

Imprints of Electro-Weak Symmetry Breaking @ Particle Colliders



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@ Wichita State University [May 1st 2024]



[made with instability.ai]



Today's Plan:

- 1 High Energy Physics
- 2 Open Questions,
- 3 → the Breaking of Symmetry,
- 4 Extended Scalar Sectors.



1

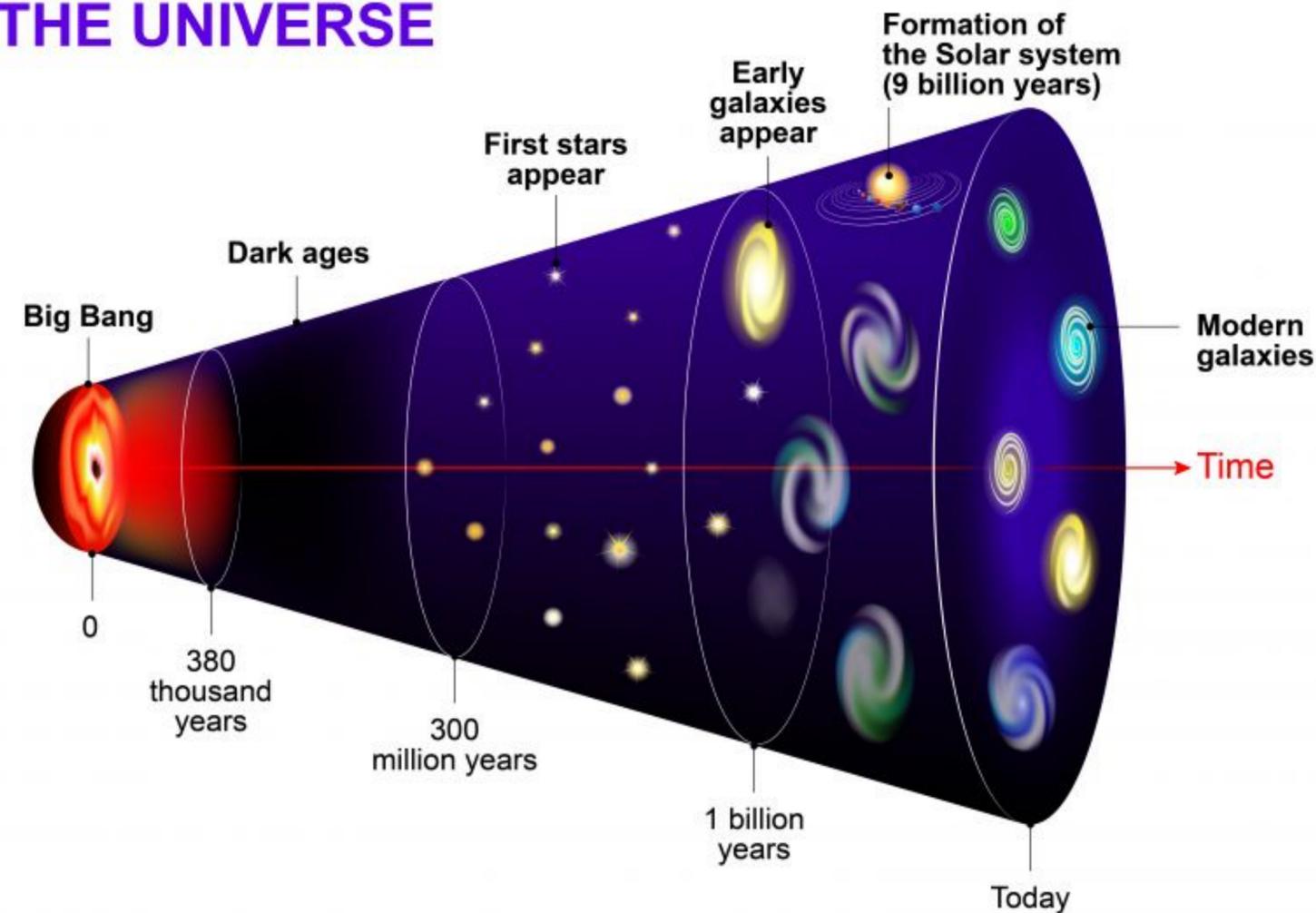
What is High Energy Physics?



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Aim: Smash particles together to “emulate” conditions closer to the Big Bang! (“Particle Archaeology”)

EVOLUTION OF THE UNIVERSE



← HIGHER ENERGY



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Aim: Smash particles together to understand the structure of matter today! (“Particle Sociology”)



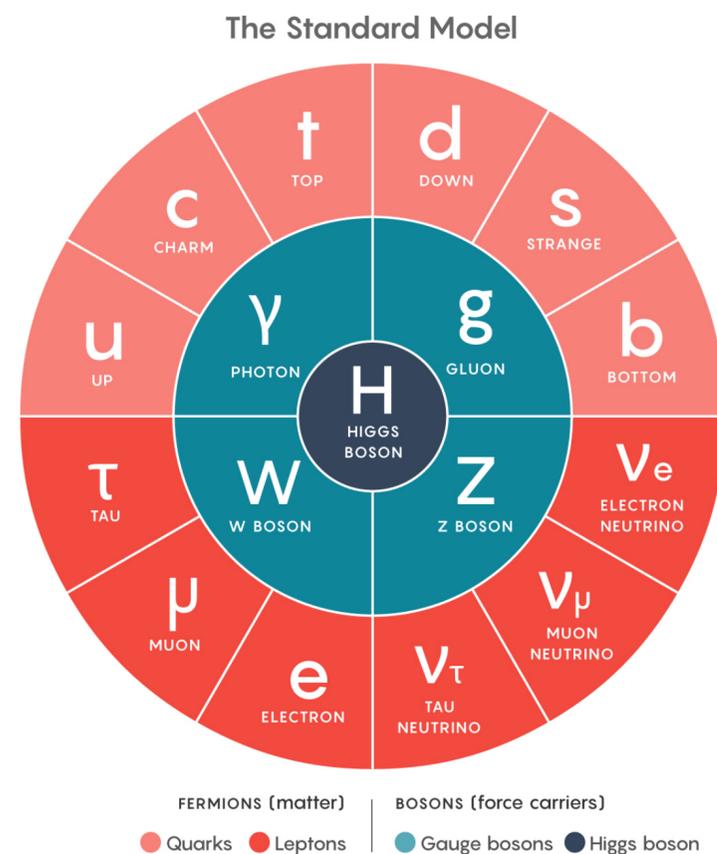
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→ understand

