## March 24 Math 2306 sec 59 Spring 2016

## Section 10: Variation of Parameters

For the equation in standard form

$$
\frac{d^{2} y}{d x^{2}}+P(x) \frac{d y}{d x}+Q(x) y=g(x)
$$

suppose $\left\{y_{1}(x), y_{2}(x)\right\}$ is a fundamental solution set for the associated homogeneous equation. We seek a particular solution of the form

$$
y_{p}(x)=u_{1}(x) y_{1}(x)+u_{2}(x) y_{2}(x)
$$

where $u_{1}$ and $u_{2}$ are functions we will determine (in terms of $y_{1}, y_{2}$ and g).

This method is called variation of parameters.

## Variation of Parameters: Derivation of $y_{p}$

$$
y^{\prime \prime}+P(x) y^{\prime}+Q(x) y=g(x)
$$

Set $y_{p}=u_{1}(x) y_{1}(x)+u_{2}(x) y_{2}(x)$

We found that the functions $u_{1}$ and $u_{2}$ were determined as

$$
u_{1}(x)=\int \frac{-y_{2}(x) g(x)}{W} d x \quad \text { and } \quad u_{2}(x)=\int \frac{y_{1}(x) g(x)}{W} d x
$$

where $W$ is the Wronskian of $y_{1}$ and $y_{2}$.

Example:
Solve the ODE

$$
x^{2} y^{\prime \prime}+x y^{\prime}-y=x^{2}
$$

given that $y_{1}=x$ is one solution to the associated homogeneous equation.

We need $y_{1}, y_{2}$ that solve $x^{2} y^{\prime \prime}+x y^{\prime}-y=0$
in standard form $y^{\prime \prime}+\frac{1}{x} y^{\prime}-\frac{1}{x^{2}} y=0$

We can get $y_{2}$ by reduction of arden.

$$
y_{2}=u(x) y_{1}(x) \text { where } u(x)=\int \frac{e^{-\int \rho(x) d x}}{\left(y_{1}(x)\right)^{2}} d x
$$

$$
\begin{aligned}
& P(x)=\frac{1}{x} \quad-\int P(x) d x=-\int \frac{1}{x} d x=-\ln |x| \\
&=\ln x^{-1} \quad \text { assuming } \\
& x>0
\end{aligned} \quad \begin{aligned}
& u=\int \frac{e^{\ln x^{-1}}}{(x)^{2}} d x=\int \frac{x^{-1}}{x^{2}} d x=\int x^{-3} d x=\frac{-1}{2} x^{-2} \\
& y_{2}=u y_{1}=\frac{-1}{2} x^{-2} \cdot x=\frac{-1}{2} x^{-1}
\end{aligned}
$$

By super position, we con take $y_{2}=x^{-1}$.

The ronkomogeneows egn is

$$
x^{2} y^{\prime \prime}+x y^{\prime}-y=x^{2}, y_{1}=x, y_{2}=x^{-1}
$$

Standerd form:

$$
y^{\prime \prime}+\frac{1}{x} y^{\prime}-\frac{1}{x^{2}} y=1 \Rightarrow g(x)=1
$$

Wronskien:

$$
W=\left|\begin{array}{cc}
x & x^{-1} \\
1 & -x^{-2}
\end{array}\right|=x\left(-x^{-2}\right)-1 \cdot x^{-1}=-x^{-1}-x^{-1}=-2 x^{-1}
$$

$$
\begin{aligned}
u_{1} & =\int \frac{-y_{2}(x) g(x)}{w} d x=\int \frac{-x^{-1} \cdot 1}{-2 x^{-1}} d x=\frac{1}{2} \int \frac{x^{-1}}{x^{-1}} d x \\
& =\frac{1}{2} \int d x=\frac{1}{2} x \\
u_{2} & =\int \frac{y_{1}(x) g(x)}{w} d x=\int \frac{x \cdot 1}{-2 x^{-1}} d x=-\frac{1}{2} \int \frac{x}{x^{-1}} d x \\
& =\frac{-1}{2} \int x \cdot x d x=\frac{-1}{2} \int x^{2} d x=\frac{-1}{2} \cdot \frac{x^{3}}{3}=\frac{-x^{3}}{6}
\end{aligned}
$$

