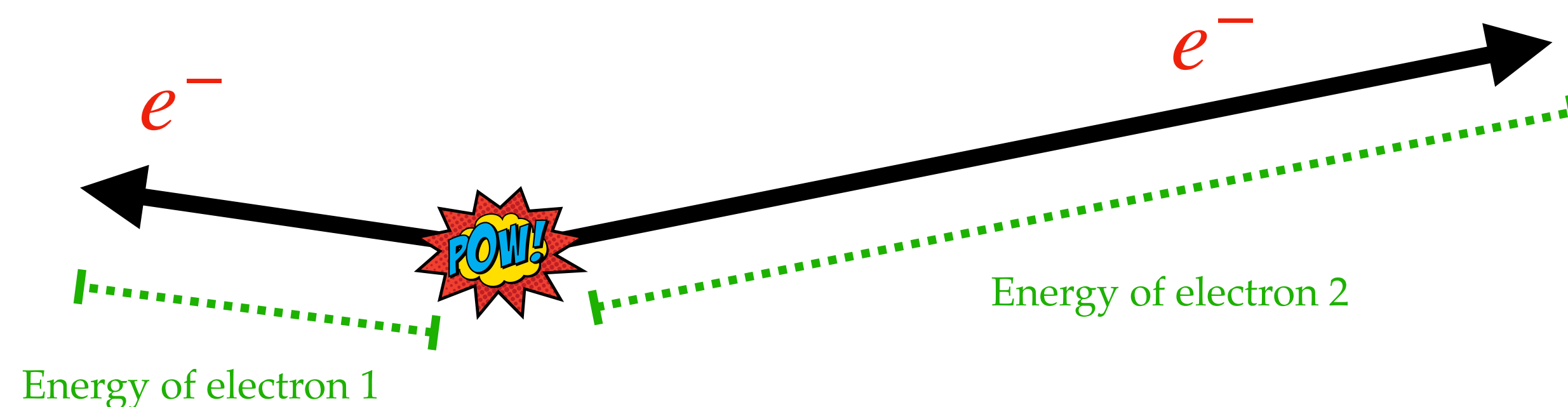
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If $c^2 < 1$ then we can simply add terms for increased precision!

In reality: $|\Psi|^2$ depends on the **energy & momentum** of the particles involved!



F

We have a way to describe what we observe at experiments!

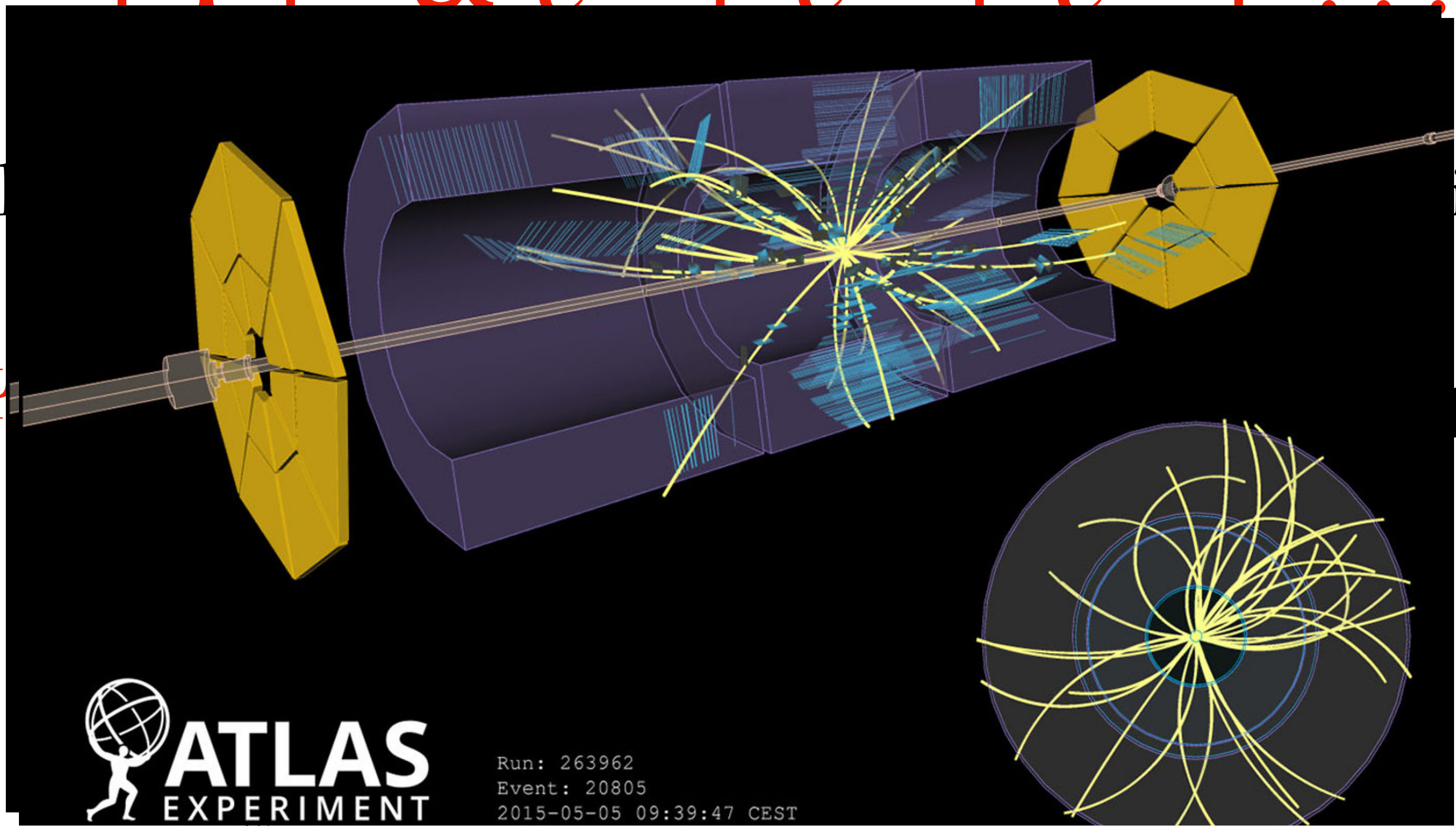
$$|\Psi| \propto c^4 + c^0 + c^0 + \dots$$

If $c^2 < 1$ then

collision!

In reality: $|\Psi|$

particles involved!



Monte Carlo Simulations

- So, can we calculate everything using pen and paper? **Not really...**

Monte Carlo Simulations

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$$|\Psi|^2 \propto c^4 + c^6 + c^8 + \dots \leftarrow$$

Adding more & more terms is extremely challenging!

Monte Carlo Simulations

- So, can we calculate everything using pen and paper? **Not really...**

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- **Solution:** Use approximate c^n for n large & model situations where $c^2 \geq 1$!
- **Monte Carlo simulations accomplish this!**



Based on **randomness** and **probability**.
Just like Quantum Mechanics!

Pictured: Casino Monte Carlo, Monaco.



How do we make sense of it all?



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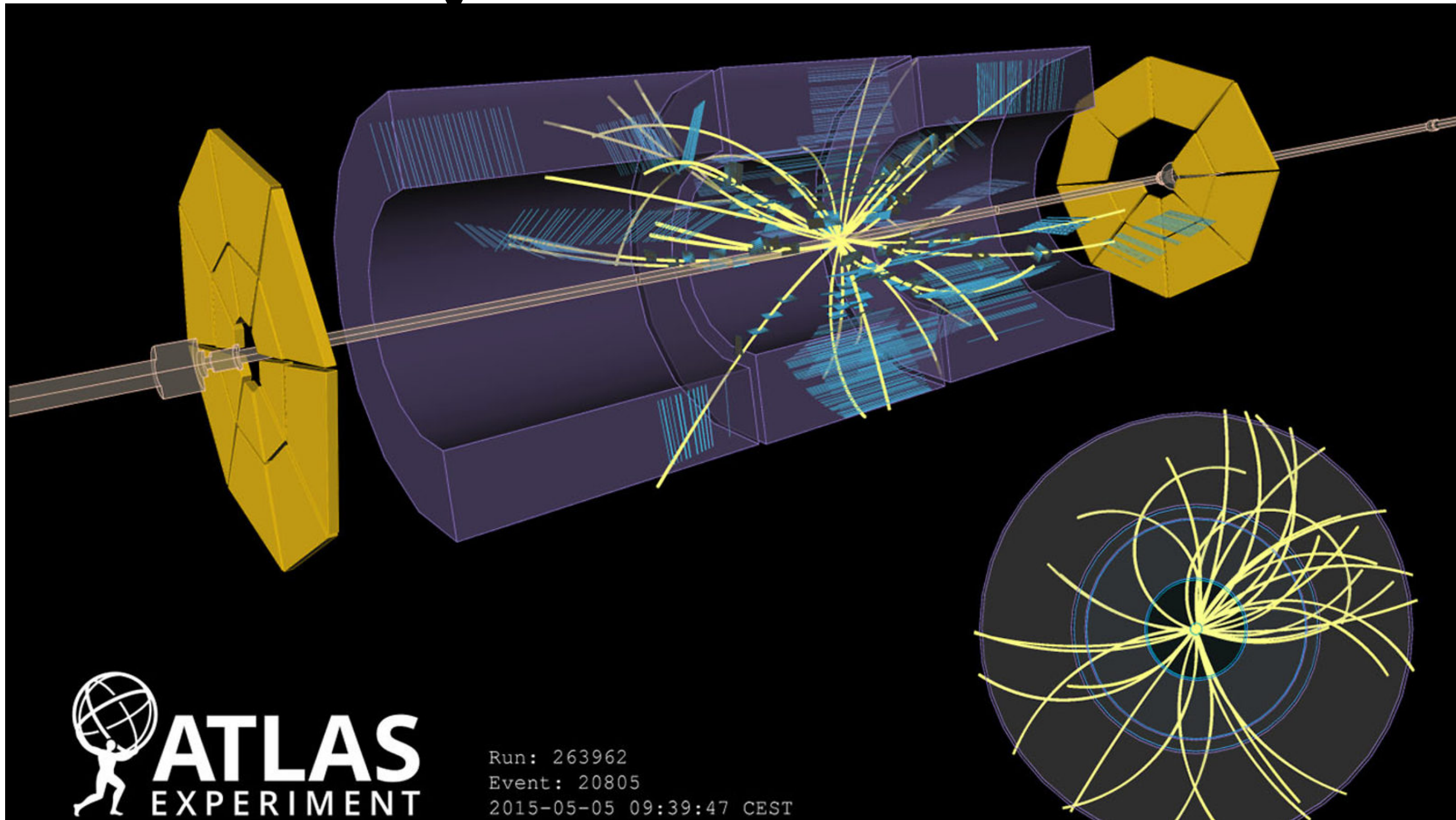
$$\begin{aligned}\mathcal{L} = & -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + i\bar{\Psi}\not{D}\Psi \\ & + \Psi_i y_{ij} \Psi_j + \text{h.c.} \\ & + |D_\mu\Phi|^2 - V(\Phi)\end{aligned}$$



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Experiment

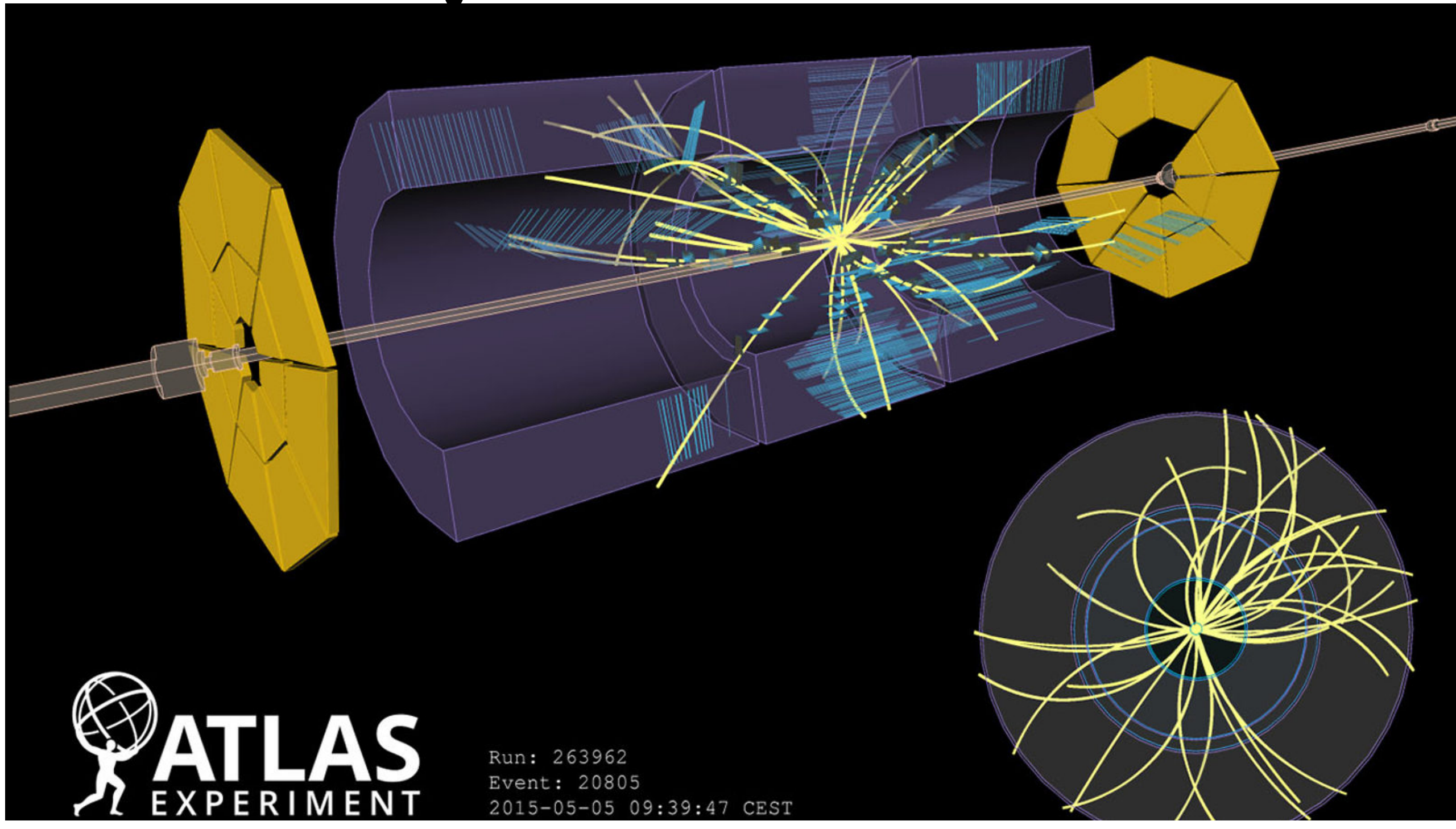


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Simulations

Experiment



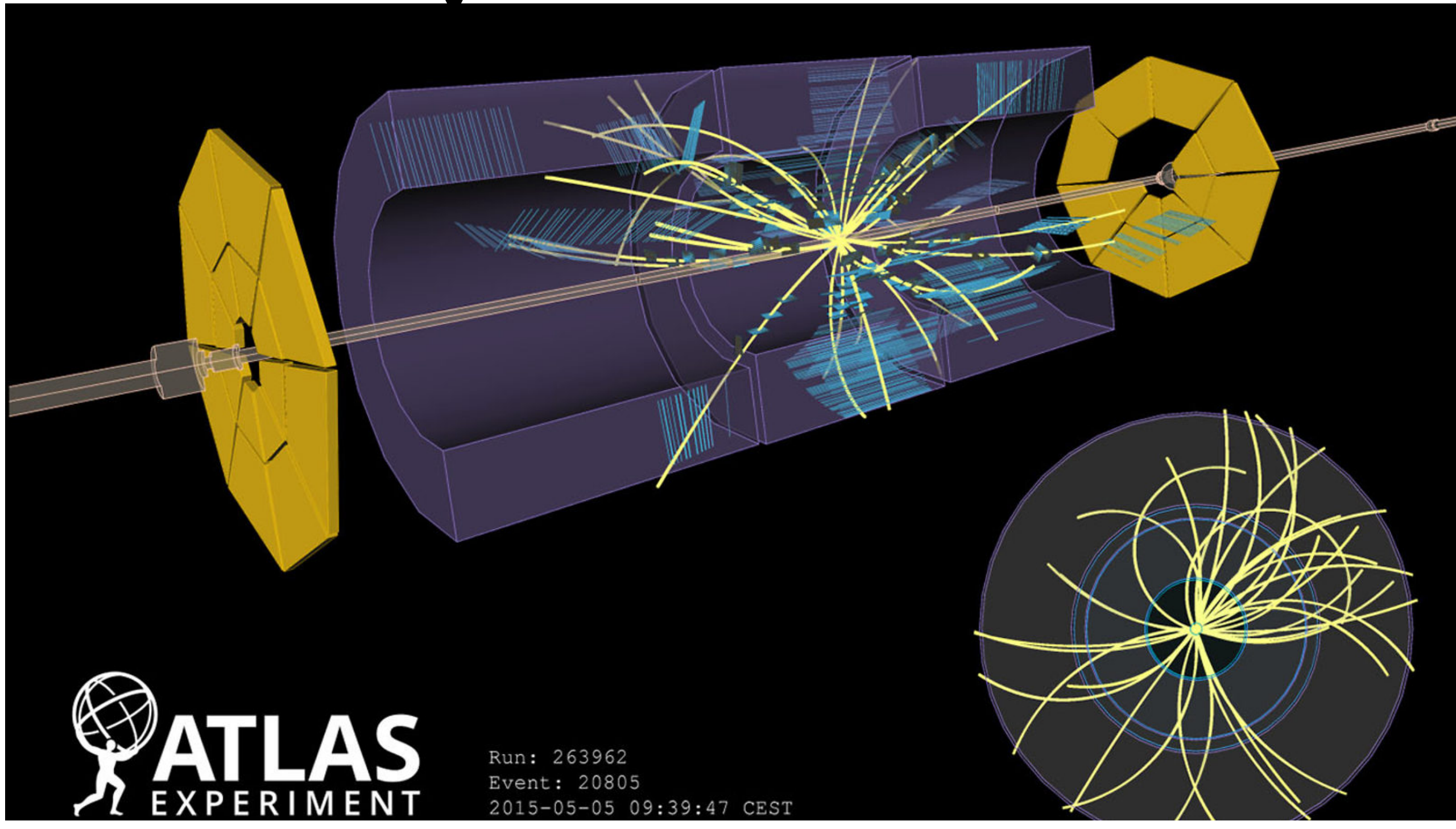
How do we make sense of it all?

Shut up and Simulate!

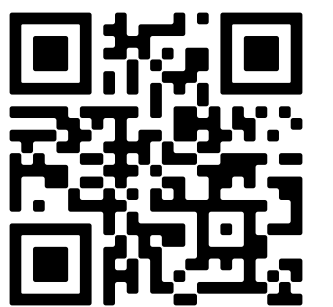
Simulations



Experiment

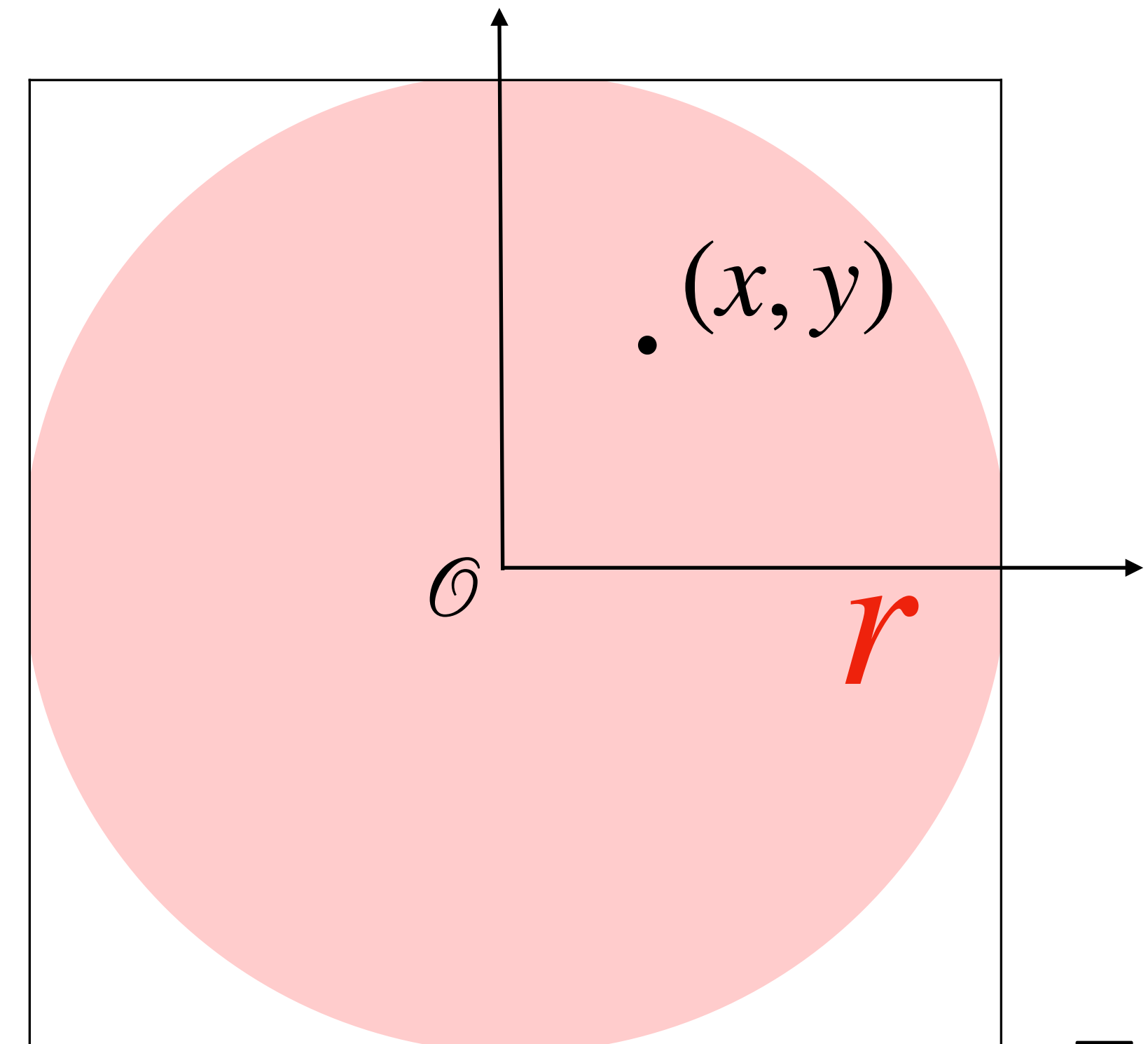


Simulations via the Monte Carlo method: Calculate π !



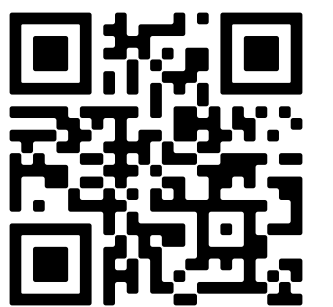
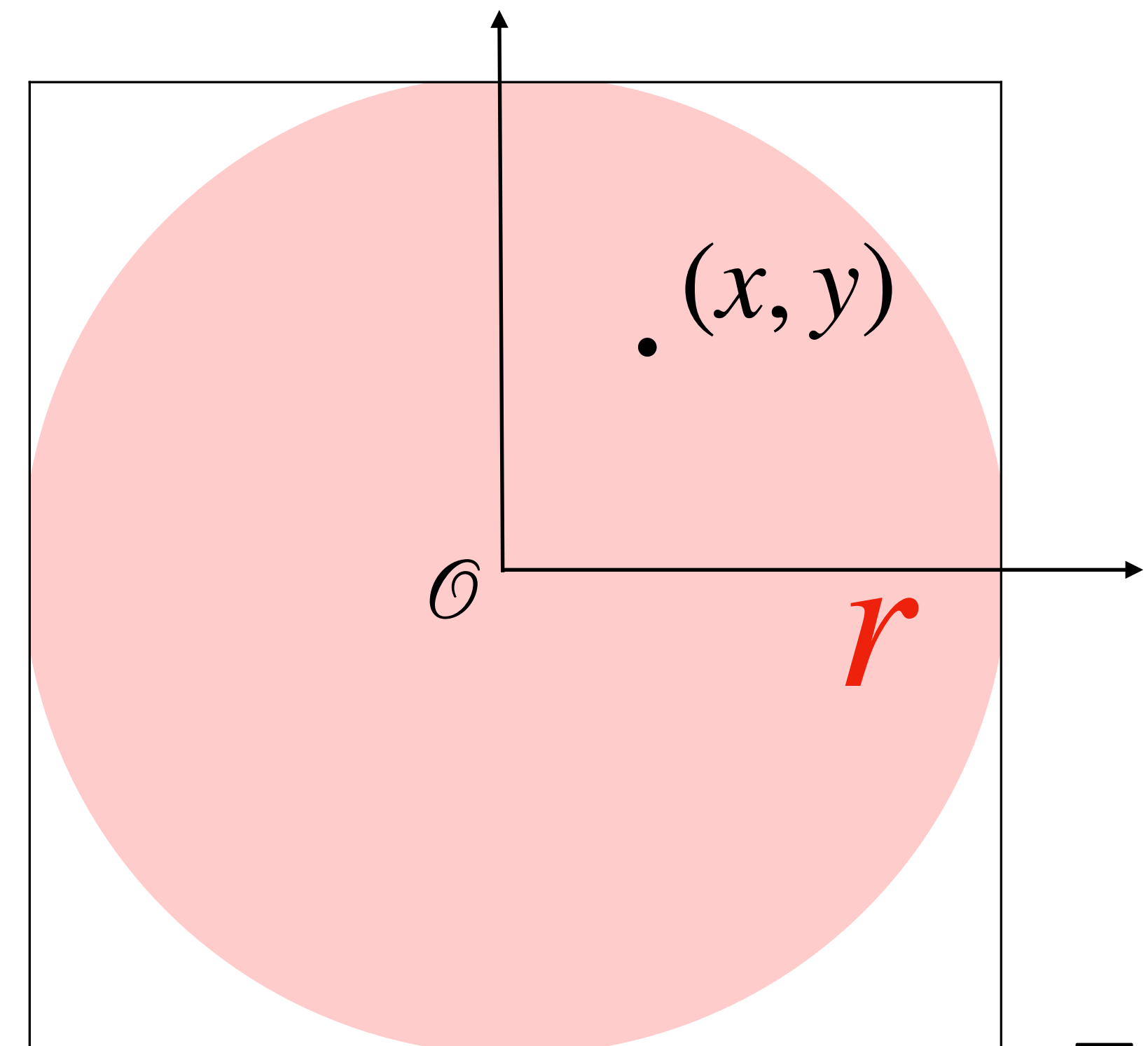
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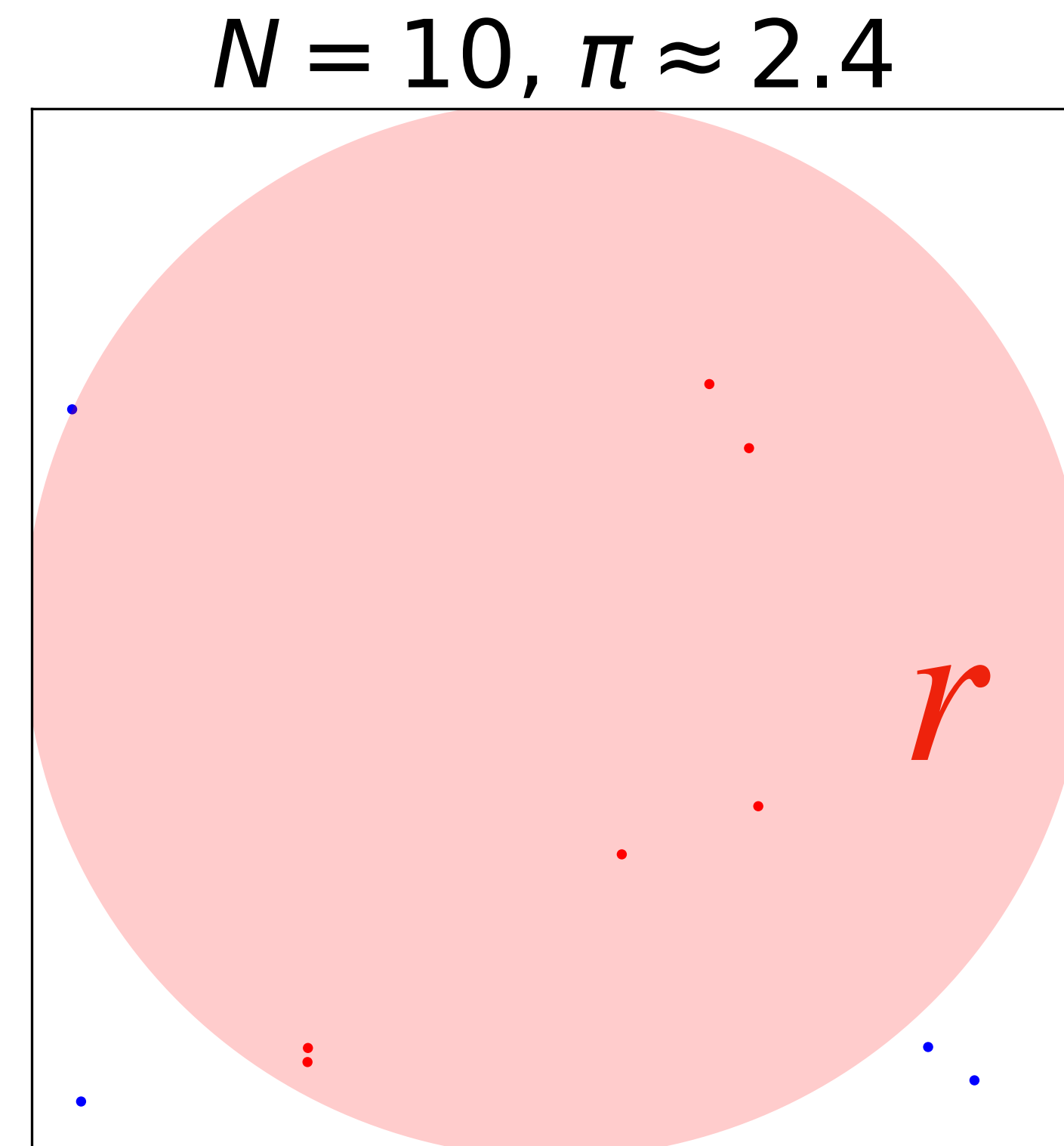
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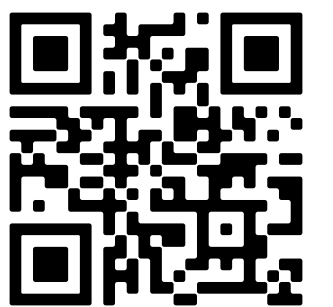
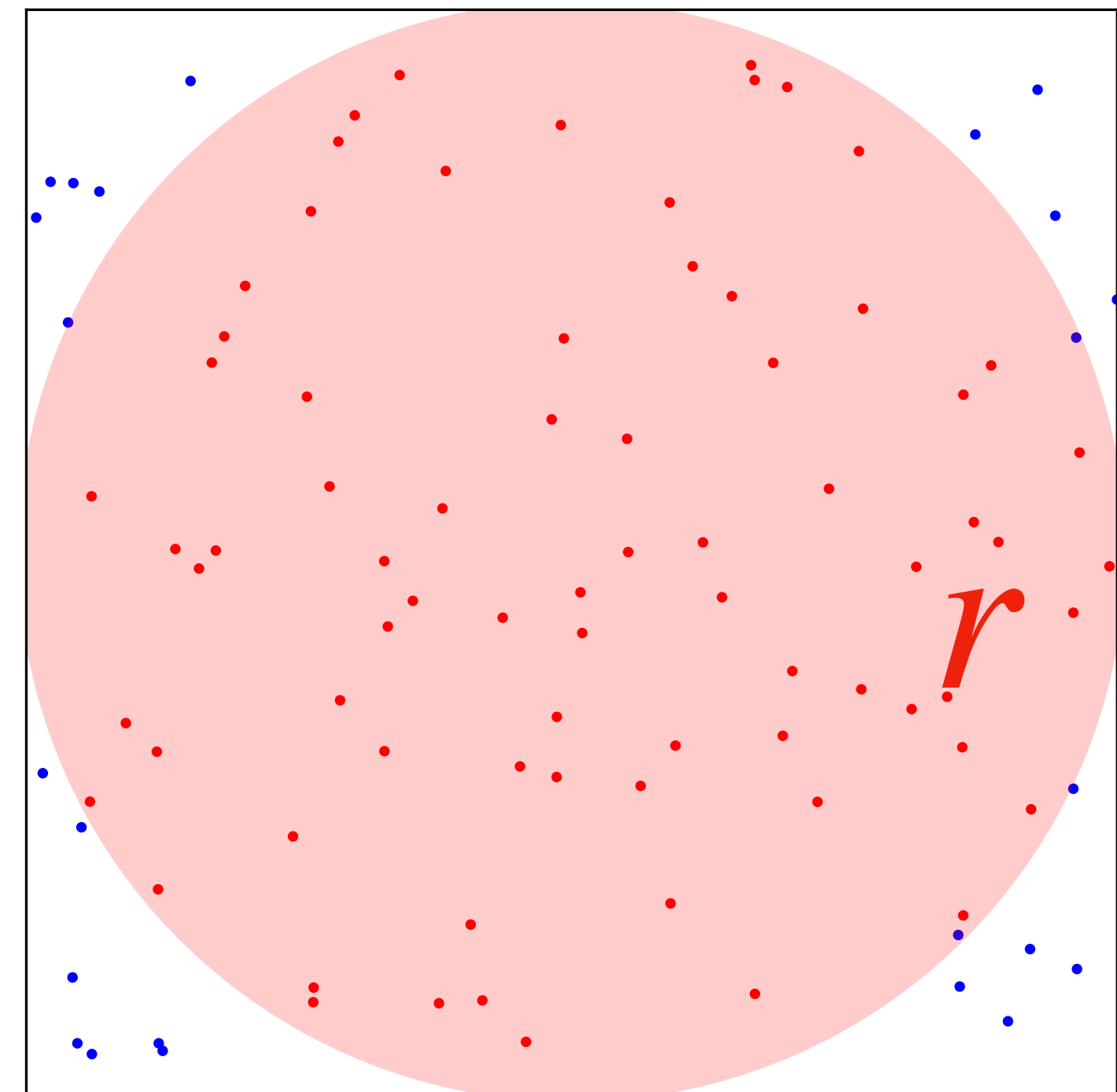
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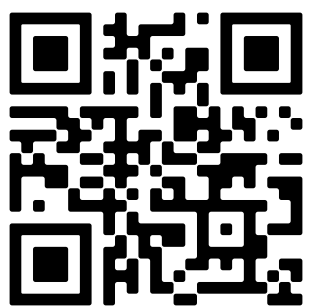
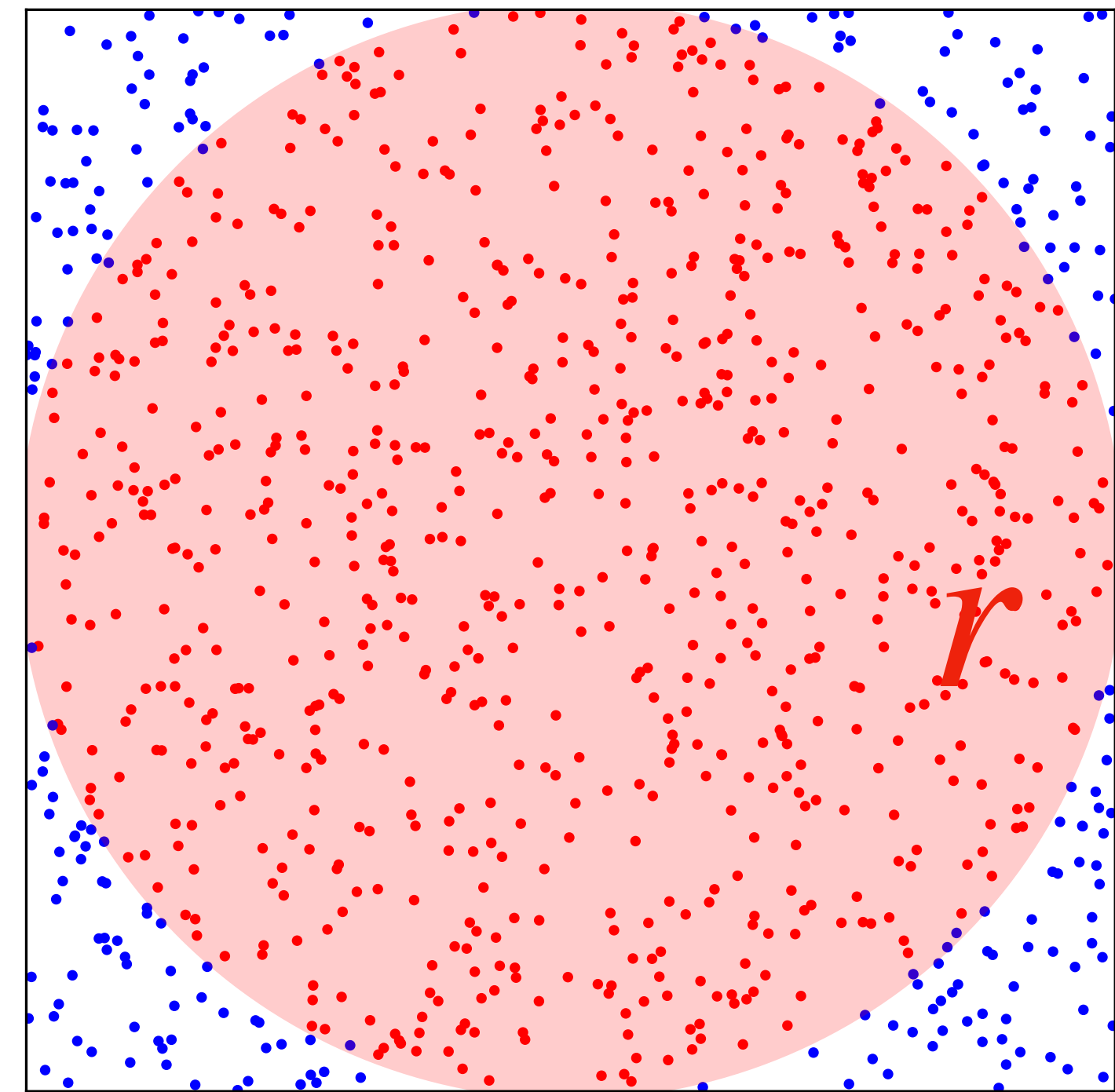
$N = 10^2, \pi \approx 3.04 \pm 0.35$



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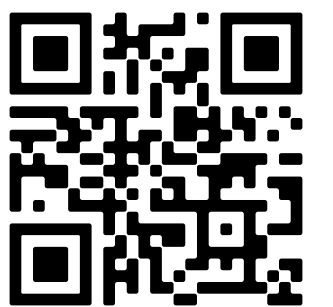
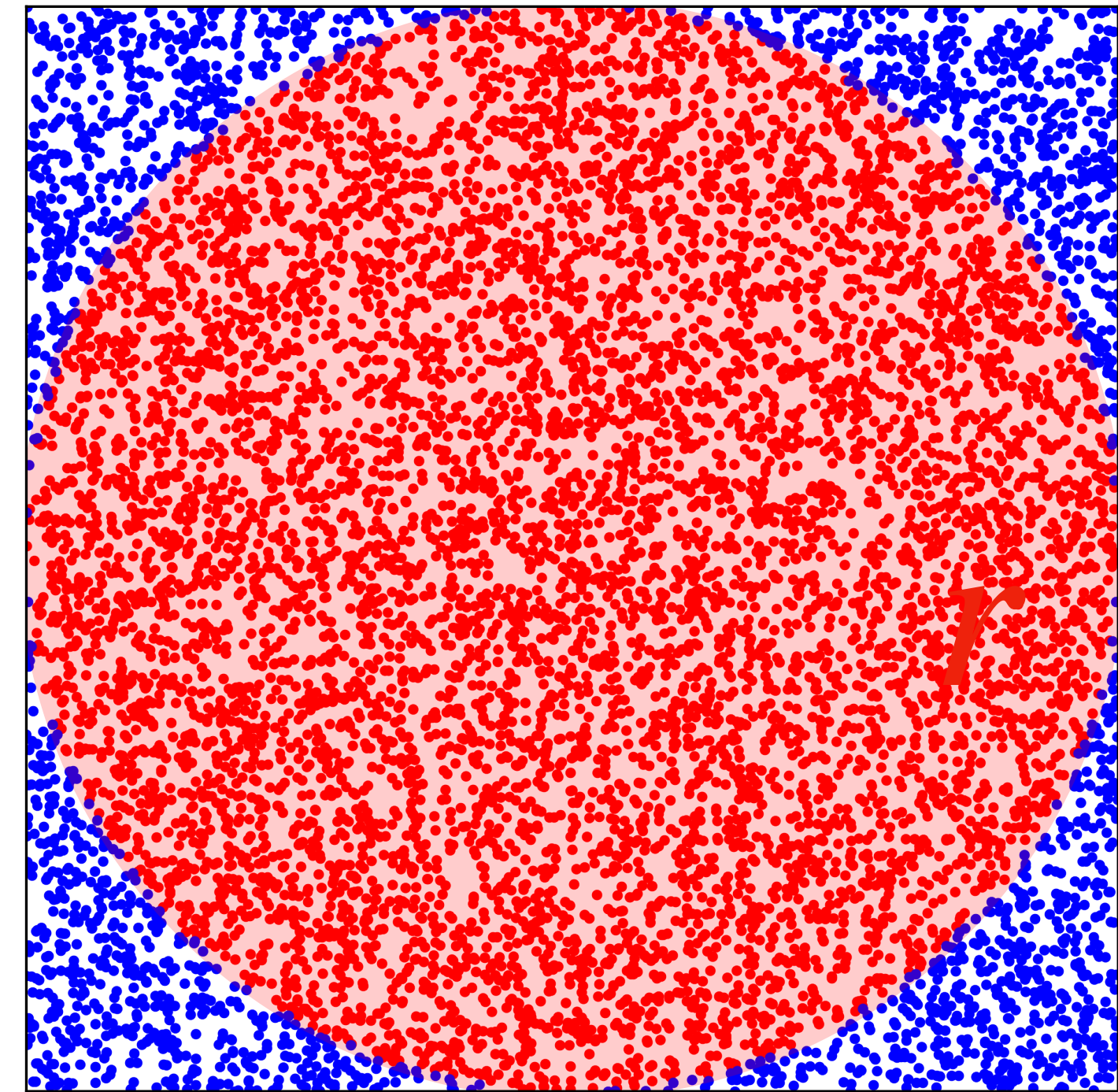
$$N = 10^3, \pi \approx 3.08 \pm 0.11$$



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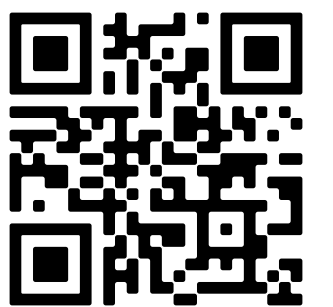
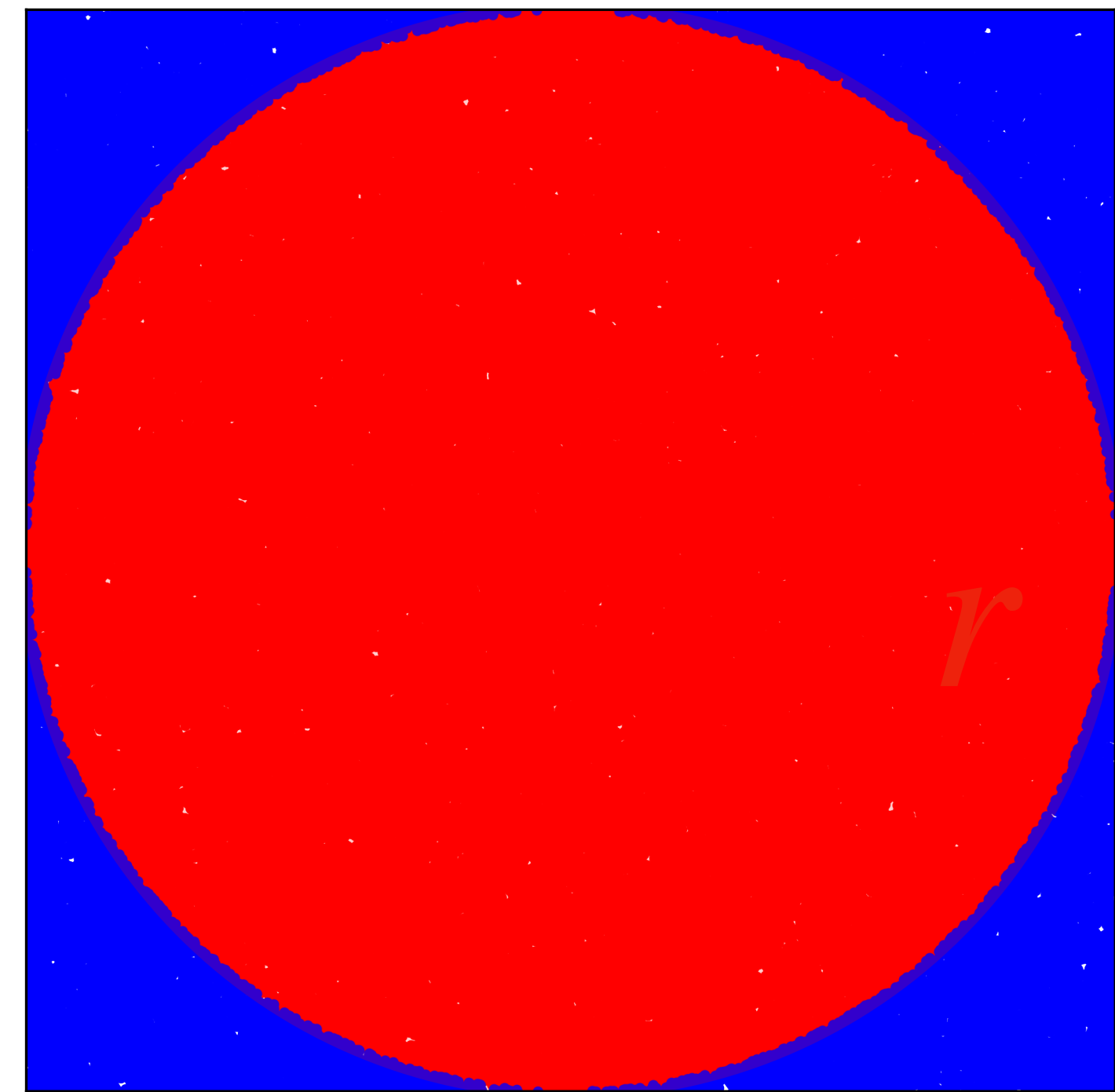
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$$N = 10^5, \pi \approx 3.14 \pm 0.01$$

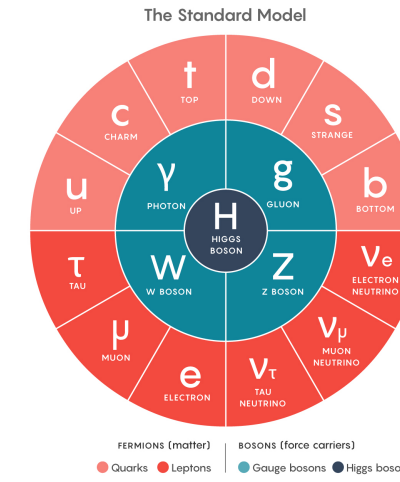


Simulations of Events via Monte Carlo

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- A set of particles and their interactions! e.g.:

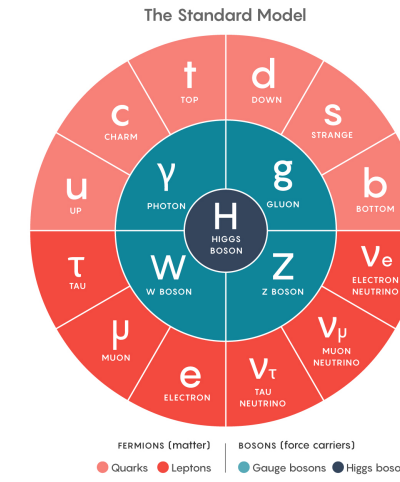


- Calculate probabilities via Quantum Field Theory → actually $|\mathcal{M}|^2$ (the matrix element).