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e.g. **The “Standard Model” (SM)**



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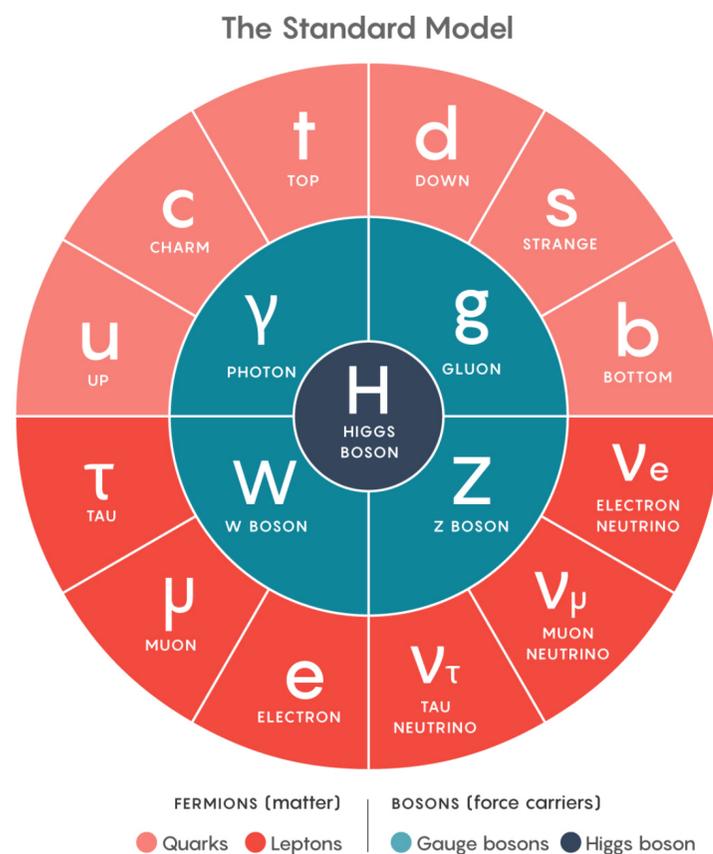
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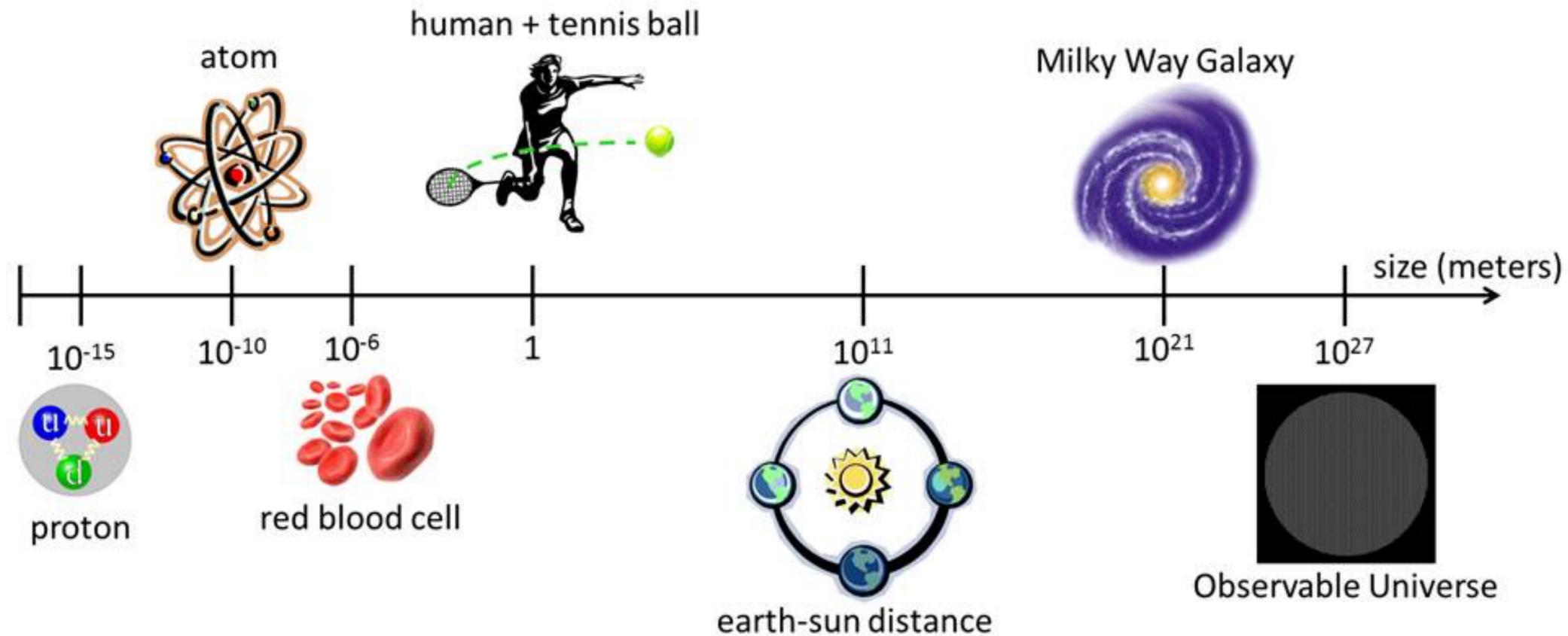
e.g. **The “Standard Model” (SM)**

& discover **exotic** phenomena!