$$
\begin{aligned}
y_{p} & =u_{1} y_{1}+u_{2} y_{2} \\
& =\frac{1}{2} x(x)+\left(\frac{-x^{3}}{6}\right) x^{-1} \\
& =\frac{1}{2} x^{2}-\frac{1}{6} x^{2}=\left(\frac{1}{2}-\frac{1}{6}\right) x^{2}=\left(\frac{3}{6}-\frac{1}{6}\right) x^{2} \\
& =\frac{2}{6} x^{2}=\frac{1}{3} x^{2}
\end{aligned}
$$

The gerard solution to the ODE is

$$
y=y_{c}+y_{1}, \quad y=c_{1} x+c_{2} x^{-1}+\frac{1}{3} x^{2}
$$

Solve the IVP

$$
x^{2} y^{\prime \prime}+x y^{\prime}-y=x^{2}, \quad y(1)=-1, \quad y^{\prime}(1)=0
$$

From the last example, the general solution to the ODE is

$$
y=c_{1} x+c_{2} x^{-1}+\frac{1}{3} x^{2} .
$$

$$
y^{\prime}=c_{1}-c_{2} x^{-2}+\frac{2}{3} x
$$

$$
\begin{aligned}
& y(1)=-1 \\
&=c_{1} \cdot 1+c_{2} \frac{1}{1}+\frac{1}{3}(1)^{2} \\
&-1=c_{1}+c_{2}+\frac{1}{3} \Rightarrow c_{1}+c_{2}=\frac{-4}{3} \\
& y^{\prime}(1)=c_{1}-c_{2} \frac{1}{1^{2}}+\frac{2}{3} \cdot 1=0 \Rightarrow c_{1}-c_{2}=\frac{-2}{3} \\
& c_{1}+c_{2}=\frac{-4}{3} \\
& c_{1}-c_{2}=\frac{-2}{3} \\
& 2 c_{1}=\frac{-6}{3}=-2 \Rightarrow c_{1}=-1 \\
& c_{2}=\frac{-4}{3}-c_{1}=\frac{-4}{3}-(-1)=\frac{-1}{3}
\end{aligned}
$$

add

The solution to the IVP is

$$
y=-x-\frac{1}{3} x^{-1}+\frac{1}{3} x^{2} .
$$

## Section 11: Linear Mechanical Equations

## Simple Harmonic Motion

We consider a flexible spring from which a mass is suspended. In the absence of any damping forces (e.g. friction, a dash pot, etc.), and free of any external driving forces, any initial displacement or velocity imparted will result in free, undamped motion-a.k.a. simple harmonic motion.

## Building an Equation: Hooke's Law



At equilibrium, displacement $x(t)=0$.
Hooke's Law: $\mathrm{F}_{\text {spring }}=k \mathrm{x}$
Figure: In the absence of any displacement, the system is at equilibrium. Displacement $x(t)$ is measured from equilibrium $x=0$.

## Building an Equation: Hooke's Law

Newton's Second Law: $F=$ ma Force $=$ mass times acceleration

$$
a=\frac{d^{2} x}{d t^{2}} \quad \Longrightarrow \quad F=m \frac{d^{2} x}{d t^{2}}
$$

Hooke's Law: $F=k x$ Force exerted by the spring is proportional to displacement
The force imparted by the spring opposes the direction of motion.

$$
\begin{gathered}
m x^{\prime \prime}+k x=0 \quad \Rightarrow x^{\prime \prime}+\frac{k}{m} x=0 \\
m \frac{d^{2} x}{d t^{2}}=-k x \quad \Longrightarrow \quad x^{\prime \prime}+\omega^{2} x=0 \quad \text { where } \quad \omega=\sqrt{\frac{k}{m}}
\end{gathered}
$$

Convention We'll Use: Up will be positive ( $x>0$ ), and down will be negative $(x<0)$. This orientation is arbitrary and follows the convention in Trench.