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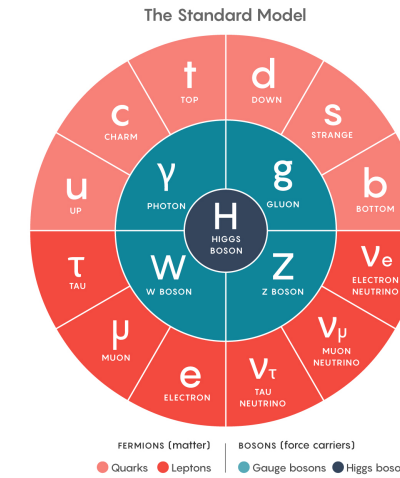
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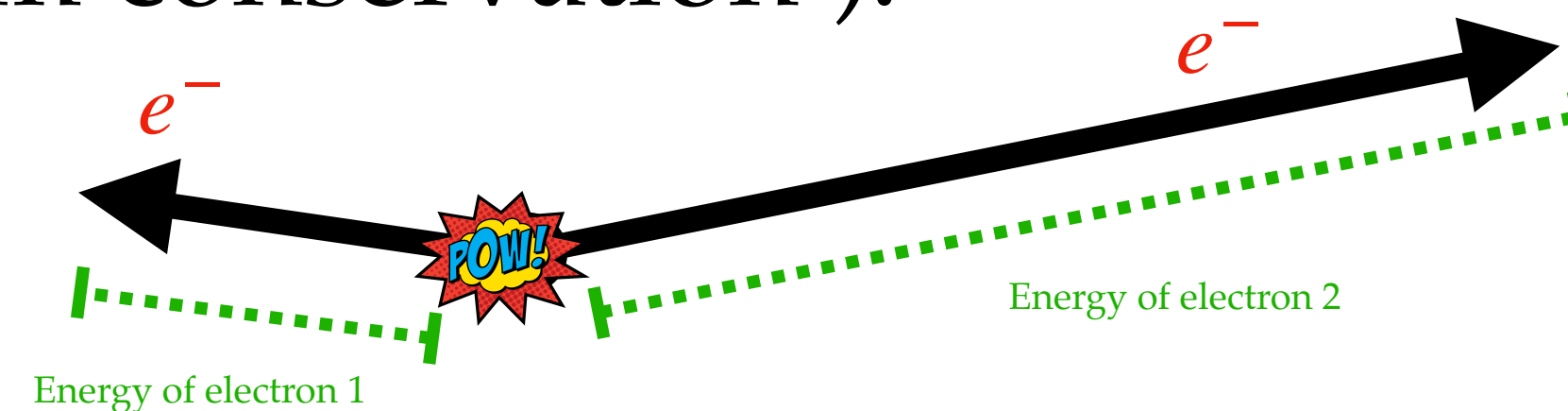
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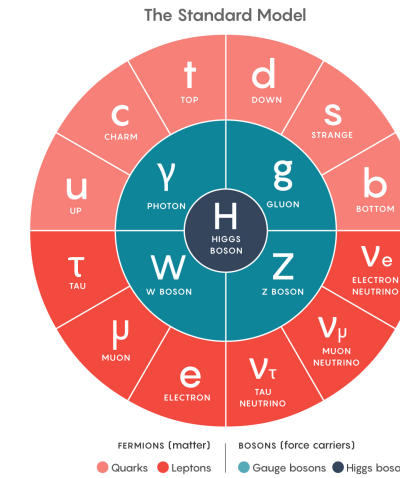
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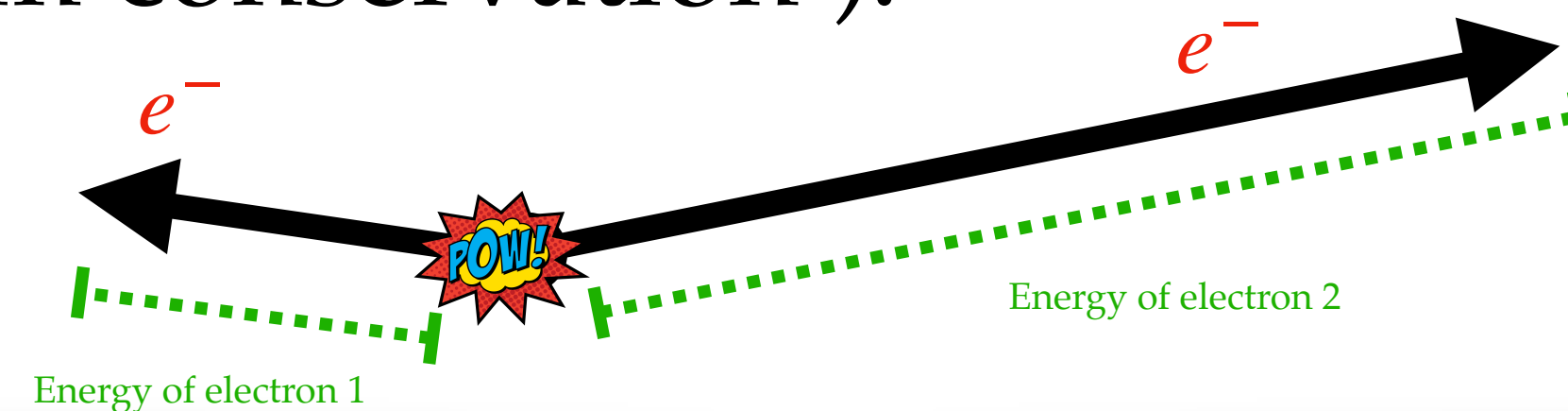
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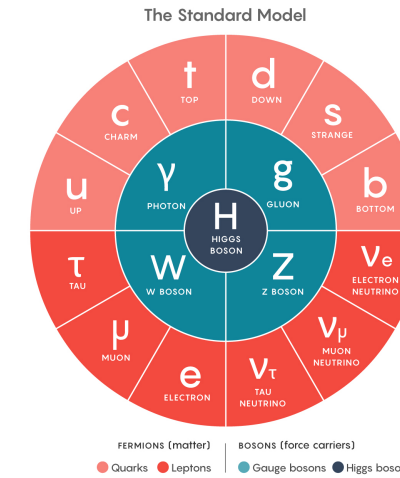


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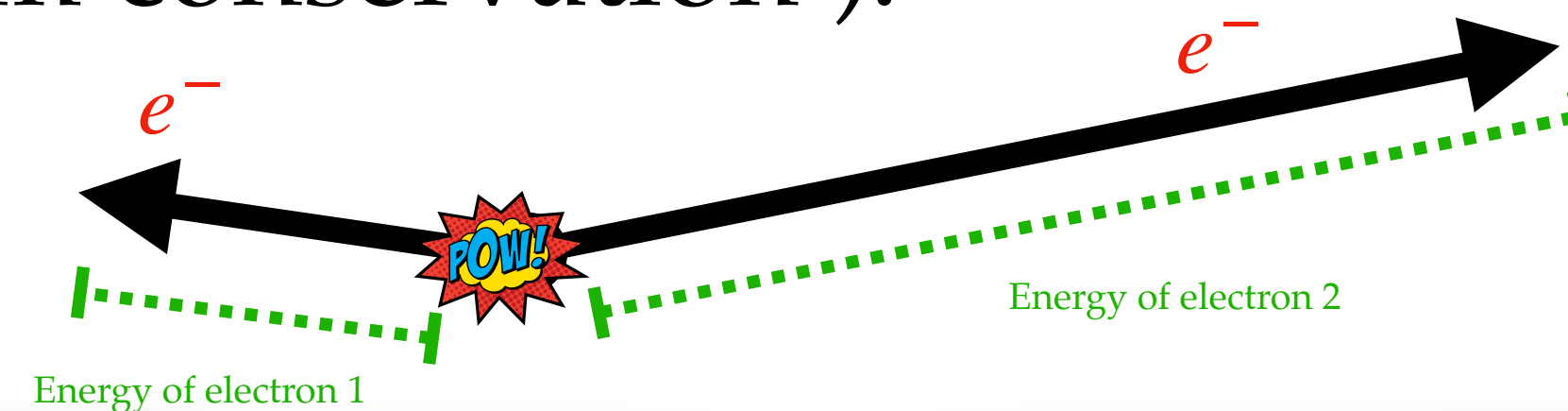
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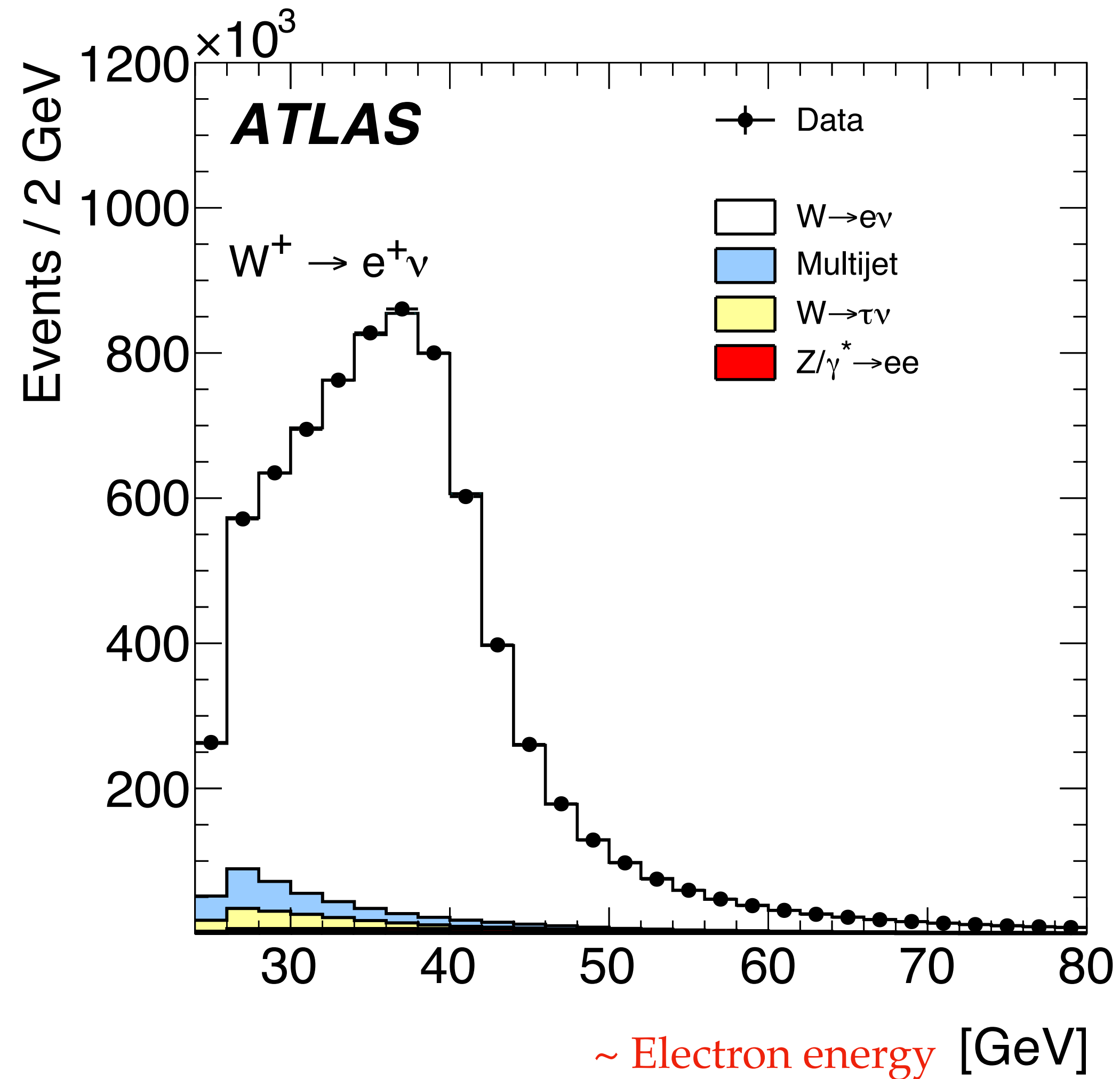
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[e.g. “How-to: write a parton-level Monte Carlo particle physics event generator”, AP, [arXiv:1412.4677](https://arxiv.org/abs/1412.4677)]

Theory **VS** Experiments

Construct histograms &
Compare to real data!

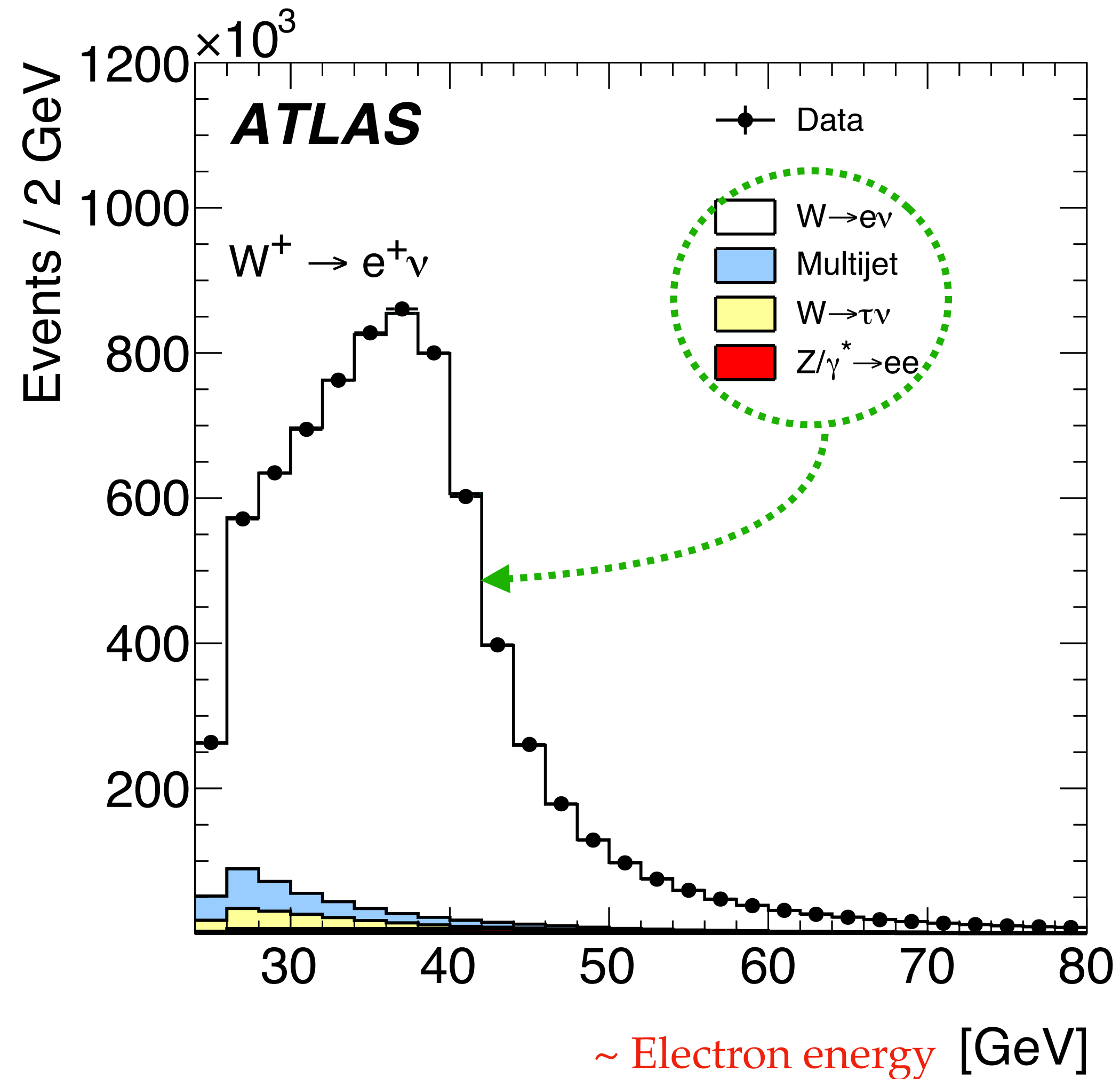
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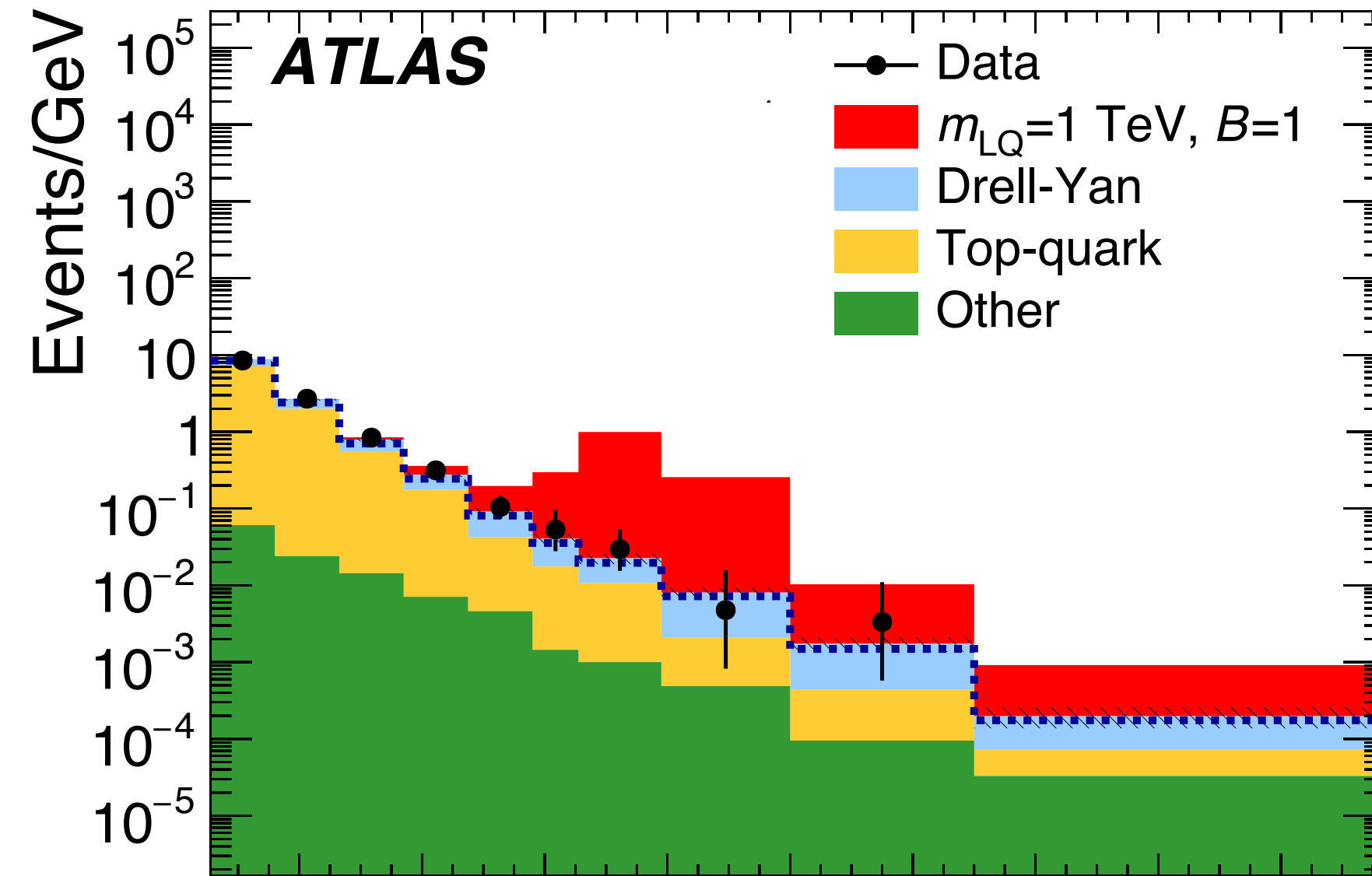


Searching for New Phenomena



Construct histograms &
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e.g.:



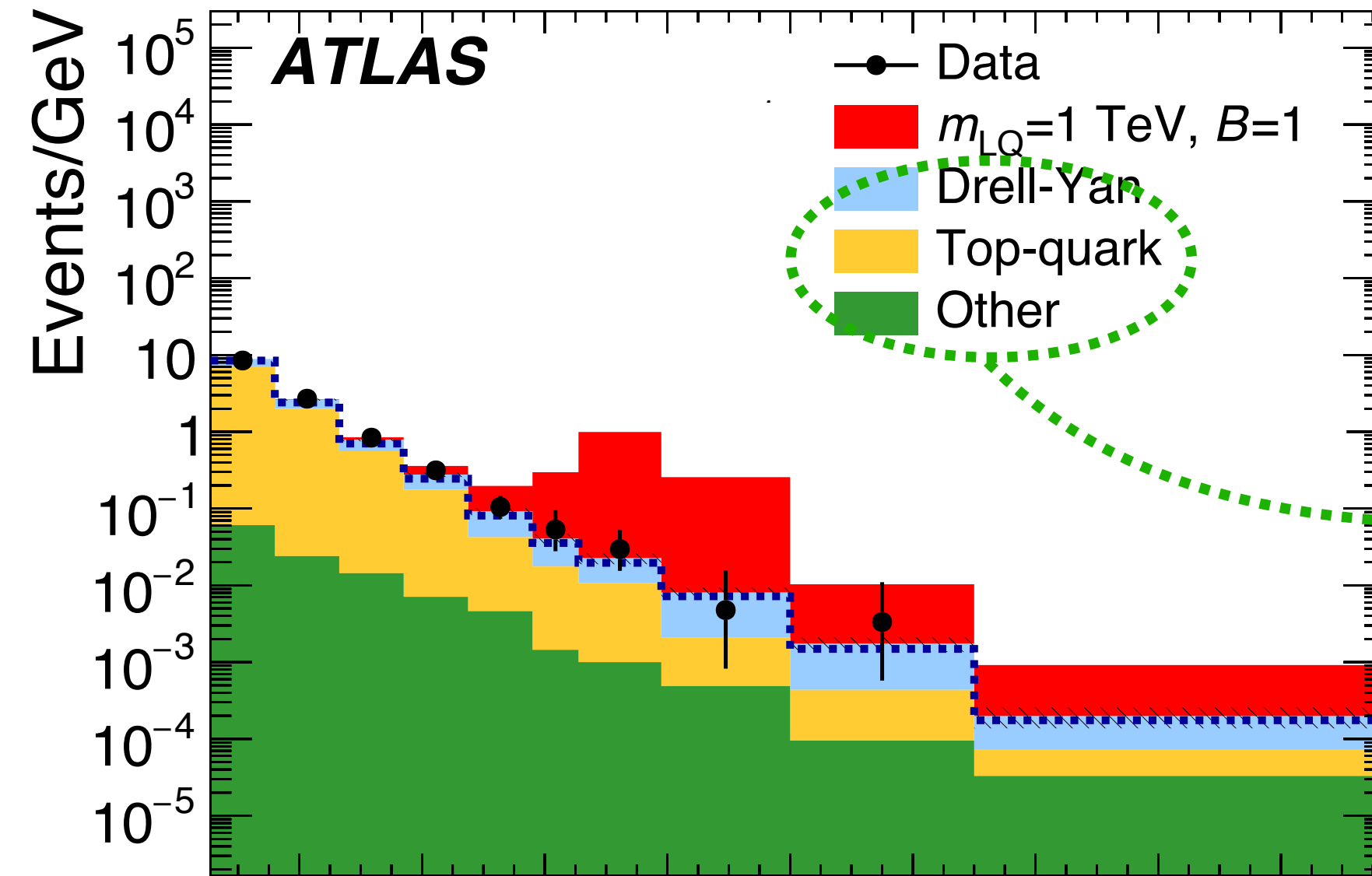
Some Observable Quantity [GeV]

Searching for New Phenomena



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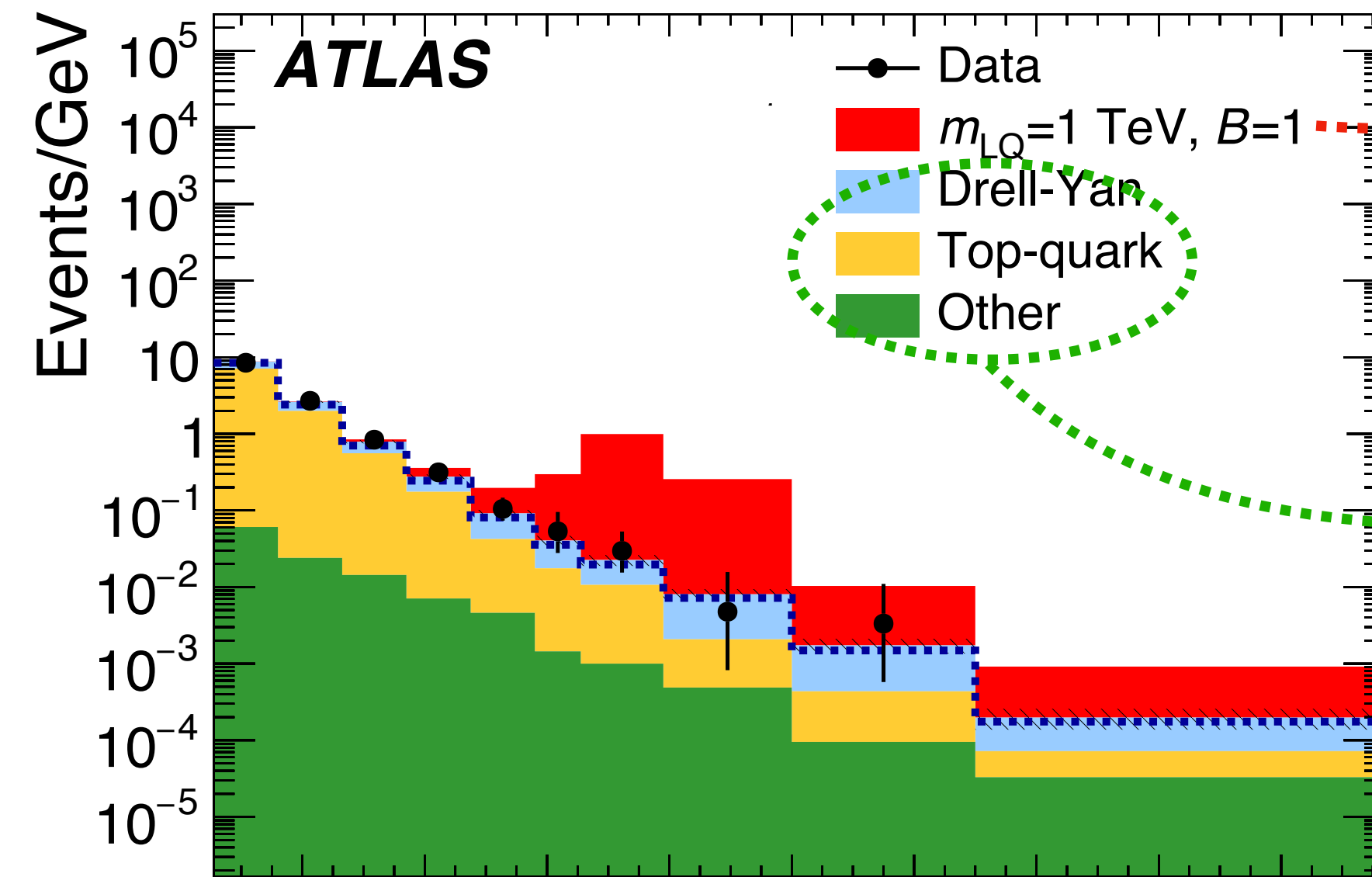
Searching for New Phenomena



NEW!

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e.g.:



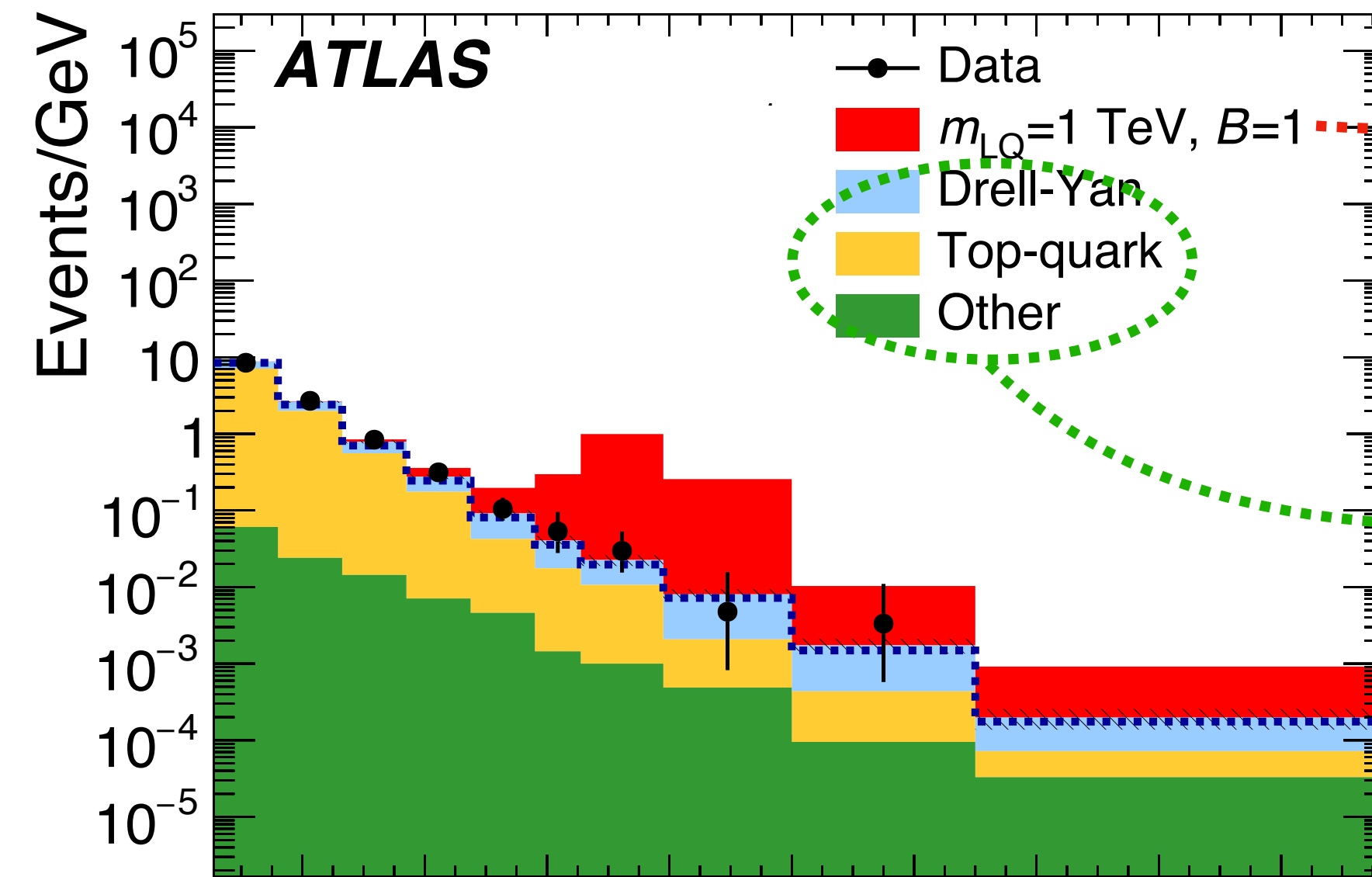
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Searching for New Phenomena



NEW!

Construct histograms &
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e.g.:



Hypothetical particle

Simulation of
stuff that we know exist!

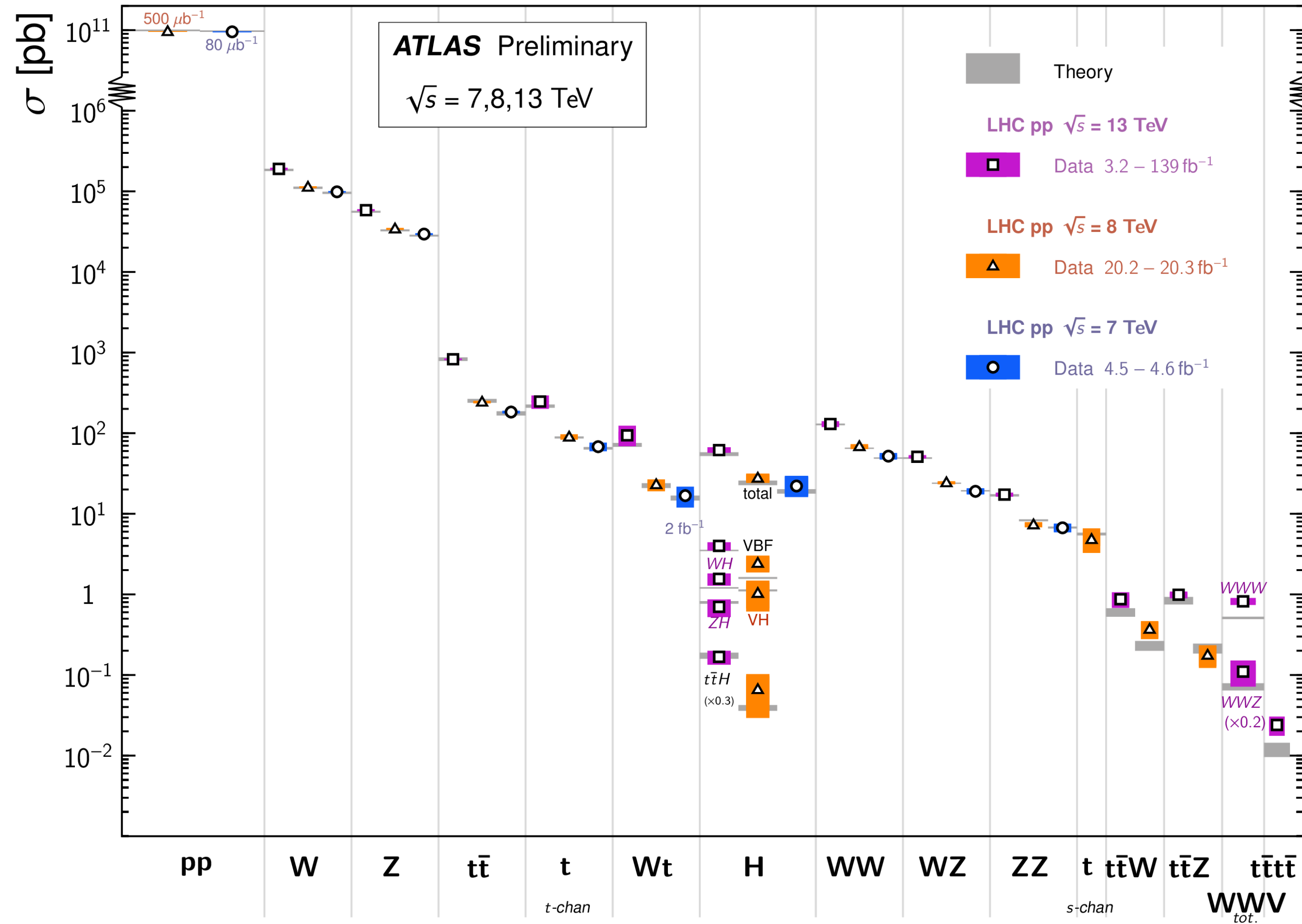
Some Observable Quantity [GeV]

⇒ Determine whether a hypothetical particle is compatible with data!

Experiment *VS* Theory

Standard Model Total Production Cross Section Measurements

Status: February 2022



Cross Section [pb]

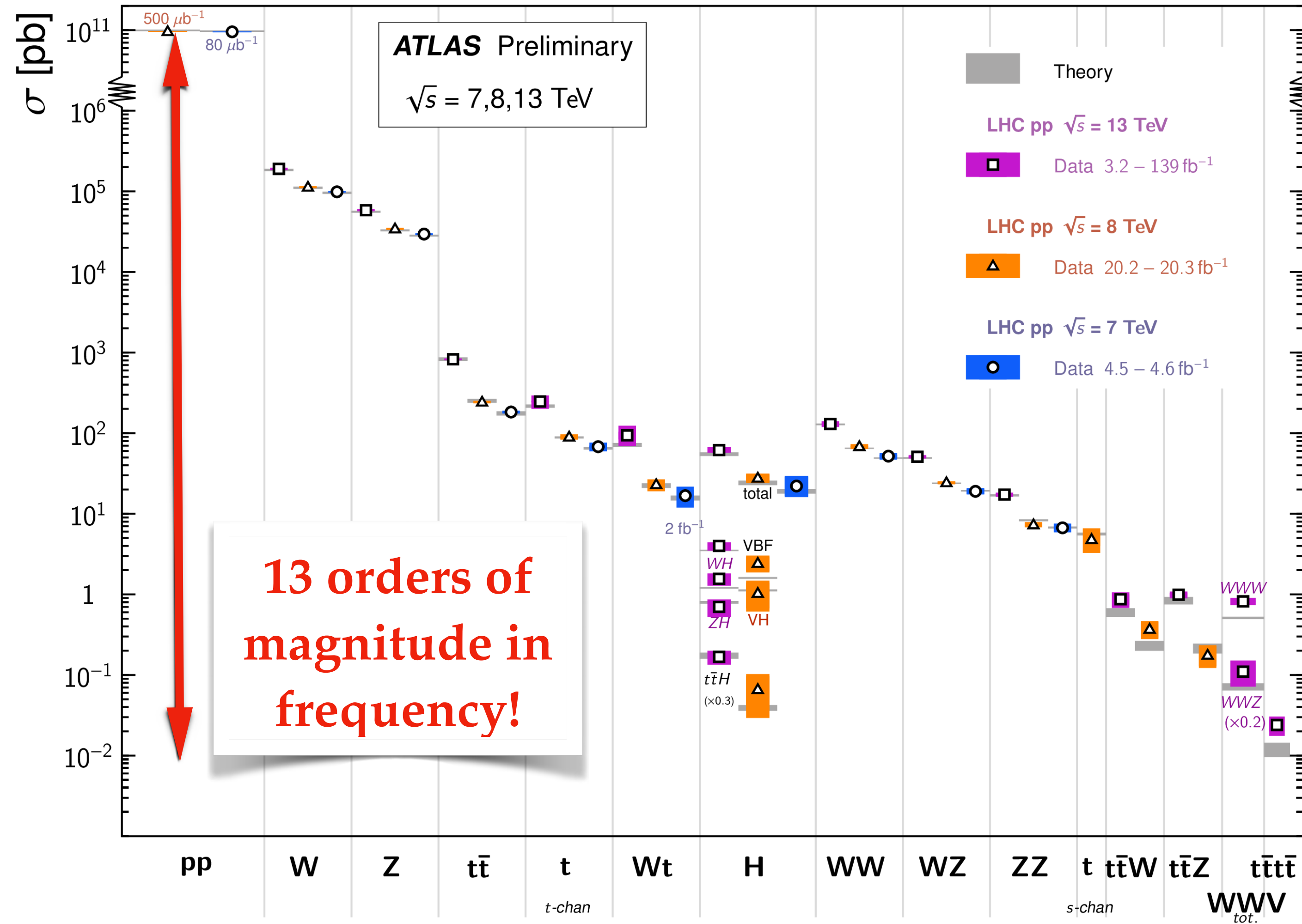
~ Frequency of
type of event

Process Type

Experiment *VS* Theory

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Cross Section [pb]

~ Frequency of
type of event

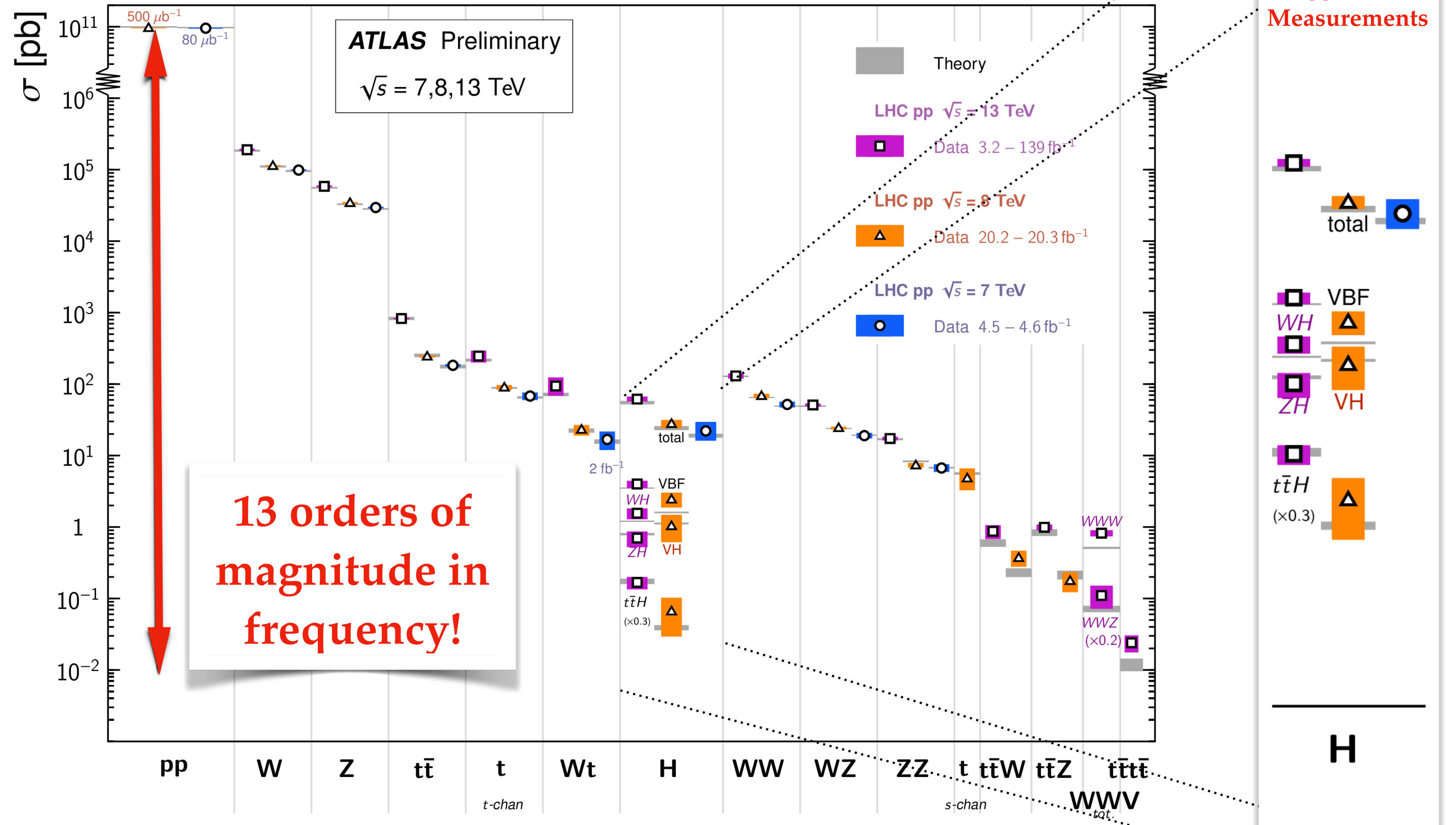
13 orders of
magnitude in
frequency!

Process Type

Experiment *VS* Theory

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Cross Section [pb]

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Process Type

Experiment VS New Phenomena

ATLAS Heavy Particle Searches* - 95% CL Upper Exclusion Limits

Status: March 2023

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.6 - 139) \text{ fb}^{-1}$$

$$\sqrt{s} = 13 \text{ TeV}$$

| Model | ℓ, γ | Jets [†] | E_T^{miss} | $\int \mathcal{L} dt [\text{fb}^{-1}]$ | Limit | Reference | |
|---|--|--|-----------------------------|--|---|--|--|
| Extra dimen. | ADD $G_{KK} + g/q$ | $0 e, \mu, \tau, \gamma$ | $1 - 4 j$ | Yes | 139 | M_D 11.2 TeV $n=2$ | 2102.10874 |
| | ADD non-resonant $\gamma\gamma$ | 2γ | - | - | 36.7 | M_S 8.6 TeV $n=3$ HLZ NLO | 1707.04147 |
| | ADD QBH | - | $2 j$ | - | 139 | M_{th} 9.4 TeV $n=6$ | 1910.08447 |
| | ADD BH multijet | - | $\geq 3 j$ | - | 3.6 | M_{th} 9.55 TeV $n=6, M_D=3 \text{ TeV}, \text{rot BH}$ | 1512.02586 |
| | RS1 $G_{KK} \rightarrow \gamma\gamma$ | 2γ | - | - | 139 | G_{KK} mass 4.5 TeV $k/\bar{M}_{Pl} = 0.1$ | 2102.13405 |
| | Bulk RS $G_{KK} \rightarrow WW/ZZ$ | multi-channel | - | - | 36.1 | G_{KK} mass 2.3 TeV $k/\bar{M}_{Pl} = 1.0$ | 1808.02380 |
| | Bulk RS $g_{KK} \rightarrow tt$ | $1 e, \mu$ | $\geq 1 b, \geq 1 J/2 j$ | Yes | 36.1 | g_{KK} mass 3.8 TeV $\Gamma/m = 15\%$ | 1804.10823 |
| | 2UED / RPP | $1 e, \mu$ | $\geq 2 b, \geq 3 j$ | Yes | 36.1 | KK mass 1.8 TeV Tier (1,1), $\mathcal{B}(A^{(1,1)} \rightarrow tt) = 1$ | 1803.09678 |
| | Gauge bosons | SSM $Z' \rightarrow \ell\ell$ | $2 e, \mu$ | - | - | 139 | Z' mass 5.1 TeV |
| SSM $Z' \rightarrow \tau\tau$ | | 2τ | - | - | 36.1 | Z' mass 2.42 TeV | 1709.07242 |
| Leptophobic $Z' \rightarrow bb$ | | - | $2 b$ | - | 36.1 | Z' mass 2.1 TeV | 1805.09299 |
| Leptophobic $Z' \rightarrow tt$ | | $0 e, \mu$ | $\geq 1 b, \geq 2 J$ | Yes | 139 | Z' mass 4.1 TeV $\Gamma/m = 1.2\%$ | 2005.05138 |
| SSM $W' \rightarrow \ell\nu$ | | $1 e, \mu$ | - | Yes | 139 | W' mass 6.0 TeV | 1906.05609 |
| SSM $W' \rightarrow \tau\nu$ | | 1τ | - | Yes | 139 | W' mass 5.0 TeV | ATLAS-CONF-2021-025 |
| SSM $W' \rightarrow tb$ | | - | $\geq 1 b, \geq 1 J$ | - | 139 | W' mass 4.4 TeV | ATLAS-CONF-2021-043 |
| HVT $W' \rightarrow WZ$ model B | | $0-2 e, \mu$ | $2 j / 1 J$ | Yes | 139 | W' mass 4.3 TeV | 2004.14636 |
| HVT $W' \rightarrow WZ \rightarrow \ell\nu \ell' \ell'$ model C | | $3 e, \mu$ | $2 j$ (VBF) | Yes | 139 | W' mass 340 GeV $g_V = 3, g_V c_H = 1, g_R = 0$ | 2207.03925 |
| HVT $Z' \rightarrow WW$ model B | | $1 e, \mu$ | $2 j / 1 J$ | Yes | 139 | Z' mass 3.9 TeV $g_V = 3$ | 2004.14636 |
| LRSB $W_R \rightarrow \mu N_R$ | | 2μ | $1 J$ | - | 80 | W_R mass 5.0 TeV $m(N_R) = 0.5 \text{ TeV}, g_L = g_R$ | 1904.12679 |
| CI | CI $qqqq$ | - | $2 j$ | - | 37.0 | Λ 21.8 TeV η_{LL} | 1703.09127 |
| | CI $\ell\ell qq$ | $2 e, \mu$ | - | - | 139 | Λ 35.8 TeV η_{LL} | 2006.12946 |
| | CI $e e b s$ | $2 e$ | $1 b$ | - | 139 | Λ 1.8 TeV $g_* = 1$ | 2105.13847 |
| | CI $\mu\mu b s$ | 2μ | $1 b$ | - | 139 | Λ 2.0 TeV $g_* = 1$ | 2105.13847 |
| | CI $tttt$ | $\geq 1 e, \mu$ | $\geq 1 b, \geq 1 j$ | Yes | 36.1 | Λ 2.57 TeV $ C_{4t} = 4\pi$ | 1811.02305 |
| | DM | Axial-vector med. (Dirac DM) | - | $2 j$ | - | 139 | m_{med} 3.8 TeV $g_a = 0.25, g_s = 1, m(\chi) = 10 \text{ TeV}$ |
| Pseudo-scalar med. (Dirac DM) | | $0 e, \mu, \tau, \gamma$ | $1 - 4 j$ | Yes | 139 | m_{med} 376 GeV $g_a = 1, g_s = 1, m(\chi) = 1 \text{ GeV}$ | 2102.10874 |
| Vector med. Z' -2HDM (Dirac DM) | | $0 e, \mu$ | $2 b$ | Yes | 139 | $m_{Z'}$ 3.0 TeV $\tan\beta = 1, g_z = 0.8, m(\chi) = 100 \text{ GeV}$ | 2108.13391 |
| Pseudo-scalar med. 2HDM+a | | multi-channel | - | - | 139 | m_a 800 GeV $\tan\beta = 1, g_s = 1, m(\chi) = 10 \text{ GeV}$ | ATLAS-CONF-2021-036 |
| LQ | Scalar LQ 1 st gen | $2 e$ | $\geq 2 j$ | Yes | 139 | LQ mass 1.8 TeV $\beta = 1$ | 2006.05872 |
| | Scalar LQ 2 nd gen | 2μ | $\geq 2 j$ | Yes | 139 | LQ mass 1.7 TeV $\beta = 1$ | 2006.05872 |
| | Scalar LQ 3 rd gen | 1τ | $2 b$ | Yes | 139 | LQ_3^u mass 1.49 TeV $\mathcal{B}(LQ_3^u \rightarrow b\tau) = 1$ | 2303.01294 |
| | Scalar LQ 3 rd gen | $0 e, \mu$ | $\geq 2 j, \geq 2 b$ | Yes | 139 | LQ_3^d mass 1.24 TeV $\mathcal{B}(LQ_3^d \rightarrow \nu\tau) = 1$ | 2004.14060 |
| | Scalar LQ 3 rd gen | $\geq 2 e, \mu, \geq 1 \tau, \geq 1 j, \geq 1 b$ | - | - | 139 | LQ_3^d mass 1.43 TeV $\mathcal{B}(LQ_3^d \rightarrow \tau\nu) = 1$ | 2101.11582 |
| | Scalar LQ 3 rd gen | $0 e, \mu, \geq 1 \tau, 0 - 2 j, 2 b$ | Yes | 139 | LQ_3^d mass 1.26 TeV $\mathcal{B}(LQ_3^d \rightarrow b\nu) = 1$ | 2101.12527 | |
| | Vector LQ mix gen | multi-channel | $\geq 1 j, \geq 1 b$ | Yes | 139 | LQ_3^u mass 2.0 TeV $\mathcal{B}(\tilde{U}_1 \rightarrow t\mu) = 1, \text{Y-M coupl.}$ | ATLAS-CONF-2022-052 |
| | Vector LQ 3 rd gen | $2 e, \mu, \tau$ | $\geq 1 b$ | Yes | 139 | LQ_3^u mass 1.96 TeV $\mathcal{B}(LQ_3^u \rightarrow b\tau) = 1, \text{Y-M coupl.}$ | 2303.01294 |
| | Vector-like fermions | VLQ $TT \rightarrow Zt + X$ | $2e/2\mu \geq 3e, \mu$ | $\geq 1 b, \geq 1 j$ | - | 139 | T mass 1.46 TeV SU(2) doublet |
| VLQ $BB \rightarrow Wt/Zb + X$ | | multi-channel | - | - | 36.1 | B mass 1.34 TeV SU(2) doublet | 1808.02343 |
| VLQ $T_{5/3} T_{5/3} / T_{5/3} \rightarrow Wt + X$ | | $2(SS) \geq 3 e, \mu$ | $\geq 1 b, \geq 1 j$ | Yes | 36.1 | $T_{5/3}$ mass 1.64 TeV $\mathcal{B}(T_{5/3} \rightarrow Wt) = 1, c(T_{5/3} Wt) = 1$ | 1807.11883 |
| VLQ $T \rightarrow Ht/Zt$ | | $1 e, \mu$ | $\geq 1 b, \geq 3 j$ | Yes | 139 | T mass 1.8 TeV SU(2) singlet, $\kappa_T = 0.5$ | ATLAS-CONF-2021-040 |
| VLQ $Y \rightarrow Wb$ | | $1 e, \mu$ | $\geq 1 b, \geq 1 j$ | Yes | 36.1 | Y mass 1.85 TeV $\mathcal{B}(Y \rightarrow Wb) = 1, c_R(Wb) = 1$ | 1812.07343 |
| VLQ $B \rightarrow Hb$ | | $0 e, \mu$ | $\geq 2b, \geq 1j, \geq 1J$ | - | 139 | B mass 2.0 TeV SU(2) doublet, $\kappa_B = 0.3$ | ATLAS-CONF-2021-018 |
| VLL $\tau' \rightarrow Z\tau/H\tau$ | | multi-channel | $\geq 1 j$ | Yes | 139 | τ' mass 898 GeV SU(2) doublet | 2303.05441 |
| Excited ferm. | Excited quark $q^* \rightarrow qg$ | - | $2 j$ | - | 139 | q^* mass 6.7 TeV only u^* and d^* , $\Lambda = m(q^*)$ | 1910.08447 |
| | Excited quark $q^* \rightarrow q\gamma$ | 1γ | $1 j$ | - | 36.7 | q^* mass 5.3 TeV only u^* and d^* , $\Lambda = m(q^*)$ | 1709.10440 |
| | Excited quark $b^* \rightarrow bg$ | - | $1 b, 1 j$ | - | 139 | b^* mass 3.2 TeV | 1910.08447 |
| | Excited lepton τ^* | 2τ | $\geq 2 j$ | - | 139 | τ^* mass 4.6 TeV $\Lambda = 4.6 \text{ TeV}$ | 2303.09444 |
| Other | Type III Seesaw | $2,3,4 e, \mu$ | $\geq 2 j$ | Yes | 139 | N^0 mass 910 GeV | 2202.02039 |
| | LRSB Majorana ν | 2μ | $2 j$ | - | 36.1 | N_R mass 3.2 TeV $m(W_R) = 4.1 \text{ TeV}, g_L = g_R$ | 1809.11105 |
| | Higgs triplet $H^{\pm\pm} \rightarrow W^\pm W^\pm$ | $2,3,4 e, \mu$ (SS) | various | Yes | 139 | $H^{\pm\pm}$ mass 350 GeV DY production | 2101.11961 |
| | Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$ | $2,3,4 e, \mu$ (SS) | - | - | 139 | $H^{\pm\pm}$ mass 1.08 TeV DY production | 2211.07505 |
| | Multi-charged particles | - | - | - | 139 | multi-charged particle mass 1.59 TeV DY production, $ q = 5e$ | ATLAS-CONF-2022-034 |
| | Magnetic monopoles | - | - | - | 34.4 | monopole mass 2.37 TeV DY production, $ g = 1g_D, \text{spin } 1/2$ | 1905.10130 |

$\sqrt{s} = 13 \text{ TeV}$
partial data

$\sqrt{s} = 13 \text{ TeV}$
full data

10⁻¹ 1 10 Mass scale [TeV]

*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).