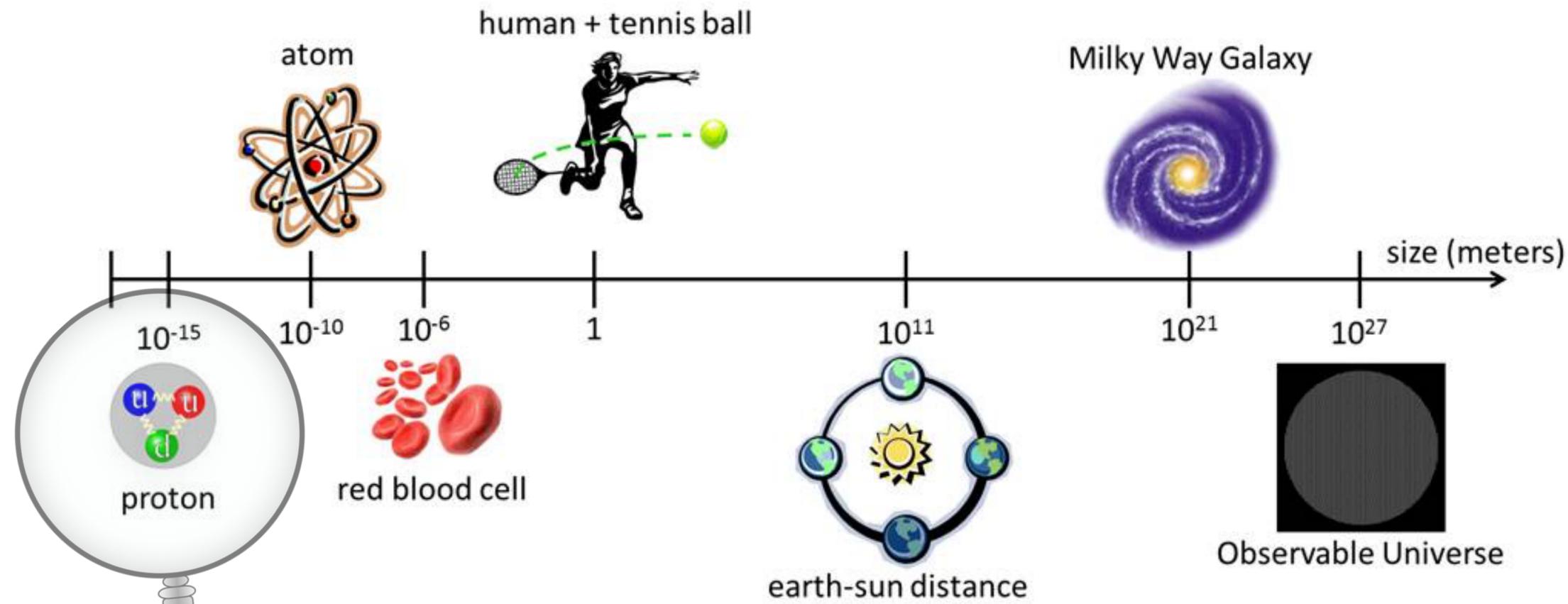
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Rule: $E \propto 1/\lambda \Rightarrow$ Higher Energy \Leftrightarrow Smaller Scales!



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$m_p \sim 1 \text{ GeV} \Rightarrow$ To probe proton structure, go to higher energies than this!



High Energy Particle Colliders

Large Hadron Collider (LHC)

@ CERN

in Geneva, Switzerland



Collides: (primarily)
Protons to Protons
Center-of-mass energy:
 $13,600 \text{ GeV} = 13.6 \text{ TeV}$