Obtaining the Spring Constant (US Customary Units)
If an object with weight $W$ pounds stretches a spring $\delta x$ feet ${ }^{1}$ from it's length with no mass attached, then by Hooke's law we compute the spring constant via the equation

$$
W=k \delta x
$$

The units for $k$ in this system of measure are $\mathrm{lb} / \mathrm{ft}$.

$$
w \left\lvert\, b=k \delta \dot{x} f t \Rightarrow k=\frac{w}{\delta x} \frac{b b}{f t}\right.
$$

${ }^{1}$ Note that $\delta x=\mathrm{w} /$ mass equilibrium - $\mathrm{w} / \mathrm{o}$ mass equilibrium.

Obtaining the Spring Constant (US Customary Units)
Note also that Weight $=$ mass $\times$ acceleration due to gravity. Hence if we know the weight of an object, we can obtain the mass via

$$
W=m g .
$$

We typically take the approximation $g=32 \mathrm{ft} / \mathrm{sec}^{2}$. The units for mass are $\mathrm{lb} \mathrm{sec}^{2} / \mathrm{ft}$ which are called slugs.

$$
w \left\lvert\, b=m g \frac{f t}{\sec ^{2}} \Rightarrow m=\frac{w}{g} \frac{\frac{b}{f-1} / \sec ^{2}}{q}=\frac{w}{g} \frac{\mathrm{lbsec}^{2}}{f t}\right.
$$

## Obtaining the Spring Constant (SI Units)

In SI units, the weight would be expressed in Newtons (N). The appropriate units for displacement would be meters ( m ). In these units, the spring constant would have units of $\mathrm{N} / \mathrm{m}$.

It is customary to describe an object by its mass in kilograms. When we encounter such a description, we deduce the weight in Newtons

$$
W=m g \text { taking the approximation } g=9.8 \mathrm{~m} / \mathrm{sec}^{2} .
$$

## Obtaining the Spring Constant: Displacment in Equilibrium

If an object stretches a spring $\delta x$ units from it's length (with no object attached), we may say that it stretches the spring $\delta x$ units in equilibrium. Applying Hooke's law with the weight as force, we have

$$
m g=k \delta x
$$

We observe that the value $\omega$ can be deduced from $\delta x$ by

$$
\omega^{2}=\frac{k}{m}=\frac{g}{\delta x}
$$

Provided that values for $\delta x$ and $g$ are used in appropriate units, $\omega$ is in units of per second.

Simple Harmonic Motion


Here, $x_{0}$ and $x_{1}$ are the initial position (relative to equilibrium) and velocity, respectively. The solution is

$$
x(t)=x_{0} \cos (\omega t)+\frac{x_{1}}{\omega} \sin (\omega t)
$$

called the equation of motion.

Caution: The phrase equation of motion is used differently by different authors. Some, including Trench, use this phrase to refer the ODE of which (1) would be the example here. Others use it to refer to the solution to the associated IVP.
$x(t)=x_{0} \cos (\omega t)+\frac{x_{1}}{\omega} \sin (\omega t)$

Characteristics of the system include

- the period $T=\frac{2 \pi}{\omega}$,
- the frequency $f=\frac{1}{T}=\frac{\omega}{2 \pi} 2$
- the circular (or angular) frequency $\omega$, and
- the amplitude or maximum displacement $A=\sqrt{x_{0}^{2}+\left(x_{1} / \omega\right)^{2}}$
${ }^{2}$ Various authors call $f$ the natural frequency and others use this term for $\omega$. $\overline{\bar{z}}$


## Amplitude and Phase Shift

We can formulate the solution in terms of a single sine (or cosine) function. Letting

$$
x(t)=x_{0} \cos (\omega t)+\frac{x_{1}}{\omega} \sin (\omega t)=A \sin (\omega t+\phi)
$$

requires

$$
A=\sqrt{x_{0}^{2}+\left(x_{1} / \omega\right)^{2}}
$$

and the phase shift $\phi$ must be defined by

$$
\sin \phi=\frac{x_{0}}{A}, \quad \text { with } \quad \cos \phi=\frac{x_{1}}{\omega A} .
$$