Experiment **VS** New Phenomena

ATLAS Heavy Particle Searches* - 95% CL Upper Exclusion Limits

Status: March 2023

ATLAS Preliminary

$\int \mathcal{L} dt = (3.6 - 139) \text{ fb}^{-1}$

$\sqrt{s} = 13 \text{ TeV}$

| Model | ℓ, γ | Jets [†] | E_T^{miss} | $\int \mathcal{L} dt [\text{fb}^{-1}]$ | Limit | Reference | |
|---|--|--|-----------------------------|--|--|---|-------------------------------|
| Extra dimen. | ADD $G_{KK} + g/q$ | $0 e, \mu, \tau, \gamma$ | $1 - 4 j$ | Yes | 139 | M_D 11.2 TeV $n=2$ | 2102.10874 |
| | ADD non-resonant $\gamma\gamma$ | 2γ | - | - | 36.7 | M_S 8.6 TeV $n=3$ HLZ NLO | 1707.04147 |
| | ADD QBH | - | $2 j$ | - | 139 | M_{th} 9.4 TeV $n=6$ | 1910.08447 |
| | ADD BH multijet | - | $\geq 3 j$ | - | 3.6 | M_{th} 9.55 TeV $n=6, M_D = 3 \text{ TeV}$, rot BH | 1512.02586 |
| | RS1 $G_{KK} \rightarrow \gamma\gamma$ | 2γ | - | - | 139 | G_{KK} mass 4.5 TeV $k/\overline{M}_{Pl} = 0.1$ | 2102.13405 |
| | Bulk RS $G_{KK} \rightarrow WW/ZZ$ | multi-channel | - | - | 36.1 | G_{KK} mass 2.3 TeV $k/\overline{M}_{Pl} = 1.0$ | 1808.02380 |
| Gauge bosons | Bulk RS $g_{KK} \rightarrow tt$ | $1 e, \mu$ | $\geq 1 b, \geq 1 J/2 j$ | Yes | 36.1 | g_{KK} mass 3.8 TeV $\Gamma/m = 15\%$ | 1804.10823 |
| | 2UED / RPP | $1 e, \mu$ | $\geq 2 b, \geq 3 j$ | Yes | 36.1 | KK mass 1.8 TeV Tier (1,1), $\mathcal{B}(A^{(1,1)} \rightarrow tt) = 1$ | 1803.09678 |
| | SSM $Z' \rightarrow \ell\ell$ | $2 e, \mu$ | - | - | 139 | Z' mass 5.1 TeV | 1903.06248 |
| | SSM $Z' \rightarrow \tau\tau$ | 2τ | - | - | 36.1 | Z' mass 2.42 TeV | 1709.07242 |
| | Leptophobic $Z' \rightarrow bb$ | - | $2 b$ | - | 36.1 | Z' mass 2.1 TeV | 1805.09299 |
| | Leptophobic $Z' \rightarrow tt$ | $0 e, \mu$ | $\geq 1 b, \geq 2 J$ | Yes | 139 | Z' mass 4.1 TeV $\Gamma/m = 1.2\%$ | 2005.05138 |
| | SSM $W' \rightarrow \ell\nu$ | $1 e, \mu$ | - | - | 139 | W' mass 6.0 TeV | 1906.05609 |
| | SSM $W' \rightarrow \tau\nu$ | 1τ | - | - | 139 | W' mass 5.0 TeV | ATLAS-CONF-2021-025 |
| | SSM $W' \rightarrow tb$ | - | $\geq 1 b, \geq 1 J$ | - | 139 | W' mass 4.4 TeV | ATLAS-CONF-2021-043 |
| | HVT $W' \rightarrow WZ$ model B | $0-2 e, \mu$ | $2 j / 1 J$ | Yes | 139 | W' mass 4.3 TeV $g_V = 3$ | 2004.14636 |
| HVT $W' \rightarrow WZ \rightarrow \ell\nu \ell' \ell'$ model C | $3 e, \mu$ | $2 j$ (VBF) | Yes | 139 | W' mass 340 GeV $g_V c_H = 1, g_R = 0$ | 2207.03925 | |
| LVT $Z' \rightarrow WW$ model B | $1 e, \mu$ | $2 j / 1 J$ | Yes | 139 | Z' mass 3.9 TeV $g_V = 3$ | 2004.14636 | |
| LRSM $W_R \rightarrow \mu N_R$ | 2μ | $1 J$ | - | 80 | W_R mass 5.0 TeV $m(N_R) = 0.5 \text{ TeV}, g_L = g_R$ | 1904.12679 | |
| CI | CI $qqqq$ | - | $2 j$ | - | 37.0 | Λ 21.8 TeV η_{LL} | 1703.09127 |
| | CI $\ell\ell qq$ | $2 e, \mu$ | - | - | 139 | Λ 35.8 TeV η_{LL} | 2006.12946 |
| | CI $e e b s$ | $2 e$ | $1 b$ | - | 139 | Λ 1.8 TeV $g_s = 1$ | 2105.13847 |
| | CI $\mu\mu b s$ | 2μ | $1 b$ | - | 139 | Λ 2.0 TeV $g_s = 1$ | 2105.13847 |
| | CI $tttt$ | $\geq 1 e, \mu$ | $\geq 1 b, \geq 1 j$ | Yes | 36.1 | Λ 2.57 TeV $ C_{4t} = 4\pi$ | 1811.02305 |
| DM | Axial-vector med. (Dirac DM) | - | $2 j$ | - | 139 | m_{med} 3.8 TeV $g_a = 0.25, g_s = 1, m(\chi) = 10 \text{ GeV}$ | ATL-PHYS-PUB-2022-036 |
| | Pseudo-scalar med. (Dirac DM) | $0 e, \mu, \tau, \gamma$ | $1 - 4 j$ | Yes | 139 | m_{med} 376 GeV $g_a = 1, g_s = 1, m(\chi) = 1 \text{ GeV}$ | 2102.10874 |
| | Vector med. Z' -2HDM (Dirac DM) | $0 e, \mu$ | $2 b$ | Yes | 139 | $m_{Z'}$ 3.0 TeV $\tan\beta = 1, g_Z = 0.8, m(\chi) = 100 \text{ GeV}$ | 2108.13391 |
| | Pseudo-scalar med. 2HDM+a | multi-channel | - | - | 139 | m_a 800 GeV $\tan\beta = 1, g_s = 1, m(\chi) = 10 \text{ GeV}$ | ATLAS-CONF-2021-036 |
| LQ | Scalar LQ 1 st gen | $2 e$ | $\geq 2 j$ | Yes | 139 | LQ mass 1.8 TeV $\beta = 1$ | 2006.05872 |
| | Scalar LQ 2 nd gen | 2μ | $\geq 2 j$ | Yes | 139 | LQ mass 1.7 TeV $\beta = 1$ | 2006.05872 |
| | Scalar LQ 3 rd gen | 1τ | $2 b$ | Yes | 139 | LQ ₃ mass 1.49 TeV $\mathcal{B}(LQ_3^u \rightarrow b\tau) = 1$ | 2303.01294 |
| | Scalar LQ 3 rd gen | $0 e, \mu$ | $\geq 2 j, \geq 2 b$ | Yes | 139 | LQ ₃ mass 1.24 TeV $\mathcal{B}(LQ_3^u \rightarrow t\nu) = 1$ | 2004.14060 |
| | Scalar LQ 3 rd gen | $\geq 2 e, \mu, \geq 1 \tau, \geq 1 j, \geq 1 b$ | - | - | 139 | LQ ₃ mass 1.43 TeV $\mathcal{B}(LQ_3^u \rightarrow t\tau) = 1$ | 2101.11582 |
| | Scalar LQ 3 rd gen | $0 e, \mu, \geq 1 \tau, 0 - 2 j, 2 b$ | Yes | 139 | LQ ₃ mass 1.26 TeV $\mathcal{B}(LQ_3^u \rightarrow b\nu) = 1$ | 2101.12527 | |
| | Vector LQ mix gen | multi-channel | $\geq 1 j, \geq 1 b$ | Yes | 139 | LQ ₃ mass 2.0 TeV $\mathcal{B}(\tilde{U}_1 \rightarrow t\mu) = 1, \text{Y-M coupl.}$ | ATLAS-CONF-2022-052 |
| | Vector LQ 3 rd gen | $2 e, \mu, \tau$ | $\geq 1 b$ | Yes | 139 | LQ ₃ mass 1.96 TeV $\mathcal{B}(LQ_3^u \rightarrow b\tau) = 1, \text{Y-M coupl.}$ | 2303.01294 |
| | Vector-like fermions | VLQ $TT \rightarrow Zt + X$ | $2e/2\mu \geq 3e, \mu$ | $\geq 1 b, \geq 1 j$ | - | 139 | T mass 1.46 TeV SU(2) doublet |
| VLQ $BB \rightarrow Wt/Zb + X$ | | multi-channel | - | - | 36.1 | B mass 1.34 TeV SU(2) doublet | 1808.02343 |
| VLQ $T_{5/3} T_{5/3} T_{5/3} \rightarrow Wt + X$ | | $2(SS) \geq 3 e, \mu$ | $\geq 1 b, \geq 1 j$ | Yes | 36.1 | $T_{5/3}$ mass 1.64 TeV $\mathcal{B}(T_{5/3} \rightarrow Wt) = 1, c(T_{5/3} Wt) = 1$ | 1807.11883 |
| VLQ $T \rightarrow Ht/Zt$ | | $1 e, \mu$ | $\geq 1 b, \geq 3 j$ | Yes | 139 | T mass 1.8 TeV SU(2) singlet, $\kappa_T = 0.5$ | ATLAS-CONF-2021-040 |
| VLQ $Y \rightarrow Wb$ | | $1 e, \mu$ | $\geq 1 b, \geq 1 j$ | Yes | 36.1 | Y mass 1.85 TeV $\mathcal{B}(Y \rightarrow Wb) = 1, c_R(Wb) = 1$ | 1812.07343 |
| VLQ $B \rightarrow Hb$ | | $0 e, \mu$ | $\geq 2b, \geq 1j, \geq 1J$ | - | 139 | B mass 2.0 TeV SU(2) doublet, $\kappa_B = 0.3$ | ATLAS-CONF-2021-018 |
| VLL $\tau' \rightarrow Z\tau/H\tau$ | | multi-channel | $\geq 1 j$ | Yes | 139 | τ' mass 898 GeV SU(2) doublet | 2303.05441 |
| Excited ferm. | Excited quark $q^* \rightarrow qg$ | - | $2 j$ | - | 139 | q^* mass 6.7 TeV only u^* and $d^*, \Lambda = m(q^*)$ | 1910.08447 |
| | Excited quark $q^* \rightarrow q\gamma$ | 1γ | $1 j$ | - | 36.7 | q^* mass 5.3 TeV only u^* and $d^*, \Lambda = m(q^*)$ | 1709.10440 |
| | Excited quark $b^* \rightarrow bg$ | - | $1 b, 1 j$ | - | 139 | b^* mass 3.2 TeV | 1910.08447 |
| | Excited lepton τ^* | 2τ | $\geq 2 j$ | - | 139 | τ^* mass 4.6 TeV $\Lambda = 4.6 \text{ TeV}$ | 2303.09444 |
| Other | Type III Seesaw | $2,3,4 e, \mu$ | $\geq 2 j$ | Yes | 139 | N^0 mass 910 GeV | 2202.02039 |
| | LRSM Majorana ν | 2μ | $2 j$ | - | 36.1 | N_R mass 3.2 TeV | 1809.11105 |
| | Higgs triplet $H^{\pm\pm} \rightarrow W^\pm W^\pm$ | $2,3,4 e, \mu$ (SS) | various | Yes | 139 | $H^{\pm\pm}$ mass 350 GeV $m(W_R) = 4.1 \text{ TeV}, g_L = g_R$ | 2101.11961 |
| | Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$ | $2,3,4 e, \mu$ (SS) | - | - | 139 | $H^{\pm\pm}$ mass 1.08 TeV DY production | 2211.07505 |
| | Multi-charged particles | - | - | - | 139 | multi-charged particle mass 1.59 TeV DY production, $ q = 5e$ | ATLAS-CONF-2022-034 |
| Magnetic monopoles | - | - | - | 34.4 | monopole mass 2.37 TeV DY production, $ g = 1g_D, \text{spin } 1/2$ | 1905.10130 | |

Mass Scale limits: 1-10 TeV

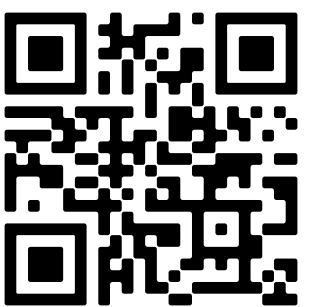
*Only a selection of the available mass limits on new states or phenomena is shown.

† Small-radius (large-radius) jets are denoted by the letter j (J).

$\sqrt{s} = 13 \text{ TeV}$
partial data

$\sqrt{s} = 13 \text{ TeV}$
full data

10⁻¹ 1 10 Mass scale [TeV]



EXPERIMENT WINS!



⇒ No new phenomena?



[credit: Alvaro Rujula,
Found in a CERN TH office]



EXPERIMENT WINS!



⇒ No new phenomena?



[credit: Alvaro Rujula,
Found in a CERN TH office]





Is that all folks?

Questions



Questions



Answers



Questions



Answers



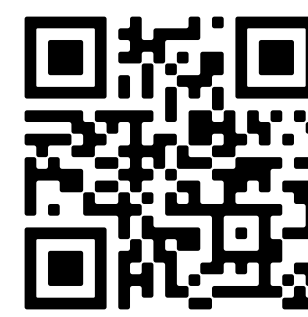
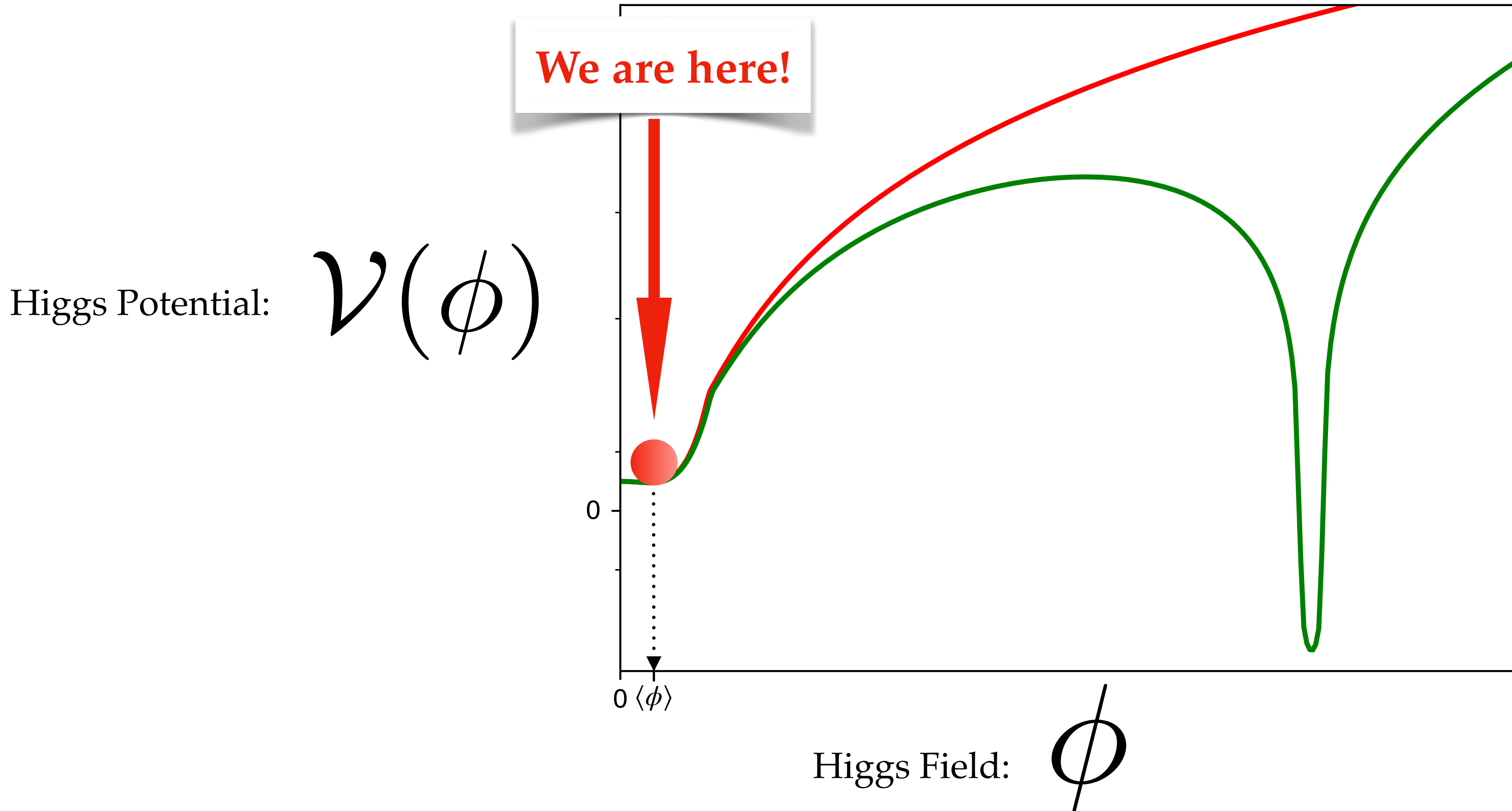
Questions



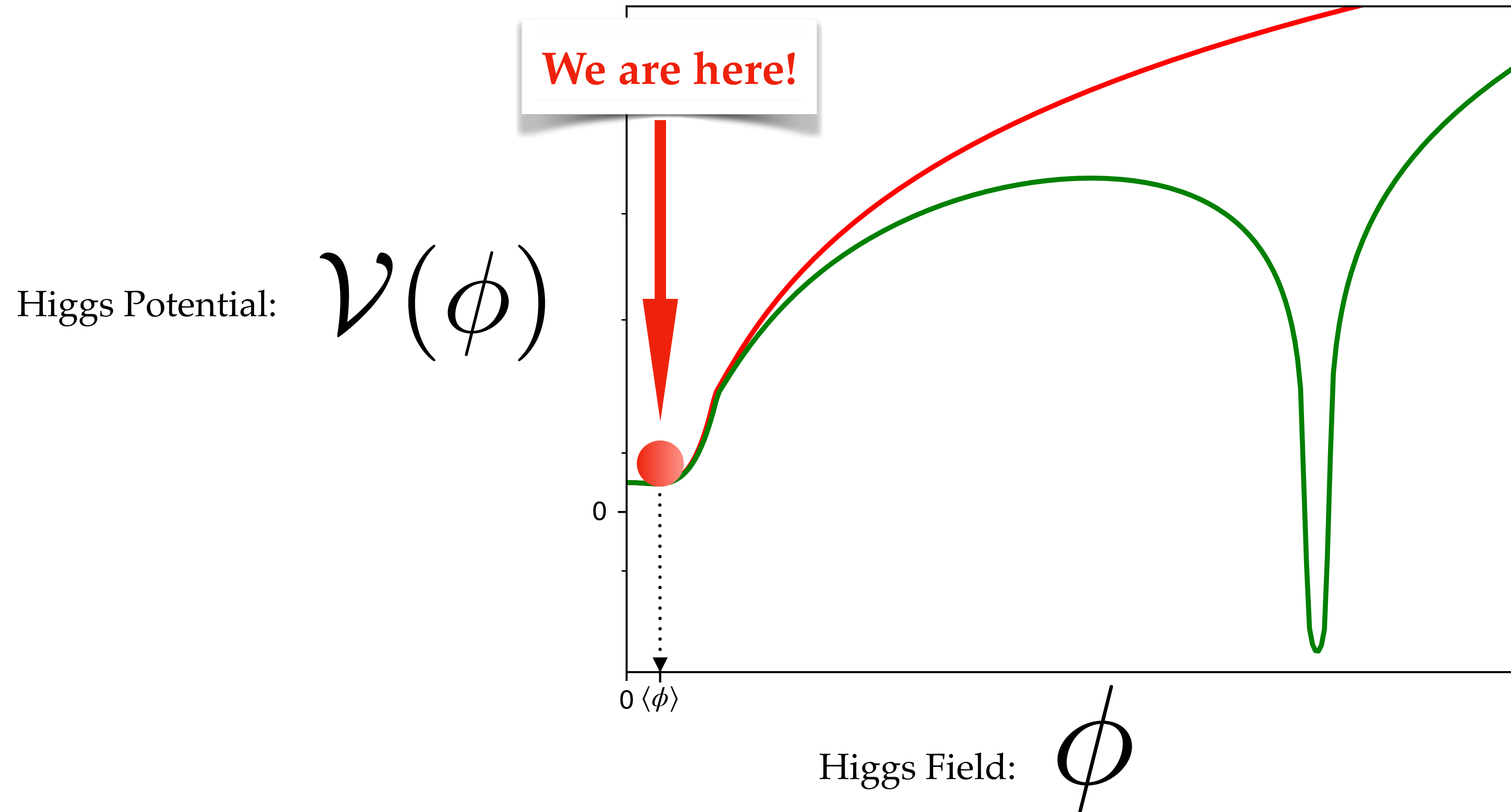
Answers



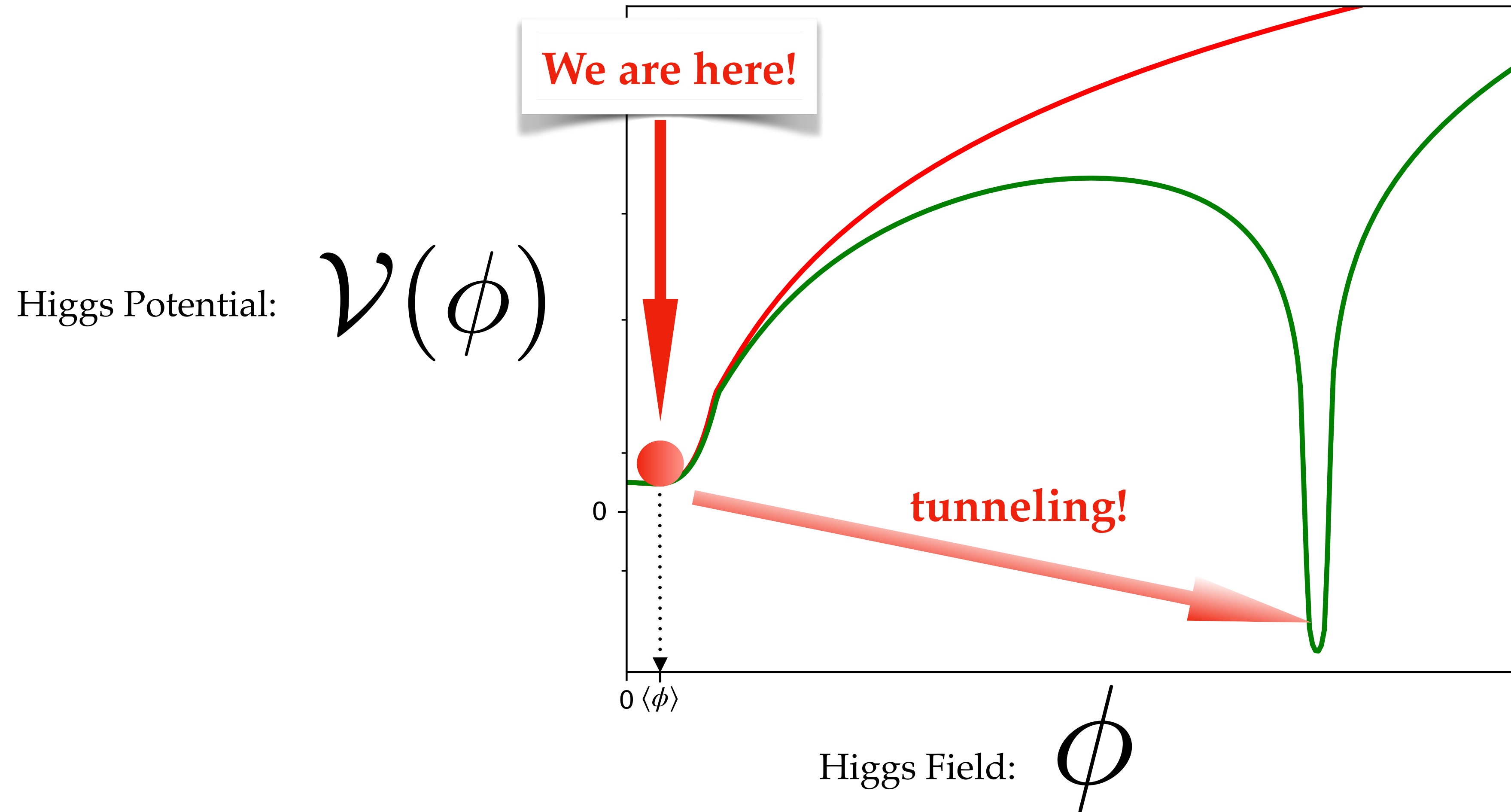
2 Q: What is the ultimate fate of our Universe?



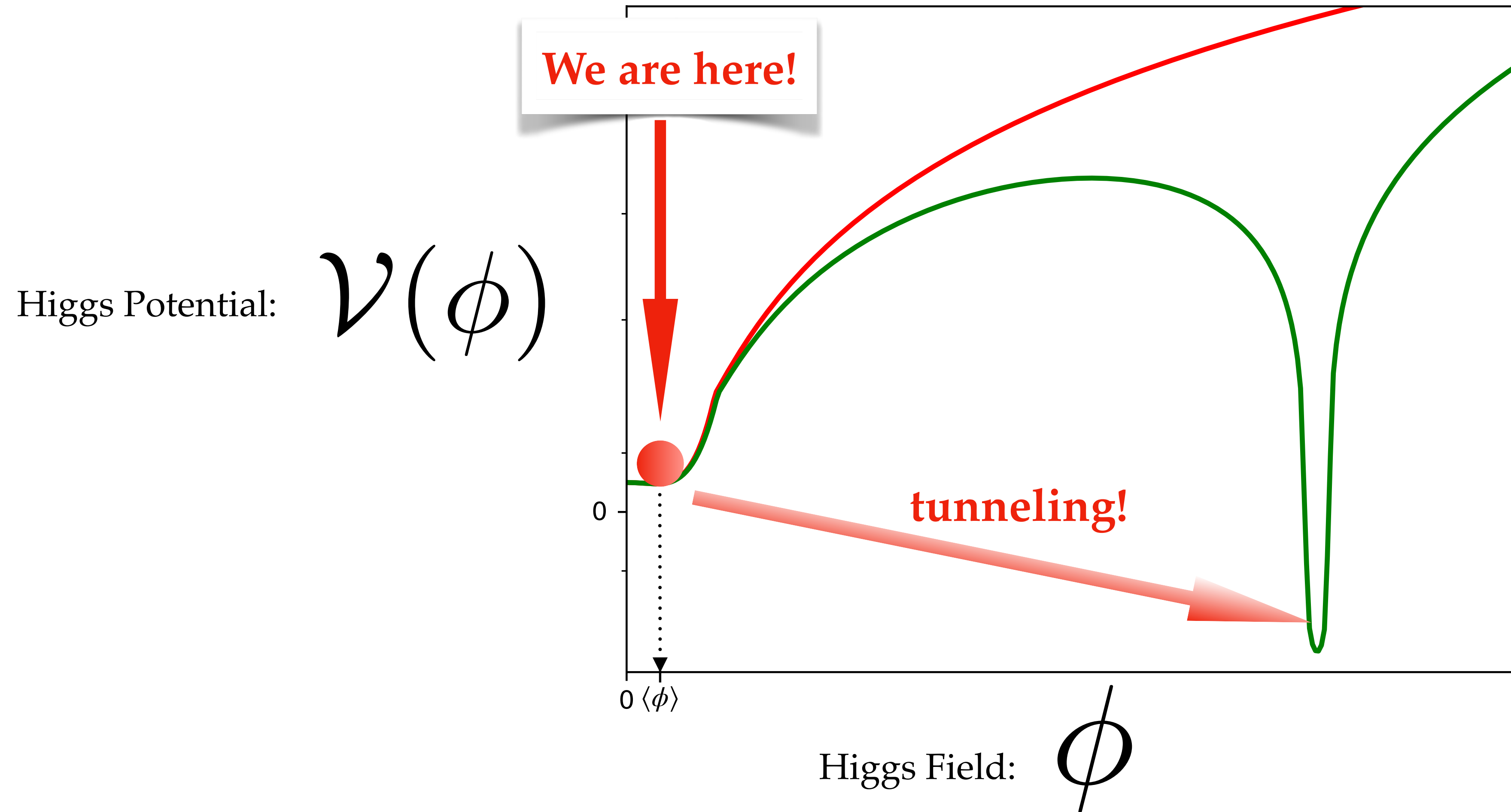
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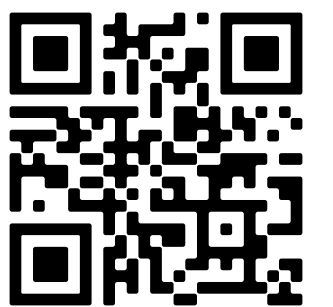
Q: What is the ultimate fate of our Universe?



Q: What is the ultimate fate of our Universe?



→ the “**Vacuum Stability Problem**”.



Q: Why does the matter we are made of exist?

[or: Why is there so much more matter than anti-matter?]



[credit: QuantaMagazine]

Q: Why does the matter we are made of exist?

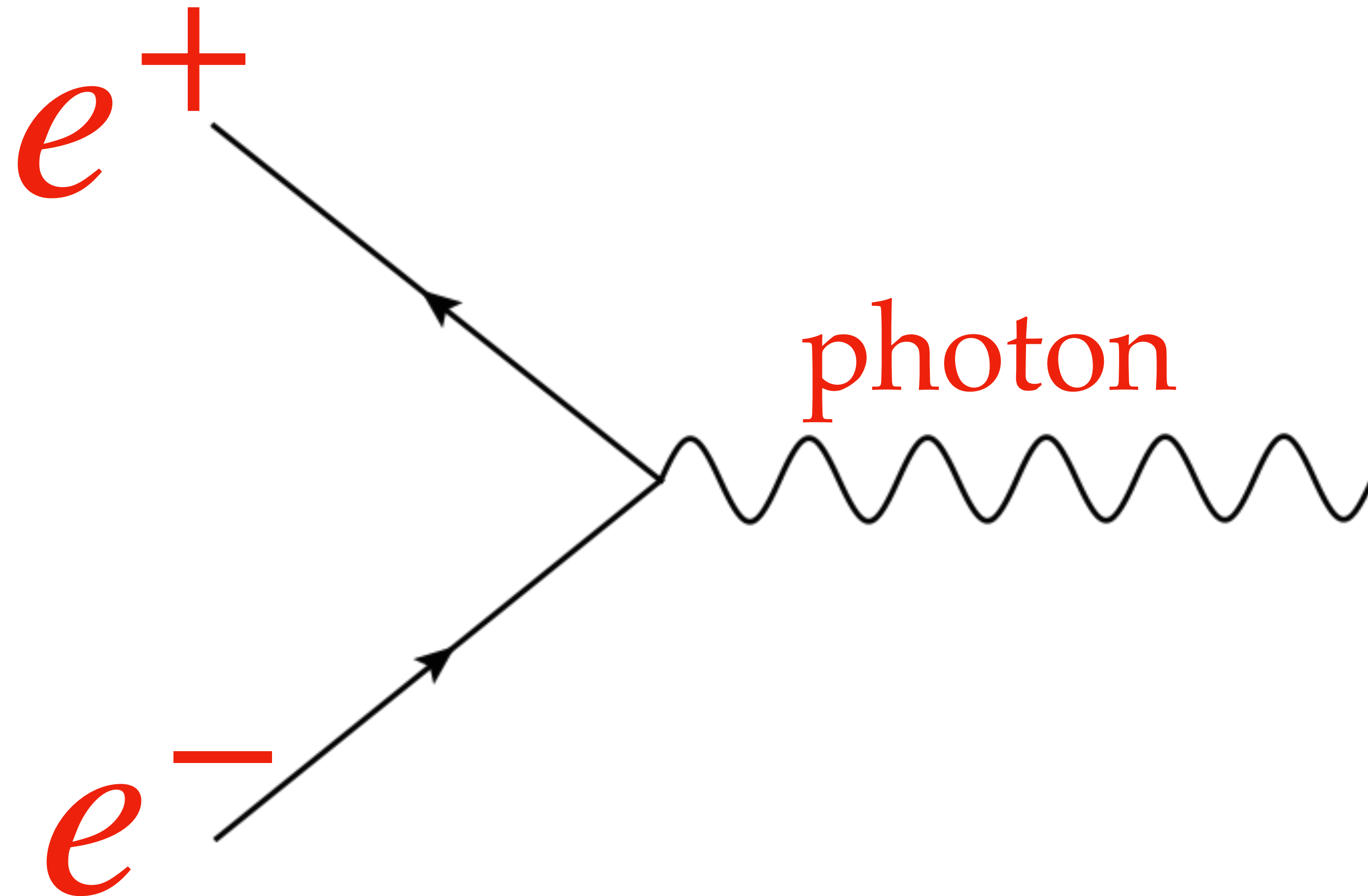
[or: Why is there so much more matter than anti-matter?]



[credit: QuantaMagazine]

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[or: Why is there so much more matter than anti-matter?]



Q: Why does the matter we are made of exist?

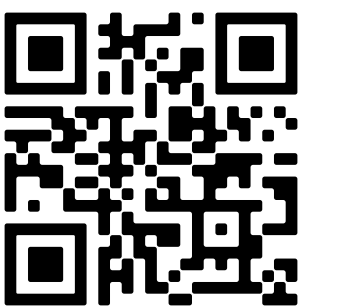
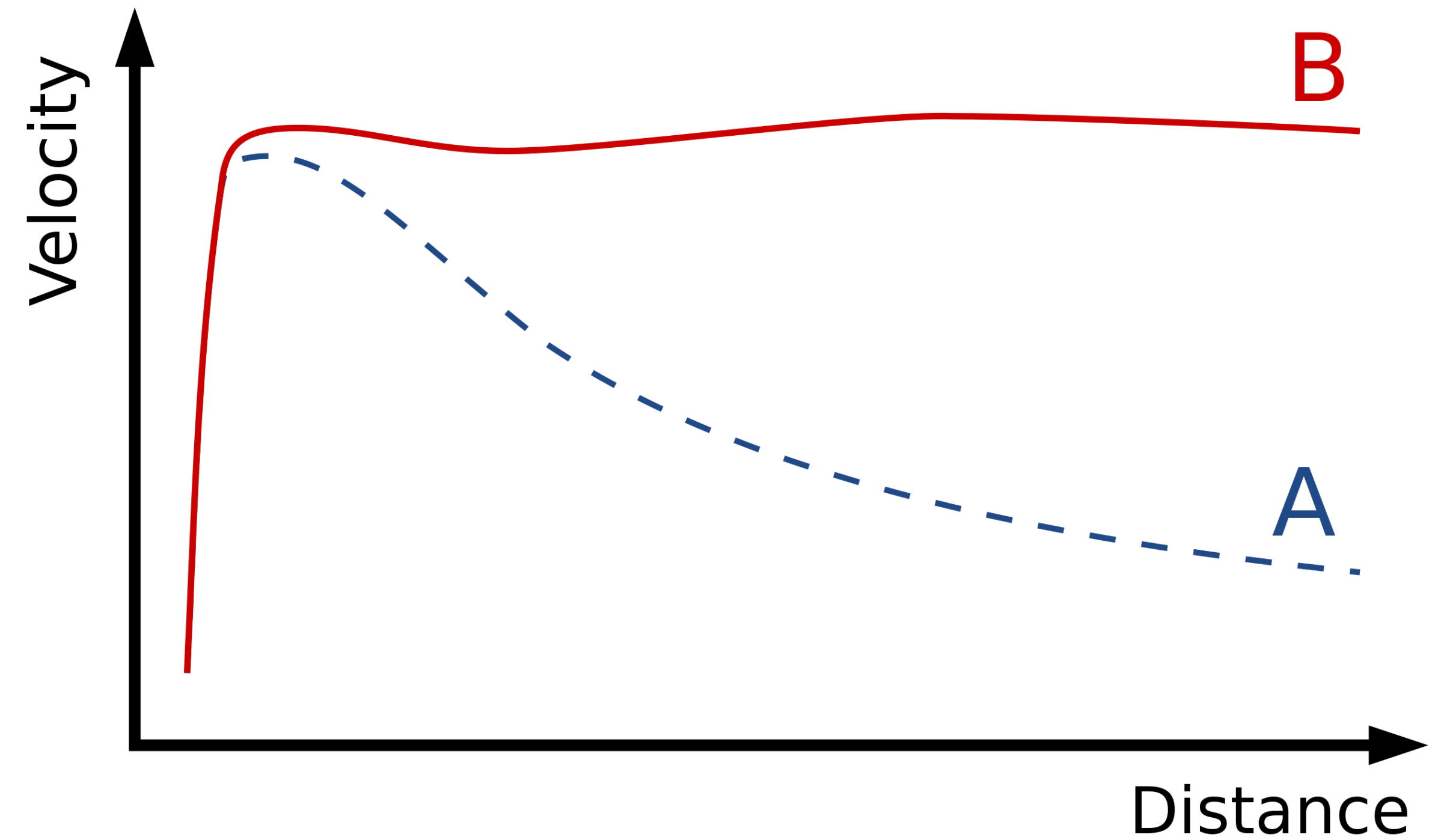
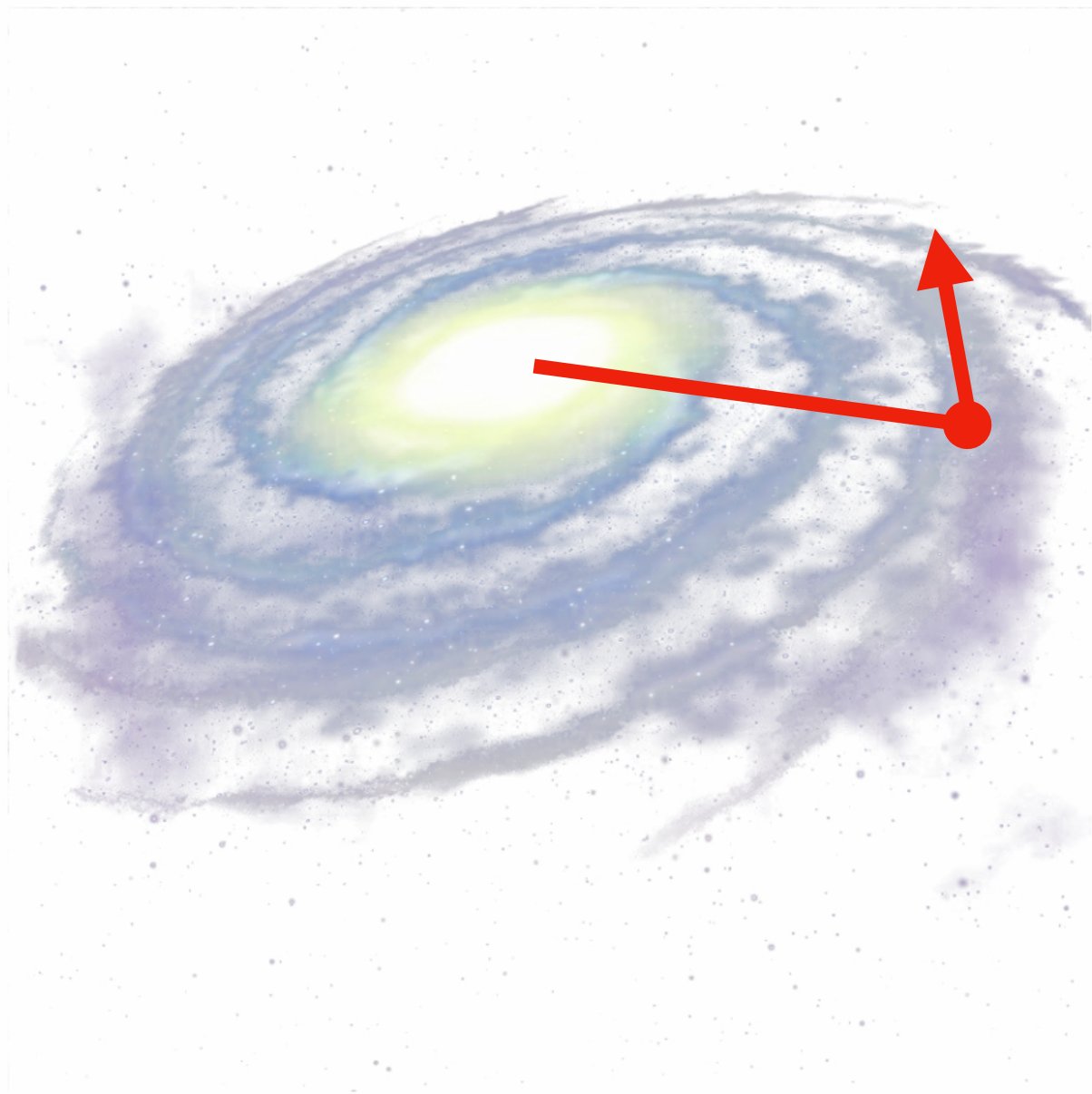
[or: Why is there so much more matter than anti-matter?]



→ the **“Matter-Anti-Matter Asymmetry”**.

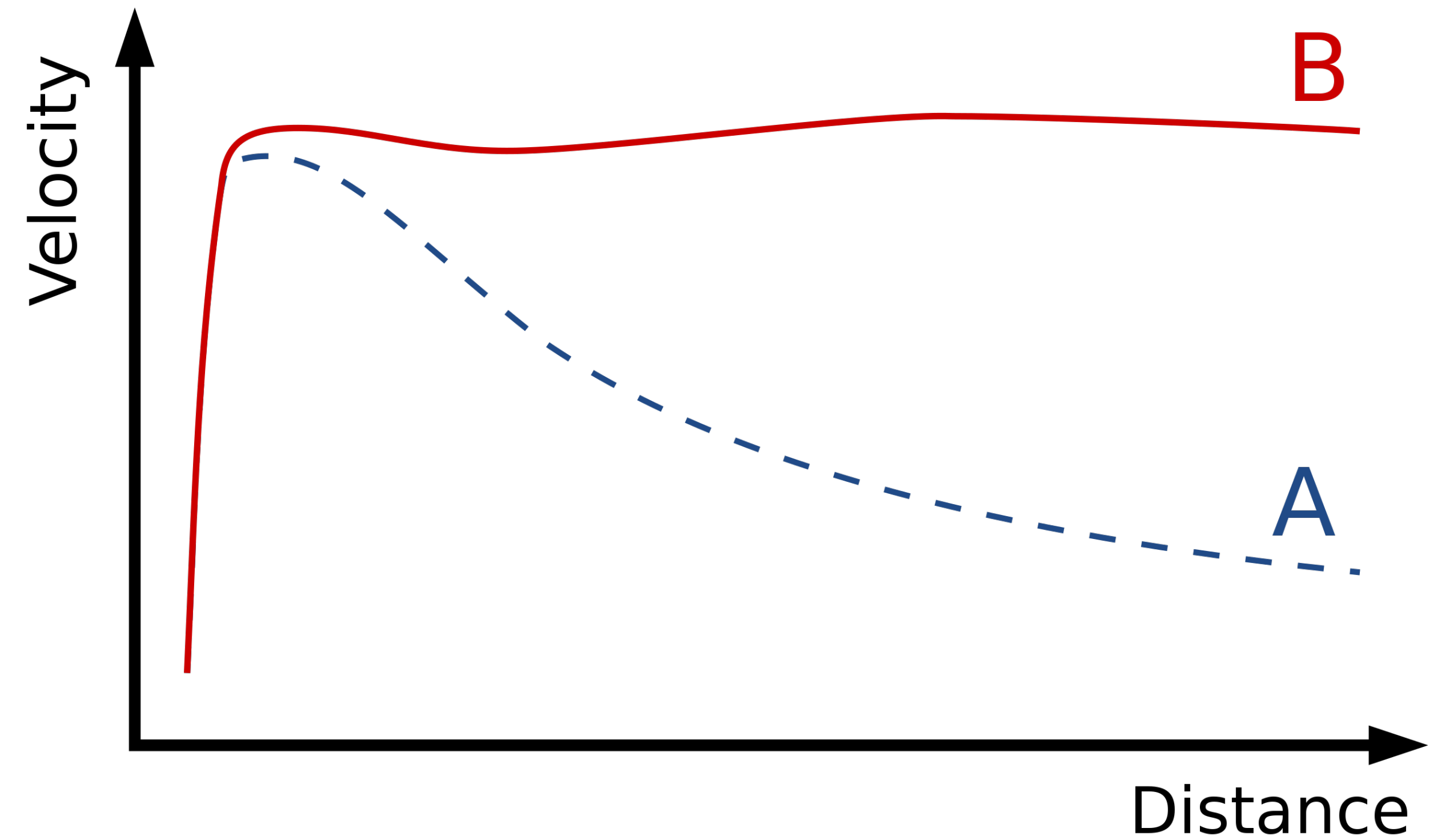
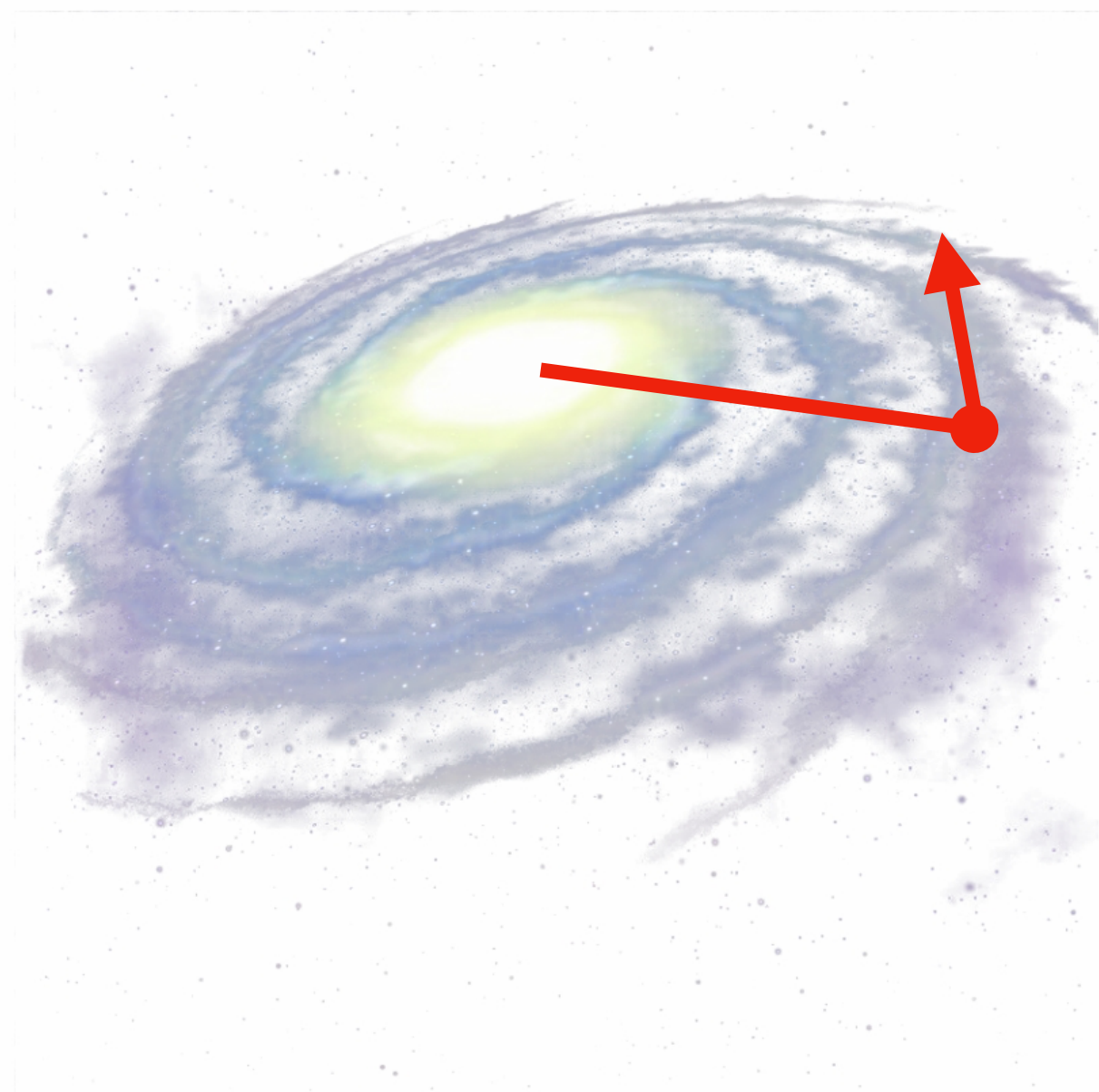
Q: What is the nature of **Dark Matter**?

Rotation curve of a typical spiral galaxy: predicted (A) and observed (B).

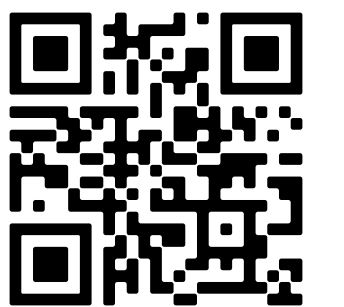


Q: What is the nature of **Dark Matter**?

Rotation curve of a typical spiral galaxy: predicted (A) and observed (B).



Q: ...





Dark Matter

Vacuum Stability

**Matter-anti-matter
Asymmetry**

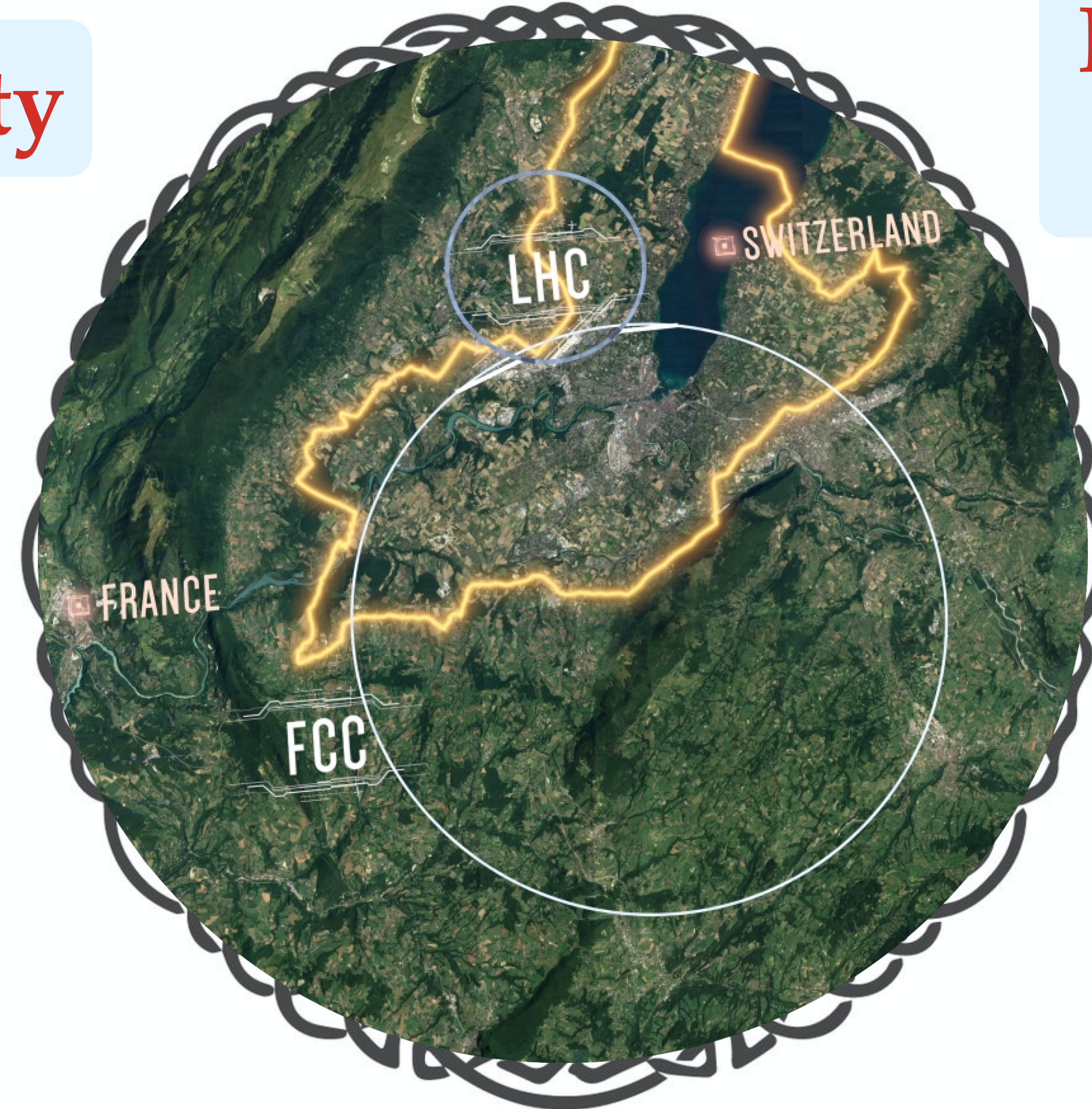


The Higgs Field & Symmetry Breaking



Dark Matter

Vacuum Stability



Matter-anti-matter
Asymmetry

e.g. Future Circular Collider:
 $pp@100 \text{ TeV}, e^+e^-$.

e.g. "High-Energy" LHC:
 $pp@27 \text{ TeV}$.

e.g. Muon Collider.

The Higgs Field & Symmetry Breaking



3

Breaking the Symmetry



Breaking the Symmetry

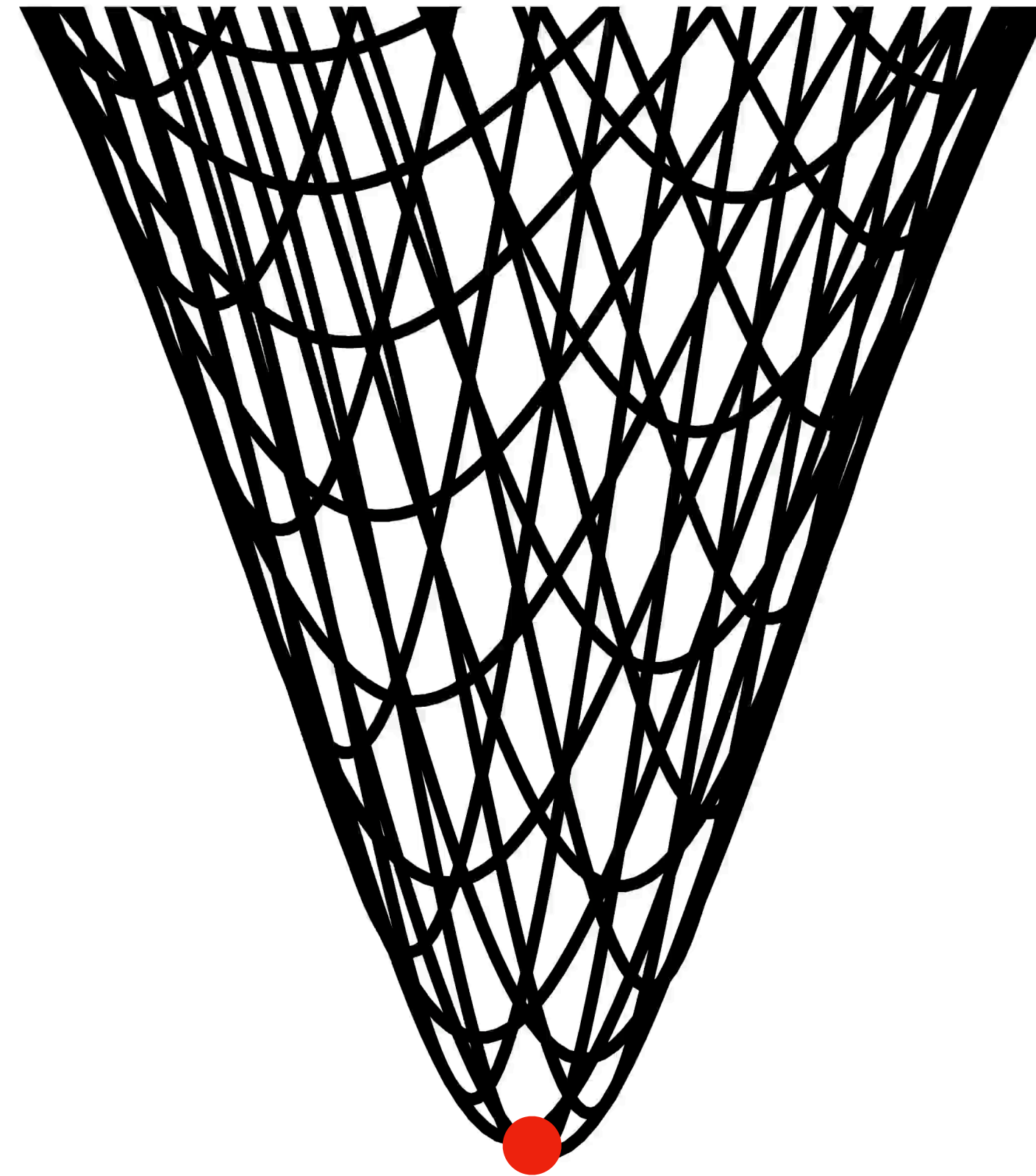


Breaking the Symmetry in the Standard Model

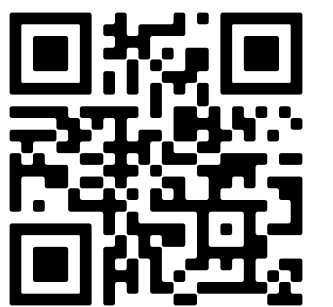
The potential of the Higgs field (ϕ), a complex doublet:

$$\phi = \begin{pmatrix} \phi_1 + i\phi_3 \\ \phi_2 + i\phi_4 \end{pmatrix}$$

(Arbitrarily) Set $\phi_3 = \phi_4 = 0$
to illustrate potential in
 (ϕ_1, ϕ_2) plane.