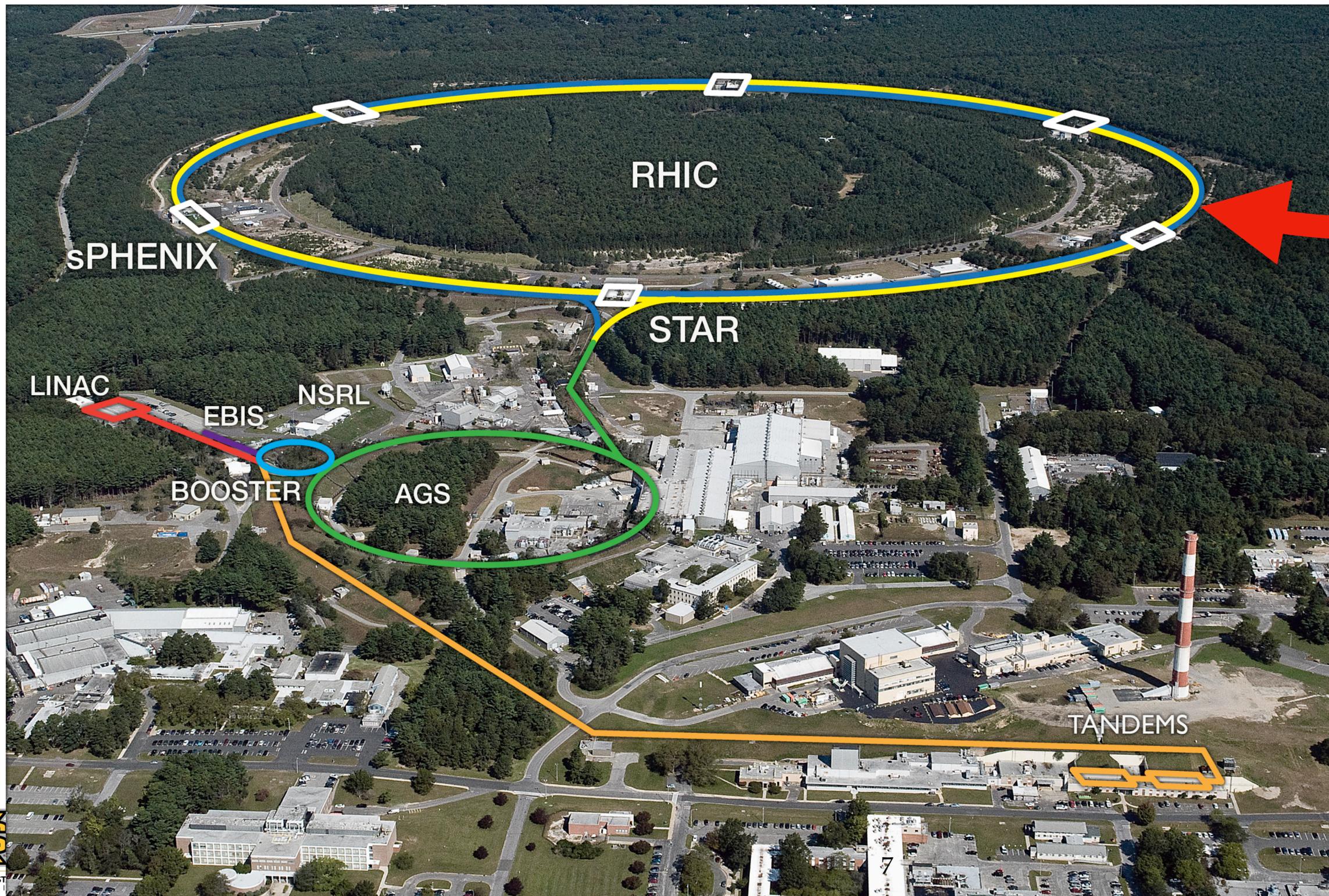


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Relativistic Heavy Ion Collider

**@ Brookhaven National Lab
in Upton, NY**

High Energy Particle Colliders



Collides: Heavy Ions

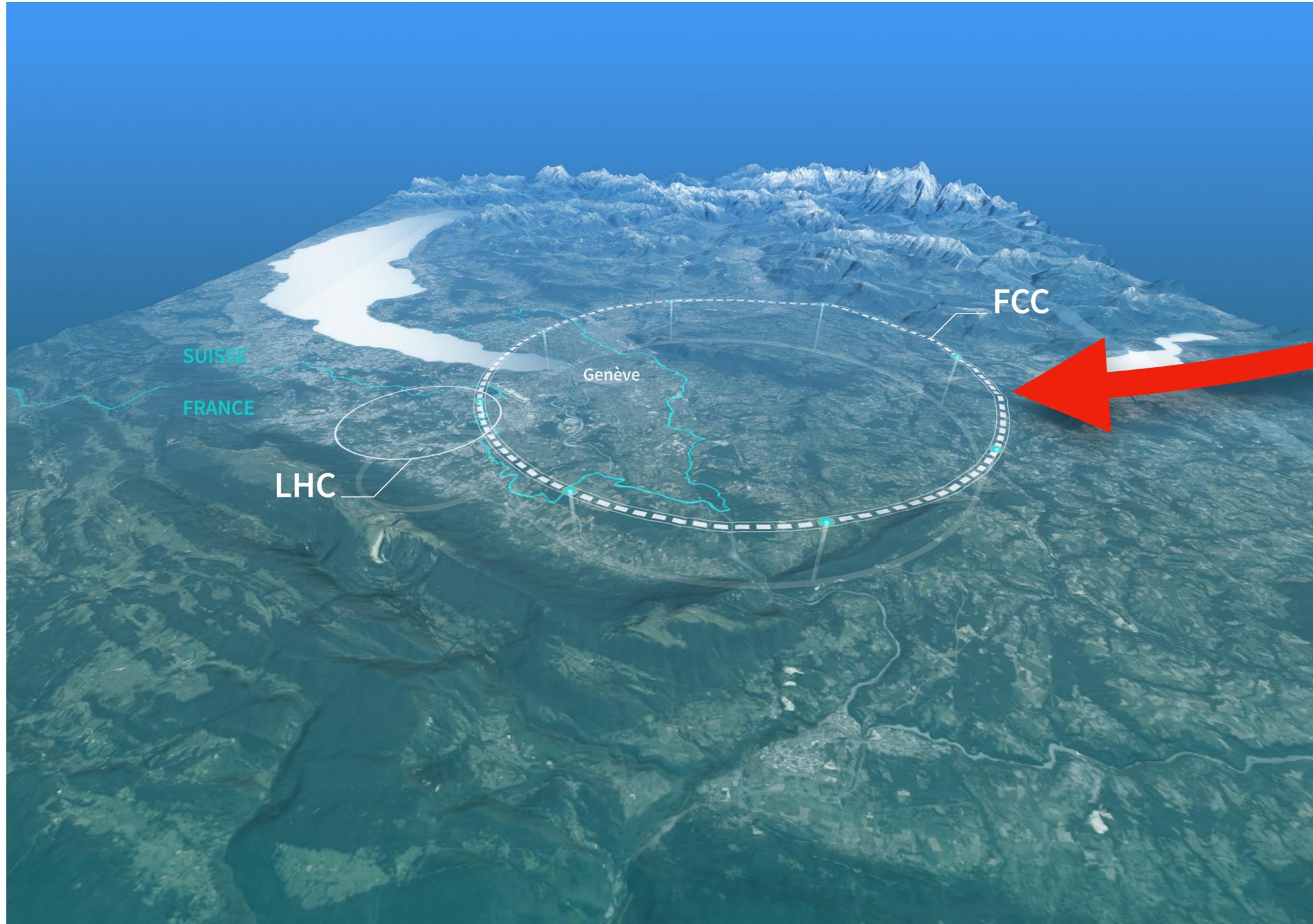
**+ Upcoming (2030s):
Electron-Ion Collider,
will collide electrons
and ions**



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High Energy Particle Colliders