$\mathcal{V}(\phi)$

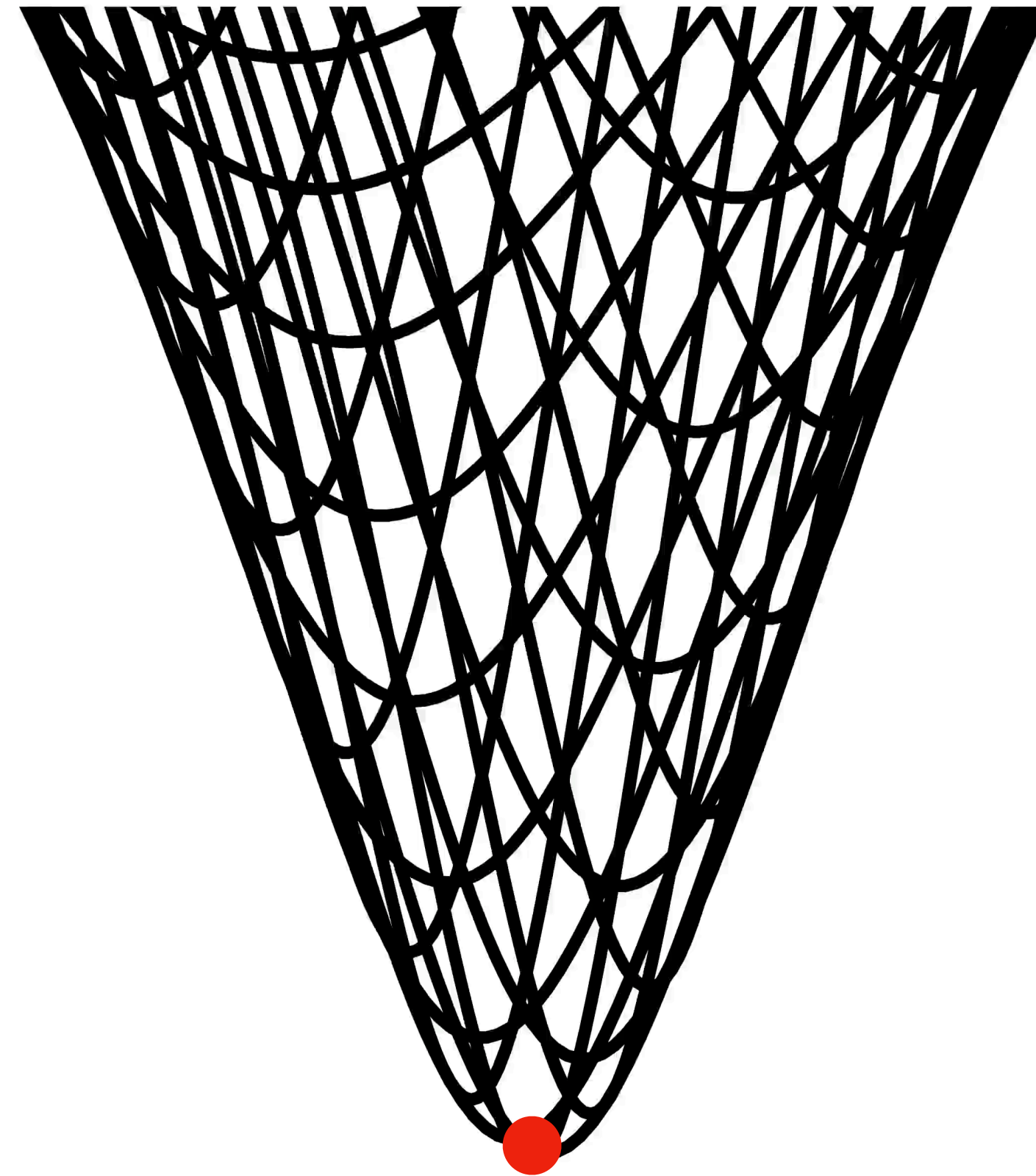


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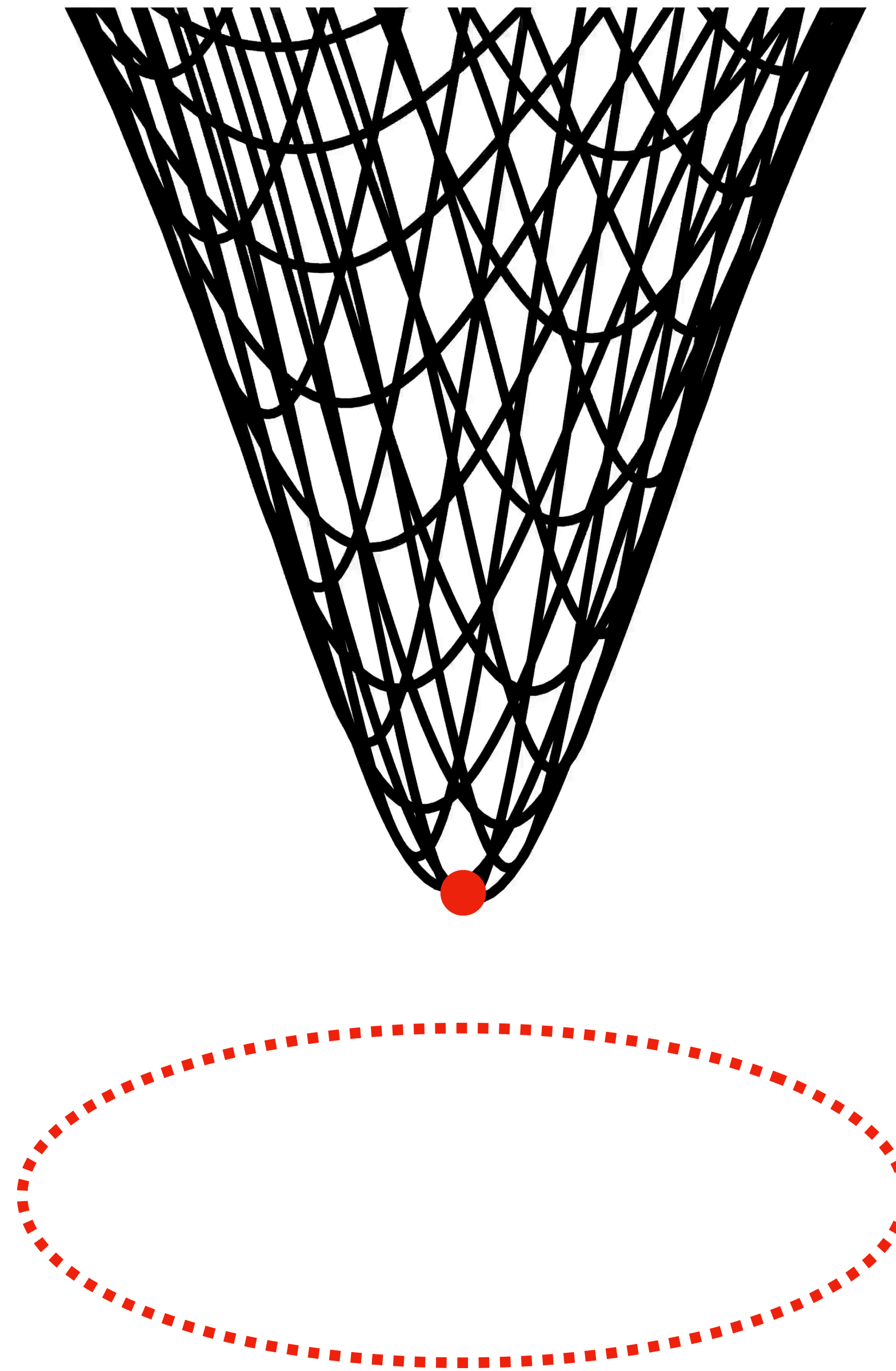


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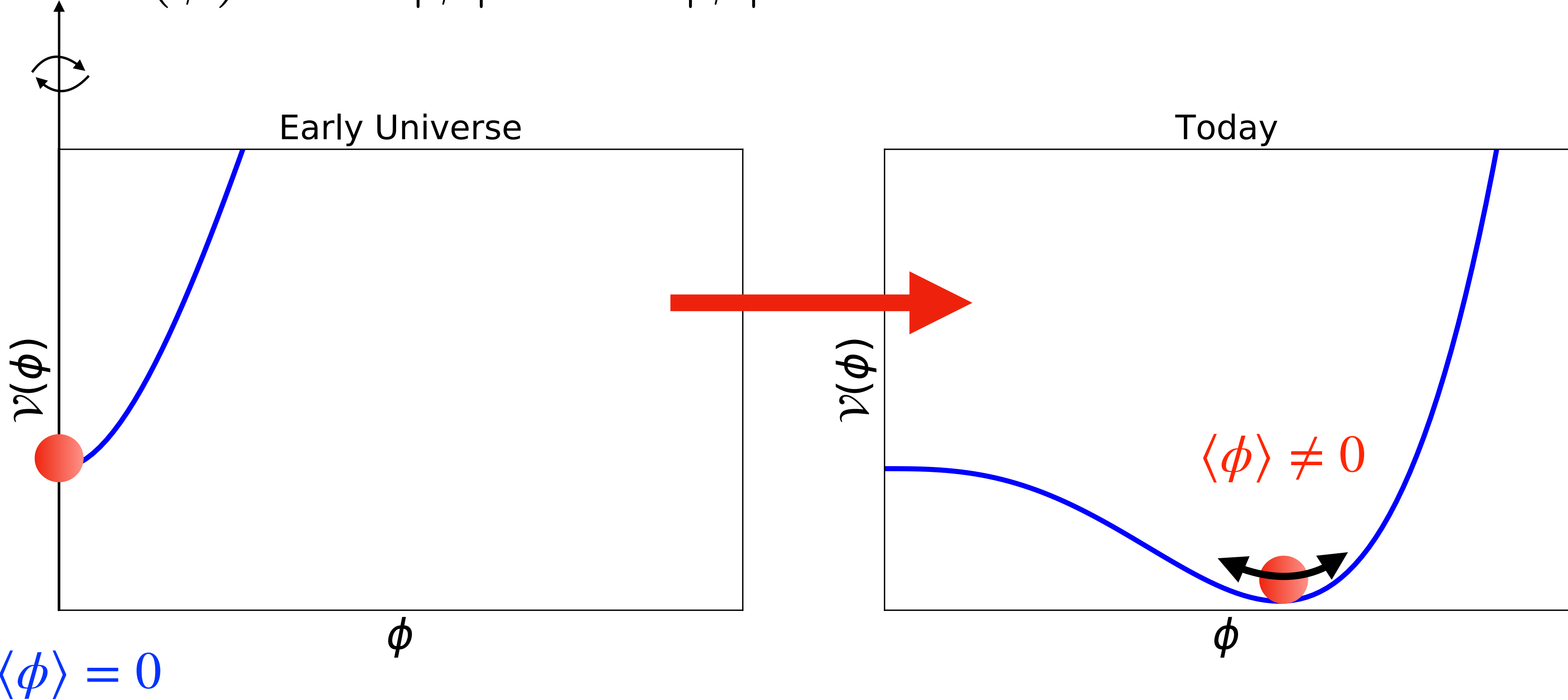
$\mathcal{V}(\phi)$



Breaking the Symmetry in the Standard Model

$$\mathcal{V}(\phi) = \text{●} |\phi|^2 + \text{■} |\phi|^4$$

Higgs field (ϕ) potential

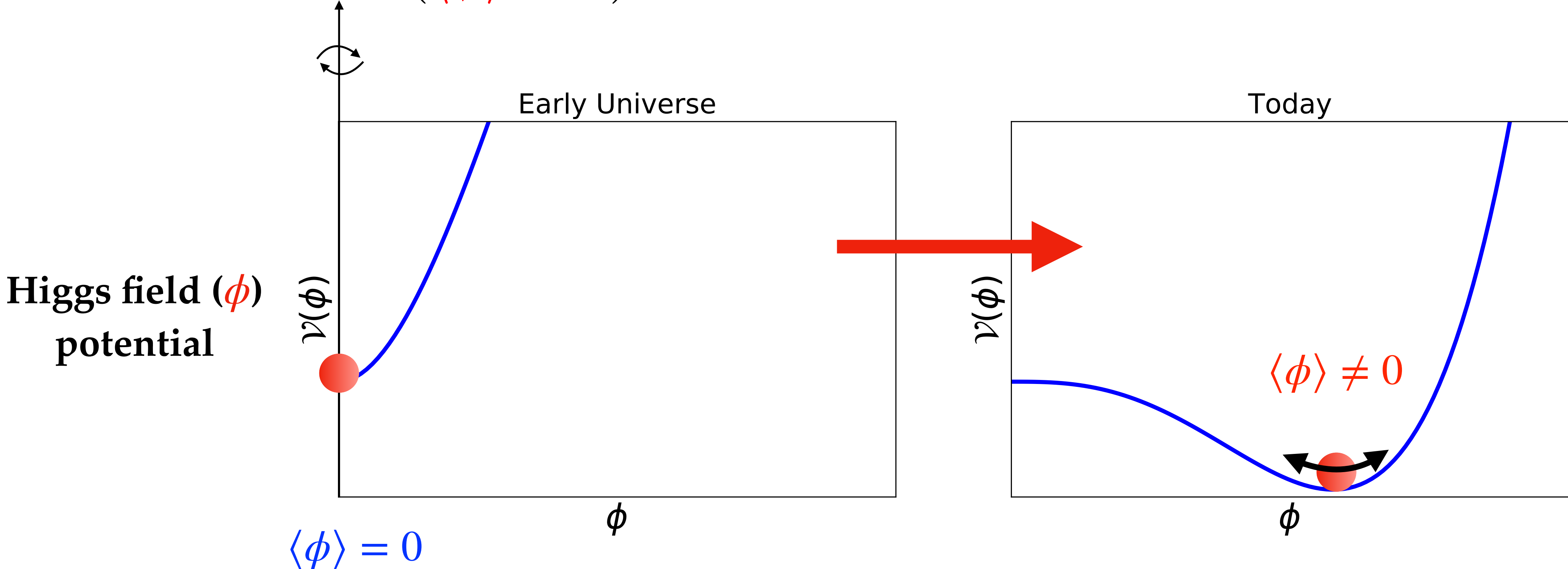


More Symmetry \rightarrow Less (obvious) Symmetry



Breaking the Symmetry in the Standard Model

$$\mathcal{V}(\langle\phi\rangle + h) = \bullet h^2 + \blacktriangle h^3 + \blacksquare h^4 \rightarrow h \text{ is the Higgs boson! (LHC, 2012)}$$



More Symmetry \rightarrow Less (obvious) Symmetry

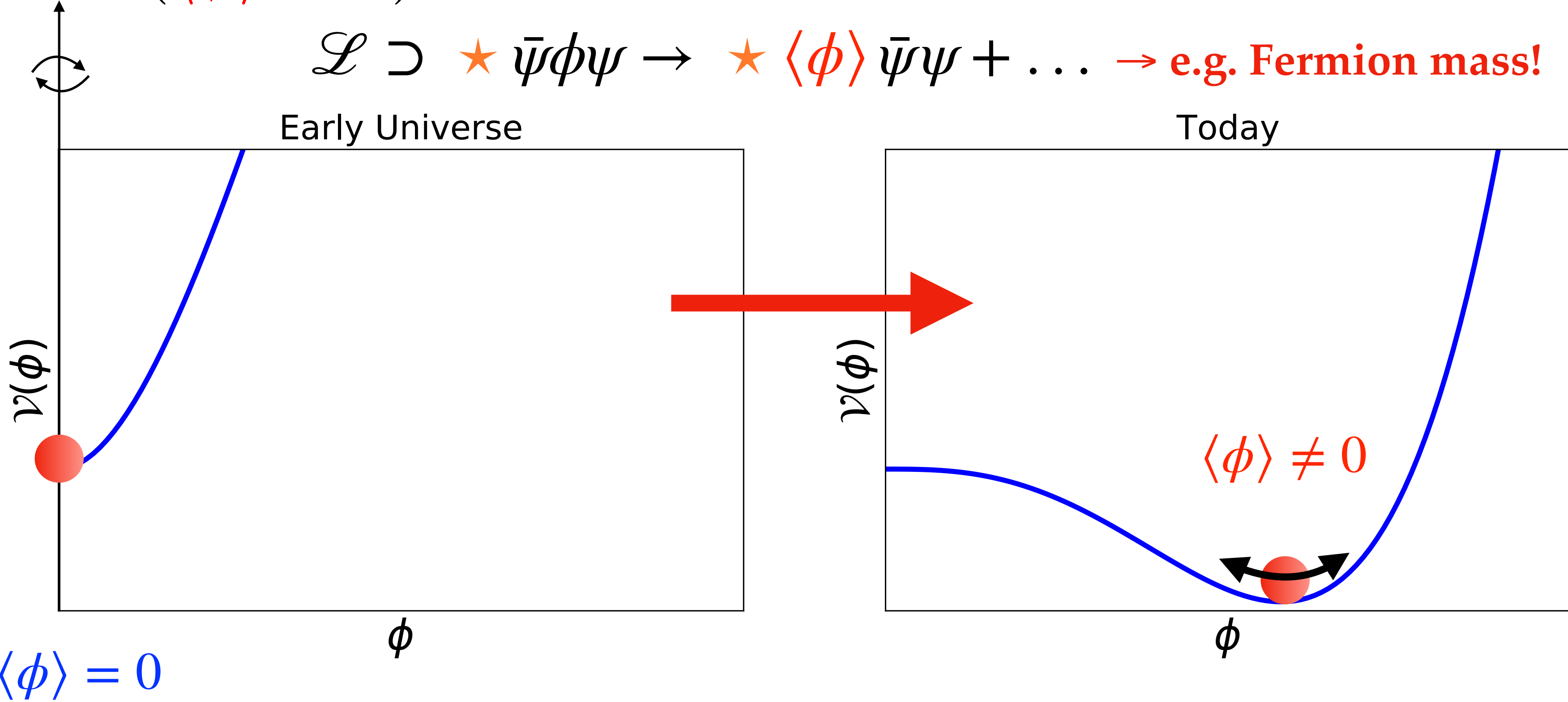


Breaking the Symmetry in the Standard Model

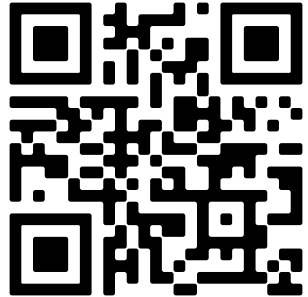
$\mathcal{V}(\langle\phi\rangle + h) = \bullet h^2 + \blacktriangle h^3 + \blacksquare h^4 \rightarrow h$ is the Higgs boson! (LHC, 2012)

$\mathcal{L} \supset \star \bar{\psi}\phi\psi \rightarrow \star \langle\phi\rangle \bar{\psi}\psi + \dots \rightarrow$ e.g. Fermion mass!

Higgs field (ϕ) potential



More Symmetry \rightarrow Less (obvious) Symmetry

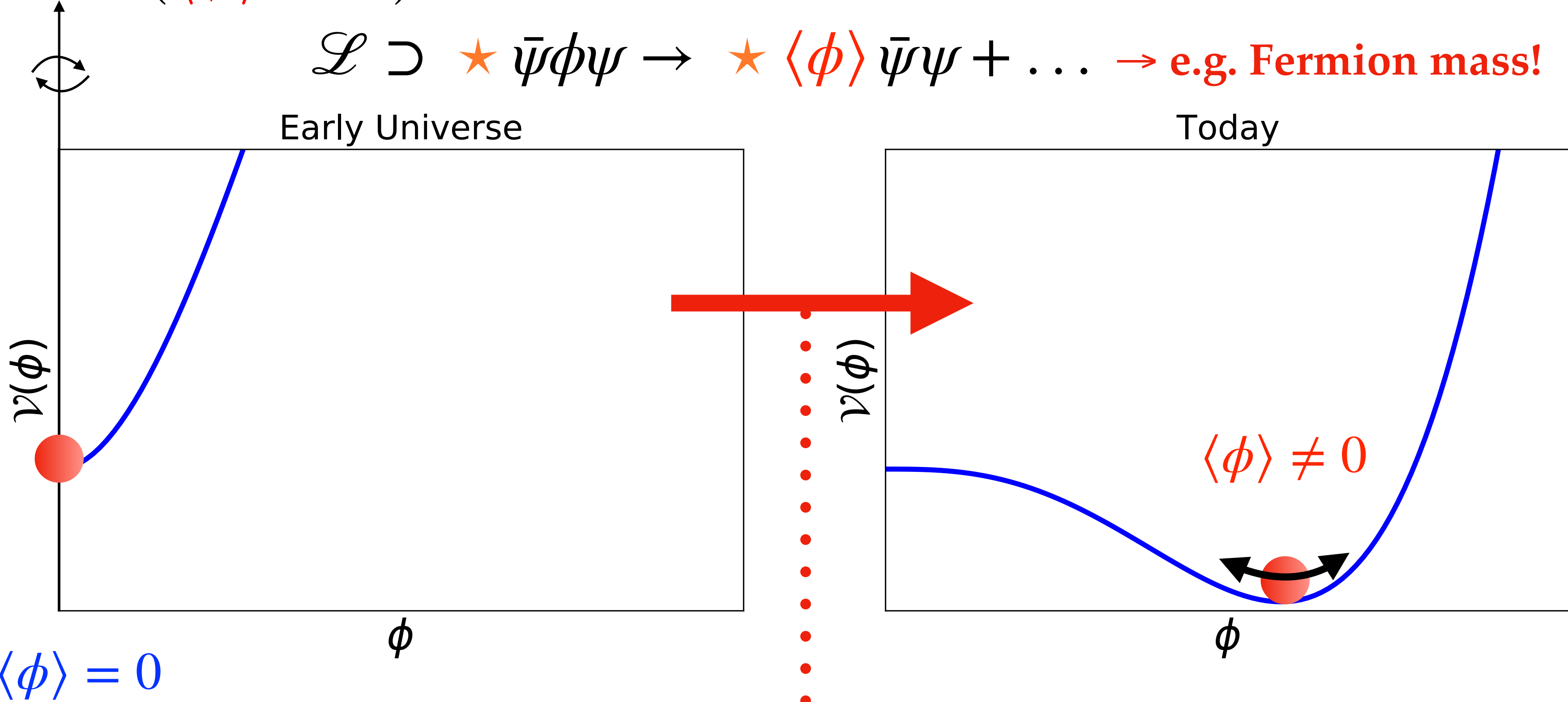


Breaking the Symmetry in the Standard Model

$\mathcal{V}(\langle\phi\rangle + h) = \bullet h^2 + \blacktriangle h^3 + \blacksquare h^4 \rightarrow h$ is the Higgs boson! (LHC, 2012)

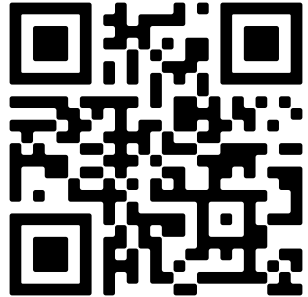
$\mathcal{L} \supset \star \bar{\psi}\phi\psi \rightarrow \star \langle\phi\rangle \bar{\psi}\psi + \dots \rightarrow$ e.g. Fermion mass!

Higgs field (ϕ) potential



More Symmetry \rightarrow Less (obvious) Symmetry

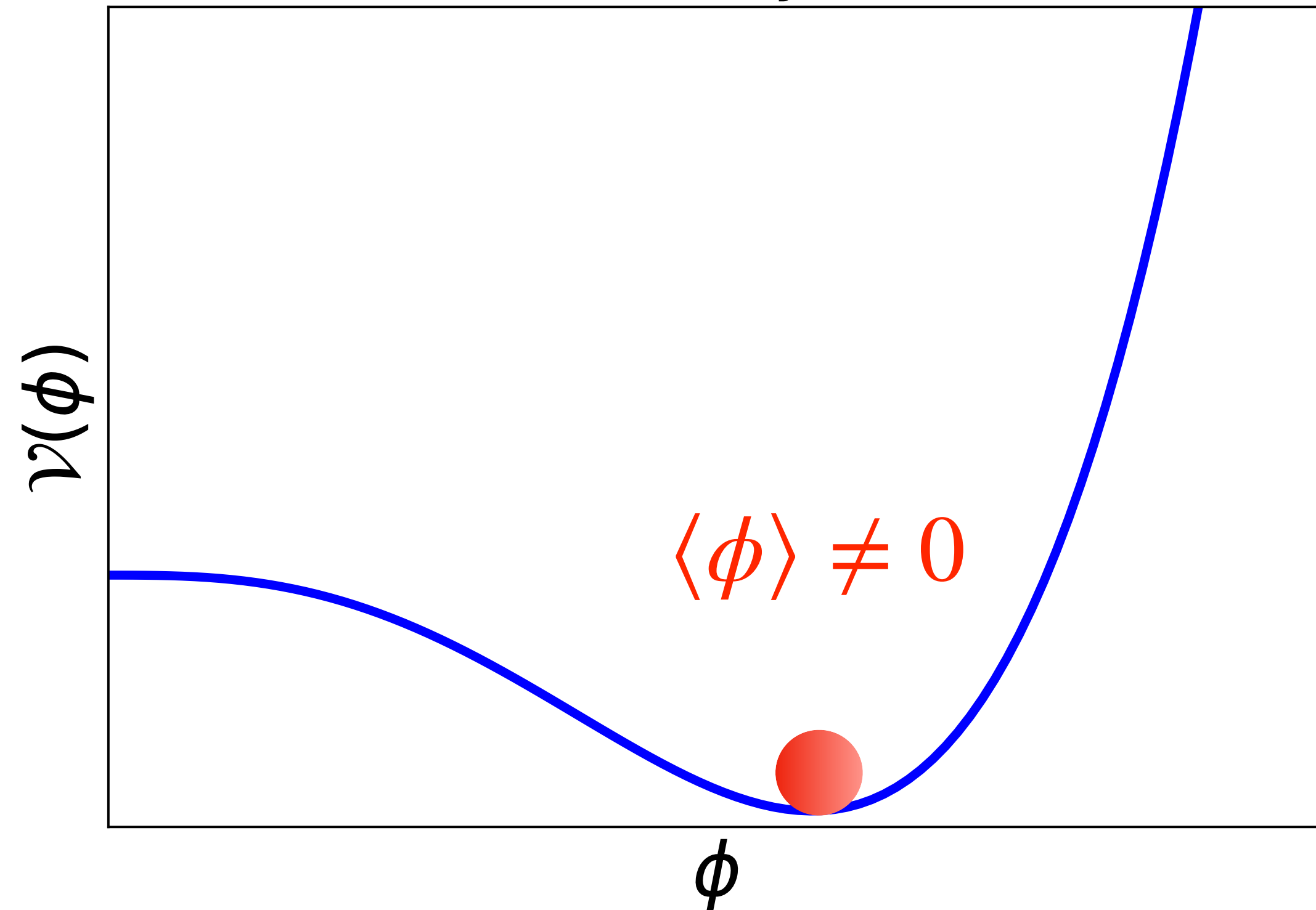
\equiv Electro-Weak Phase Transition (EWPT)



The Higgs Boson's Potential

$$\mathcal{V}(\langle\phi\rangle + h) = \bullet h^2 + \blacktriangle h^3 + \blacksquare h^4 \rightarrow \text{the Higgs boson's self-interactions.}$$

Today



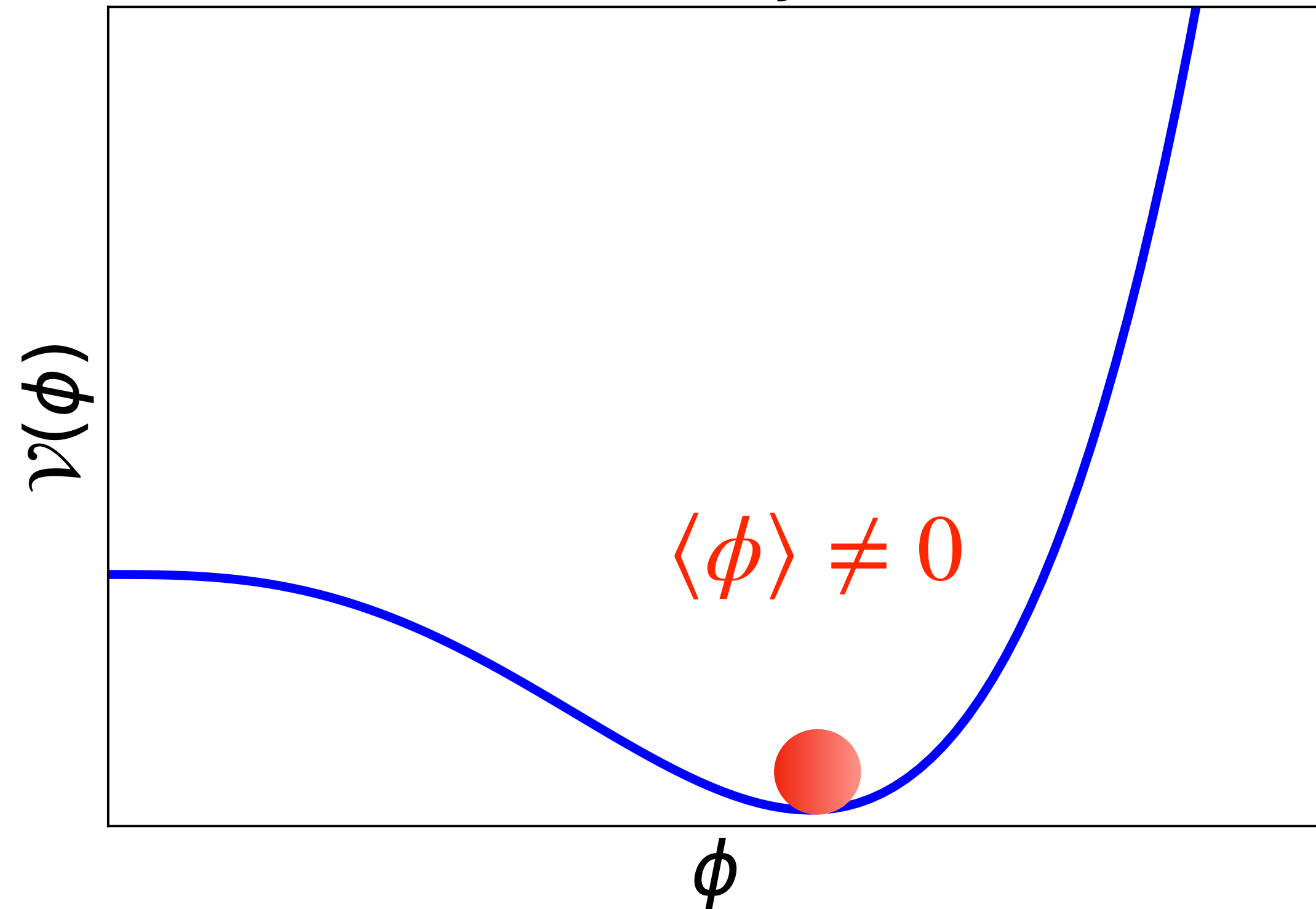
→ Determine shape of potential by measuring:

$$\{ \bullet, \blacktriangle, \blacksquare \}$$

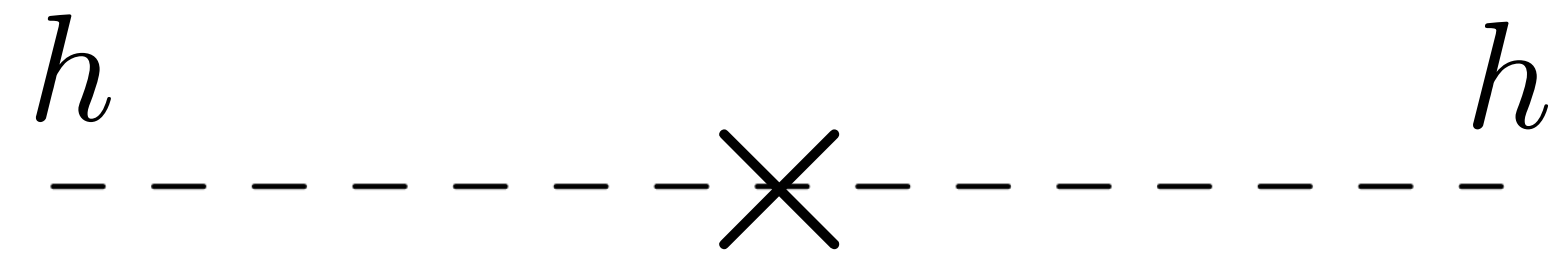
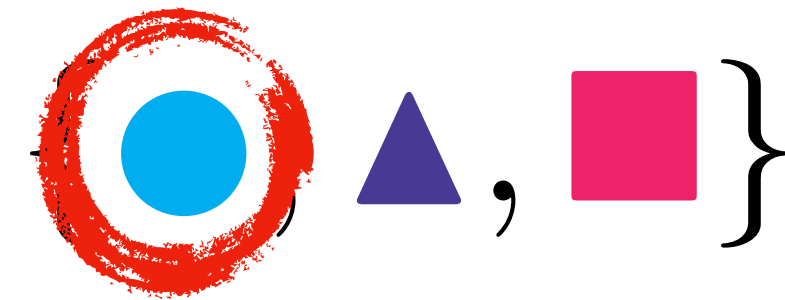
The Higgs Boson's Potential

$$\mathcal{V}(\langle\phi\rangle + h) = \bullet h^2 + \blacktriangle h^3 + \blacksquare h^4 \rightarrow \text{the Higgs boson's self-interactions.}$$

Today



→ Determine shape of potential by measuring:



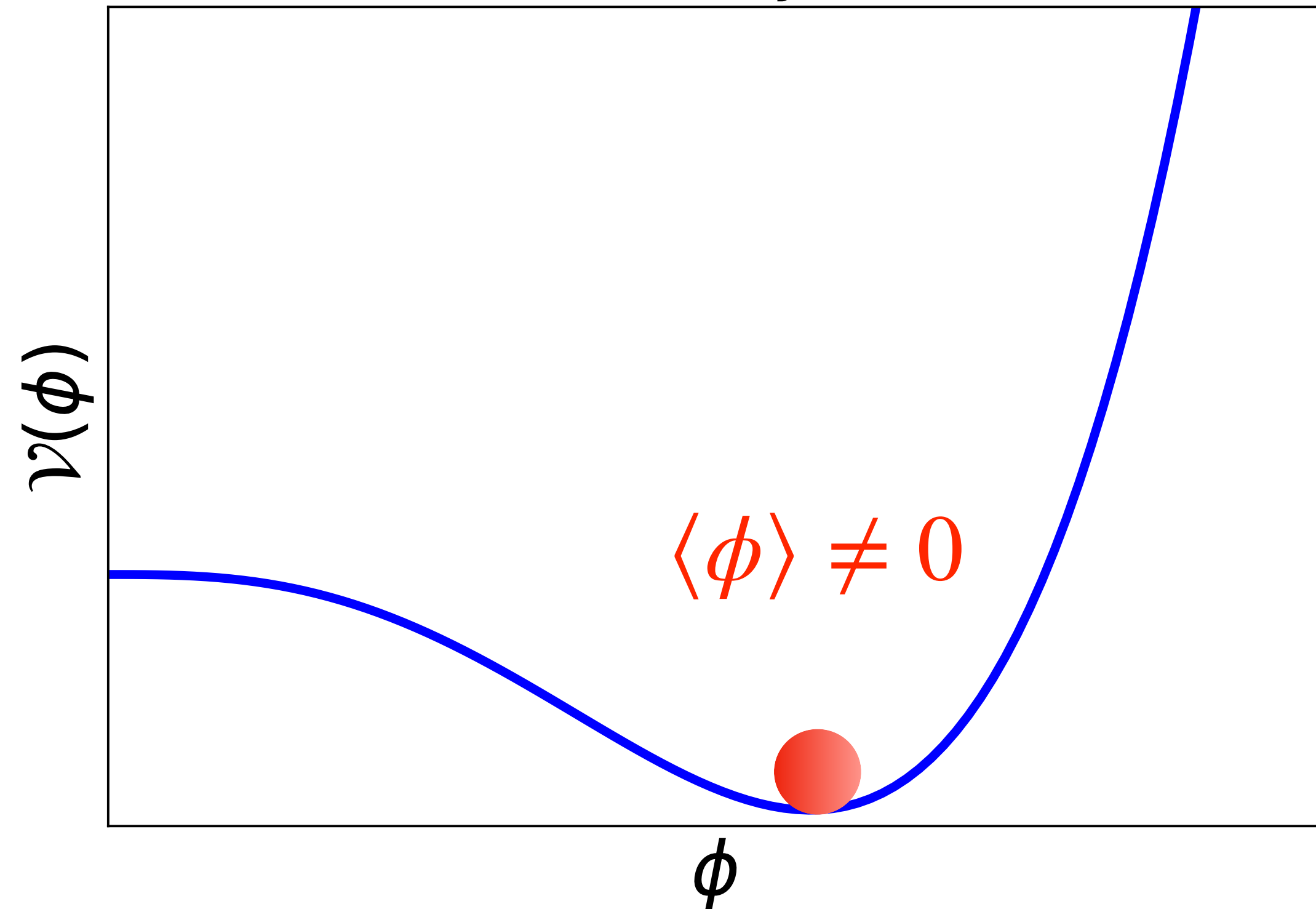
Higgs boson discovery @ LHC, 2012



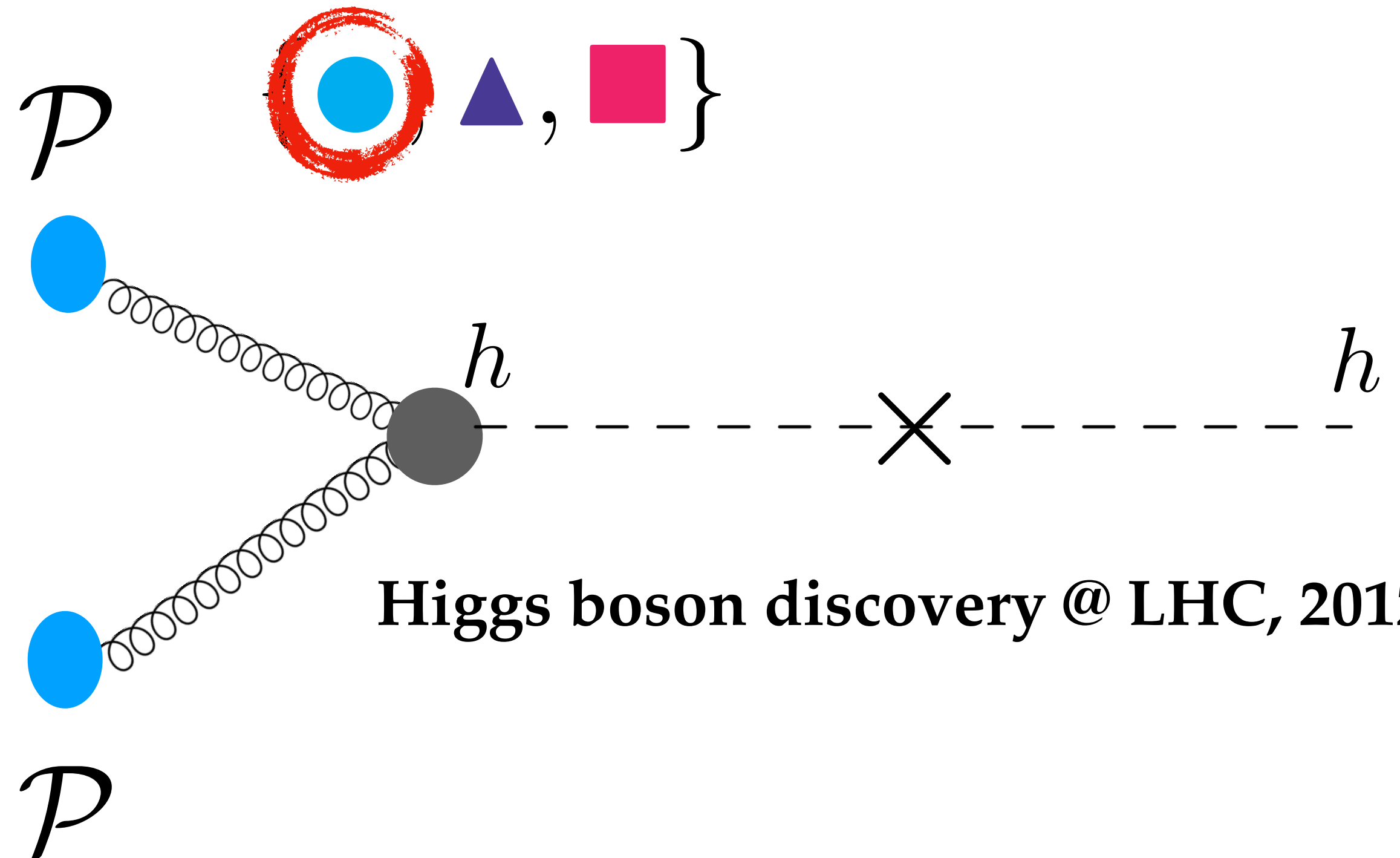
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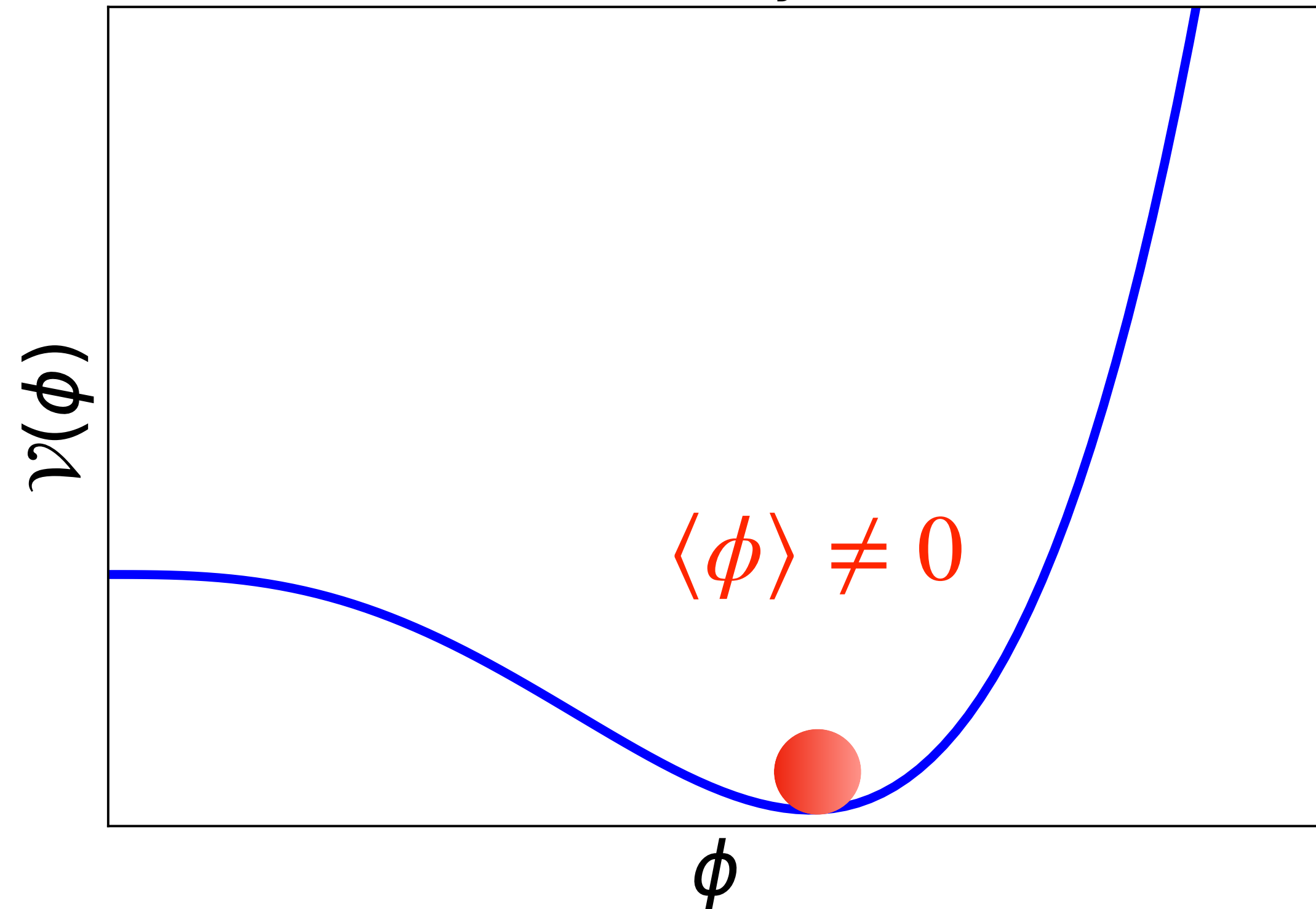
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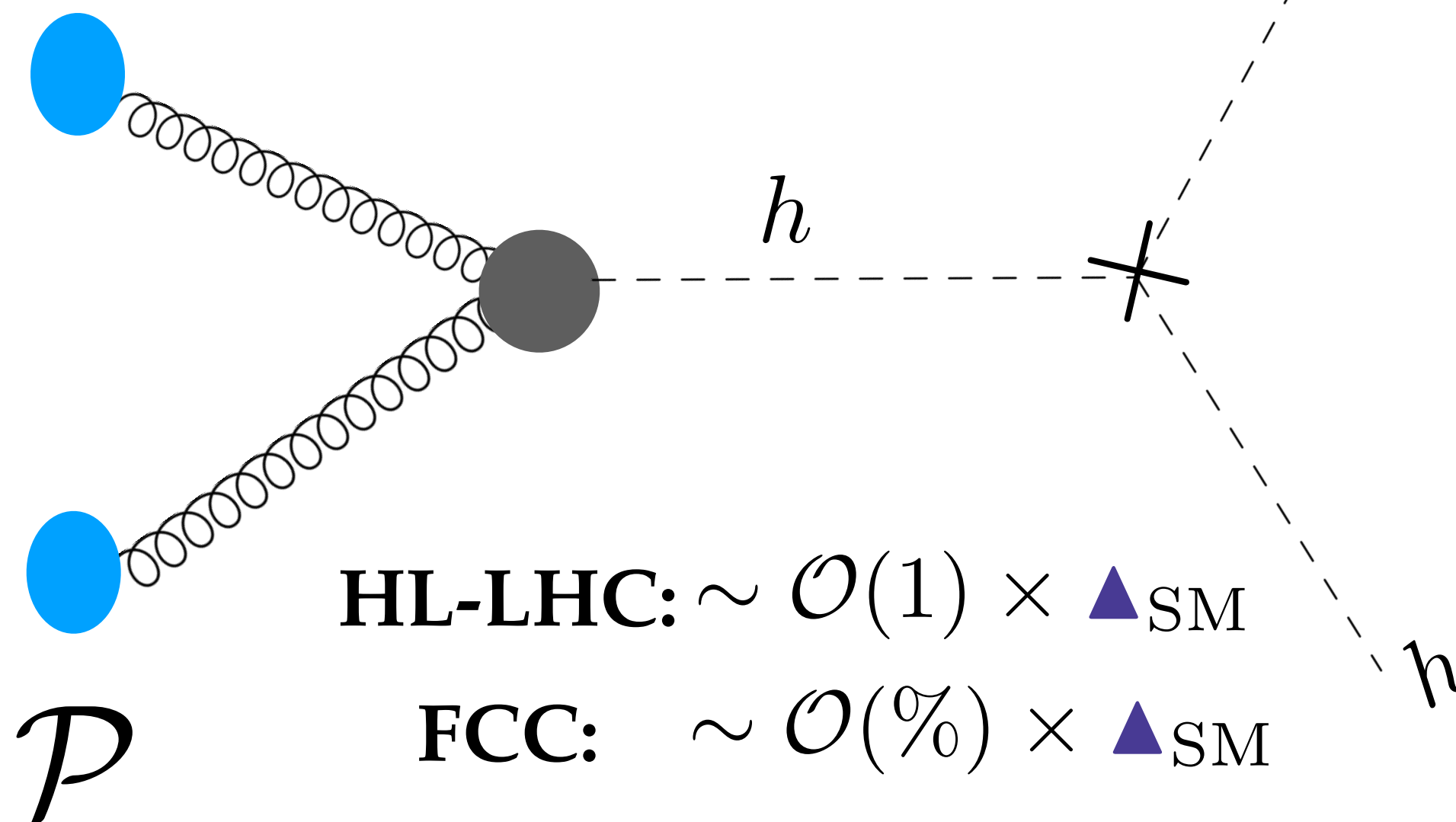
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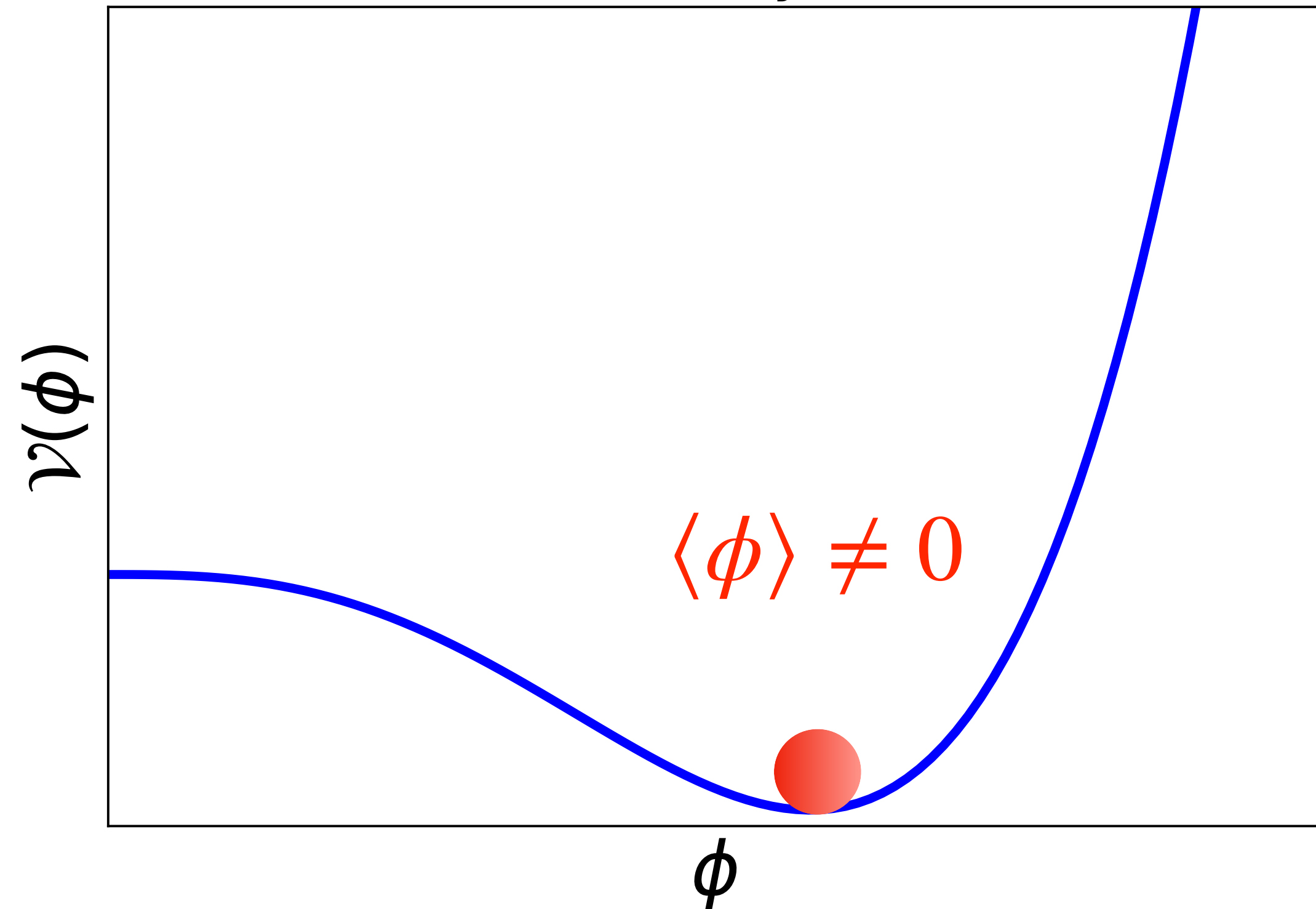
$$\mathcal{P} \quad \{ \bullet, \blacktriangle, \blacksquare \}$$