**“Future Circular Collider”
@ CERN**



**Collides: electrons to
positrons (2040s)**

&

**protons to protons
(2070s) @ Center-of-
mass energy: 100,000**

GeV = 100 TeV

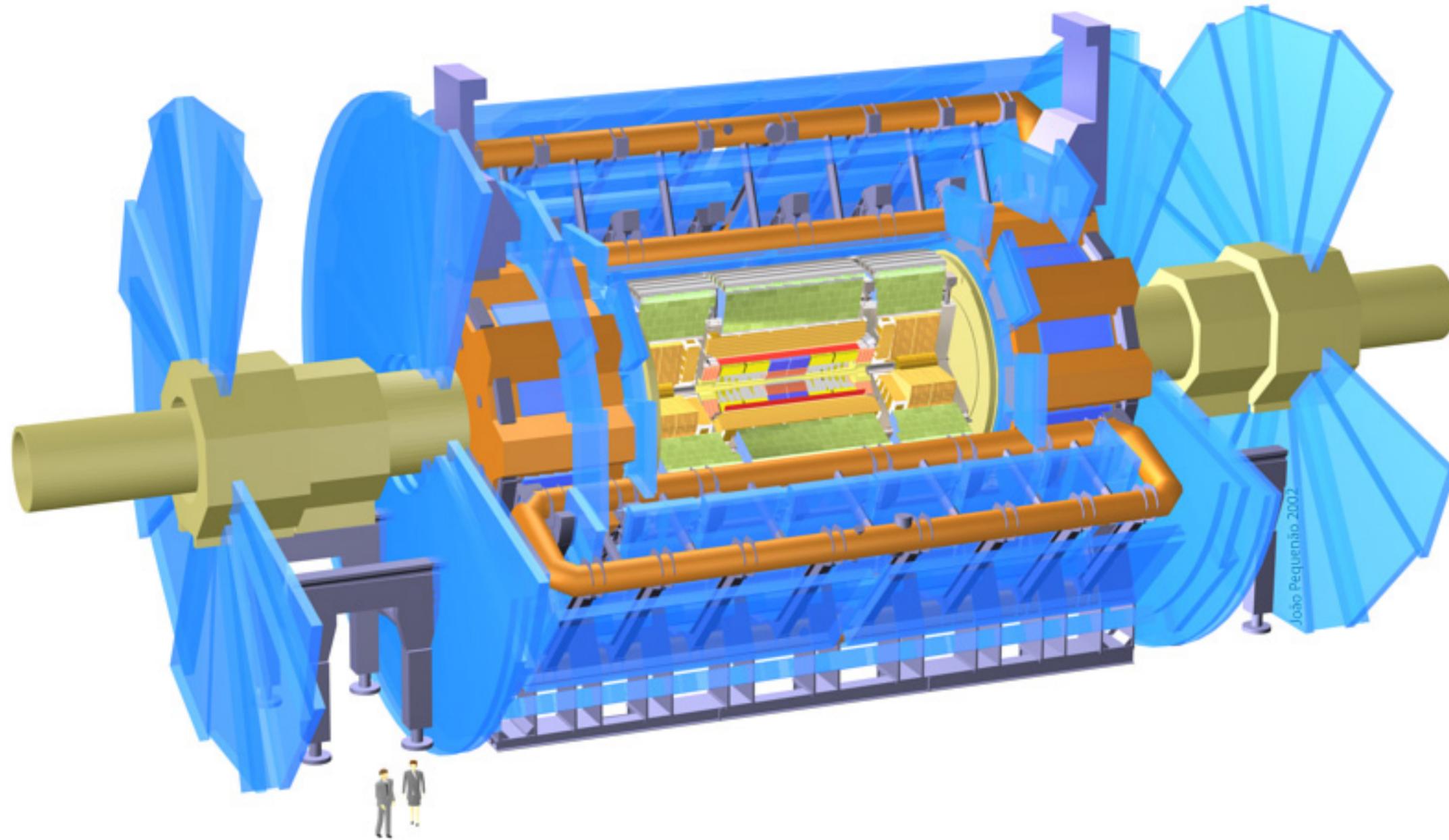


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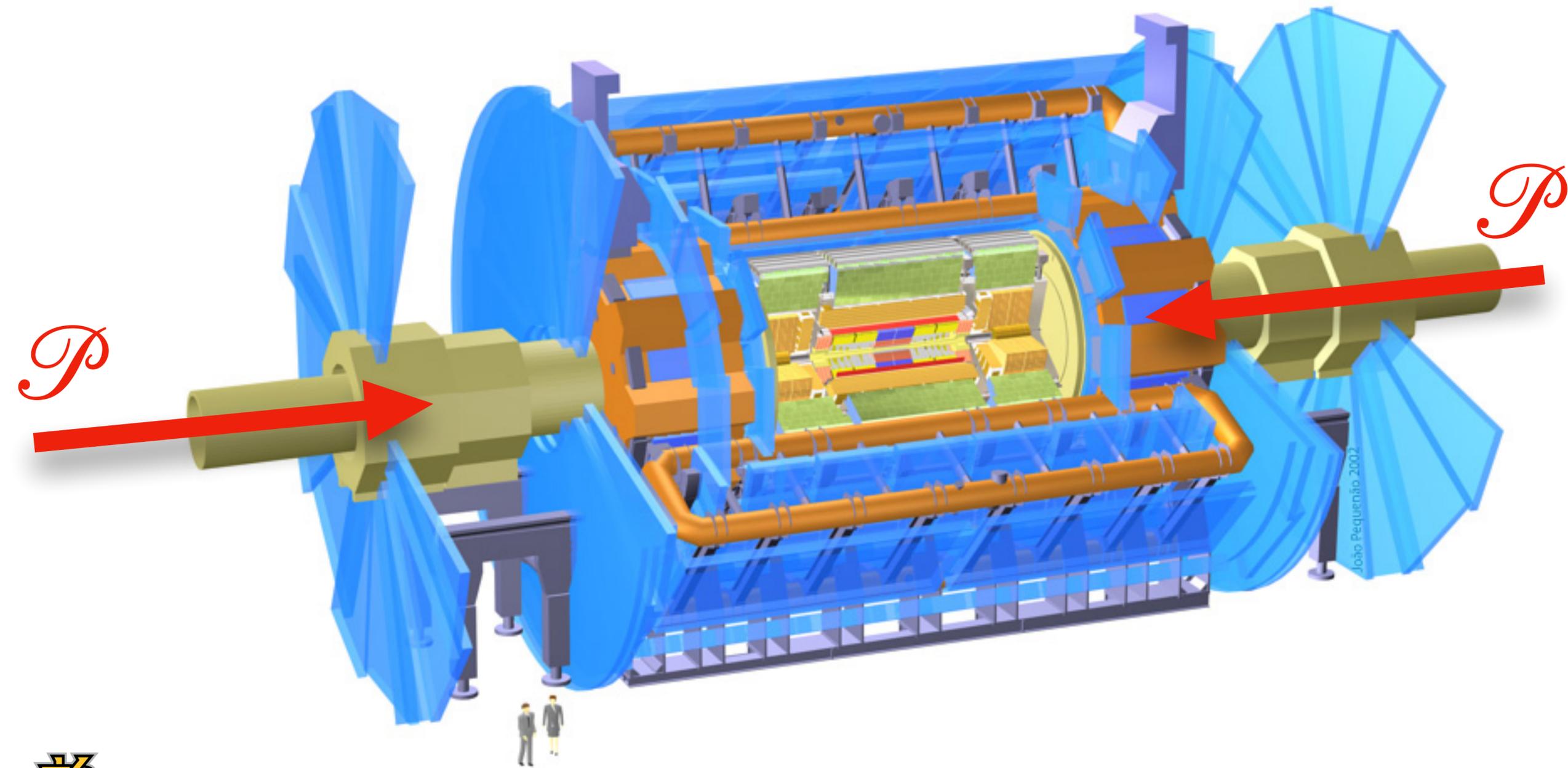
Happening Now: Particle Collisions at the LHC