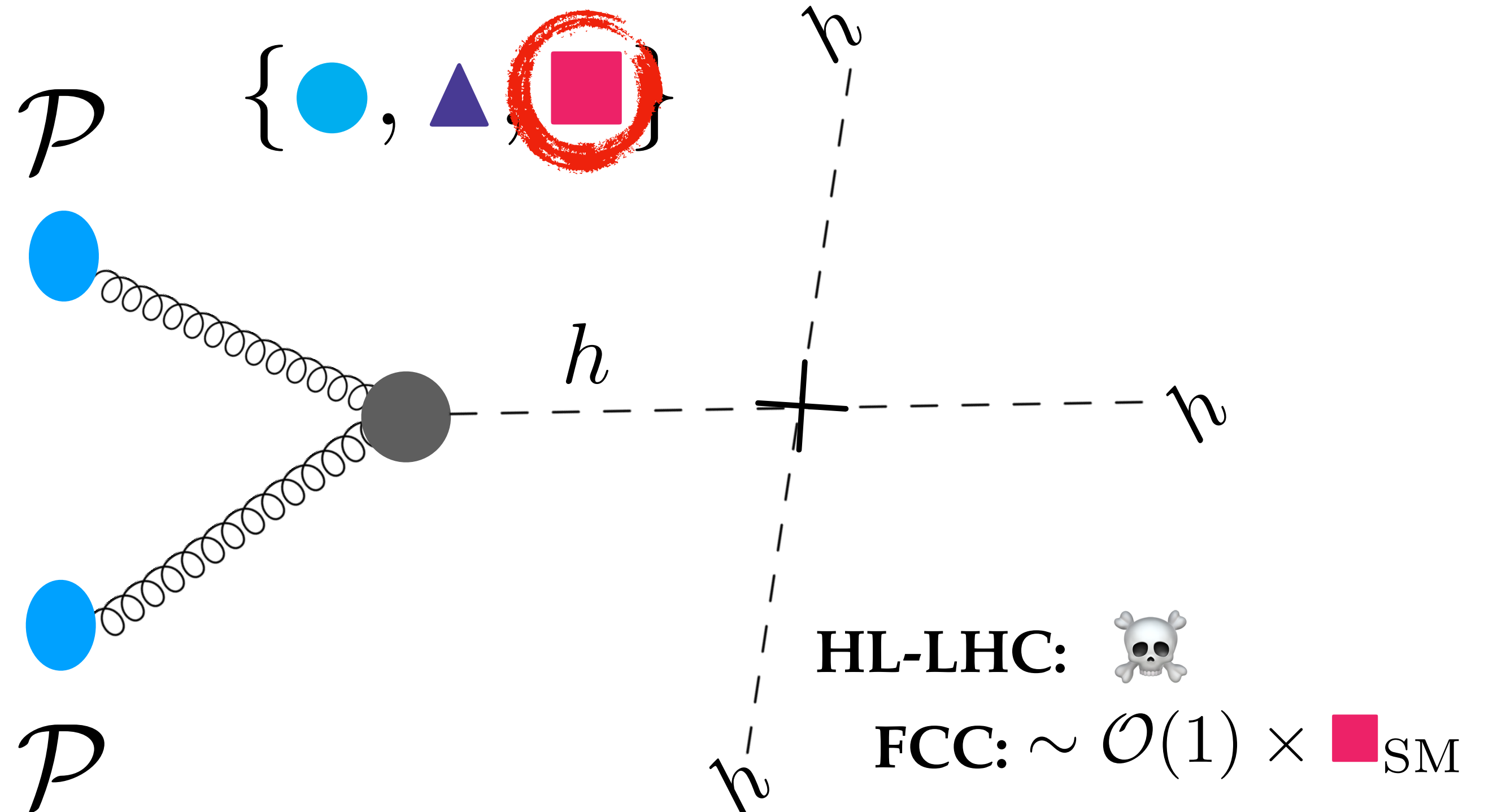
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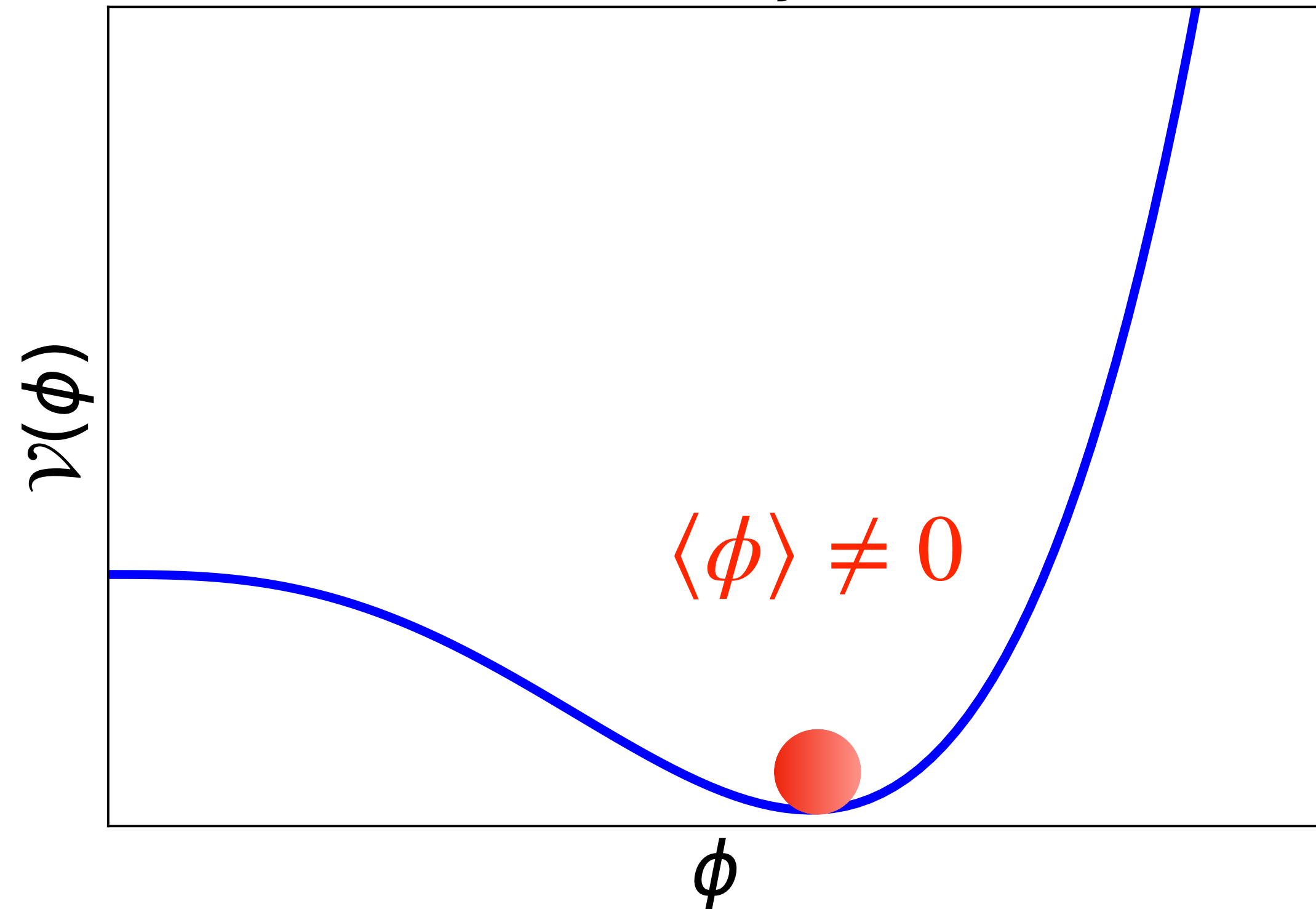


[e.g. [AP](#), Sakurai, arXiv:1508.06524, [AP](#),
Tetlalmatzi-Xolocotzi, Zaro,
arXiv:1909.09166]

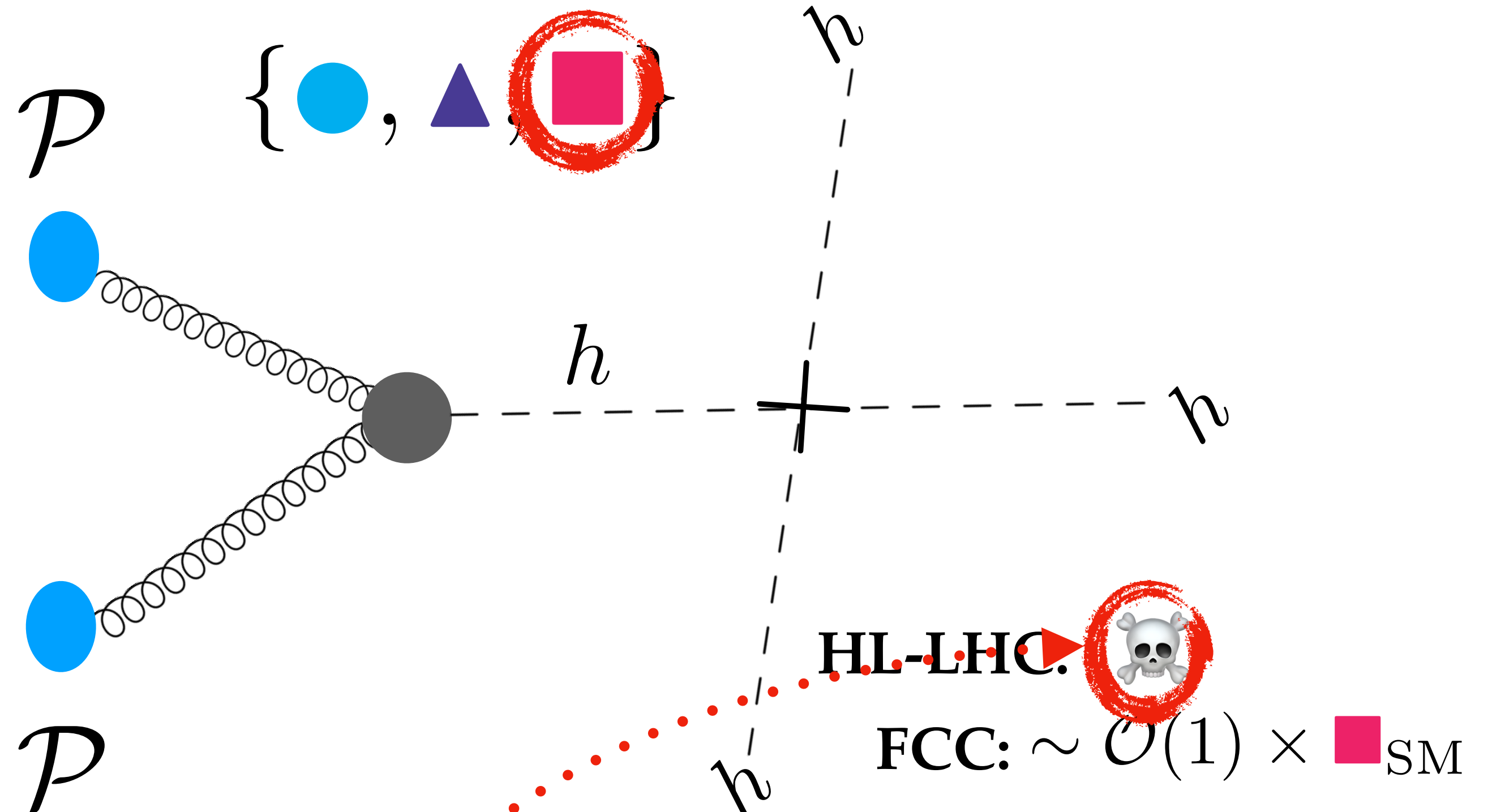
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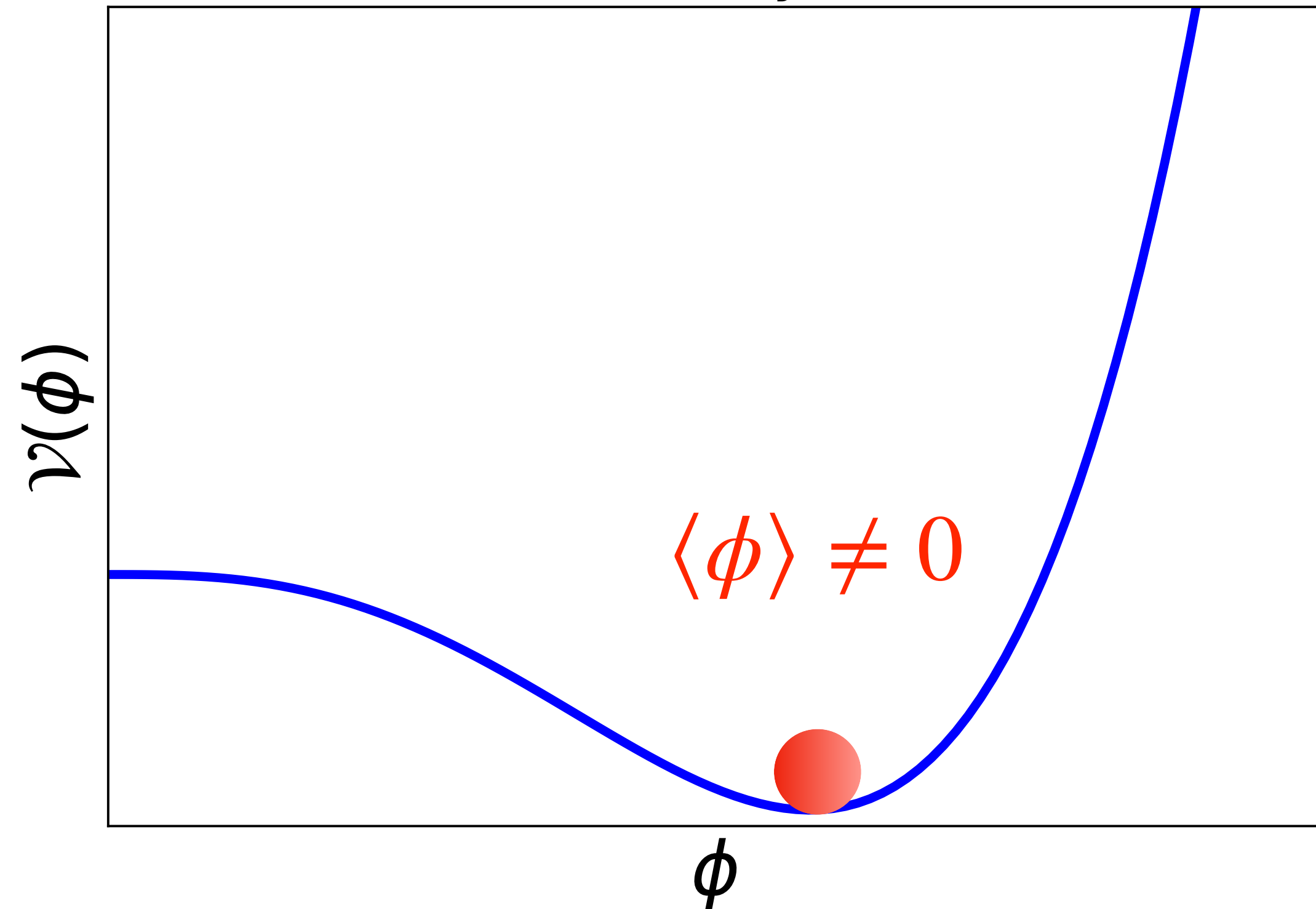
[[AP](#), Robens, Tetlalmatzi-Xolocotzi, 2101.00037 + in progress
with Osama Karkout, Carlo Pandini, [AP](#), Marieke Postma, Tristan du Pree,
Gilberto Tetlalmatzi-Xolocotzi, Jorinde van de Vis [...]]

[e.g. [AP](#), Sakurai, arXiv:1508.06524, [AP](#),
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Alexandra Carvalho, [AP](#), Marko Stamenkovic,
Gilberto Tetlalmatzi-Xolocotzi, Alberto Tonerio [...]

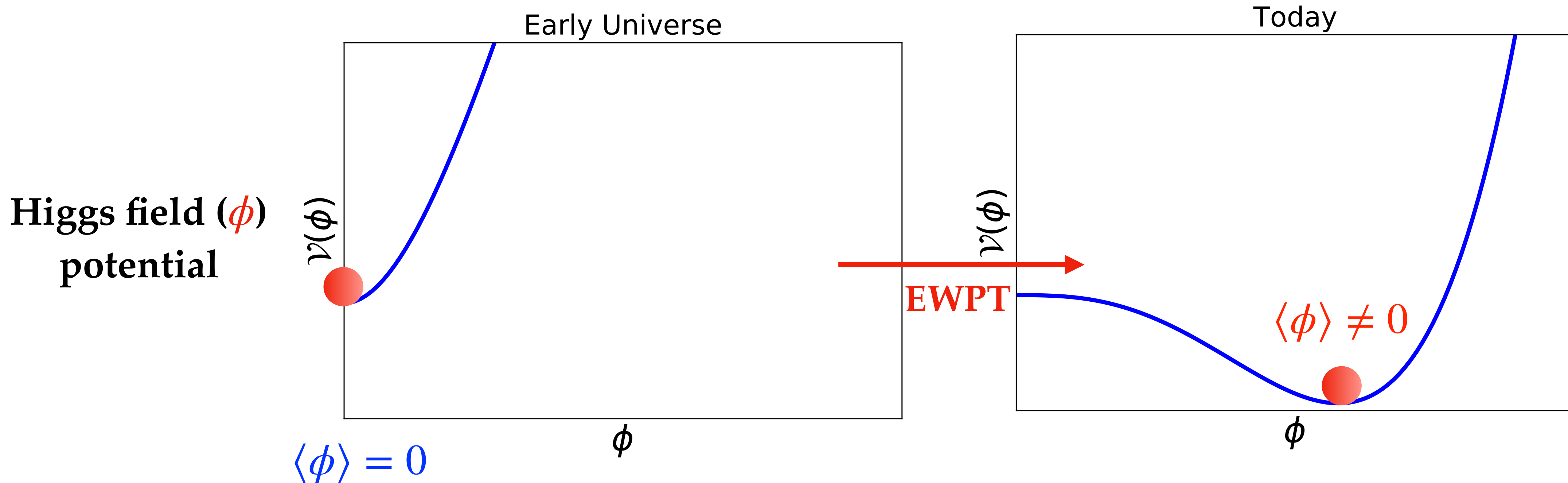
[hhh with Anomalous Couplings]

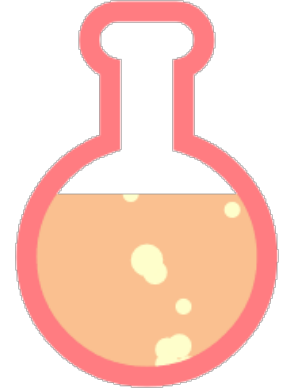
Osama Karkout, Carlo Pandini, [AP](#), Marieke Postma, Tristan du
Pree, Gilberto Tetlalmatzi-Xolocotzi, Jorinde van de Vis [...]

[hhh in TRSM + Cosmology]

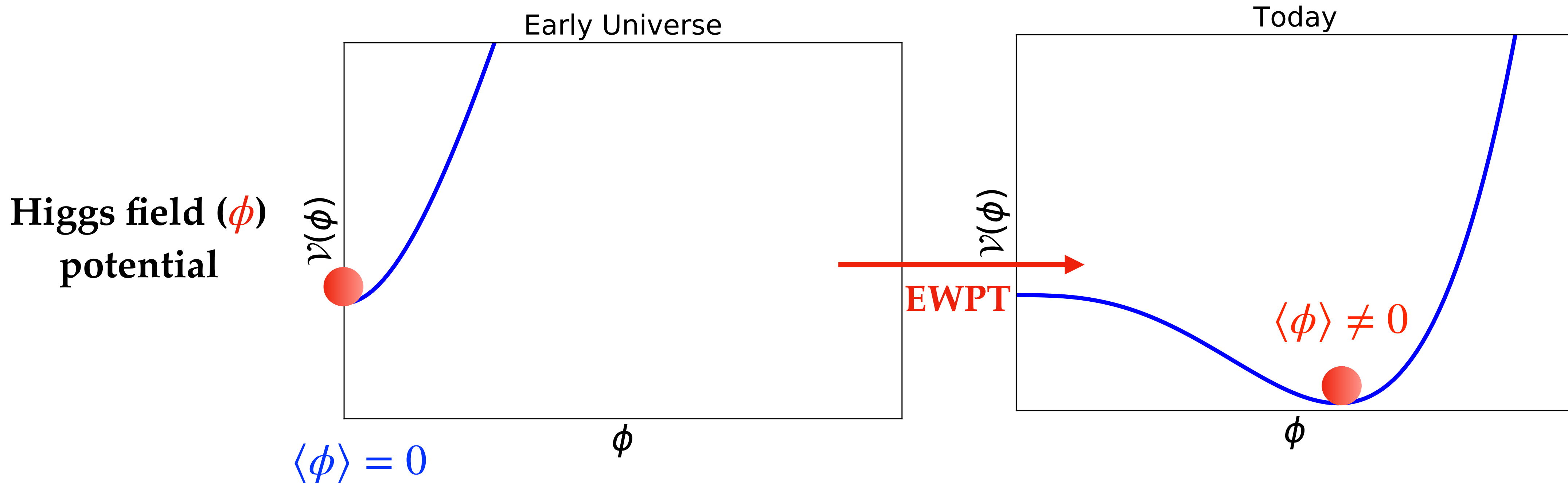


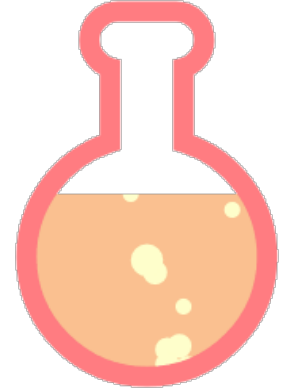
Breaking the Symmetry in the SM



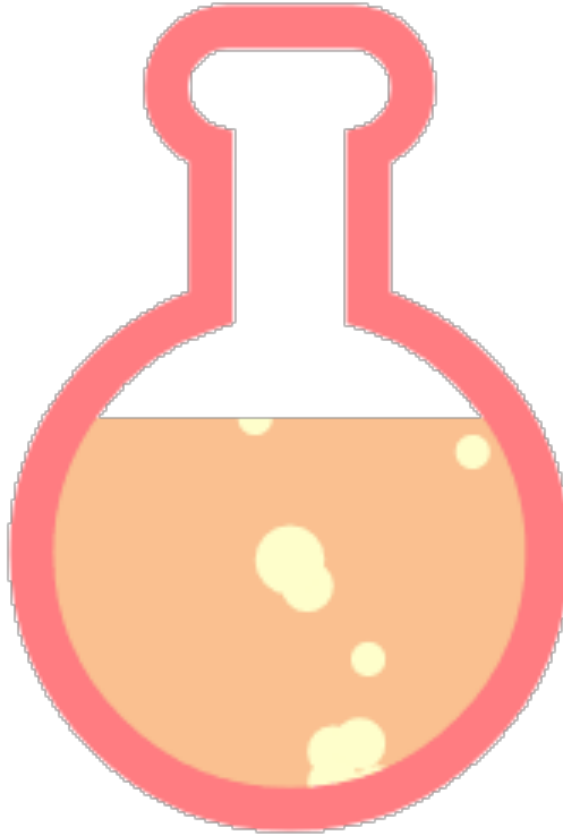
- **Nature** of EWPT \rightarrow Important open question, e.g. **its order**:
 - A **First-Order transition** (e.g. the boiling of water)? 
 - or a **Second-Order transition** (e.g. the superfluid transition) or **cross-over**?

Breaking the Symmetry in the SM

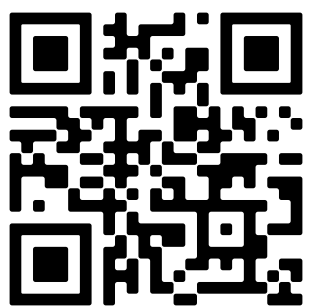


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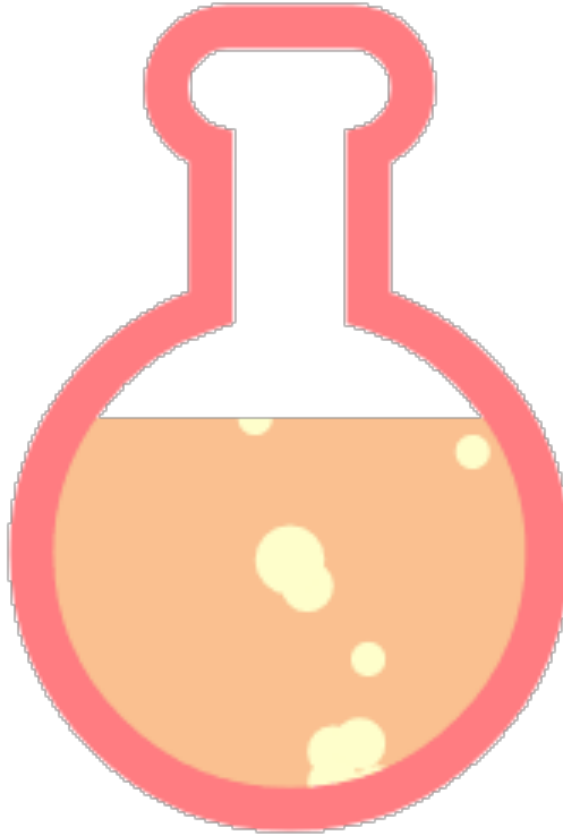
The Nature of the Phase Transition

- Clues to the origin of **matter-anti-matter asymmetry**.
- Was the asymmetry created **during the EWPT?**
 - “Electro-Weak Baryogenesis” (**EWBG**).
- Pre-requisite: **a First-Order transition**. → 
- *Note:* This **does not** occur in the SM!

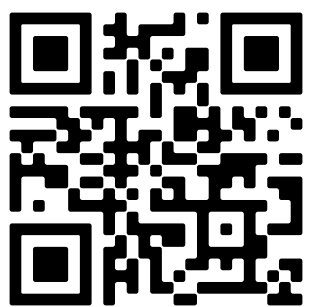
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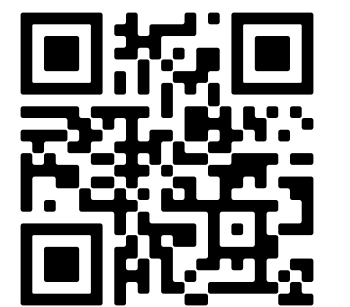
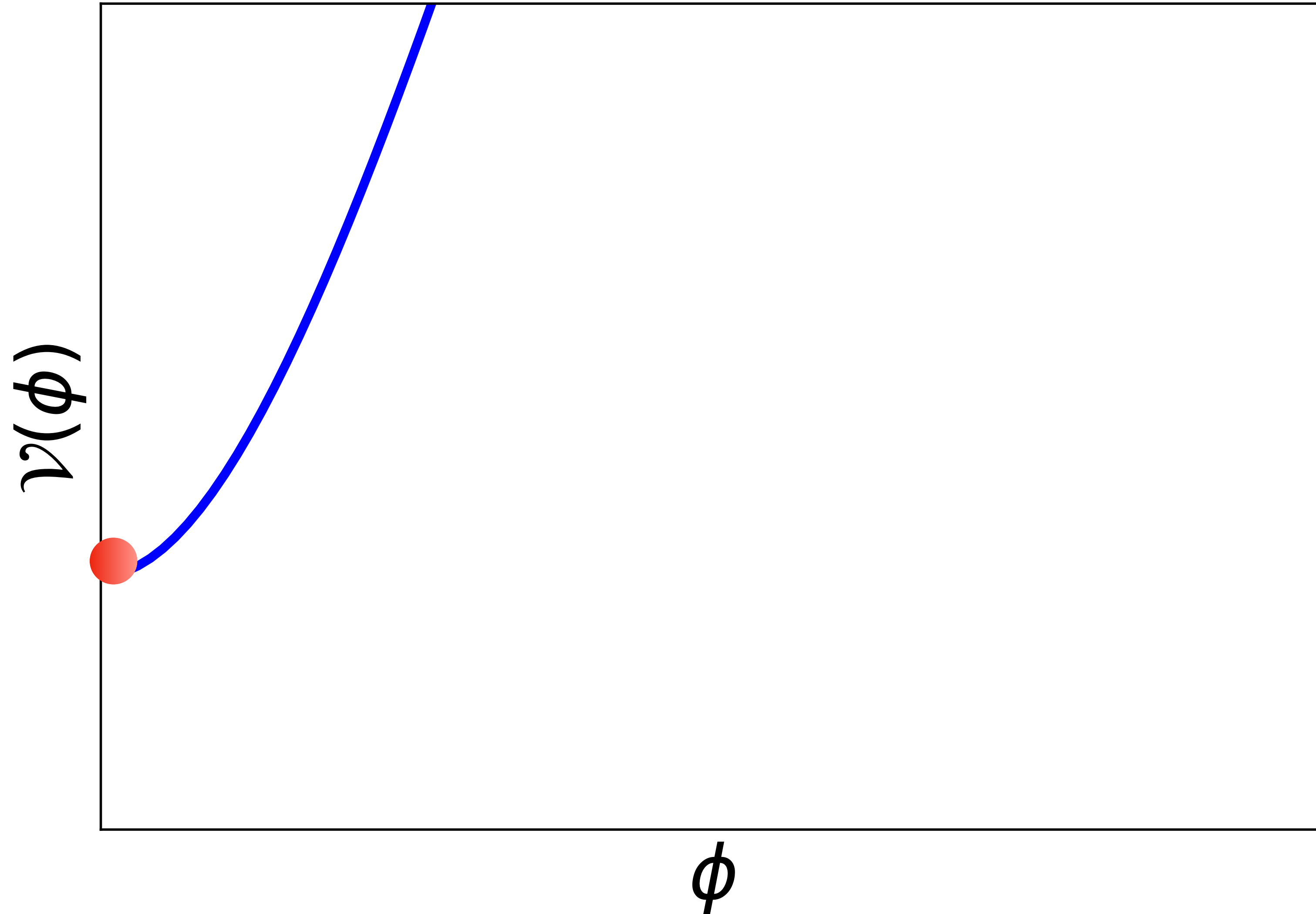
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A First-Order Electro-Weak Phase Transition: As the Universe Cools Down!

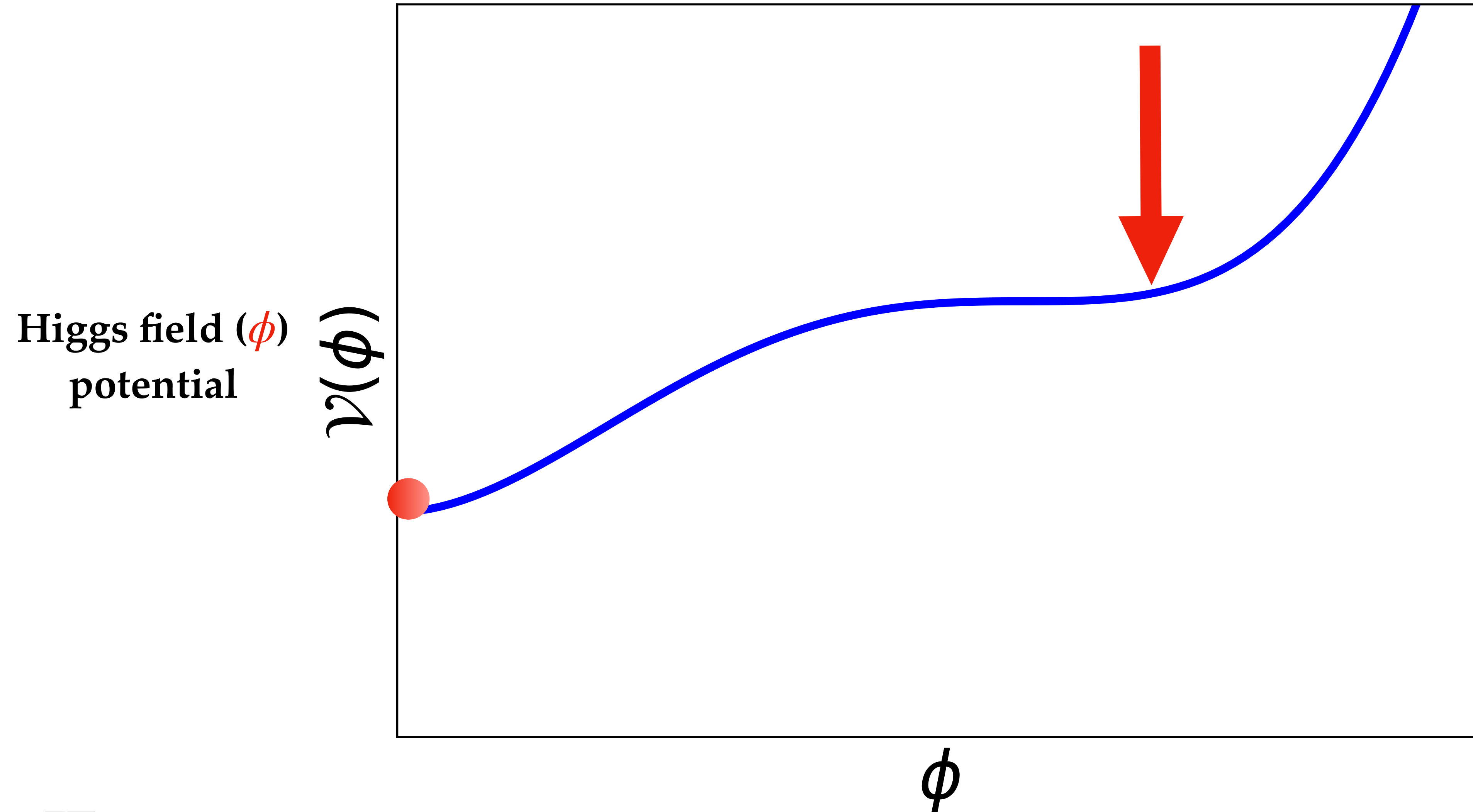
Early Universe

Higgs field (ϕ)
potential



A First-Order Electro-Weak Phase Transition: As the Universe Cools Down!

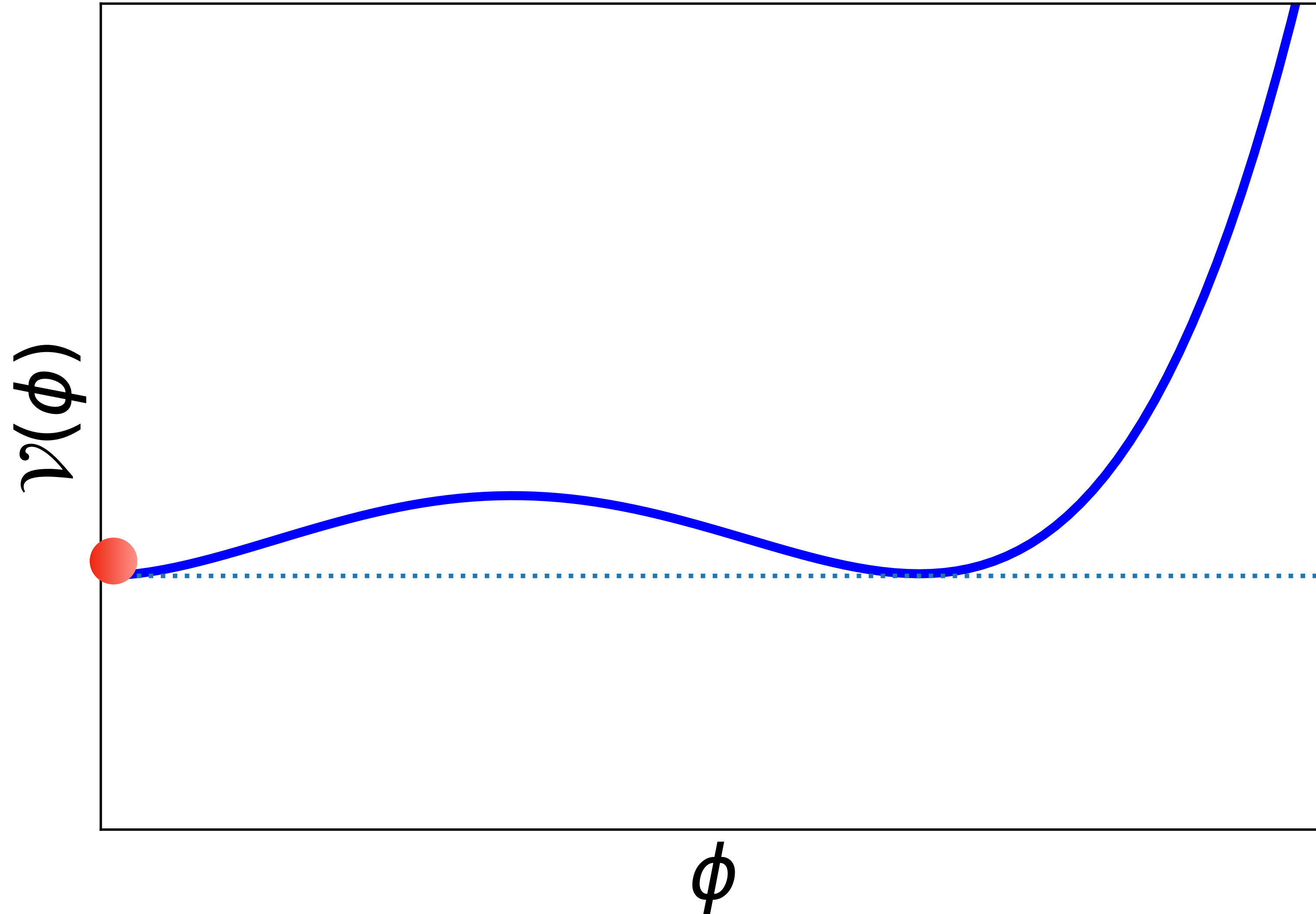
Second minimum appears



A First-Order Electro-Weak Phase Transition: As the Universe Cools Down!

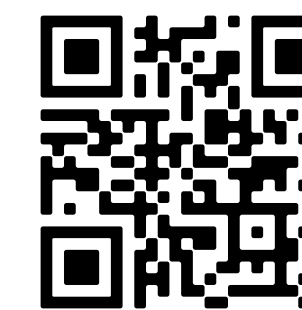
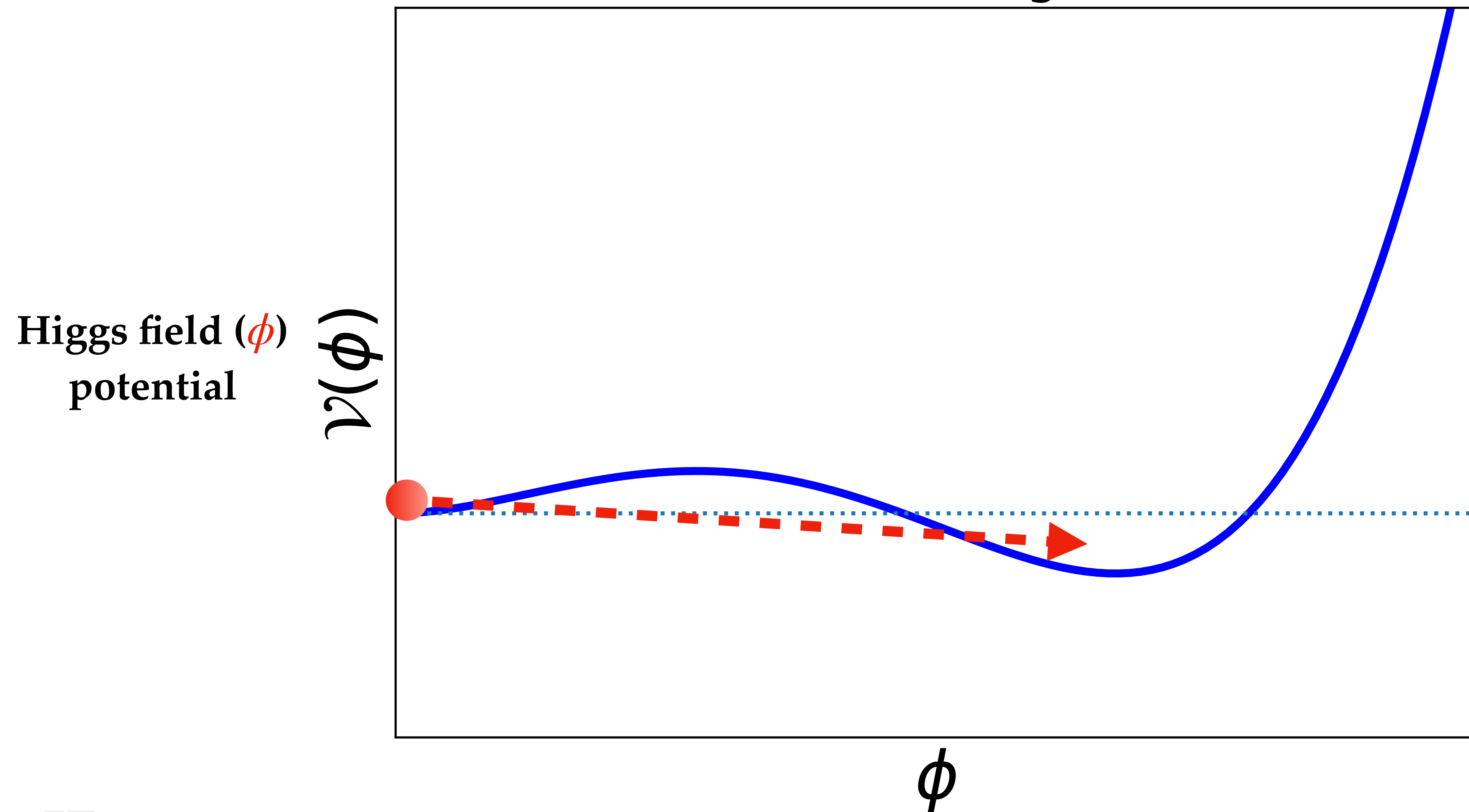
Critical temperature reached

Higgs field (ϕ)
potential



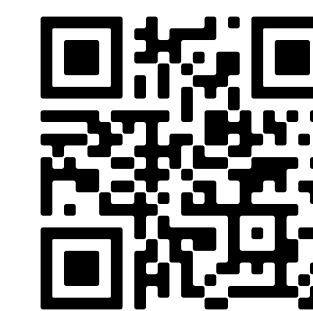
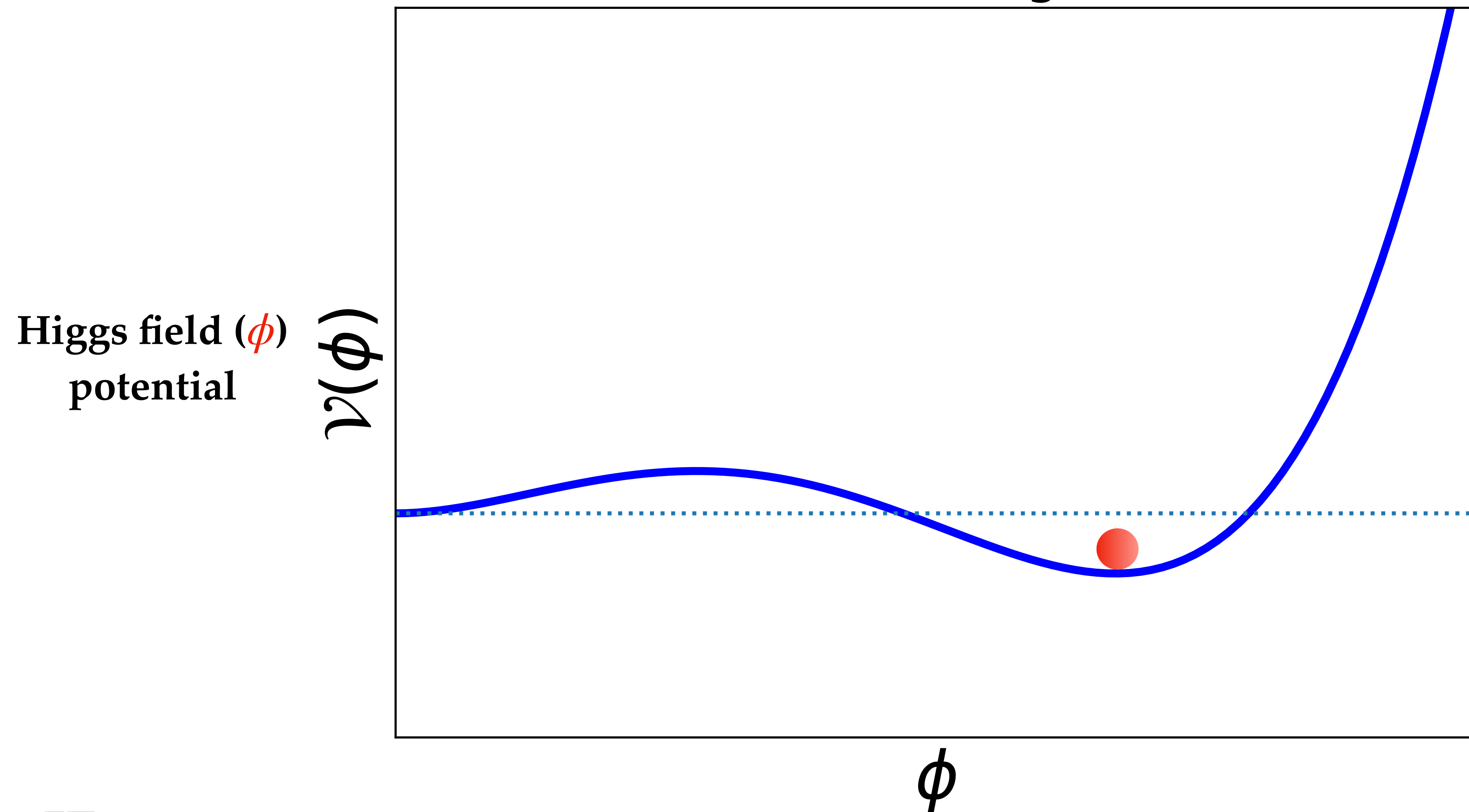
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Tunneling!



A First-Order Electro-Weak Phase Transition: As the Universe Cools Down!