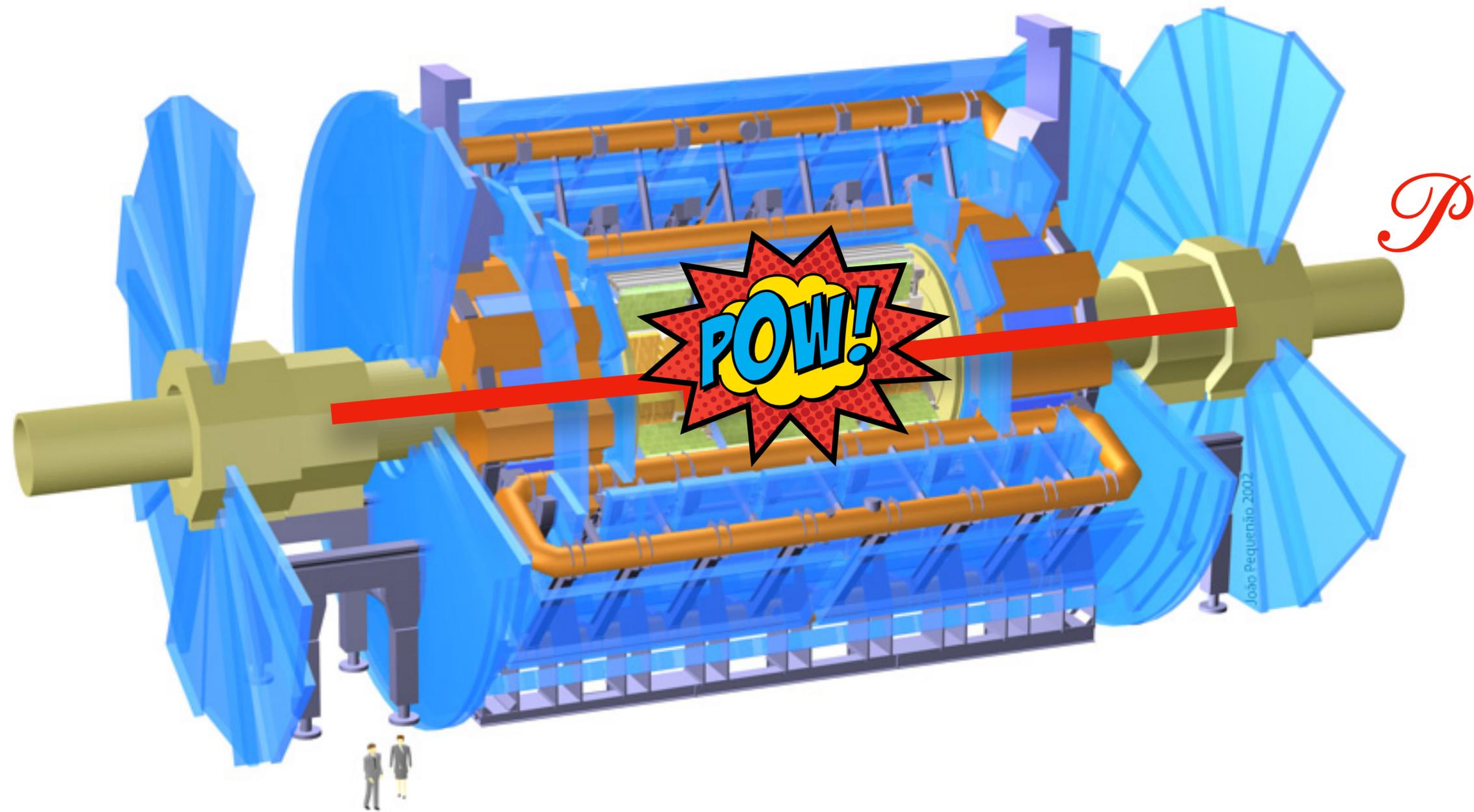
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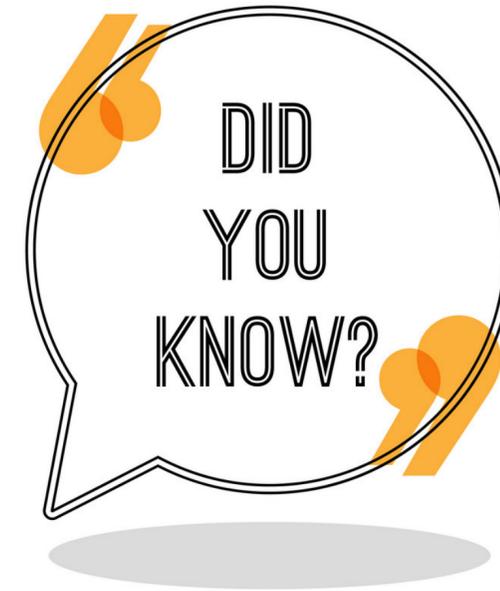