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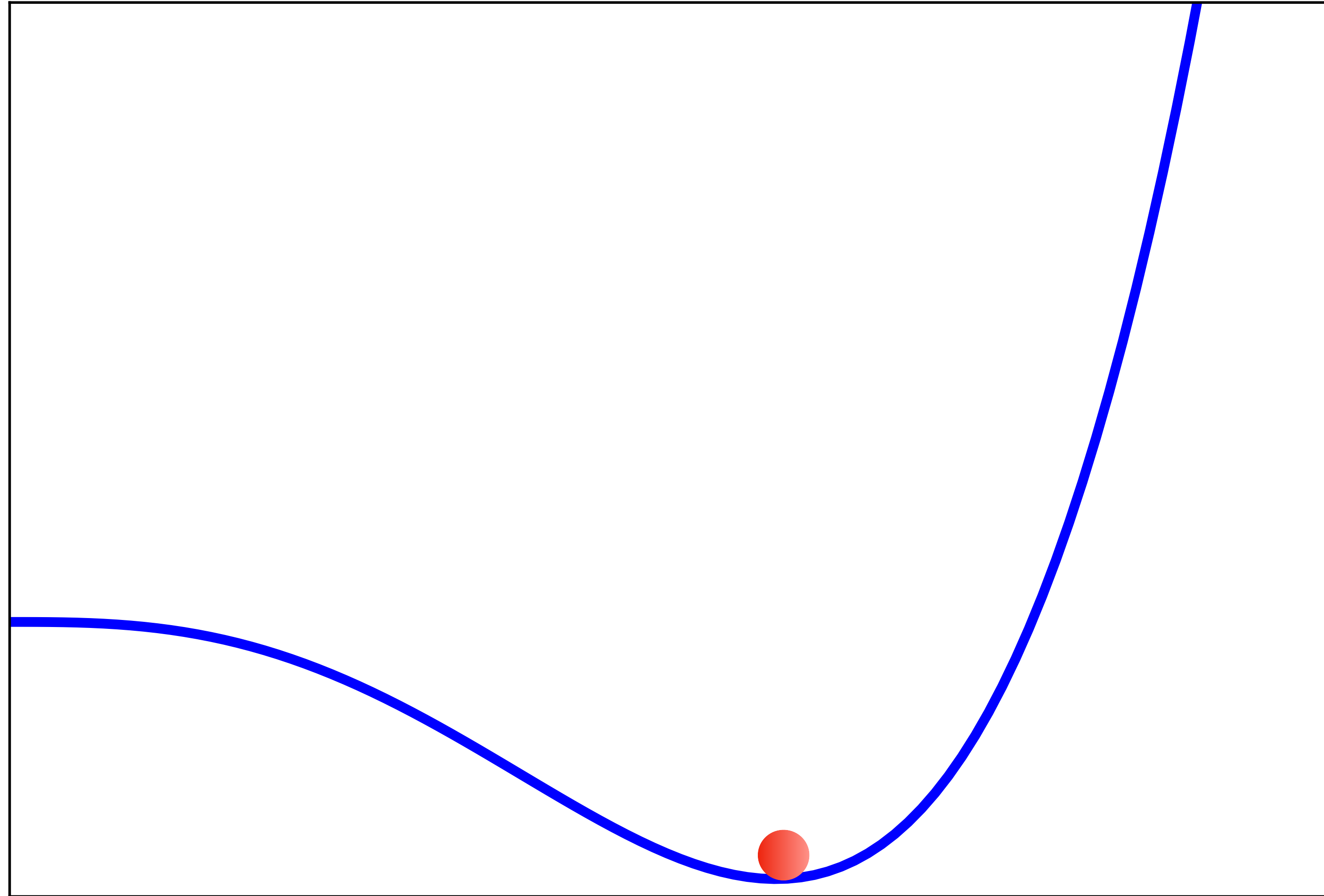


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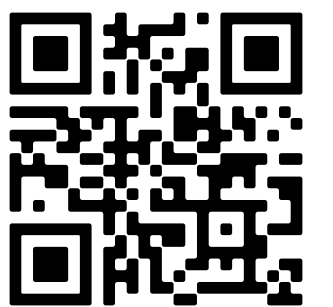
Today

Higgs field (ϕ)
potential

$V(\phi)$

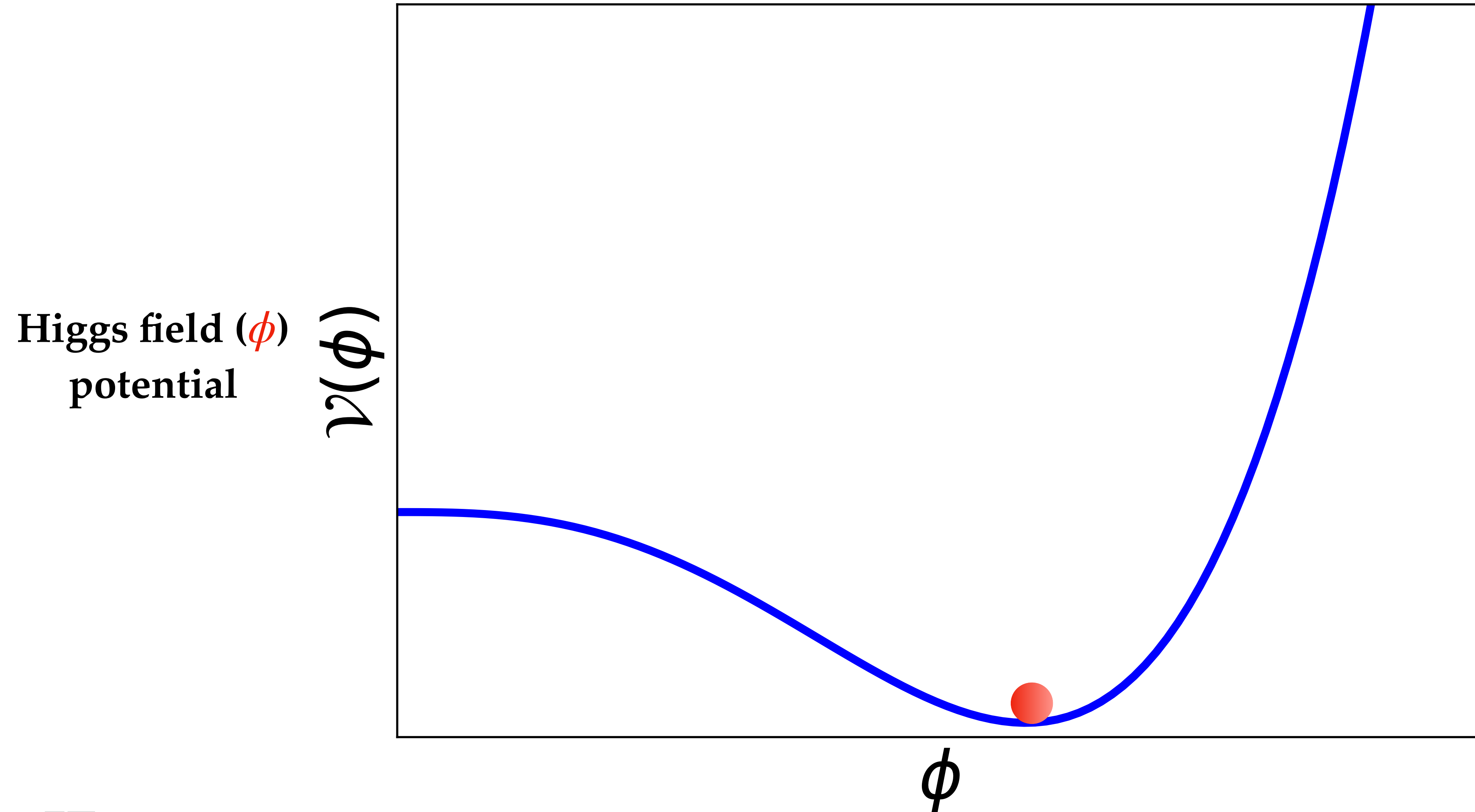


ϕ

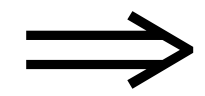


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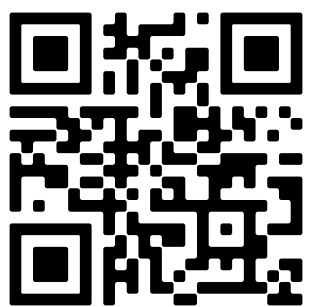
Today



$$\langle \phi \rangle \neq 0$$



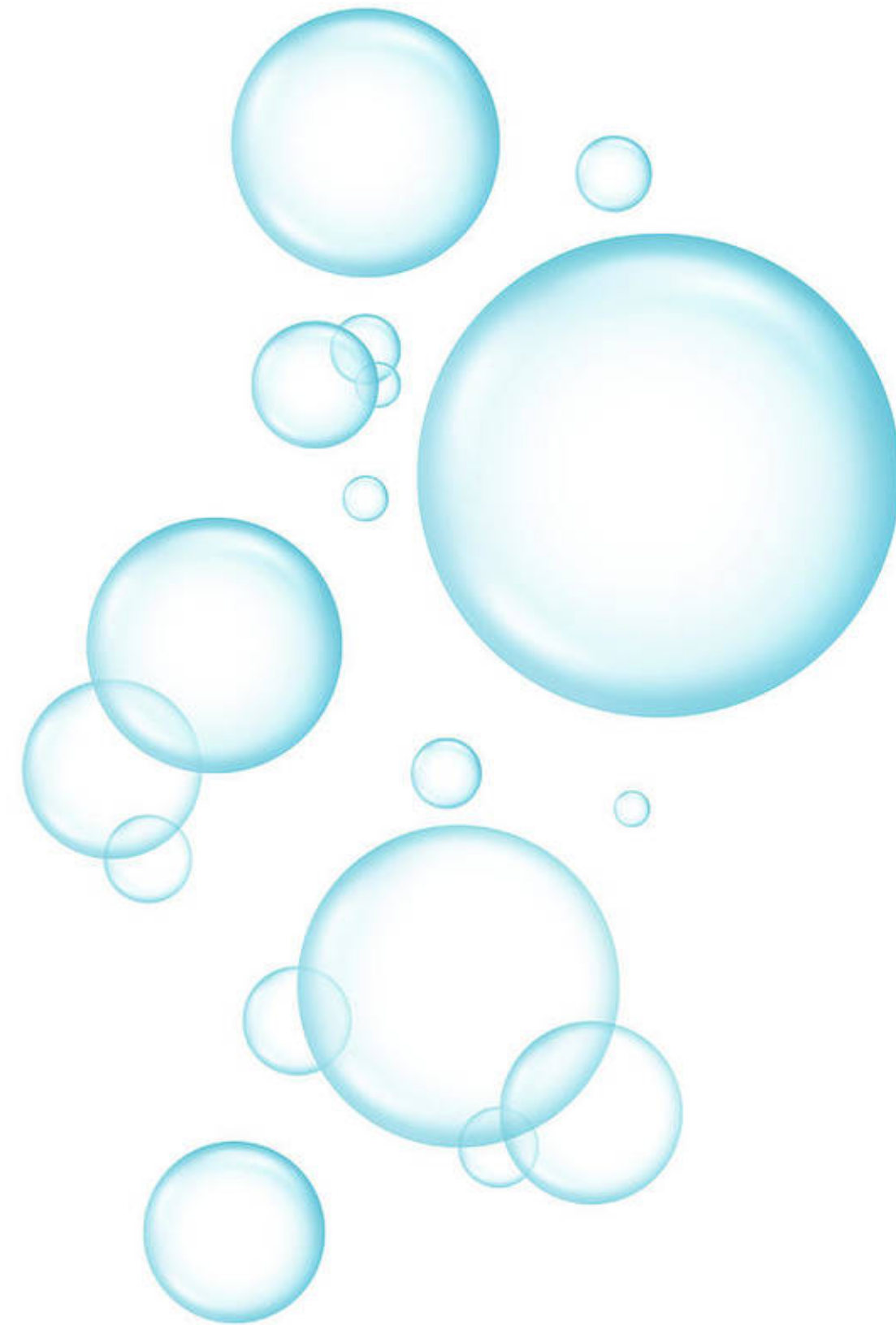
**Symmetry
Breaking.**



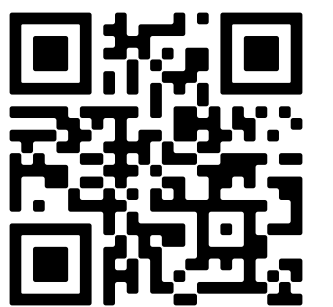
A First-Order EWPT

Some time after critical temperature is reached:

→ **Bubbles** of the broken phase nucleate and expand.*



**Note:* Gravitational Waves can form during this period!

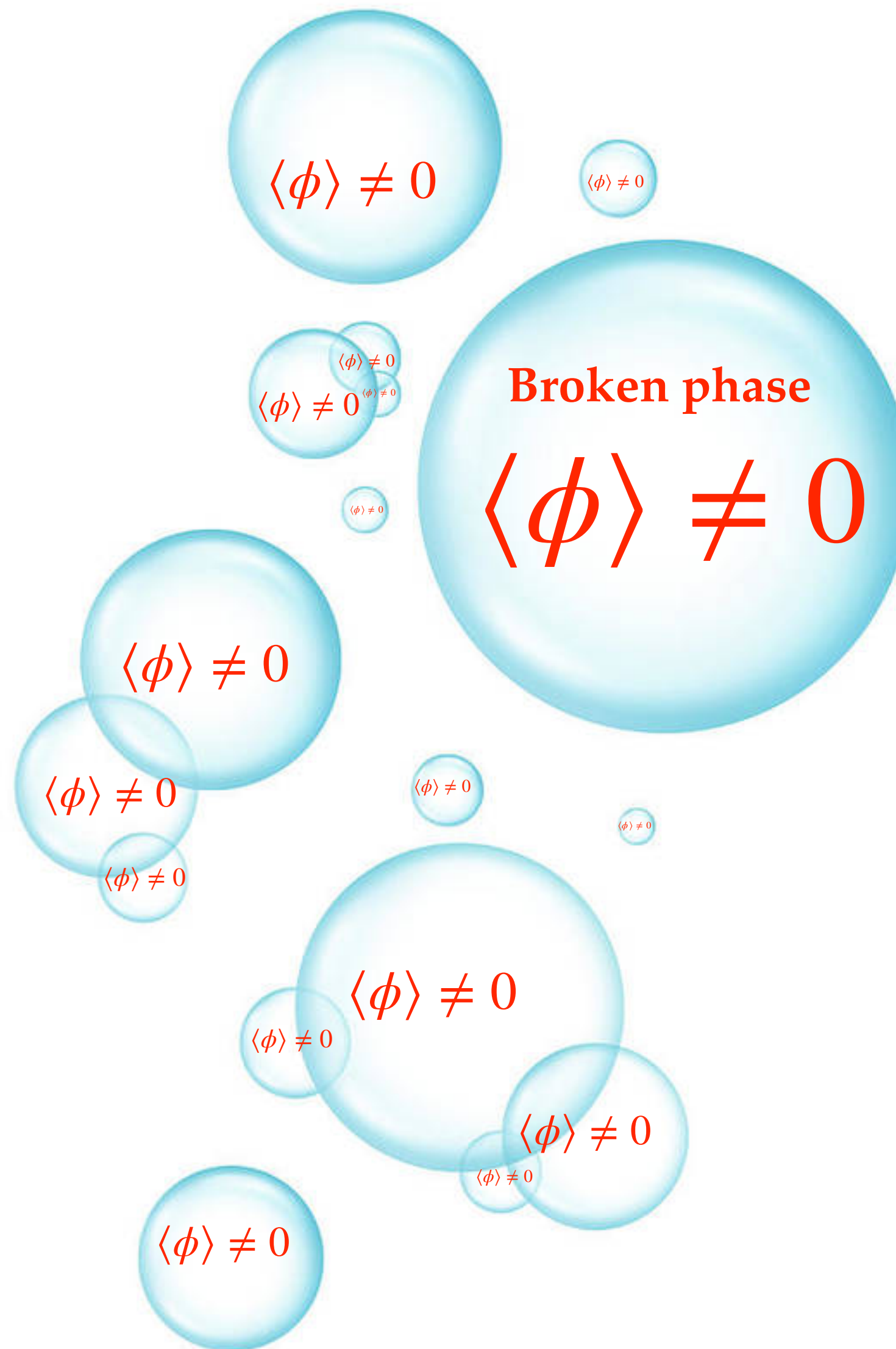


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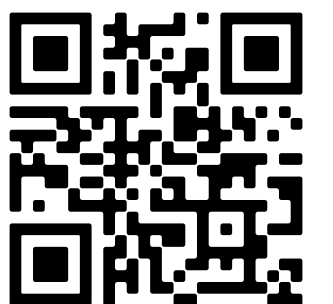
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$$\langle \phi \rangle = 0$$

Symmetric phase

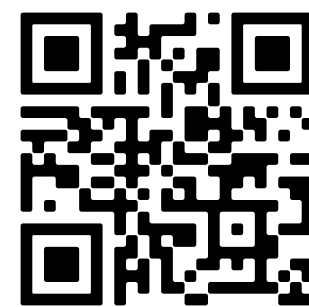
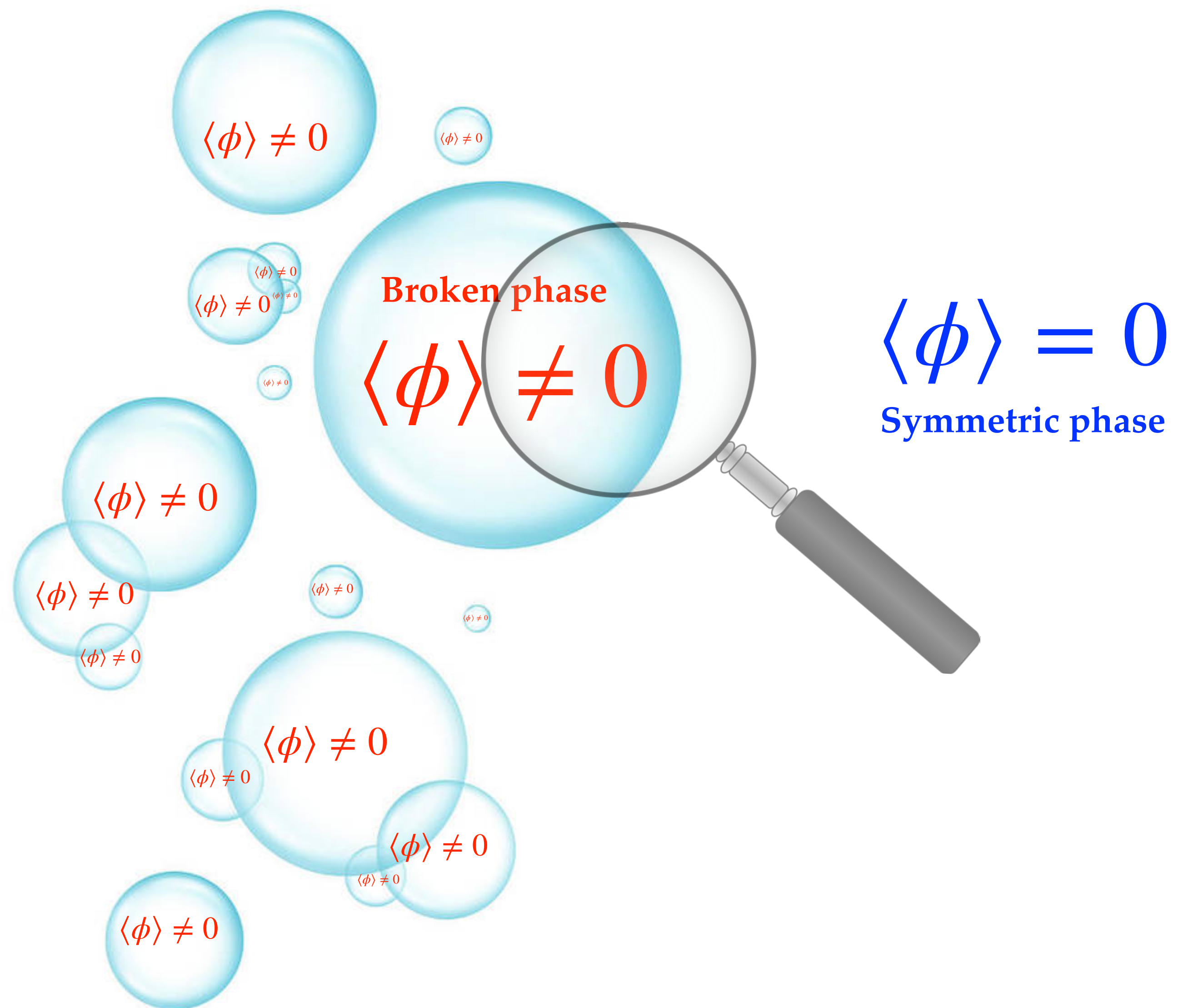


A First-Order EWPT

Some time after critical temperature is reached:

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Electro-Weak Baryogenesis

$\langle \phi \rangle \neq 0$
Broken phase

$\langle \phi \rangle = 0$
Symmetric phase



Electro-Weak Baryogenesis

Left/Right-Handed Fermions

$$\psi_L + \psi_R$$

$$\langle \phi \rangle \neq 0$$

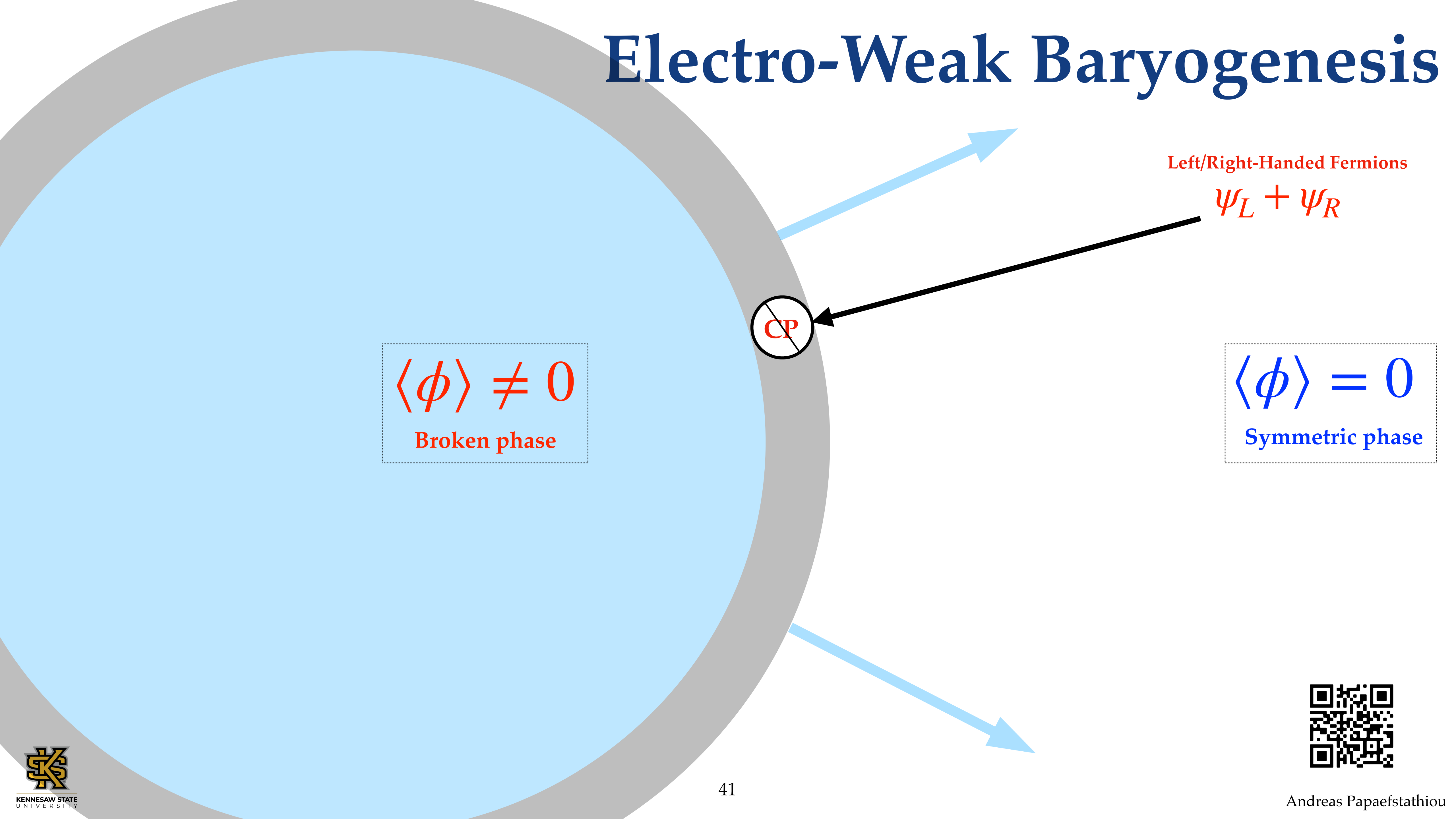
Broken phase

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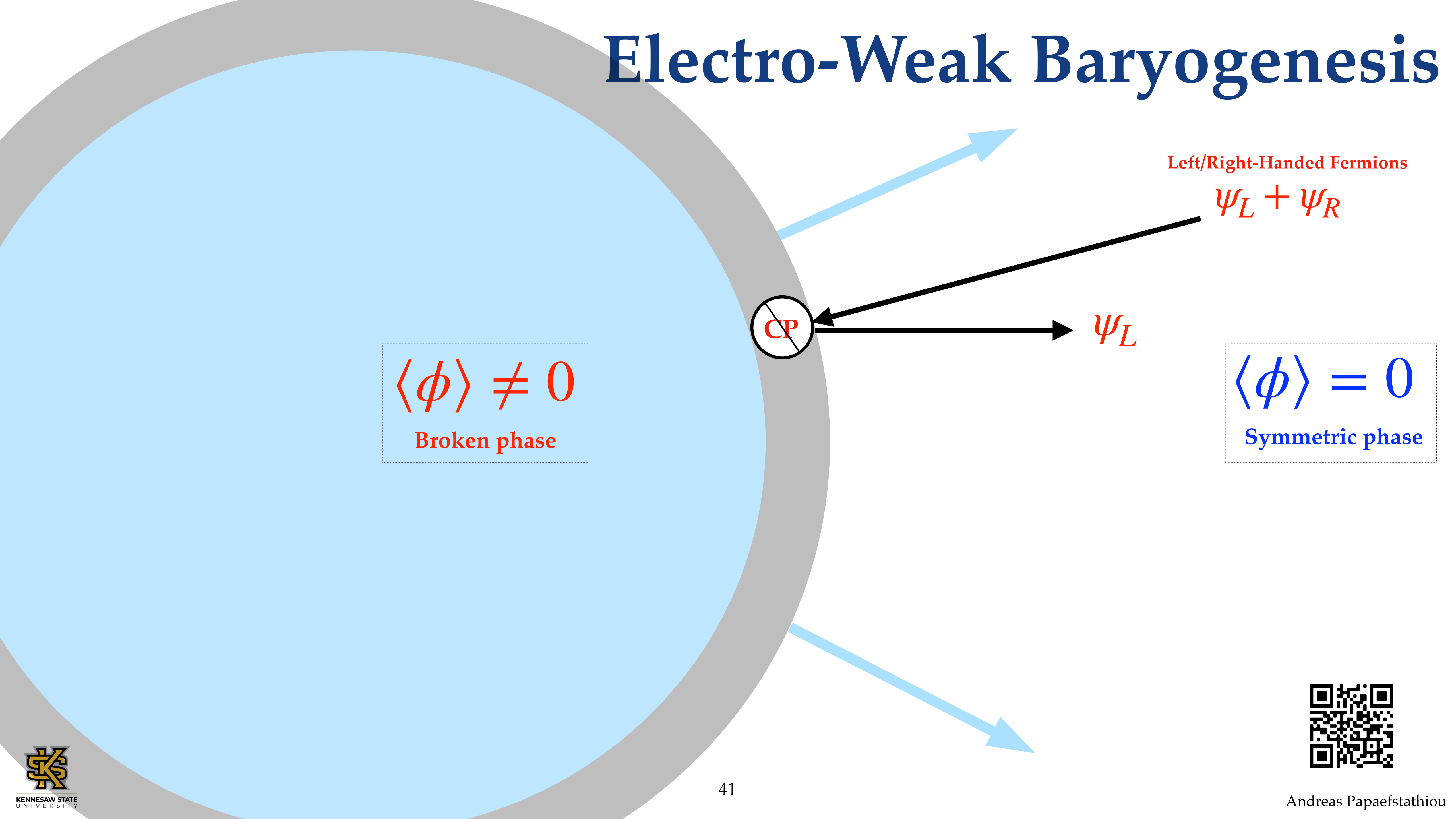
Symmetric phase



Electro-Weak Baryogenesis



Electro-Weak Baryogenesis



Left/Right-Handed Fermions

$$\psi_L + \psi_R$$

CP

$$\psi_L$$

$$\langle \phi \rangle \neq 0$$

Broken phase

$$\langle \phi \rangle = 0$$

Symmetric phase



Electro-Weak Baryogenesis

~~Sphaleron~~

$\langle \phi \rangle \neq 0$
Broken phase

~~CP~~

Left/Right-Handed Fermions

$\psi_L + \psi_R$

ψ_L

$\langle \phi \rangle = 0$
Symmetric phase

t_L, b_L

τ_L, ν_τ

e_L, ν_e

Sphaleron

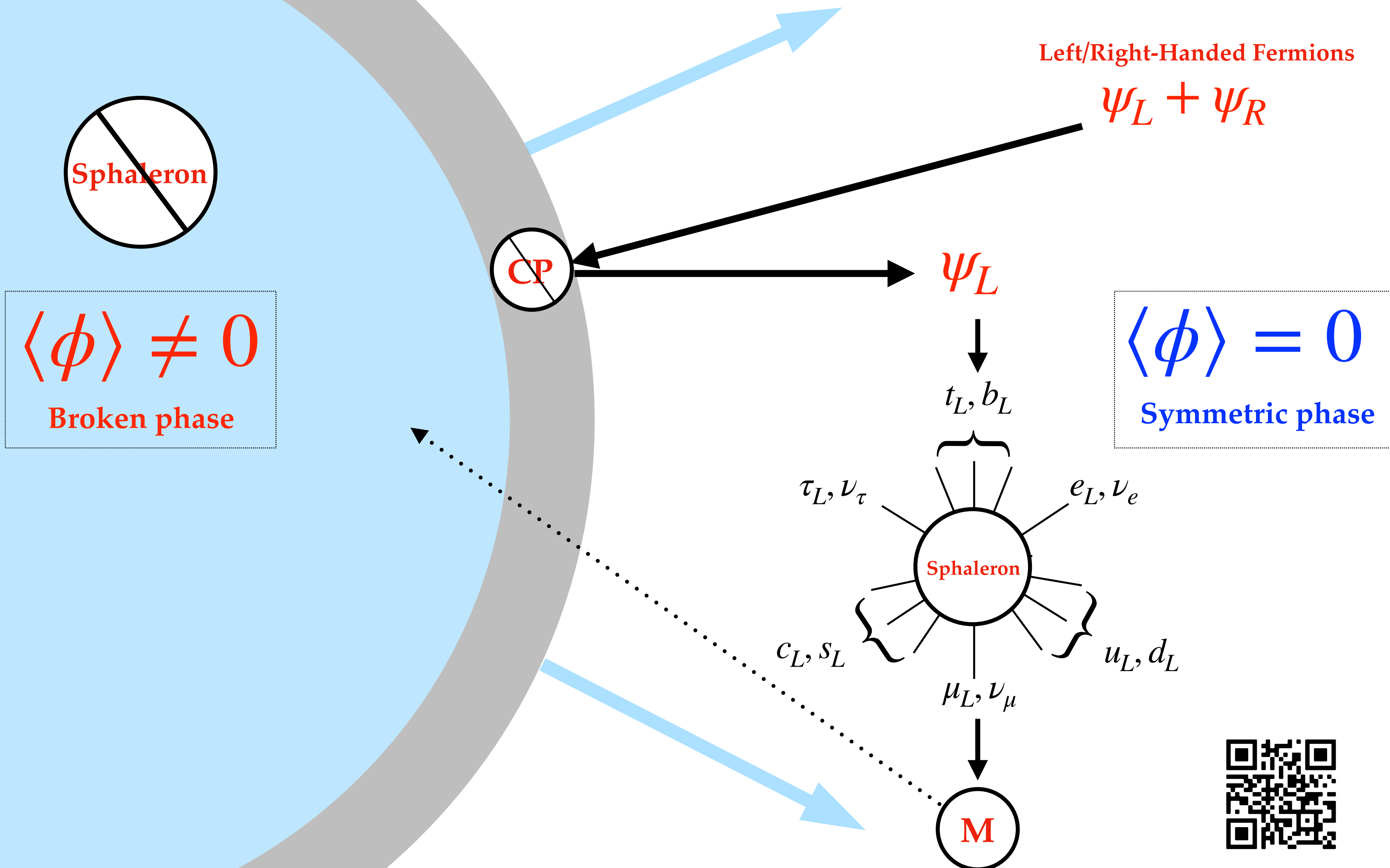
c_L, s_L

u_L, d_L

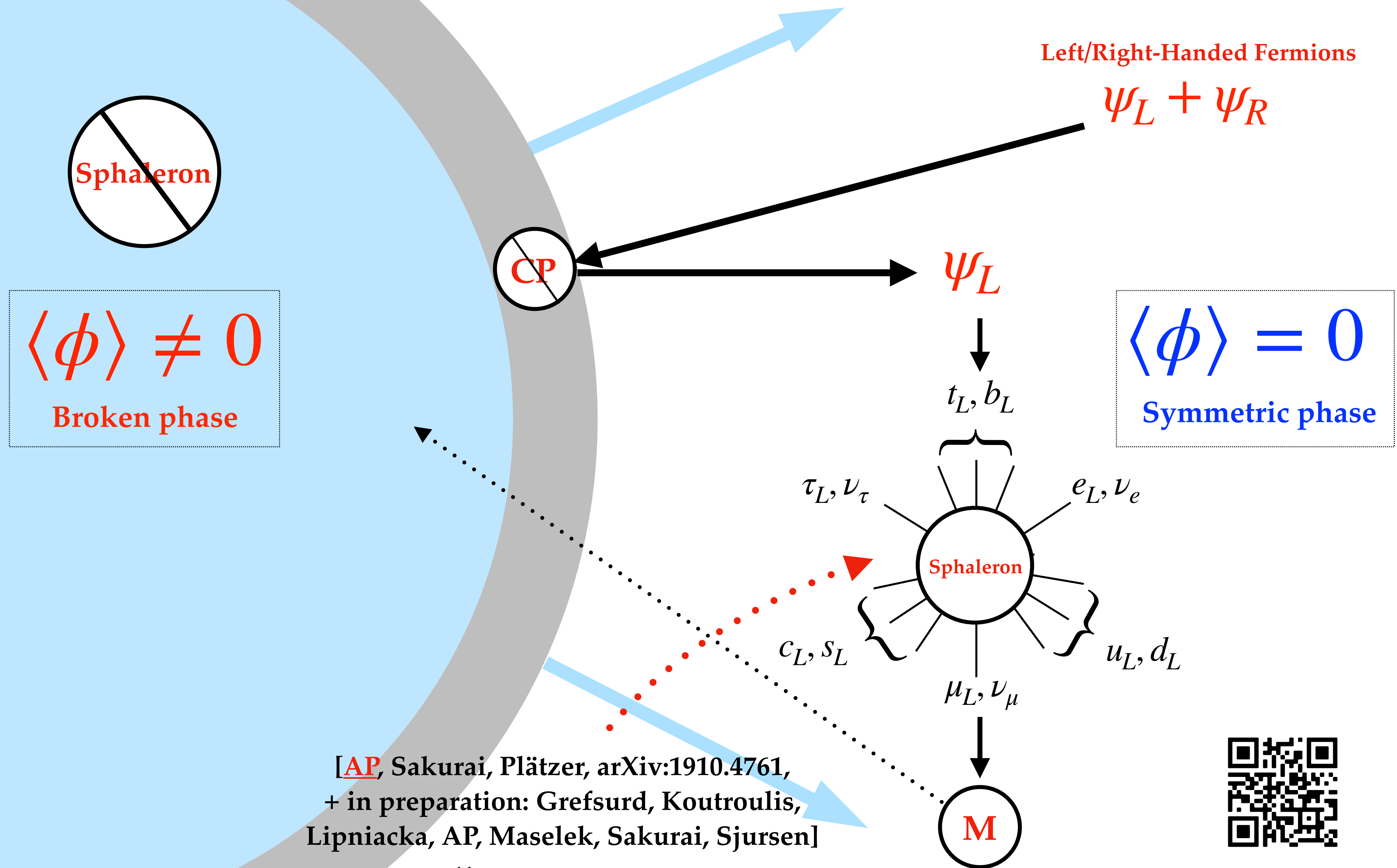
μ_L, ν_μ



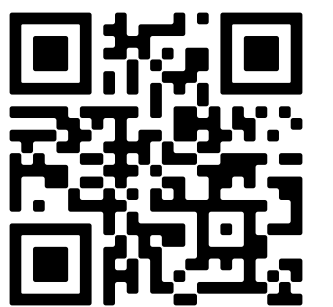
Electro-Weak Baryogenesis



Electro-Weak Baryogenesis



[AP, Sakurai, Plätzer, arXiv:1910.4761,
 + in preparation: Grefsurd, Koutroulis,
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Electro-Weak Baryogenesis

~~Sphaleron~~

~~CP~~

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$$\psi_L$$

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Symmetric phase

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$$\tau_L, \nu_\tau$$

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Sphaleron

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$$u_L, d_L$$

$$\mu_L, \nu_\mu$$

M

**A First-Order
Transition requires
New Phenomena
beyond the SM!**

[AP, Sakurai, Plätzer, arXiv:1910.4761,
+ in preparation: Grefsurd, Koutroulis,
Lipniacka, AP, Maselek, Sakurai, Sjursen]



Electro-Weak Archaeology



We live here!



$$\langle \phi \rangle \neq 0$$

Broken phase



Electro-Weak Archaeology



We live here!

$$\langle \phi \rangle \neq 0$$

Broken phase



→ What are the implications
of Electro-Weak
Baryogenesis at Colliders?



4 Extending the Scalar Sector [AP, White, arXiv:2010.00597]

- A First-Order EWPT dictates new phenomena. [Kajantie, Laine, Rummukainen, Shaposhnikov hep-ph/9605288]
- Consider the **simplest possible extension!**

$$\mathcal{V}(\phi, S) = \bullet |\phi|^2 + \blacksquare |\phi|^4$$

Add: S , a new scalar field,
No SM “charges” \equiv Singlet.

Extending the Scalar Sector [[AP](#), White, arXiv:2010.00597]

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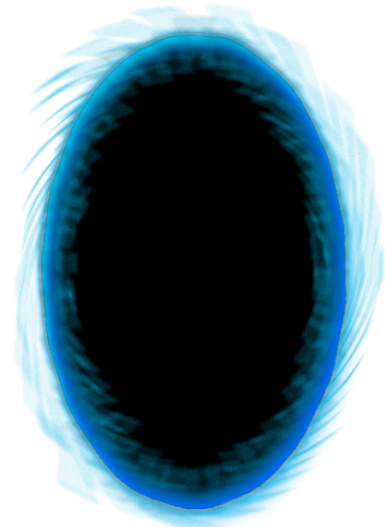
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$$+ \text{●} S^2 + \text{▲} S^3 + \text{■} S^4$$

$$+ \text{▲} |\phi|^2 S + \text{■} |\phi|^2 S^2 \leftarrow \text{“Portal” interactions.}$$



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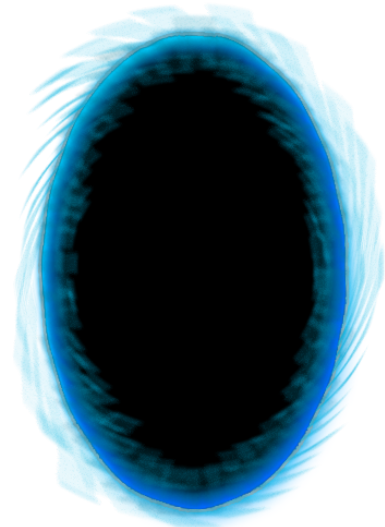
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$$+ S \times (\text{Hidden Sector}) + \dots \leftarrow \text{Dark Matter?}$$



Extending the Scalar Sector [[AP](#), White, arXiv:2010.00597]

$$\mathcal{V}(\phi, S) = \color{green}{\bullet} |\phi|^2 + \color{blue}{\blacksquare} |\phi|^4 + \color{magenta}{\bullet} S^2 + \color{cyan}{\blacktriangle} S^3 + \color{maroon}{\blacksquare} S^4 + \color{red}{\blacktriangle} |\phi|^2 S + \color{purple}{\blacksquare} |\phi|^2 S^2$$

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EWSB \leftrightarrow VEVs:

$$\phi \rightarrow \langle \phi \rangle + h$$

$$S \rightarrow \langle S \rangle + \chi$$

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EWSB \leftrightarrow VEVs:

$$\begin{aligned} \phi &\rightarrow \langle \phi \rangle + h \\ S &\rightarrow \langle S \rangle + \chi \end{aligned} \quad \color{blue}{\rightarrow} \quad \mathcal{V} \supset \color{red}{\circ} h^2 + \color{cyan}{\circ} h\chi + \color{blue}{\circ} \chi^2$$

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$$\phi \rightarrow \langle \phi \rangle + h \quad \longrightarrow \quad \mathcal{V} \supset \color{red}{\circ} h^2 + \color{cyan}{\circ} h\chi + \color{blue}{\circ} \chi^2$$

$$S \rightarrow \langle S \rangle + \chi$$

\Rightarrow Mass (squared) matrix:

$$M^2 = \begin{pmatrix} \frac{\partial^2 \mathcal{V}}{\partial h^2} & \frac{\partial^2 \mathcal{V}}{\partial h \partial \chi} \\ \frac{\partial^2 \mathcal{V}}{\partial h \partial \chi} & \frac{\partial^2 \mathcal{V}}{\partial \chi^2} \end{pmatrix}$$

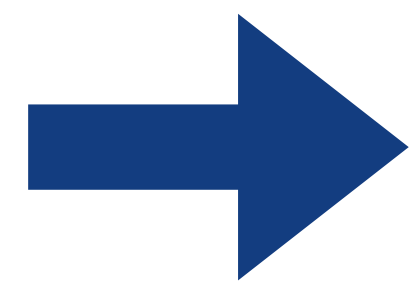
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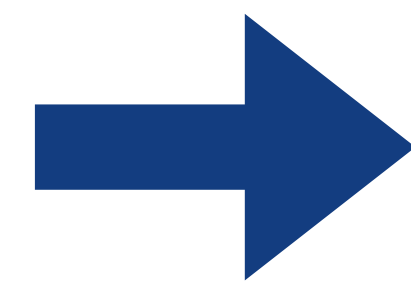
EWSB \leftrightarrow VEVs:

$$\phi \rightarrow \langle \phi \rangle + h$$

$$S \rightarrow \langle S \rangle + \chi$$



$$\mathcal{V} \supset \color{red}{\bullet} h^2 + \color{cyan}{\bullet} h\chi + \color{blue}{\bullet} \chi^2$$



Diagonalize!

\Rightarrow Mass (squared) matrix:

Mass Eigenstates

$$M^2 = \begin{pmatrix} \frac{\partial^2 \mathcal{V}}{\partial h^2} & \frac{\partial^2 \mathcal{V}}{\partial h \partial \chi} \\ \frac{\partial^2 \mathcal{V}}{\partial h \partial \chi} & \frac{\partial^2 \mathcal{V}}{\partial \chi^2} \end{pmatrix}$$

$$\begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} h \\ \chi \end{pmatrix}$$

θ : mixing angle

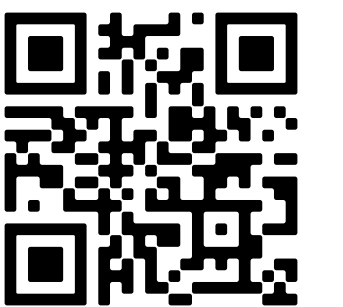
Extending the Scalar Sector [[AP](#), White, arXiv:2010.00597]

$$\mathcal{V}(\phi, S) = \color{green}{\bullet} |\phi|^2 + \color{blue}{\blacksquare} |\phi|^4 + \color{magenta}{\bullet} S^2 + \color{cyan}{\blacktriangle} S^3 + \color{maroon}{\blacksquare} S^4 + \color{red}{\blacktriangle} |\phi|^2 S + \color{purple}{\blacksquare} |\phi|^2 S^2$$

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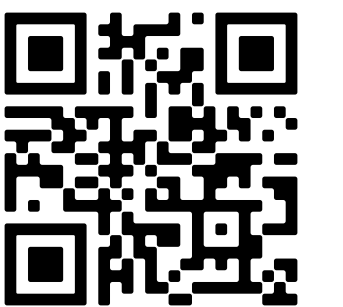
$h_1 \rightarrow$ "SM-like" Higgs boson.

$h_2 \rightarrow$ new scalar resonance.

i.e. choose: $\theta \sim 0$, and:

$$h_1 = h \cos \theta + \chi \sin \theta$$

$$h_2 = -h \sin \theta + \chi \cos \theta$$



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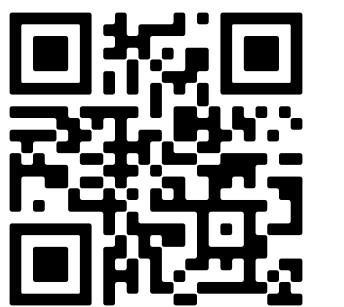
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Primary targets for collider studies!

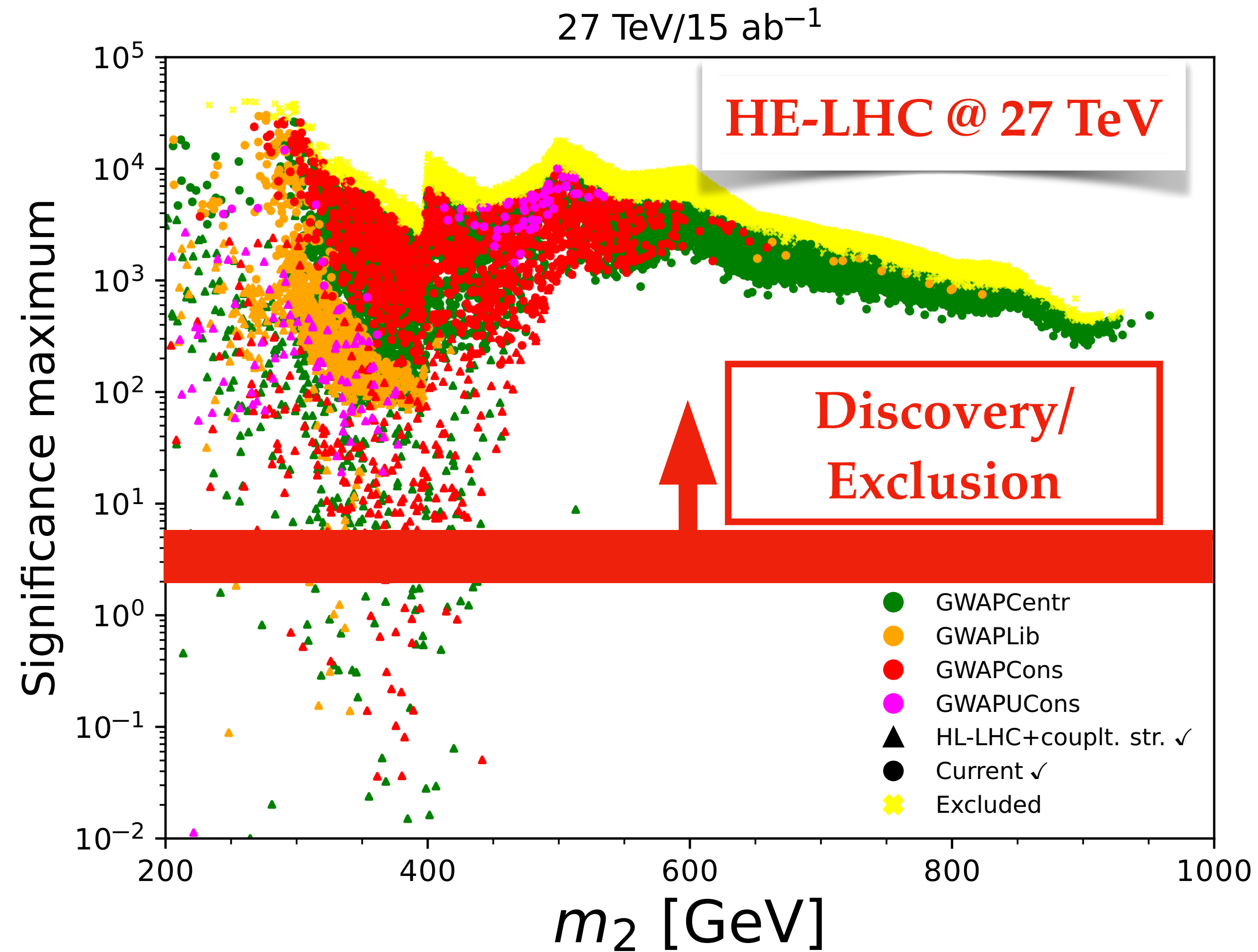
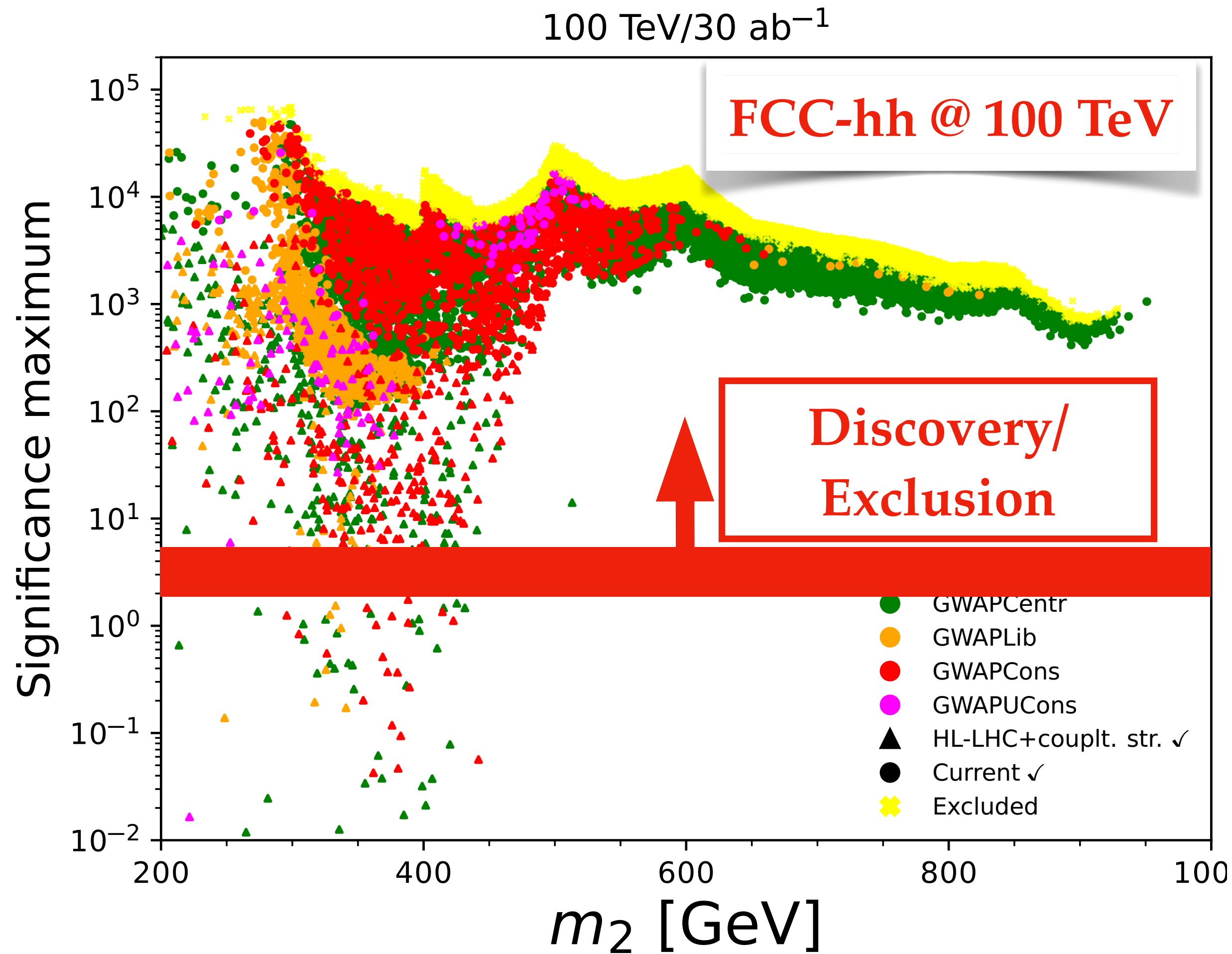
$h_1 \Rightarrow$ Reductions in Higgs boson rates.

$h_2 \Rightarrow$ New resonance searches.



Significance @ Future Colliders

Color-coded parameter-space points denotes **theoretical uncertainty**.

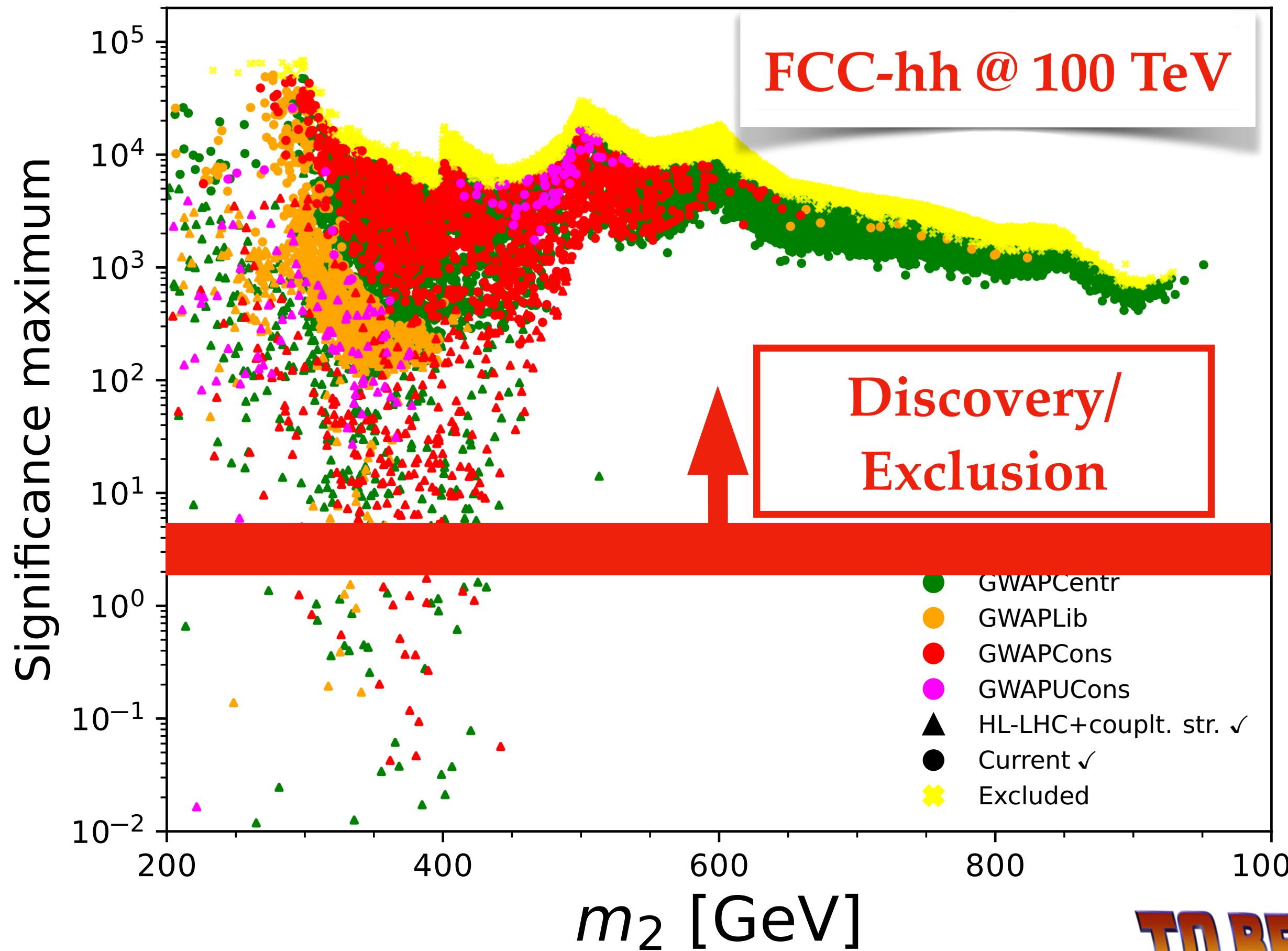


✕ = Point is excluded today.

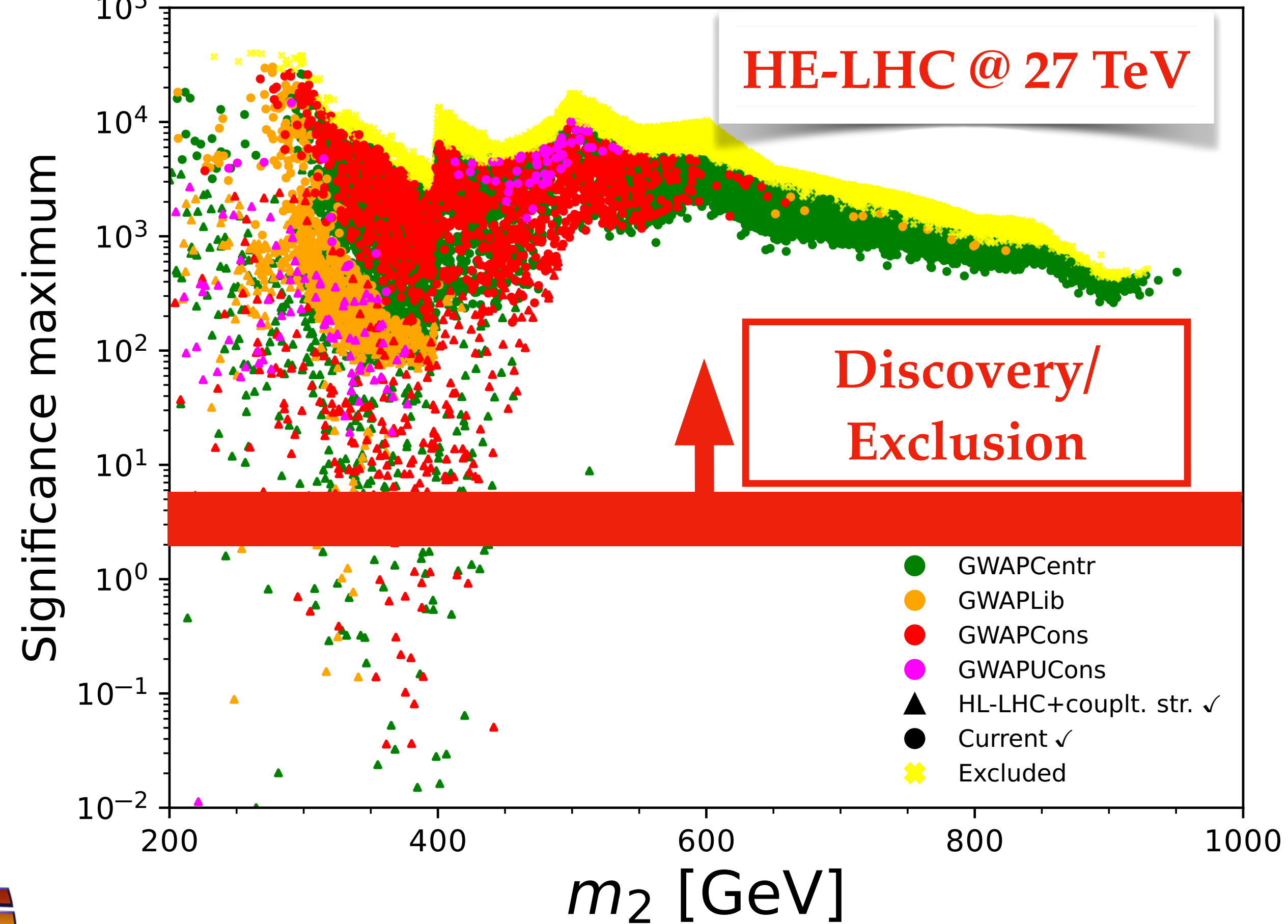
Significance @ Future Colliders

Color-coded parameter-space points denotes **theoretical uncertainty**.