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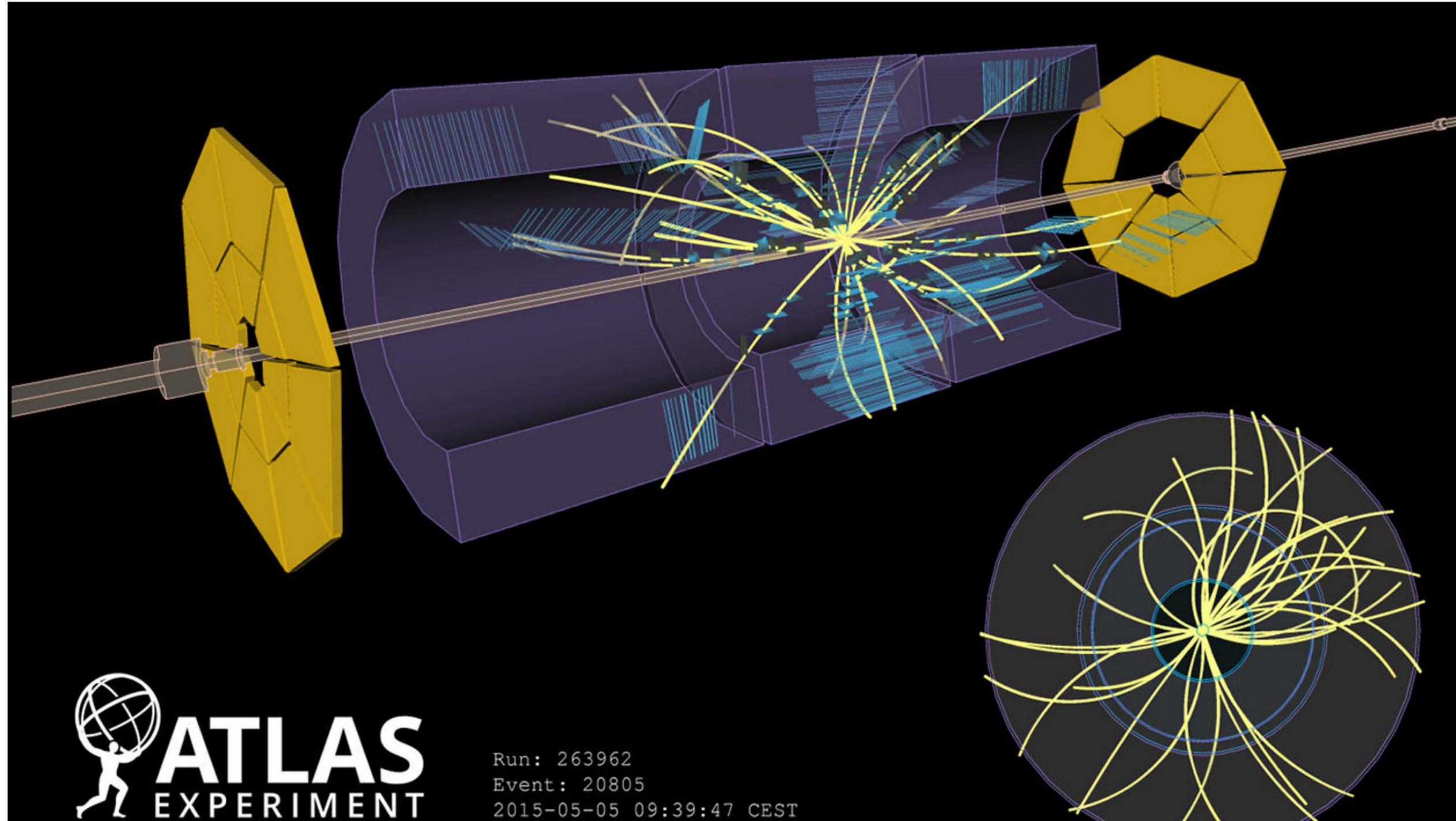


LHC: Fun Facts

- Circumference of **27 km** (17 miles).
- Lies **100 m** (330 feet) underground.
- Produces **hundreds of millions of particle collisions/second** (=“events”).
- **Energy stored** in particle beam = Energy of **1000 kg car @ 3000 km/h** (~1900 mph)!
- The data recorded can fill around **100,000 DVDs** each year!



This is a real “event”!



How do we make sense of it all?



Pictured: Richard Feynman playing the bongos.



How do we make sense of it all?

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + i\bar{\Psi}\not{D}\Psi \\ & + \Psi_i y_{ij} \Psi_j + \text{h.c.} \\ & + |D_\mu\Phi|^2 - V(\Phi)\end{aligned}$$