100 TeV/30 ab⁻¹



27 TeV/15 ab⁻¹



TO BE CONTINUED...

✱ = Point is excluded today.

Discovery Post-Mortem

“With 4 parameters I can fit an elephant and with 5 I can make him wiggle his trunk.”
— John Von Neumann

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Can we fit the shape of an elephant with 4 parameters?

Discovery Post-Mortem

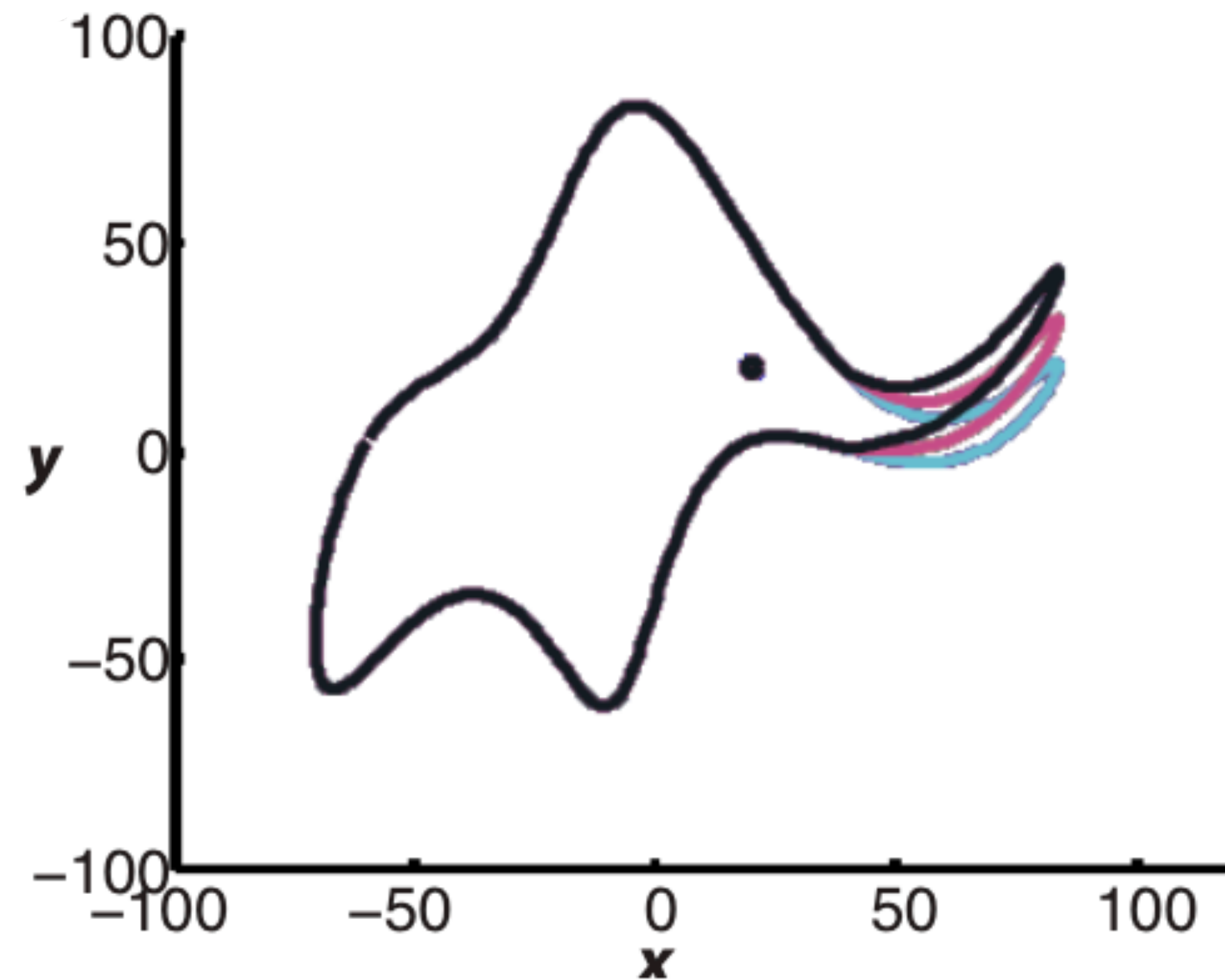
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Can we fit the shape of an elephant with 4 parameters?

→ **Yes!** With **four** complex parameters,

[and with 5 we can make it wiggle its trunk.]

[Mayer, Khairy, Howard, Am. J. Phys., Vol. 78, No. 6, June 2010]

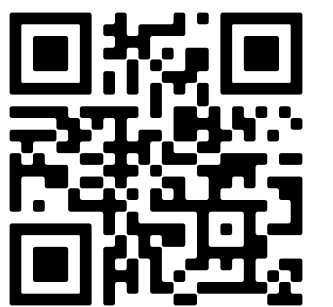


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If we discover e.g. a new scalar particle at colliders,

→ Can we verify that it is indeed the remnant of a singlet field that generates a first-order electro-weak phase transition?



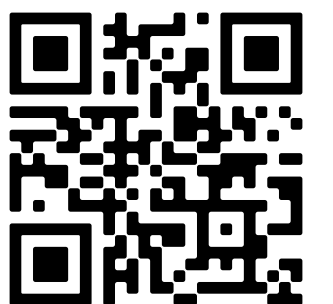
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⇒ The “**Inverse Problem**”.



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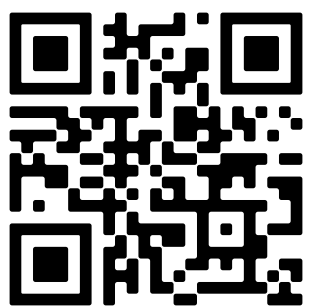
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Experimental observations → { ●, ▲, ■, ▲, ■ }

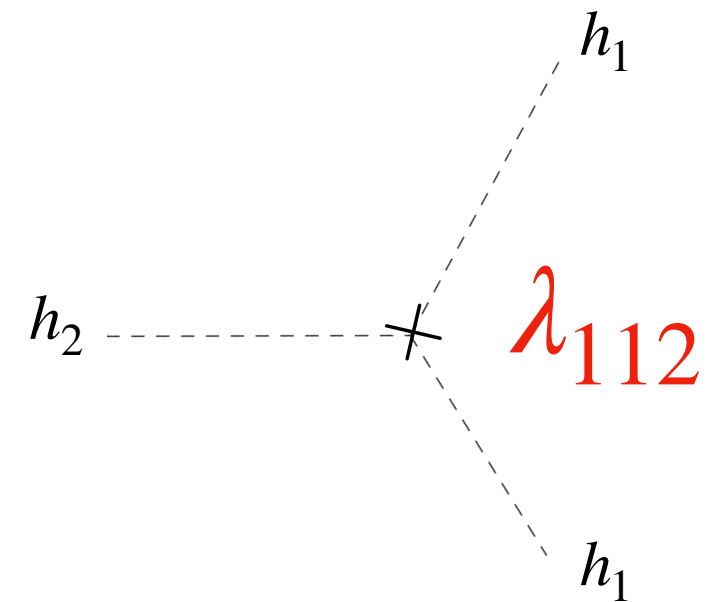


Discovery Post-Mortem Example [AP, White, arXiv:2108.11394]






Combine possible measurements:

$$pp \rightarrow h_2 \rightarrow ZZ$$

$$pp \rightarrow h_2 \rightarrow h_1 h_1$$

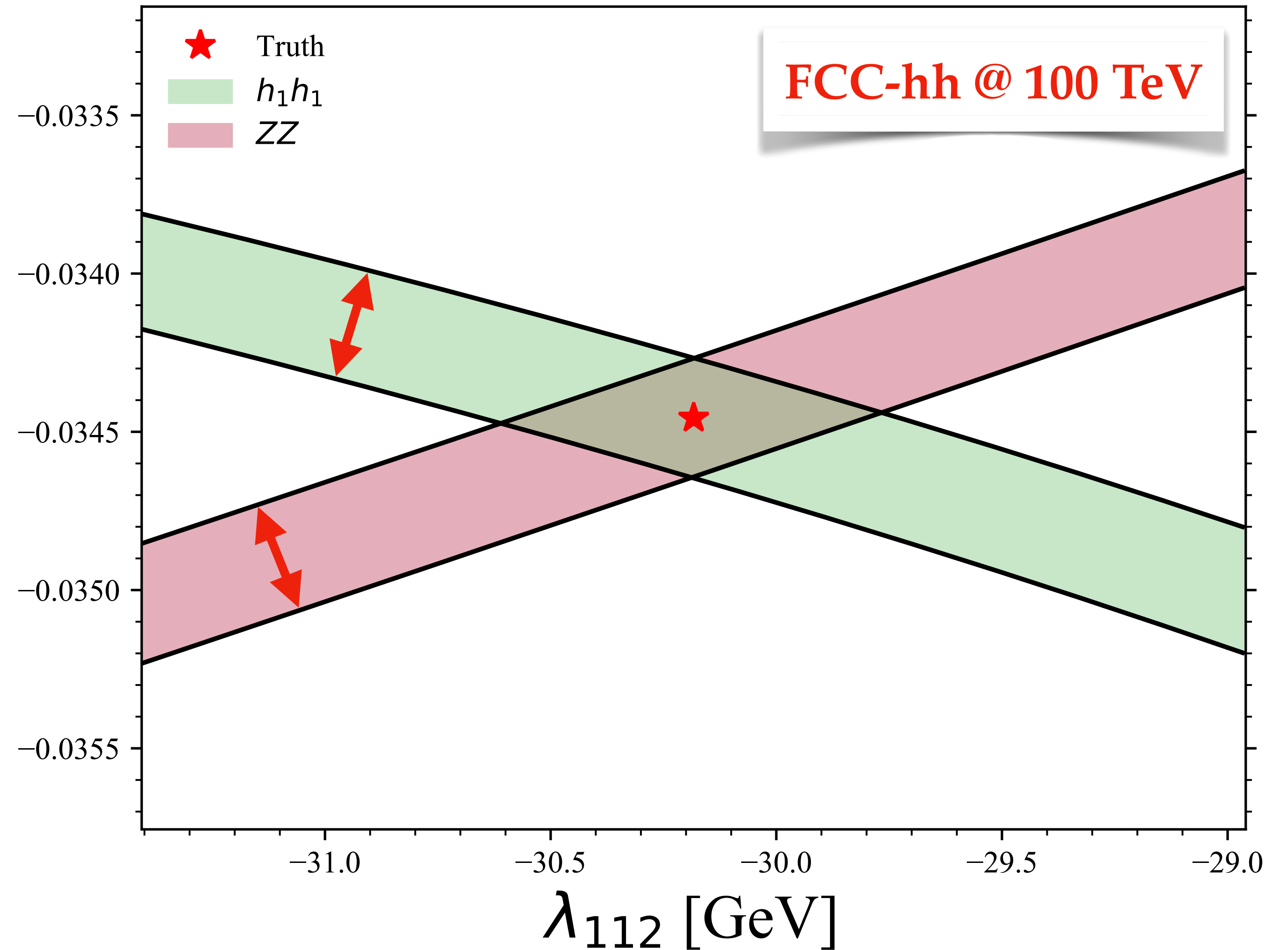


→ both functions of $\sin \theta$ & λ_{112} !

→ in turn f^{ns} of { , , , ,  }.

[Width of bands represents expected measurement uncertainty].

pp@100 TeV/30000 fb⁻¹, UCons1

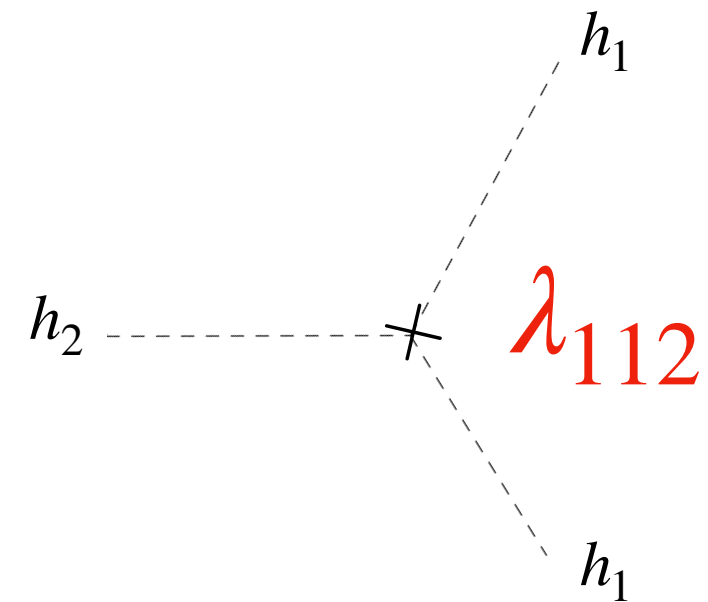


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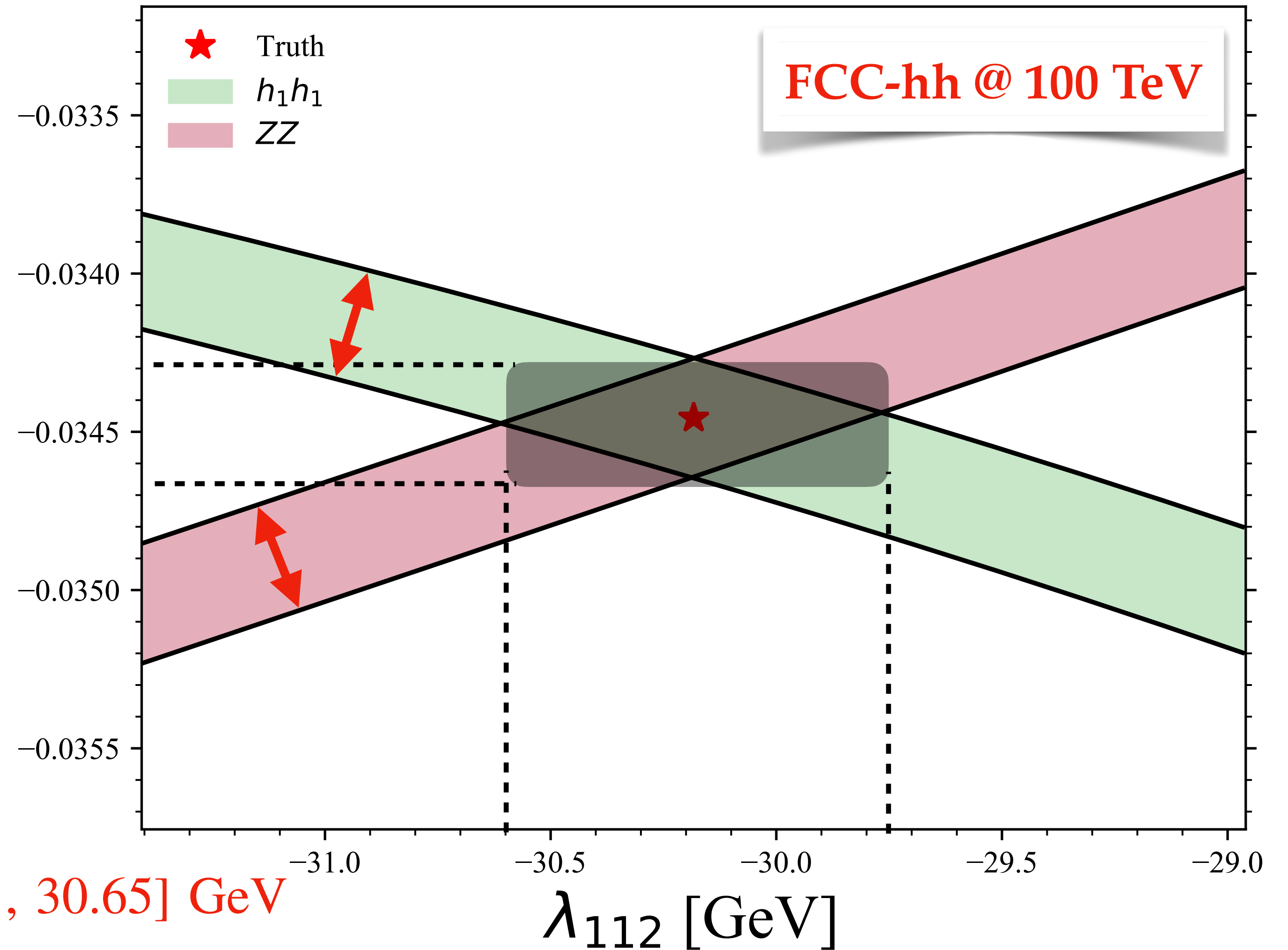


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$$\lambda_{112} \in [29.81, 30.65] \text{ GeV}$$

$$\sin \theta \in [0.0343, 0.0347]$$

The Inverse Problem in Extended Scalar Sectors:

Multi-scalar processes should play a **crucial rôle**:

$$pp \rightarrow h_1 h_2 \quad | \mathcal{M} |^2 \sim \lambda_{122}^2, \lambda_{112}^2 + \dots$$

$$pp \rightarrow h_2 h_2 \quad | \mathcal{M} |^2 \sim \lambda_{222}^2, \lambda_{122}^2 + \dots$$

$$pp \rightarrow h_1 h_1 h_1 \quad | \mathcal{M} |^2 \sim f[\lambda_{1111}, \lambda_{1112}, \lambda_{111}, \lambda_{112}]$$

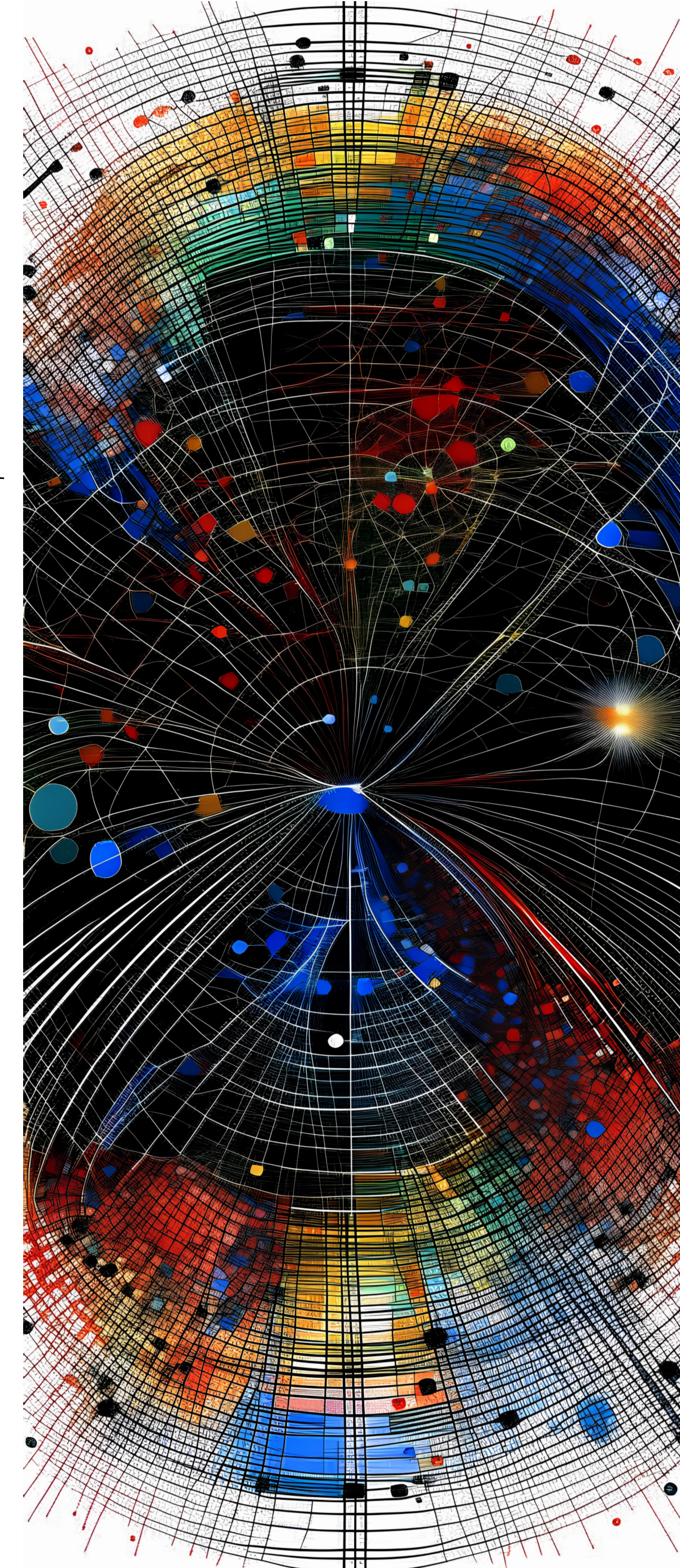
[...]

**TO BE
CONTINUED** 



Summary

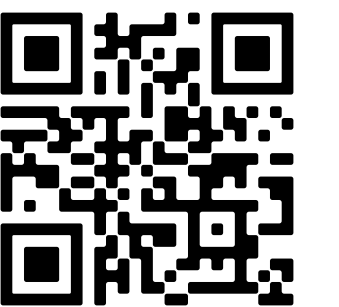
- Within the realm of the Standard Model, we can describe a whole of range of particle interactions with extraordinary precision.
- **Monte Carlo Event Generators** are invaluable tools of communication between theory and experiment.
- However, there remains a multitude of **open questions**, that may be linked via the **Electro-Weak Phase Transition**.
- The **Nature of the Electro-Weak Phase Transition** is an important scientific enquiry.
- **(Strong) First-Order EWPT [not in SM!] → Matter-Anti-Matter asymmetry.**
- Adding a **singlet scalar field to the SM can catalyze this!**
- Future particle colliders have the **potential to probe this mechanism.**



[made with stablecog] Andreas Papaefstathiou

Outlook

- Following any discovery, **solving the inverse problem** would be the **crucial** next step.
- I discussed **first steps in this exercise**, following the discovery of a new scalar.
- **Multi-scalar production processes** will play a crucial role in this endeavor.
- [Could they also be **discovery channels** themselves?]
- Questions merit investigation both at the LHC and other future colliders (e.g. a Muon Collider!).

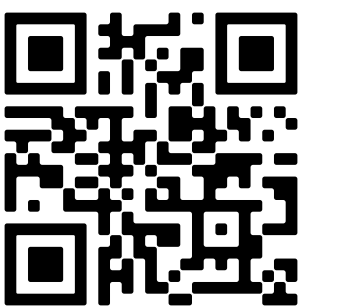


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Thanks!
Questions?



Appendix

Theoretical Uncertainties in SM+Singlet

- **Uncertainties** → Can affect e.g. the strength of the transition, $\langle \phi(T) \rangle / T$.
- Due to:
 - **gauge dependence**, [Patel, Ramsey-Musolf, arXiv:1101.4665]
 - **scale dependence** → Linde's IR problem: expansion parameter is $gn_B \sim gT/m$, (n_B mode occupation), diverges as $m \rightarrow 0 \Rightarrow$ perturbativity breaks down.
[Linde, Phys. Lett. 96B (1980) 289.]
- \Rightarrow **To make reliable and sensible statements on colliders prospects:**
 - **Crucial to take uncertainties into account.**

Theoretical Uncertainty Bands

- Define “uncertainty band” by:

1. Deriving 1-loop effective potential in the covariant gauge,

[Arnold, Espinosa, hep-ph/9212235], [Andreassen, MSc, Norwegian U. Sci. Tech., 2013]

2. Run couplings $\lambda \rightarrow \lambda(\mu)$, μ is RGE scale, [SARAH, Staub, arXiv:0806.0538]

3. Scan parameter space of Lagrangian,

4. Vary $\mu \in [\frac{1}{2} \times m_Z, 5 \times m_Z]$ & gauge params. $\xi_i \in [0,3] \rightarrow$ band of 8 pts.

5. Use **PhaseTracer** for each point in band \rightarrow Get phase transitions, $\langle \phi(T_c) \rangle / T_c$.

[Athron, Balázs, Fowlie, Zhang, arXiv:2003.02859]

Parameter-space Categories

1. Define two conditions:
 - i. VEV at 1-loop: $\langle \phi(T = 0) \rangle = 246 \pm 30 \text{ GeV}$ & deepest minimum.
 - ii. $\langle \phi(T_c) \rangle / T_c > 1$ & no other transition with higher T_c .
2. Define four mutually-exclusive categories*:

*An alternative classification appears in our article: see Appendix.

SFO-EWPT more certain



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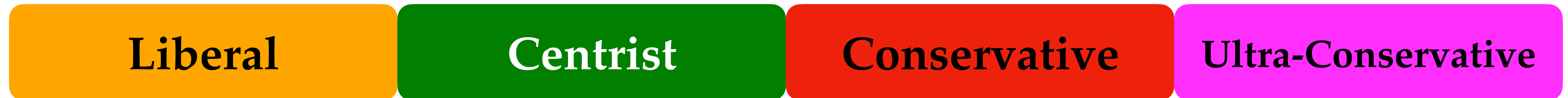
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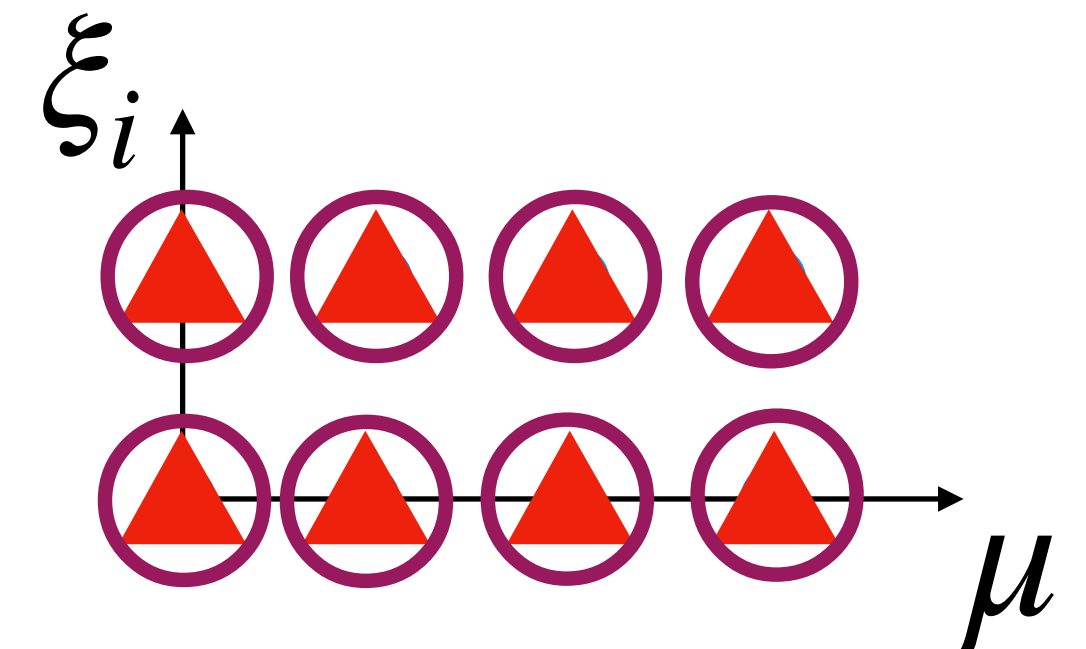
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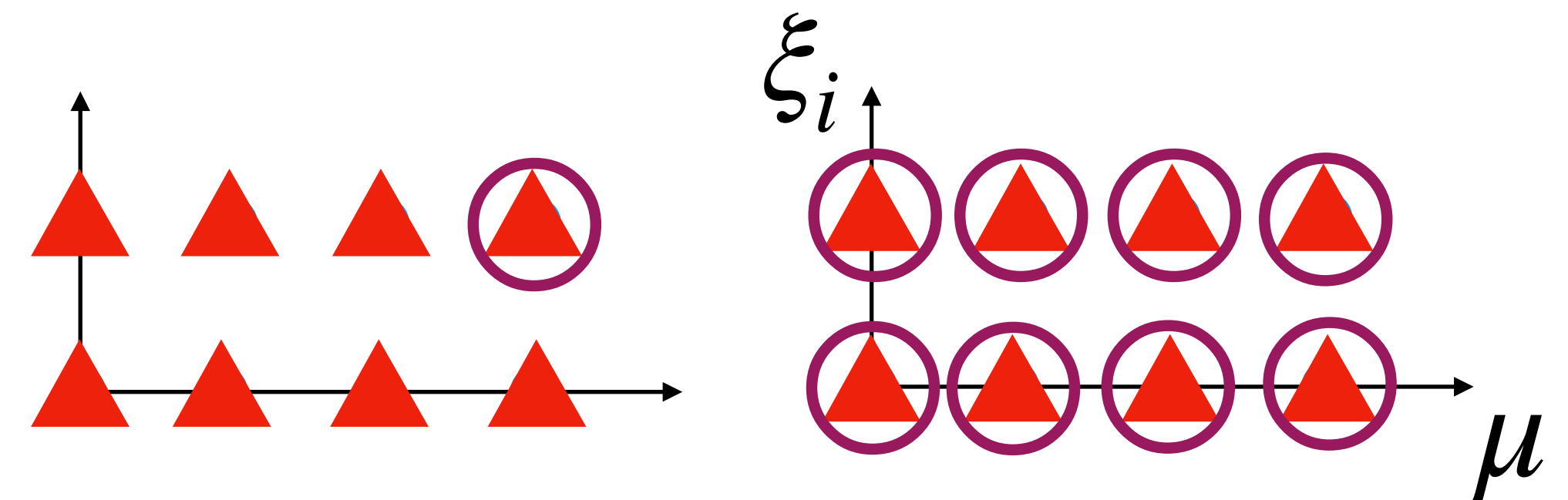
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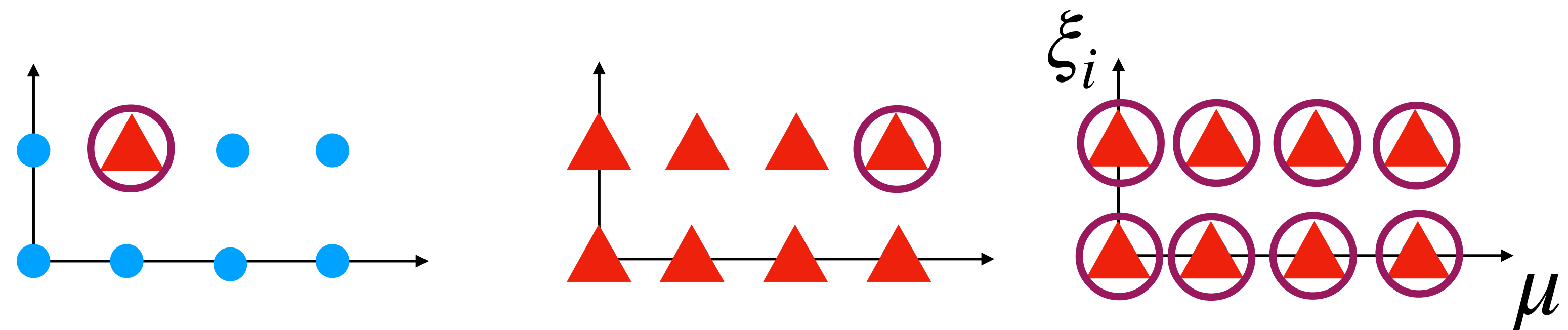
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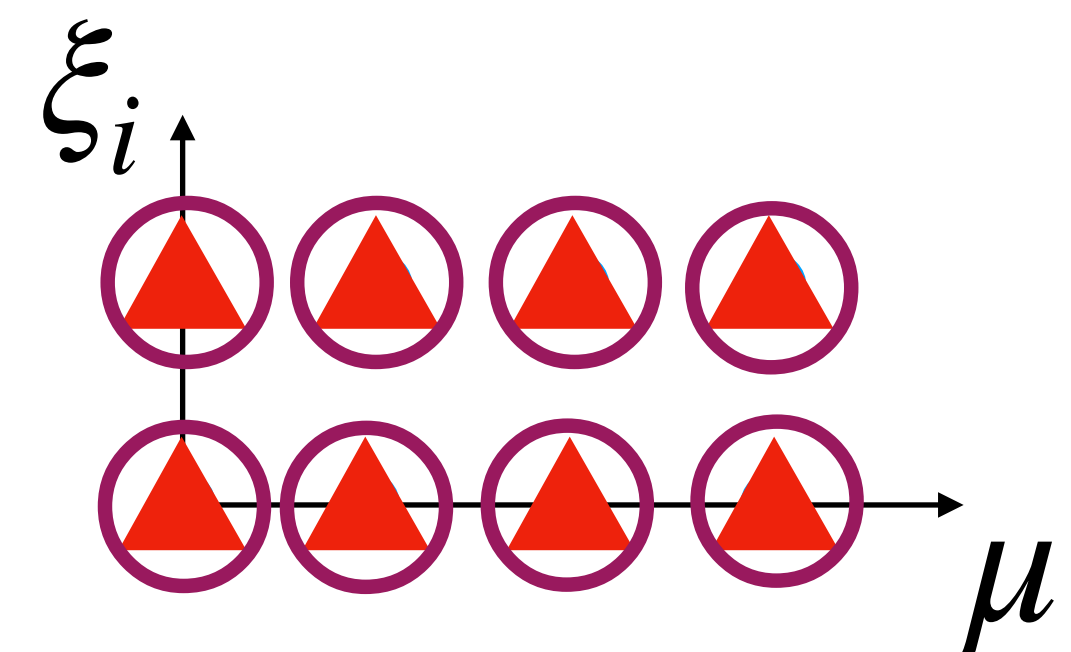
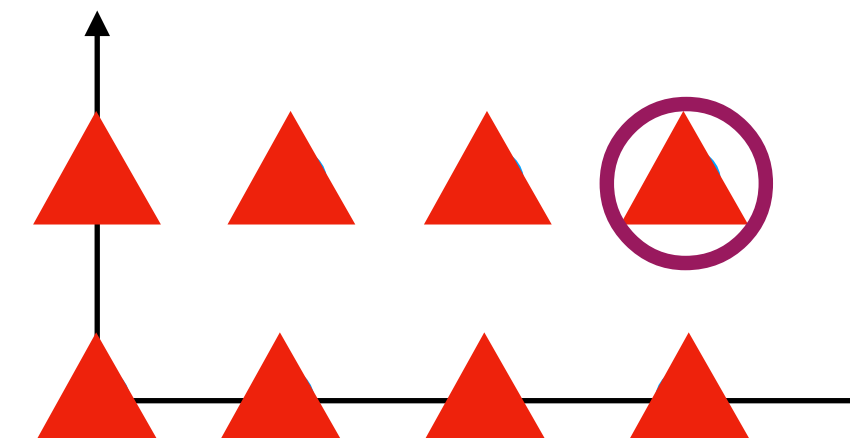
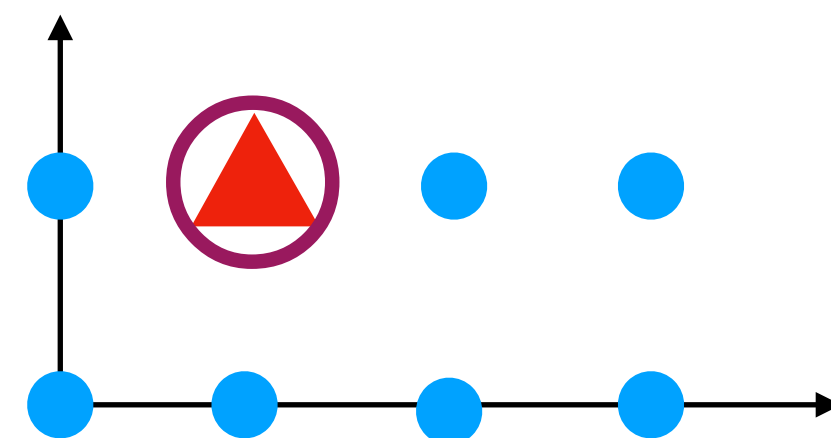
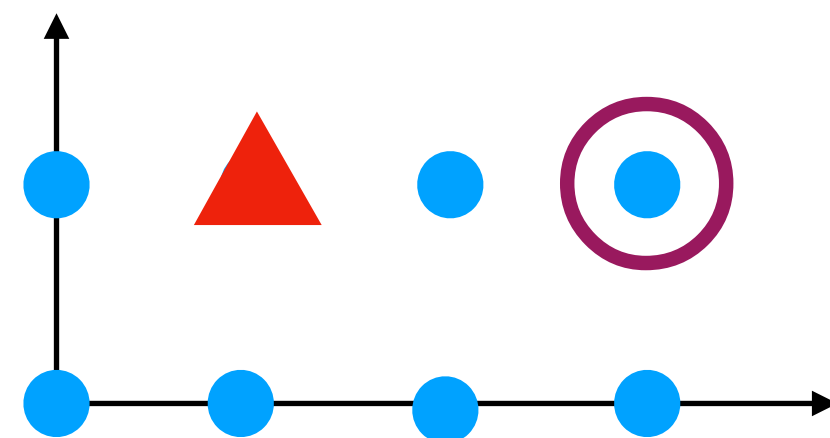
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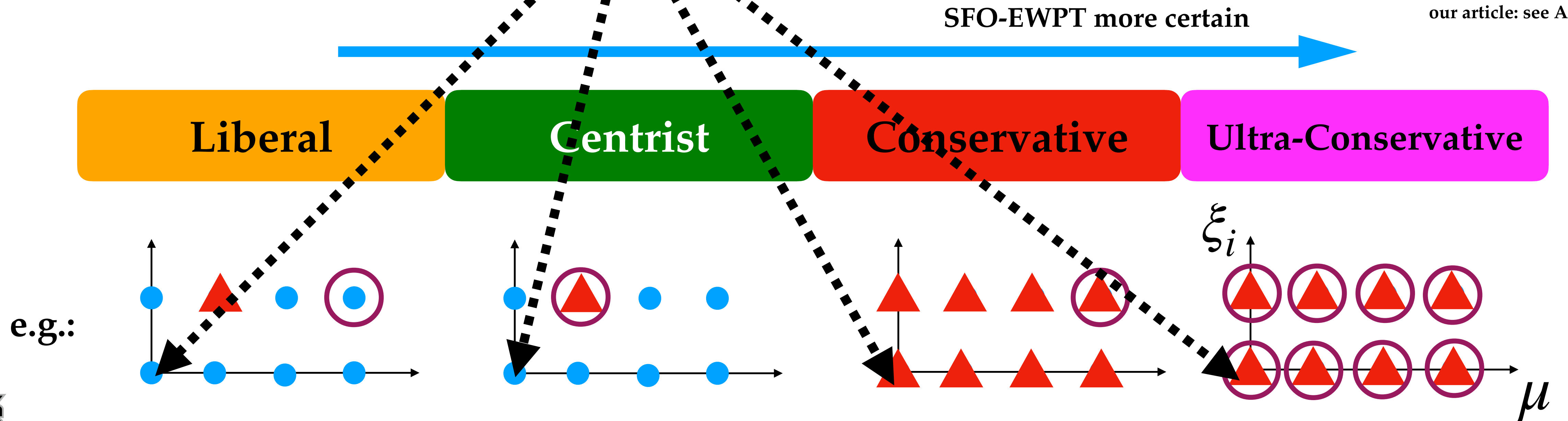
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1. Define two conditions:

Note: For phenomenological analyses, take “central” μ and ξ_i .

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Electro-Weak Precision Observables

- Real singlet scalar field
 - \rightarrow modifies Higgs contributions to diagonal weak gauge boson vacuum polarisation diagrams,
 - & introduces additional contributions.
- Quantify via S, T, U parameters. [Hagiwara, Matsumoto, Haidt, Kim, hep-ph/9409380]
- Change in EWPO \mathcal{O} ($= S, T, U$):

$$\Delta\mathcal{O} = (\mathcal{O}(m_2^2) - \mathcal{O}(m_1^2)) \times \sin^2 \theta$$

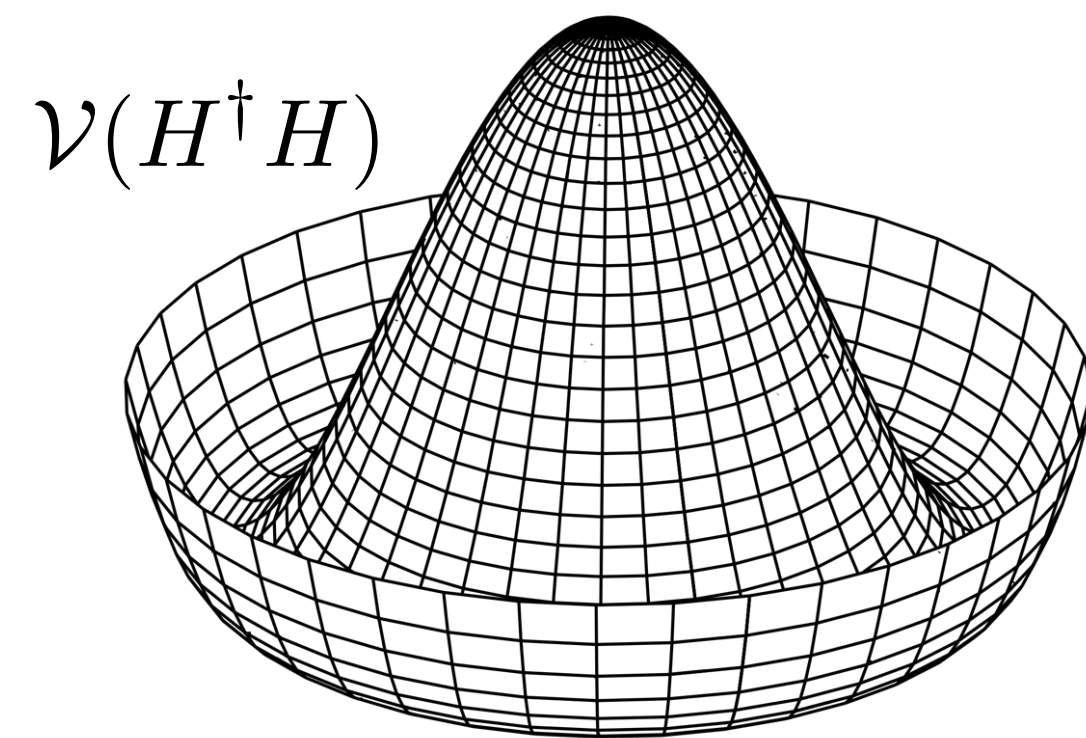
\Rightarrow calculate compatibility with experimental measurement $\Delta\mathcal{O}^{\text{EXP}}$.

The Higgs Potential & Vacuum Stability

The Importance of the Higgs sector

- the Higgs boson: the central protagonist of EWSB:

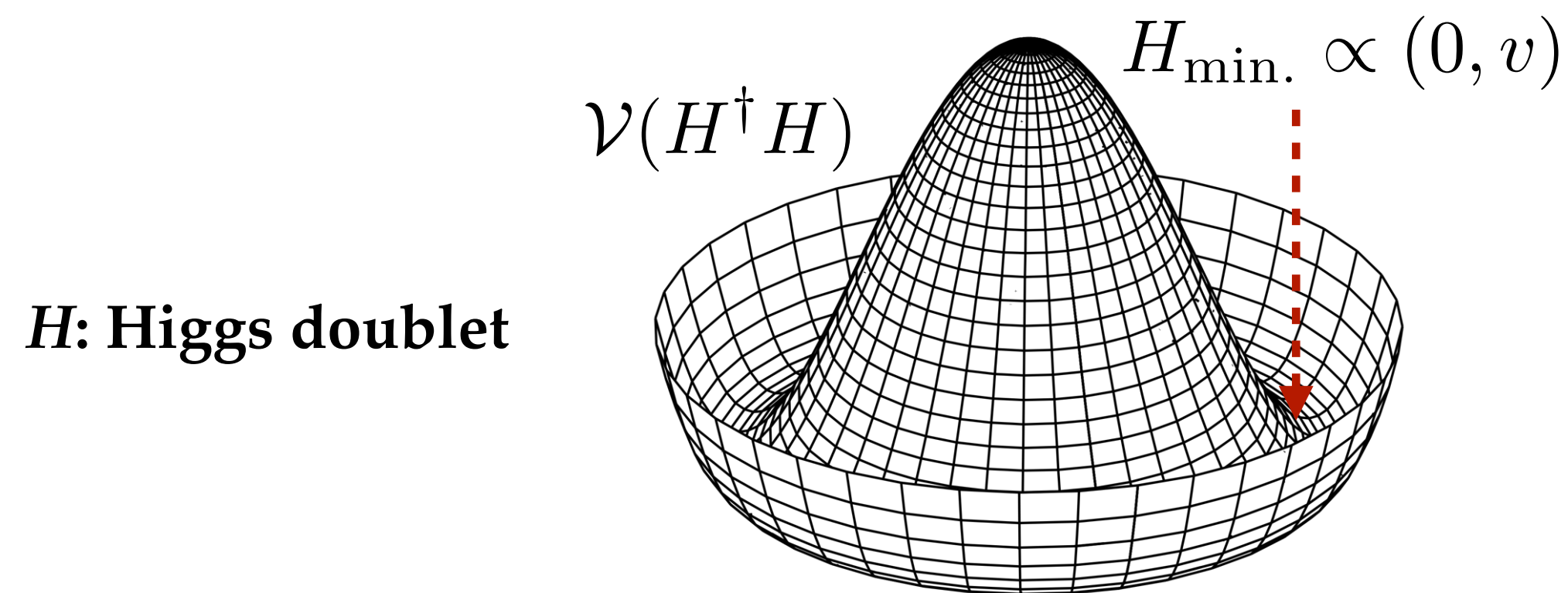
H: Higgs doublet



- an important characteristic of the Higgs boson is the way it couples to itself:

The Importance of the Higgs sector

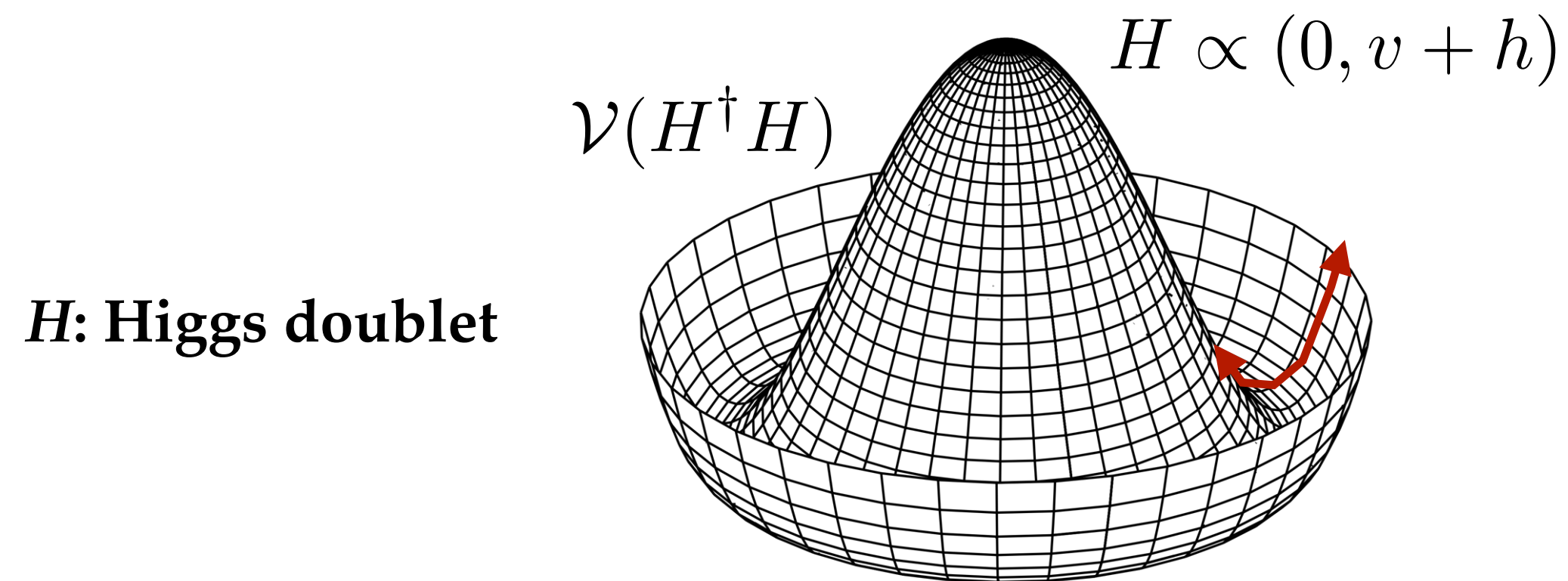
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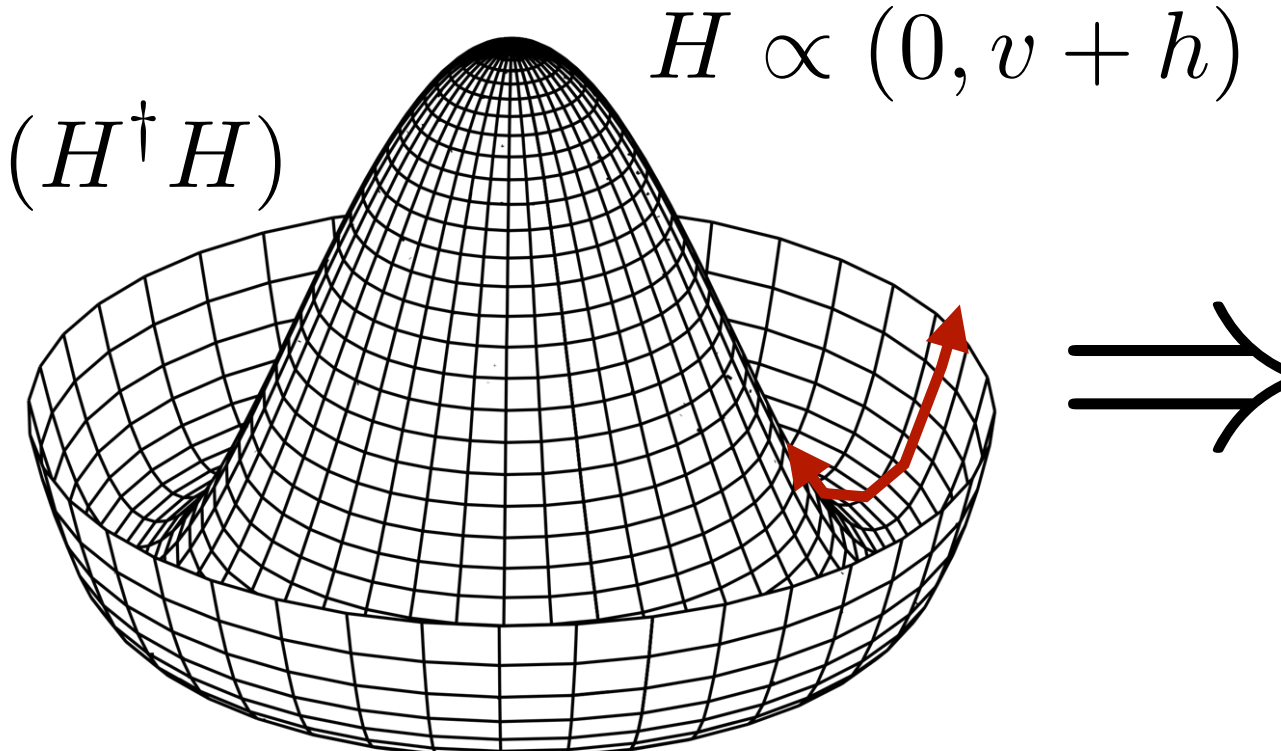
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$\mathcal{V}(H^\dagger H)$

$H \propto (0, v + h)$



e.g. fermion masses
& interactions:

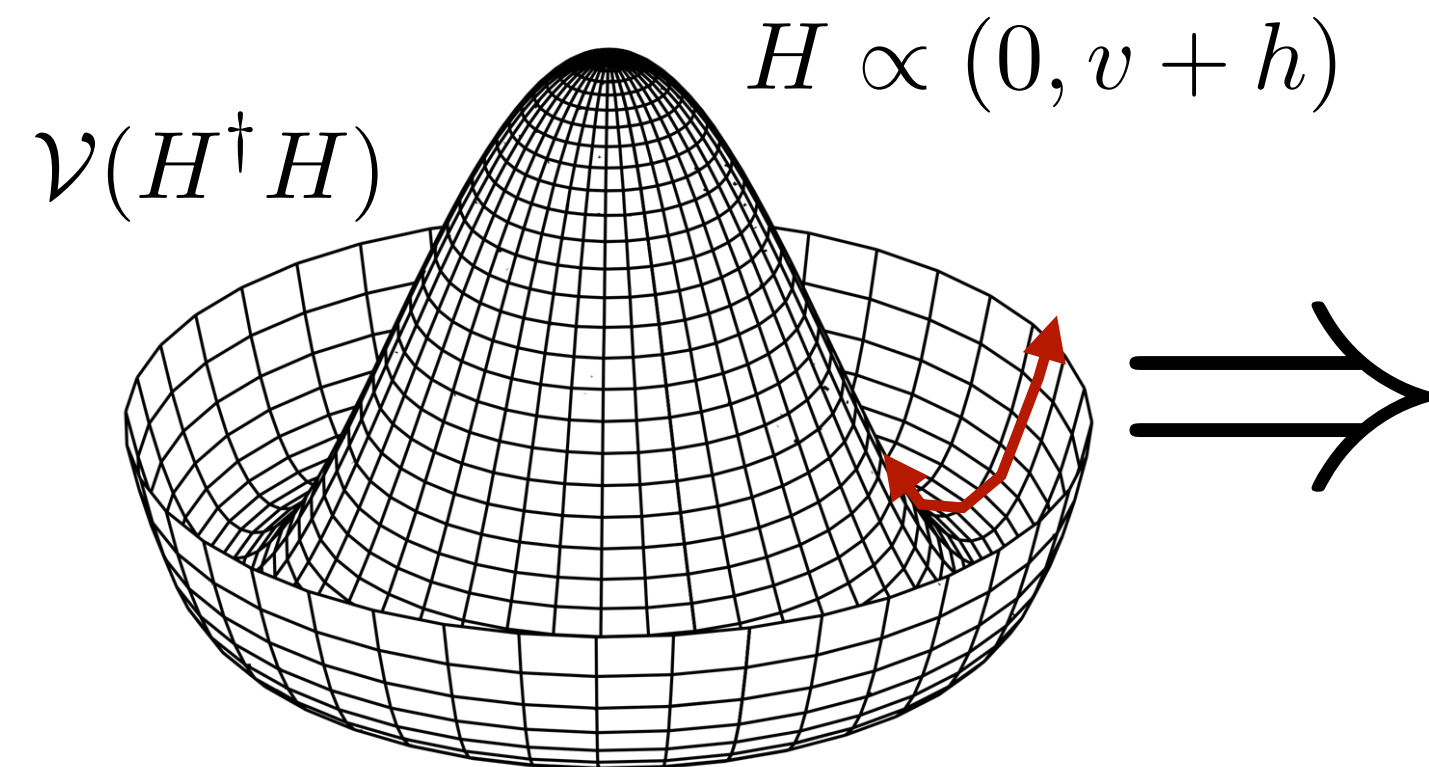
$$\mathcal{L} \supset -m_f \bar{f}_L f_R - \frac{m_f}{v} h \bar{f}_L f_R + \text{h.c.}$$

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e.g. gauge boson masses
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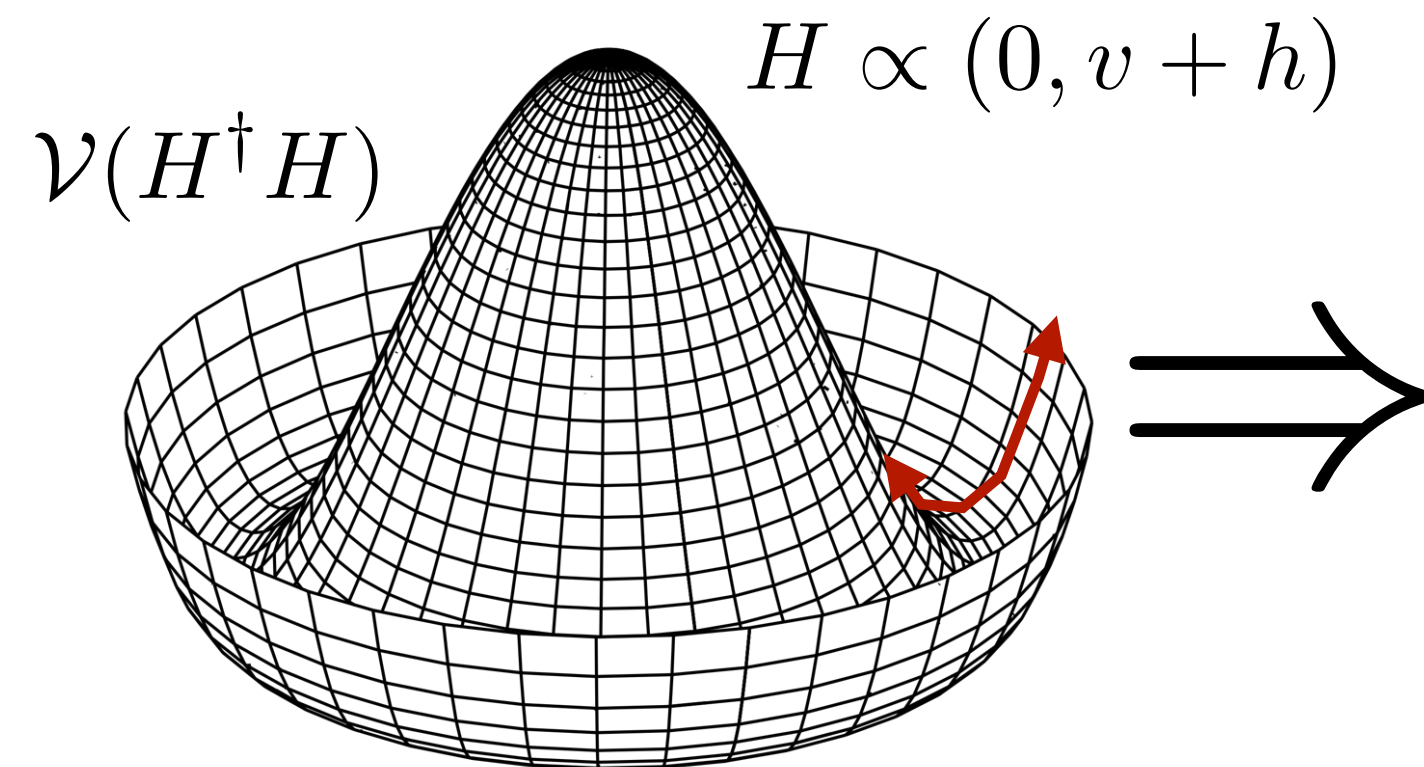
$$\mathcal{L} \supset [m_W^2 W^{\mu+} W_{\mu}^- + \frac{1}{2} m_Z^2 Z^{\mu} Z_{\mu}] \times \left(1 + \frac{h}{v}\right)^2$$

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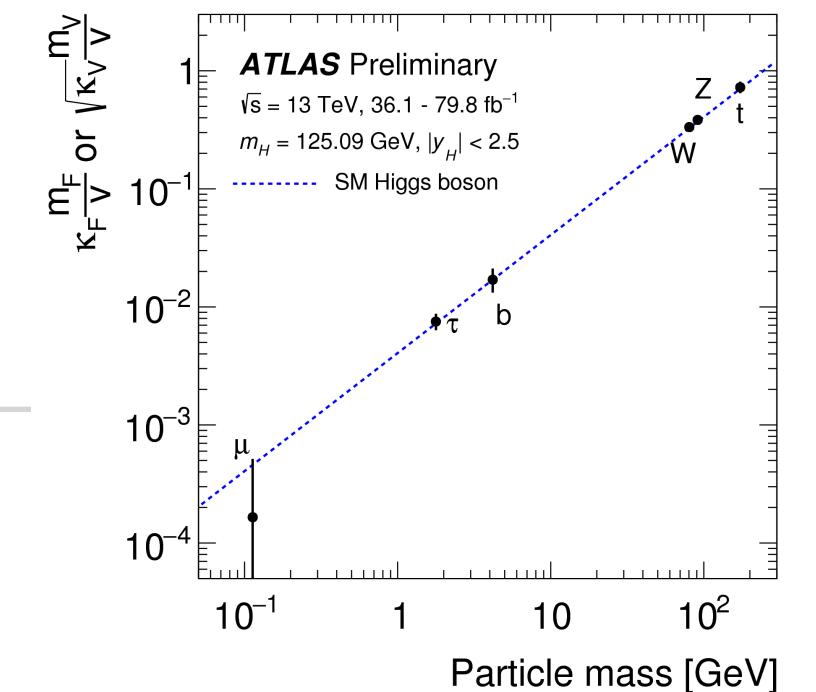
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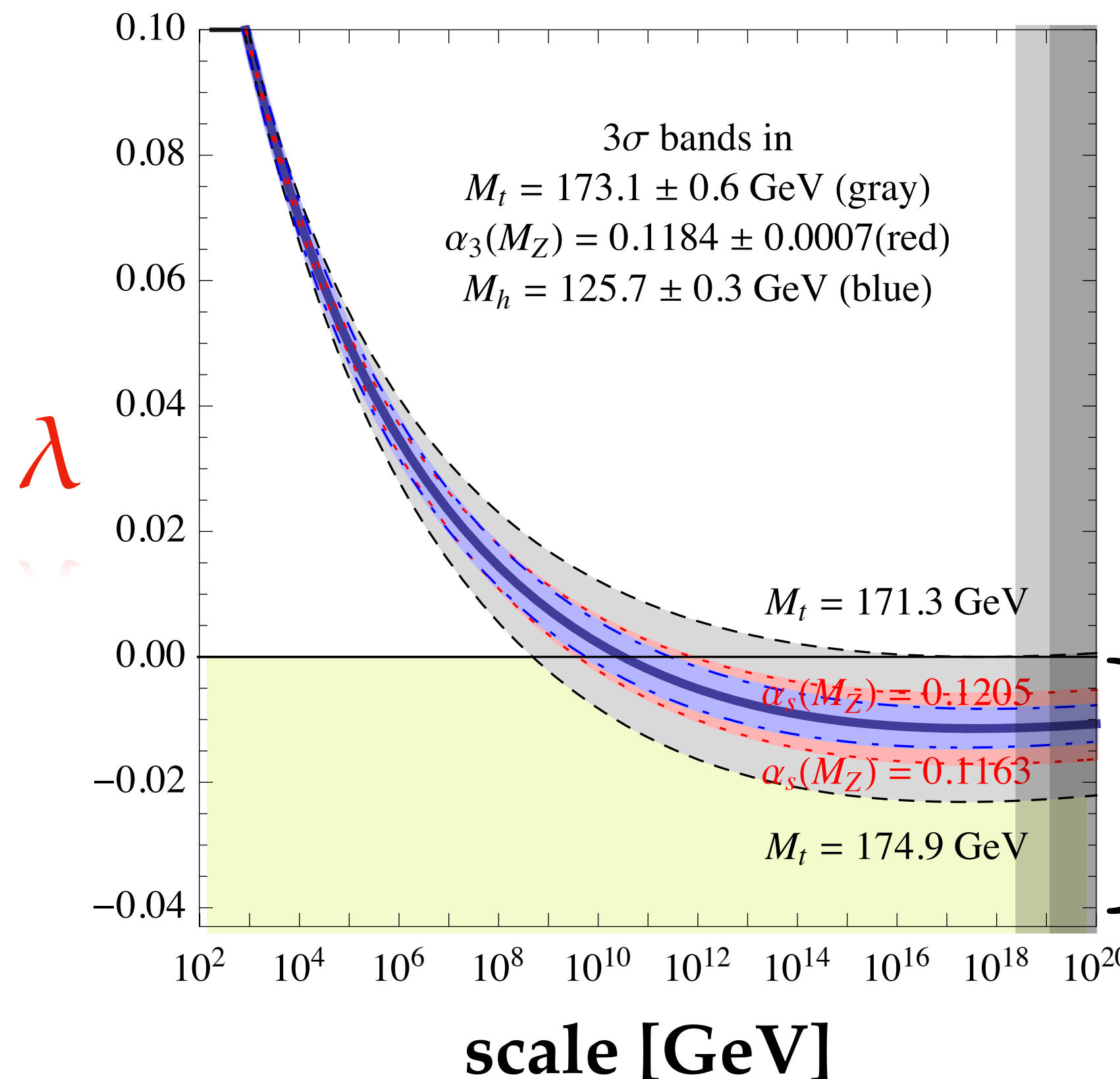
- an important characteristic of the Higgs boson is the way it couples to itself:

Vacuum Stability

- SM potential for the Higgs doublet:

$$\mathcal{V}(H^\dagger H) = -m^2(H^\dagger H) + \lambda(H^\dagger H)^2$$

- renormalisation group evolution of the coupling λ :

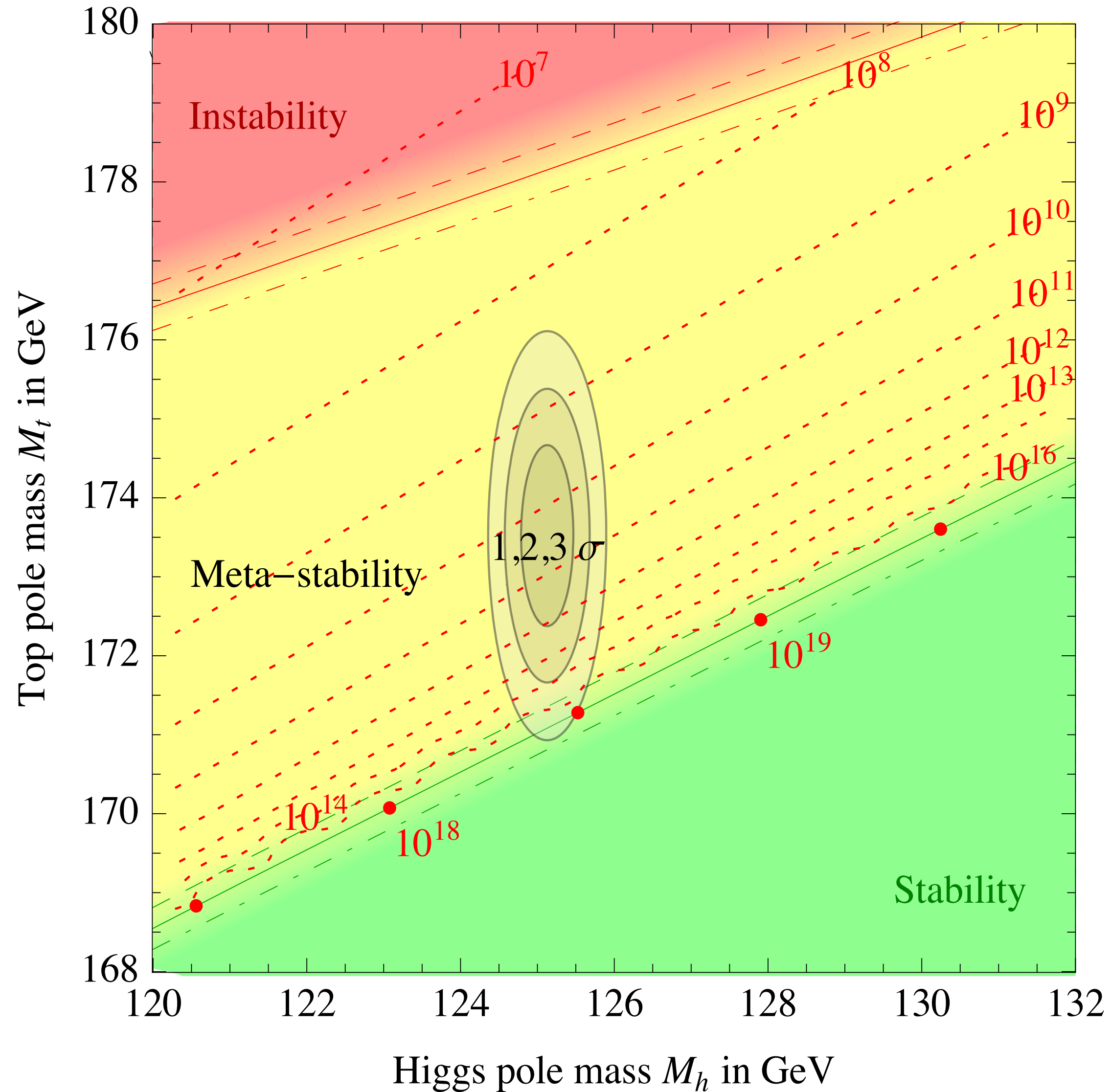


[Cabibbo, Maiani, Parisi, Petronzio, 1979, Hung, 1979, ..., Degrassi, Di Vita, Elias-Miró, Giudice, Isidori, Strumia, 1205.6497, Buttazzo, Degrassi, Giardino, Giudice, Sala, Salvio, Strumia, 1307.3536 ..., Espinosa, 1512.01222]

**potentially
unstable or
meta-stable
vacuum!**

vacuum stability

[Buttazzo, Degrassi, Giardino, Giudice, Sala, Salvio, Strumia, 1307.3536, Espinosa, 1512.01222]



in deriving this: **assumed Standard Model.**

➡ a hint for a non-standard Higgs sector?

➡ further investigation necessary.

Sphaleron/Instanton Processes

Instantons and Baryon-# Violation

- toy model:

(1+1)-dimensions, Abelian gauge field A^μ , complex scalar Φ , Dirac fermion of unit charge Ψ .

- Euclidean space action:

$$S = \int d^2x \left[\frac{1}{4} F_{\mu\nu}^2 + |(\partial_\mu - ieA_\mu)\Phi|^2 + V(\Phi) + i\bar{\Psi}(\partial_\mu - ieA_\mu)\gamma^\mu\Psi \right]$$

“Higgs potential”: $V(\Phi) = \lambda(\Phi^*\Phi - v^2)^2 \implies$ “EWSB” $\implies M_A, M_h$

Instantons and Baryon-# Violation

- consider the current:

$$K_\mu = \frac{e}{2\pi} \epsilon_{\mu\nu} A_\nu$$

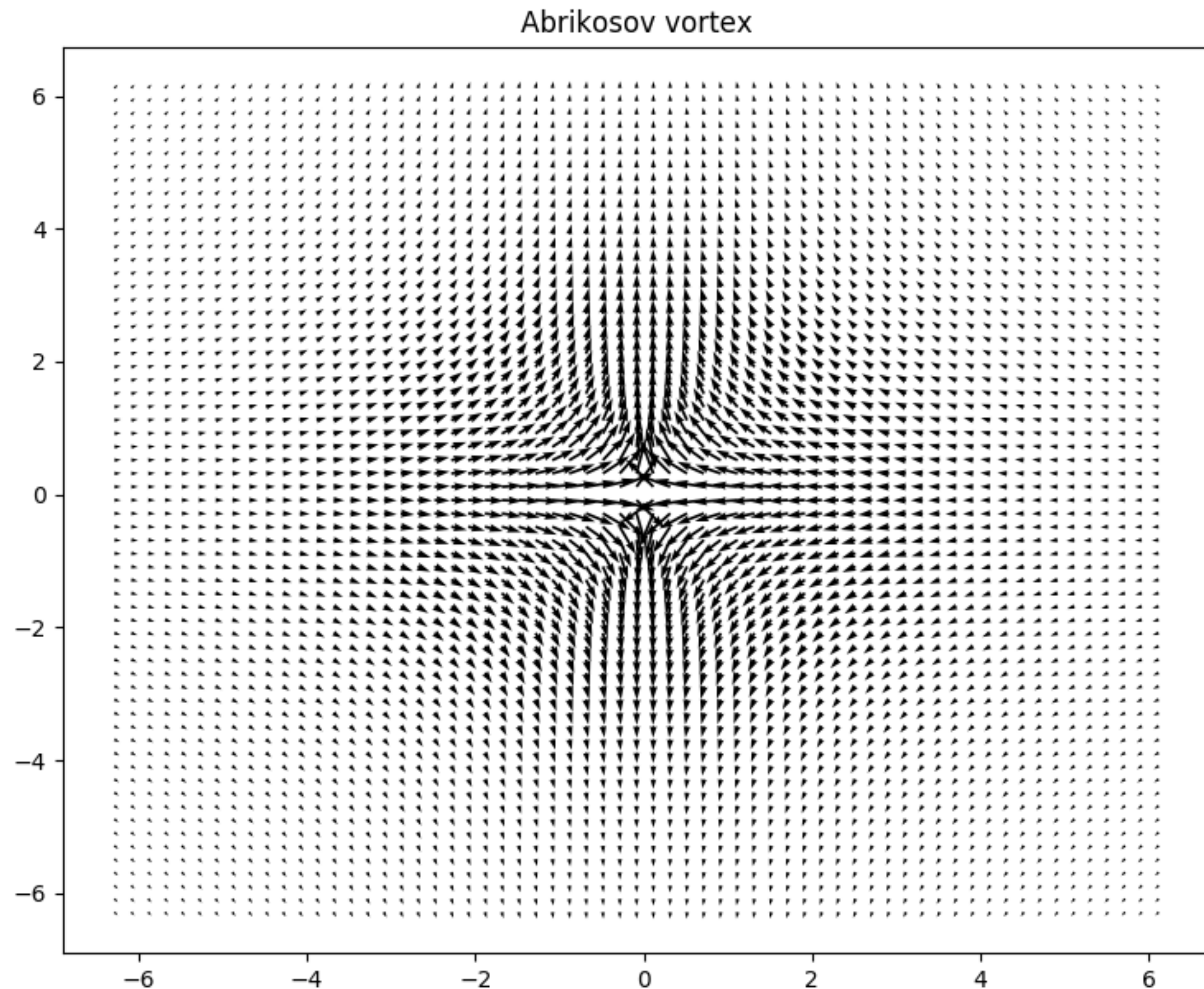
- corresponds to “charge density”:

$$N_{\text{CS}} = \int dx K_0 = \frac{e}{2\pi} \int dx A_1$$

- known as the “winding” or “Chern-Simons” number.

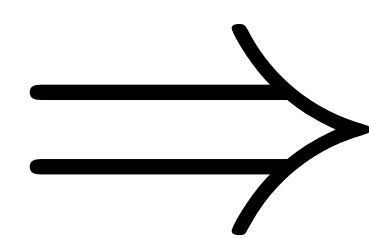
Instantons and Baryon-# Violation

- a classical solution to equations of motion is the “Abrikosov vortex”:



$$A_r = 0, \quad A_\theta = \frac{1}{er} f(r),$$

$$f(0) = 0, \quad 1 - f(r) \sim e^{-M_A r}$$



changes the Chern-Simons number by one unit:

$$\Delta N_{\text{CS}} = \int d^2x \partial_\mu K^\mu = \frac{e}{4\pi} \int d^2x \epsilon_{\mu\nu} F^{\mu\nu} = 1$$

Instantons and Baryon-# Violation

- “instanton” transition necessarily accompanied by change of chirality of fermions by two units:

$$j_{\mu}^5 = \bar{\Psi} \gamma_{\mu} \gamma_5 \Psi$$

$$\frac{1}{2} \partial_{\mu} j^{5\mu} = \frac{e}{4\pi} \epsilon_{\mu\nu} F^{\mu\nu} \quad \text{anomalous divergence of the axial-vector current.}$$

$$\Delta N_{\text{CS}} = \int d^2x \partial_{\mu} K^{\mu} = \frac{e}{4\pi} \int d^2x \epsilon_{\mu\nu} F^{\mu\nu} = 1$$

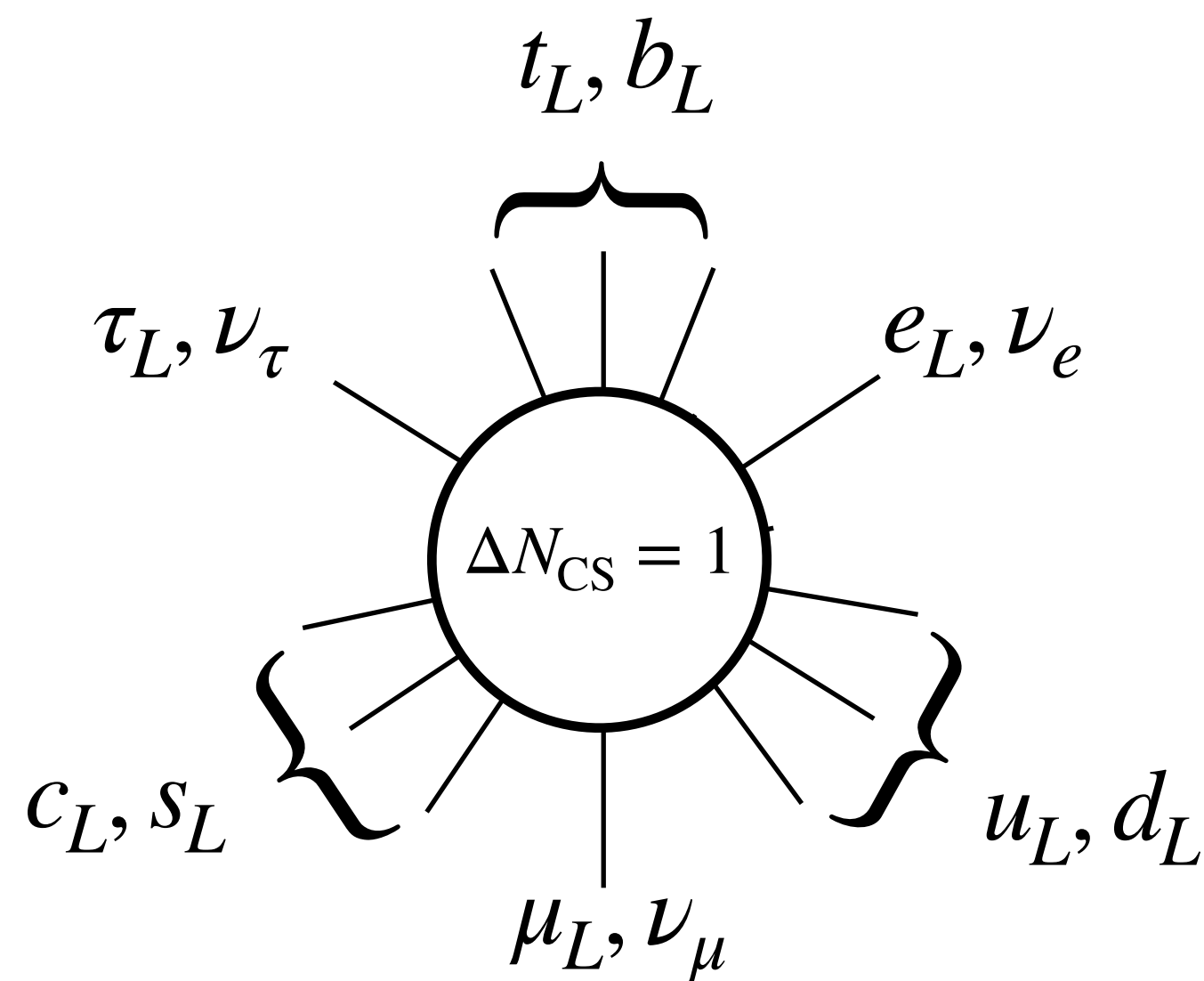
$$\Rightarrow \Delta Q_5 = \Delta \int dx j_0^5 = 2$$

EW Sphalerons at colliders?

- Rate and observability of sphaleron processes at colliders debated.

e.g. [Bezrukov, Levkov, Rebbi, Rybakov, Tinyakov, hep-ph/0304180] VS. [Tye, Wong, 1505.0360, 1710.07223].

- Ponder: **Sphaleron-induced interactions at hadron colliders:**

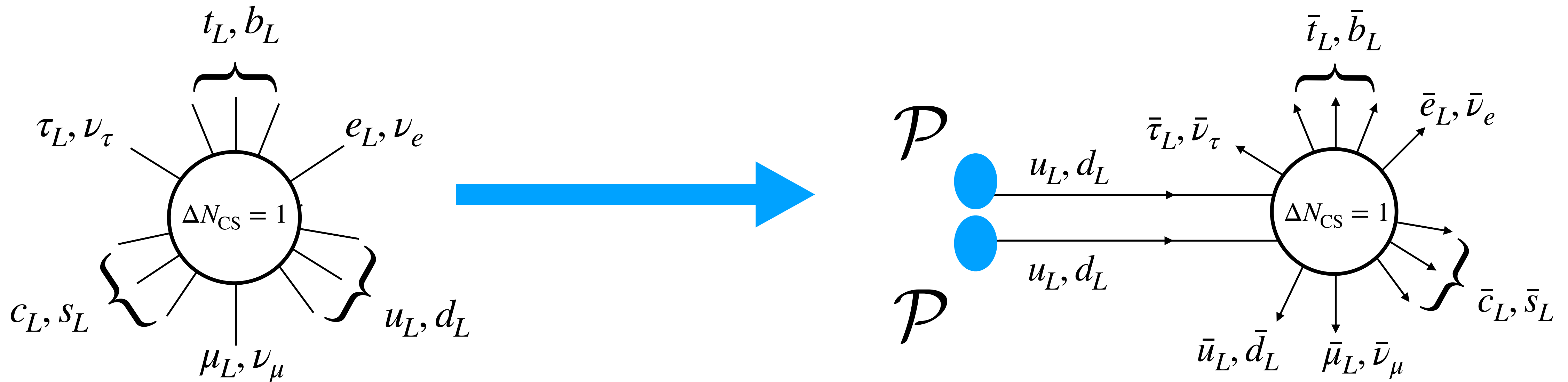


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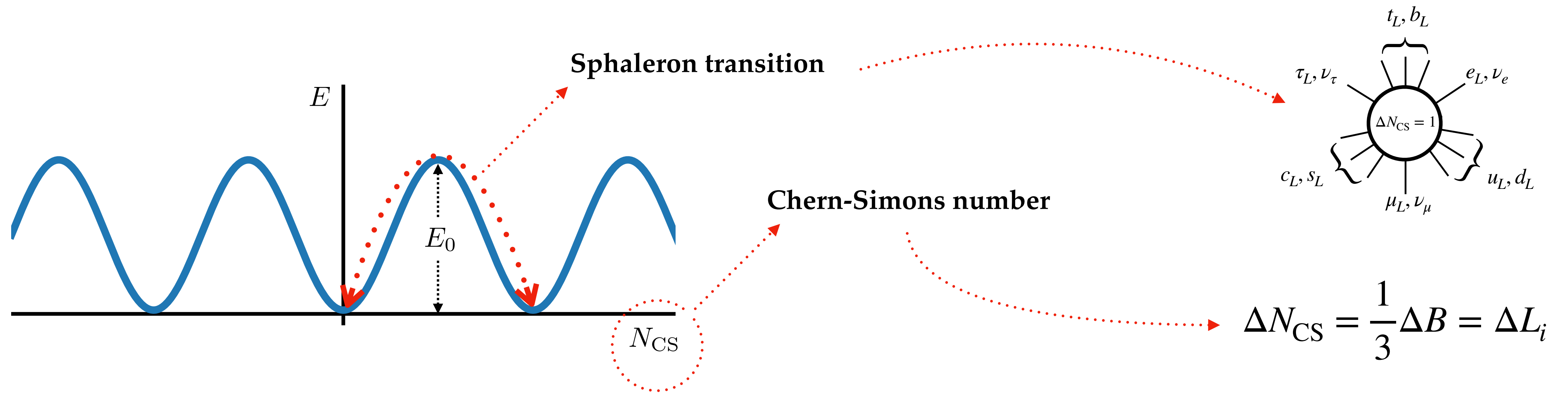
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What is the Sphaleron?



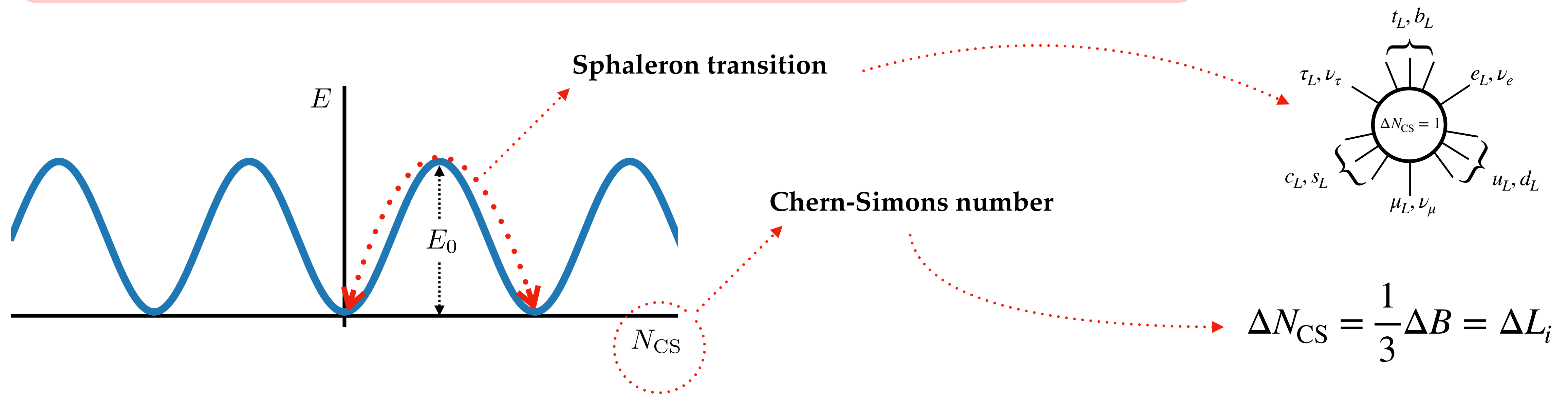
- $E_0 \sim \mathcal{O}(10)$ TeV, separates degenerate Electro-Weak vacua.



What is the Sphaleron?



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The Sphaleron energy depends crucially on the Higgs sector!

A Note on Sphaleron Suppression

- **Suppression of sphaleron rate inside bubble**

⇒ Baryon Asymmetry “swept in” broken phase and “frozen in”.

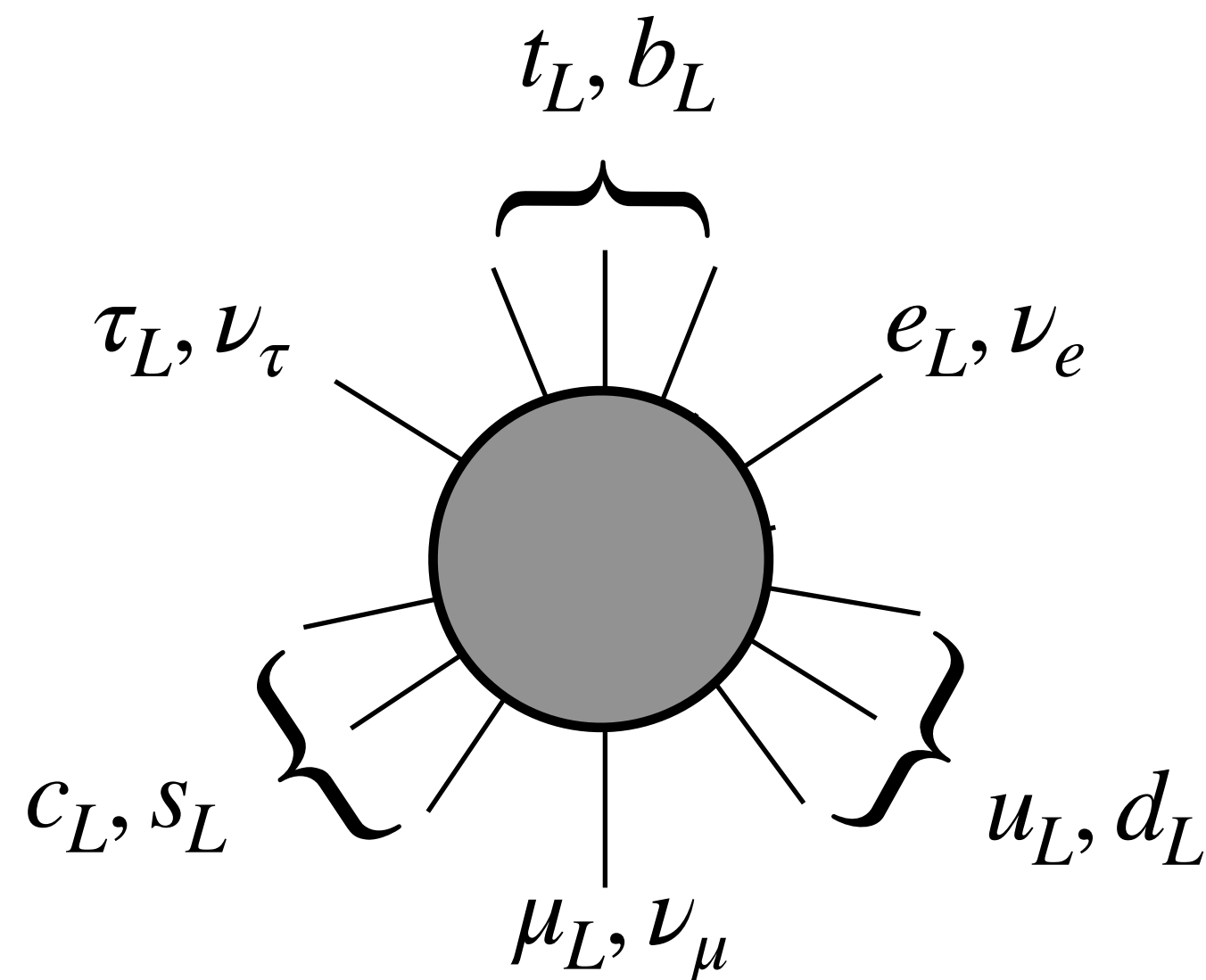
- **Rate** $\sim \exp[-\langle \phi(T_C) \rangle / T_C \times \dots]$,

[T_C : the critical temperature.]

- ⇒ Require: $\langle \phi(T_C) \rangle / T_C \geq 1 \Rightarrow$ a “**Strong**” First-Order EWPT (**SFO-EWPT**).

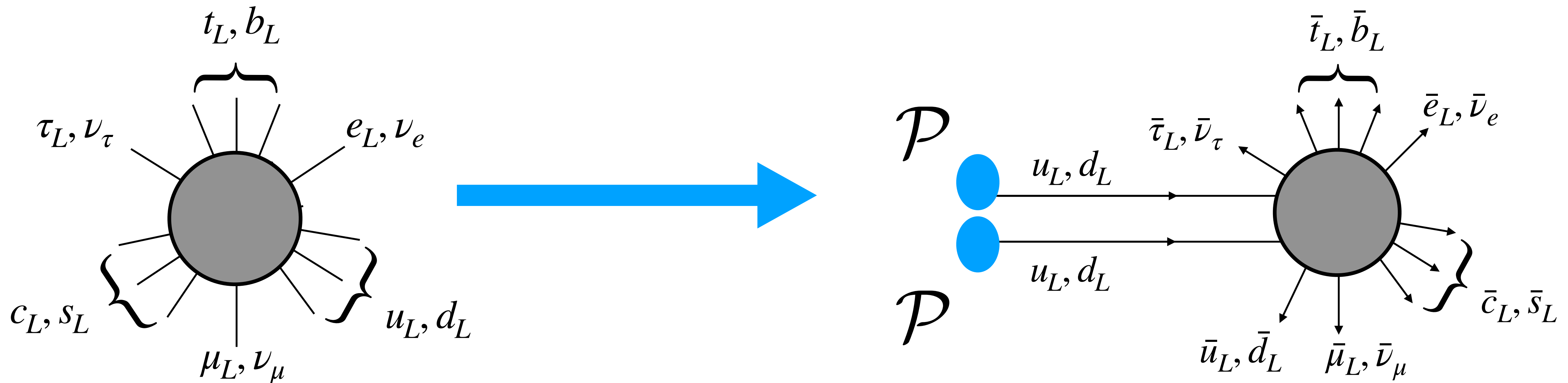
Sphaleron Suppression

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- Suppression requires “**Strong**” First-Order EWPT (**SFO-EWPT**).
- **Despite suppression: Sphalerons @ colliders?**

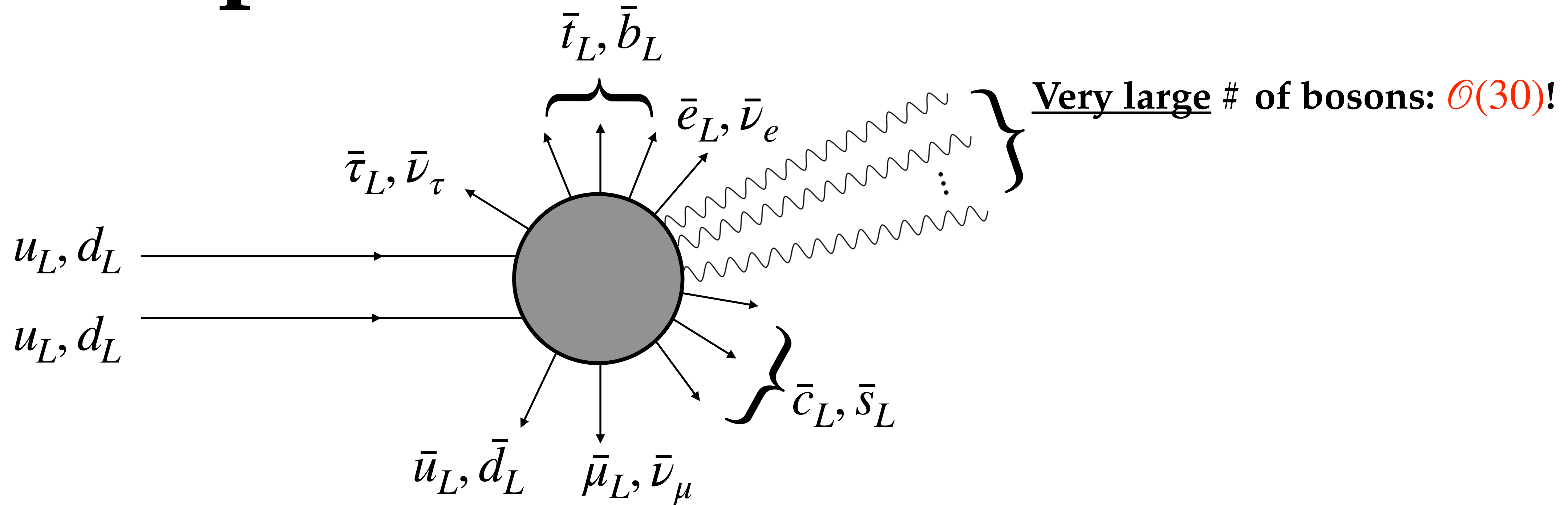


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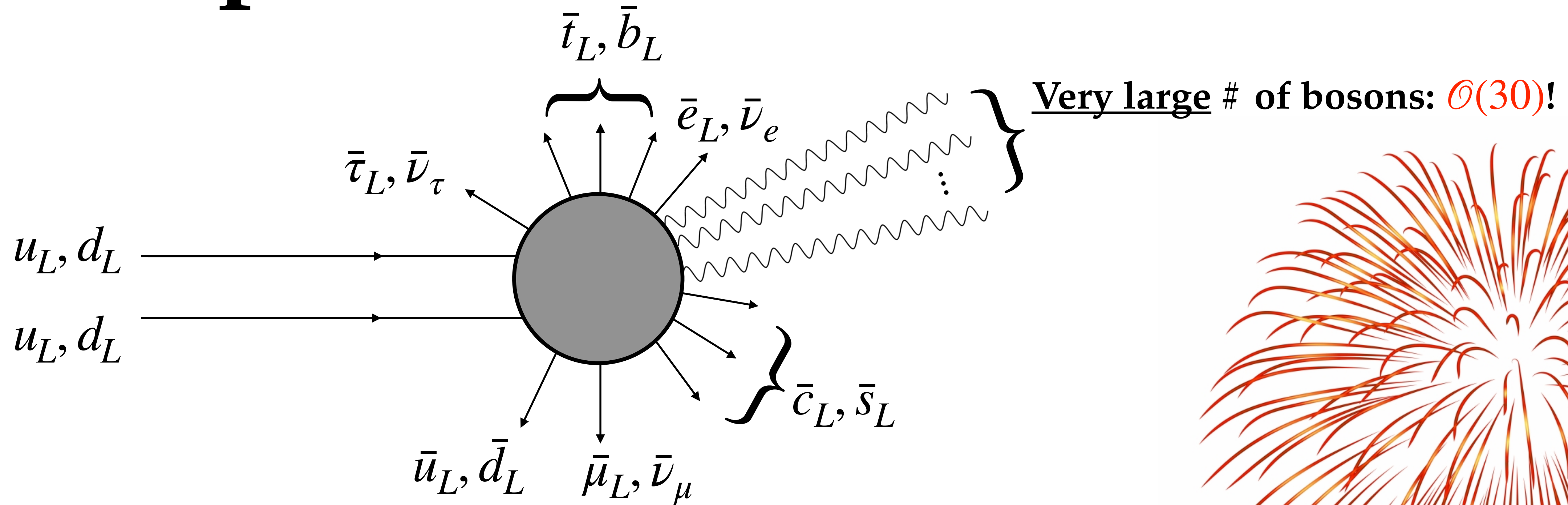
EW Sphalerons at Colliders?



- Possible enhancement if large number of bosons,

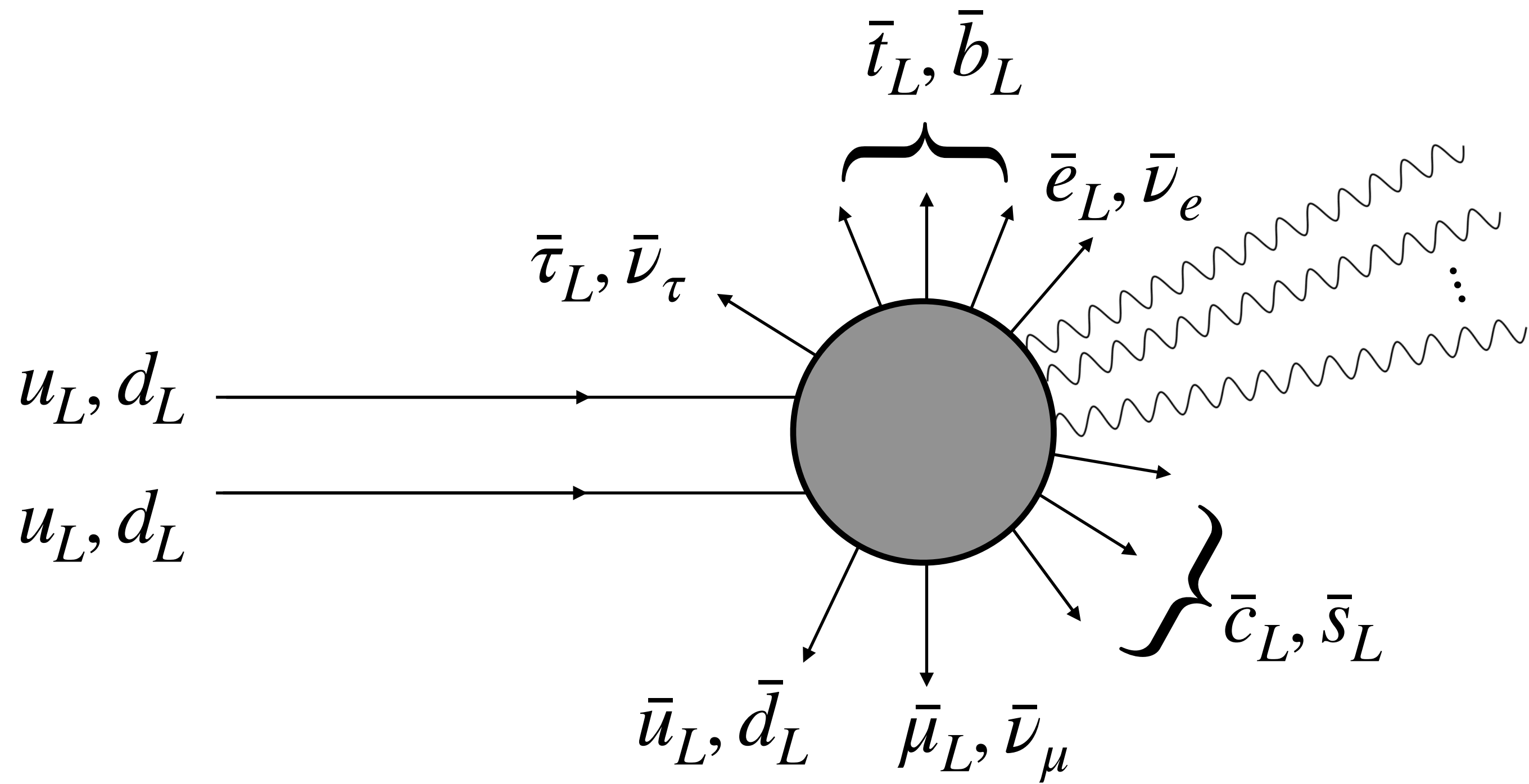
⇒ Events would **spectacularly light up detectors** at experiments!

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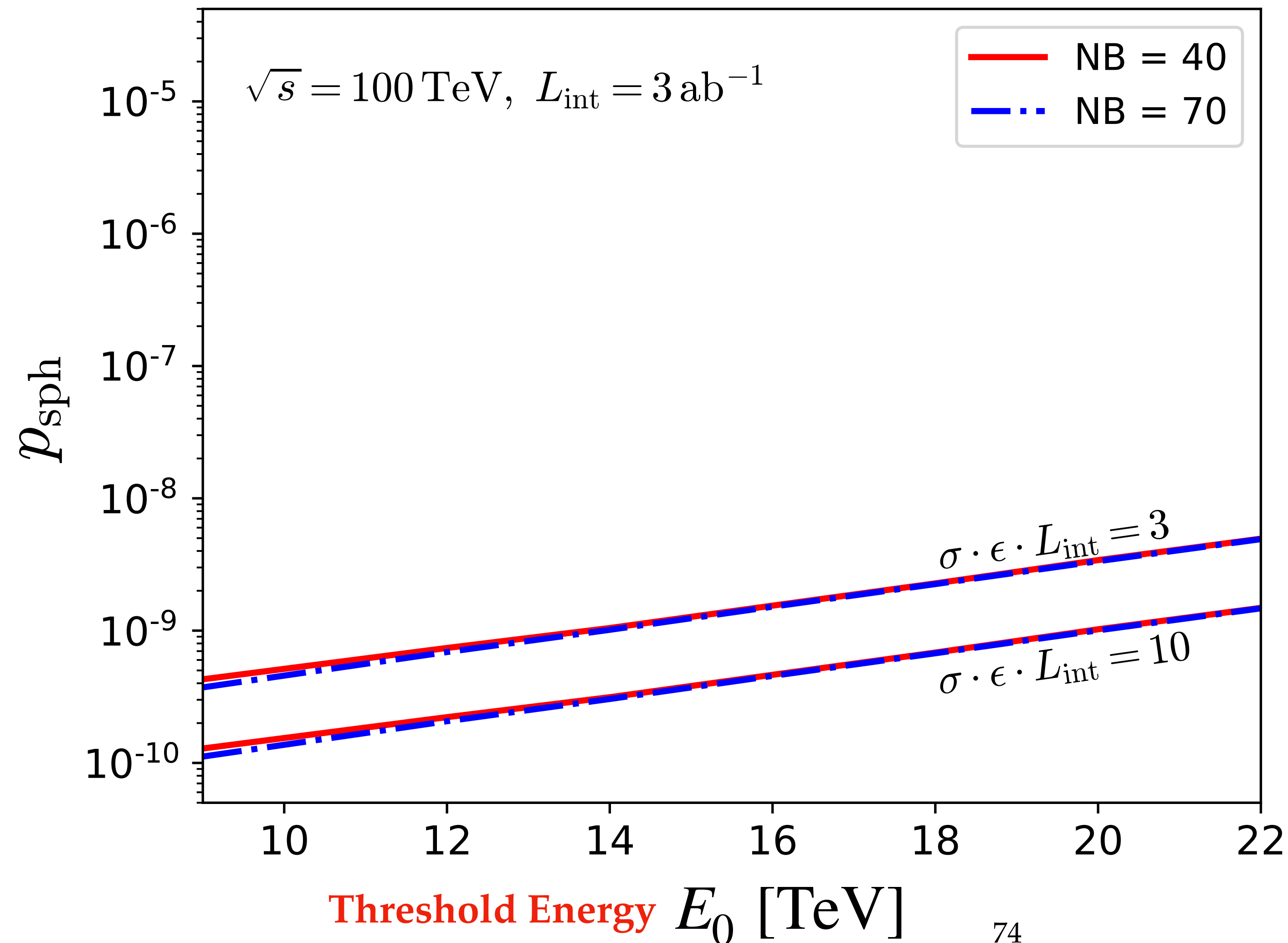
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Sphalerons at the FCC

- Parametrise parton-parton cross section by P_{sph} :

$$\hat{\sigma}(E) = \frac{P_{\text{sph}}}{m_W^2} \Theta(E - E_0)$$



→ Event Generator within HERWIG 7.

[[AP](#), Sakurai, Plätzer, [arXiv:1910.4761](#)]

EW Sphalerons at Colliders?

[[AP](#), Sakurai, Plätzer, [arXiv:1910.4761](#)]

⋮

EW Sphalerons at Colliders?

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Homework:

- (i) What can we learn about the Higgs sector and EWBG?
- (ii) New theoretical features in Sphaleron MC.
- (iii) Model discrimination, e.g. VS micro-black holes.
- (iv) Collaboration with experimentalists for measurements.