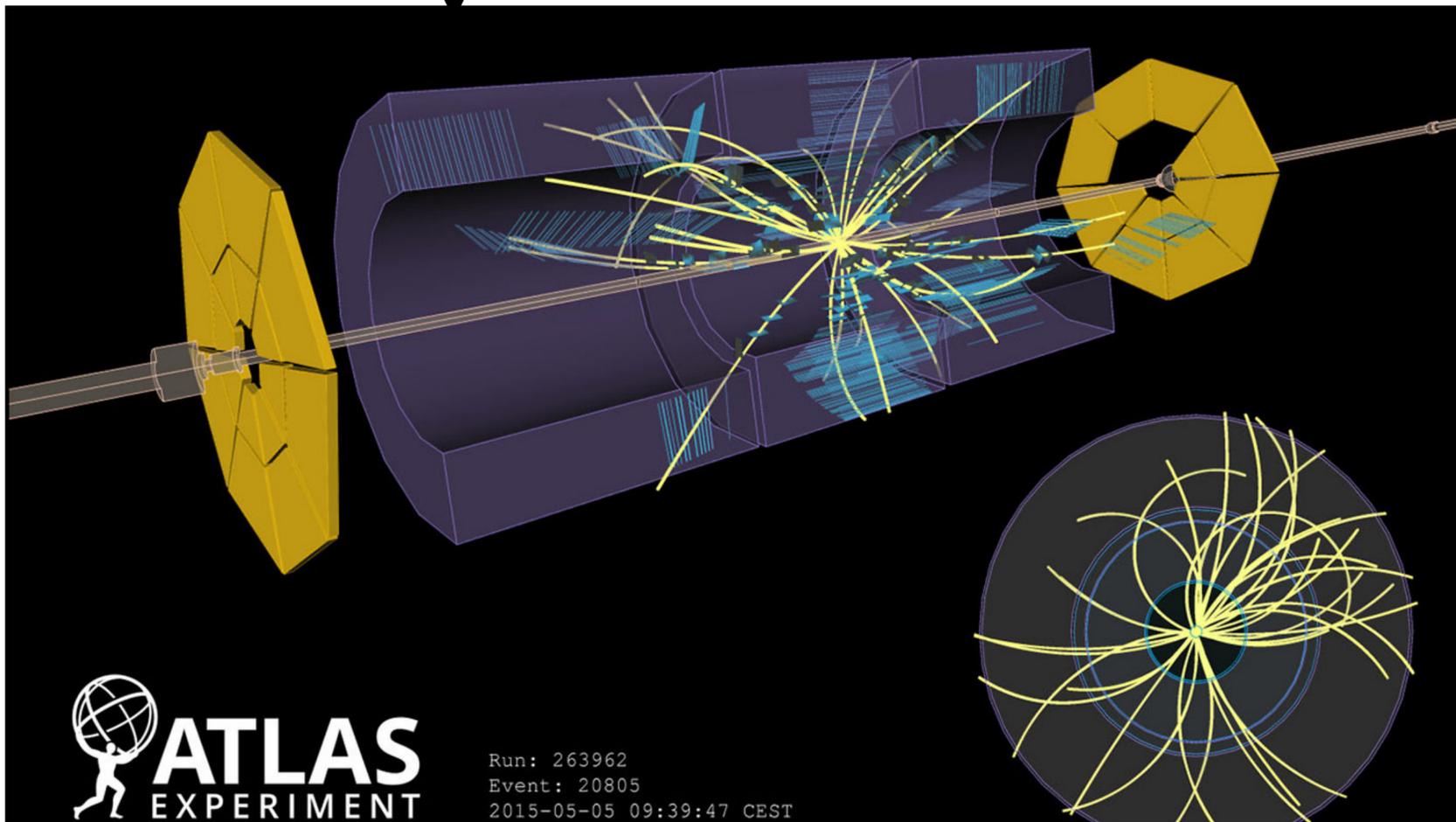
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Experiment



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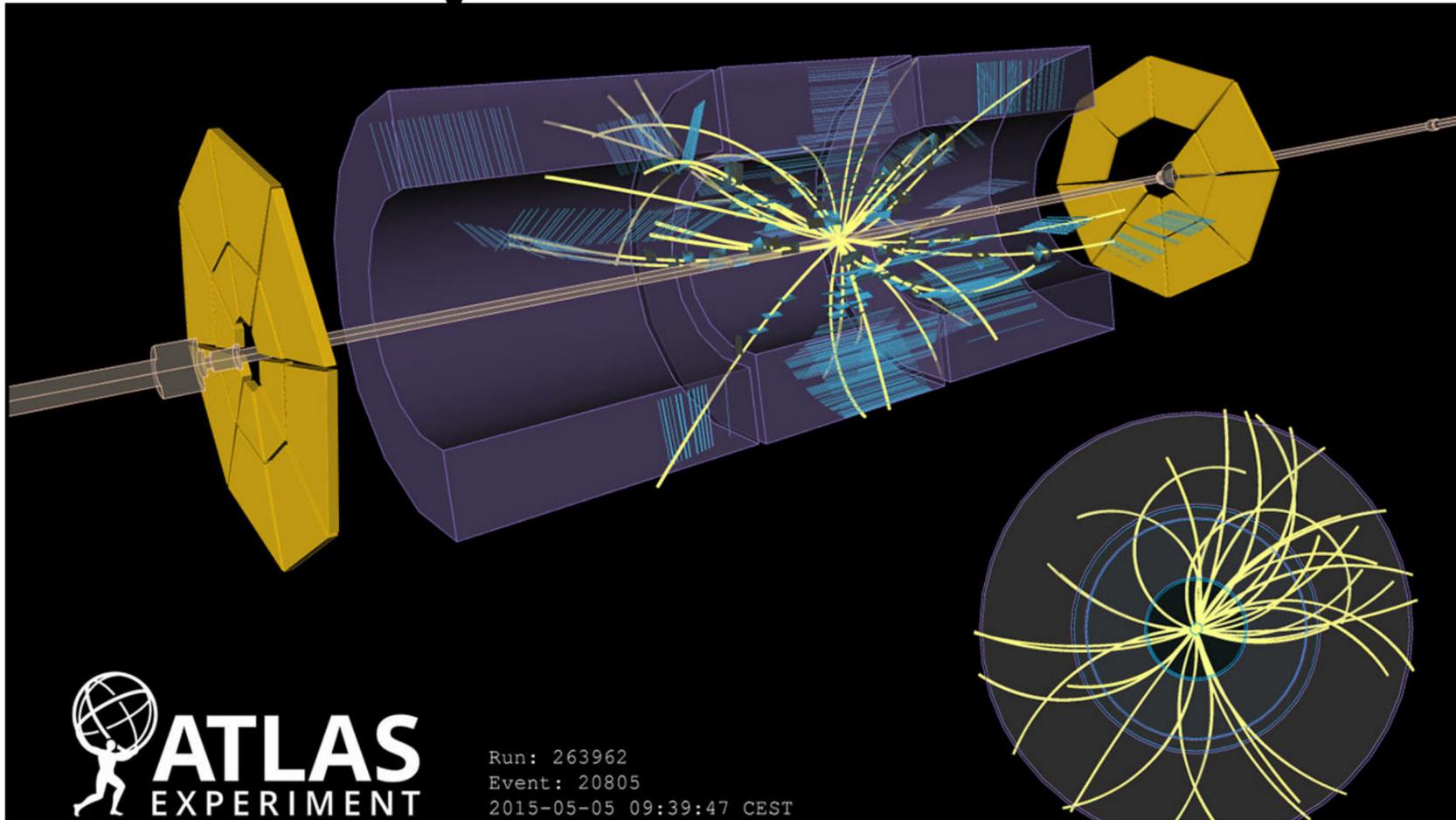


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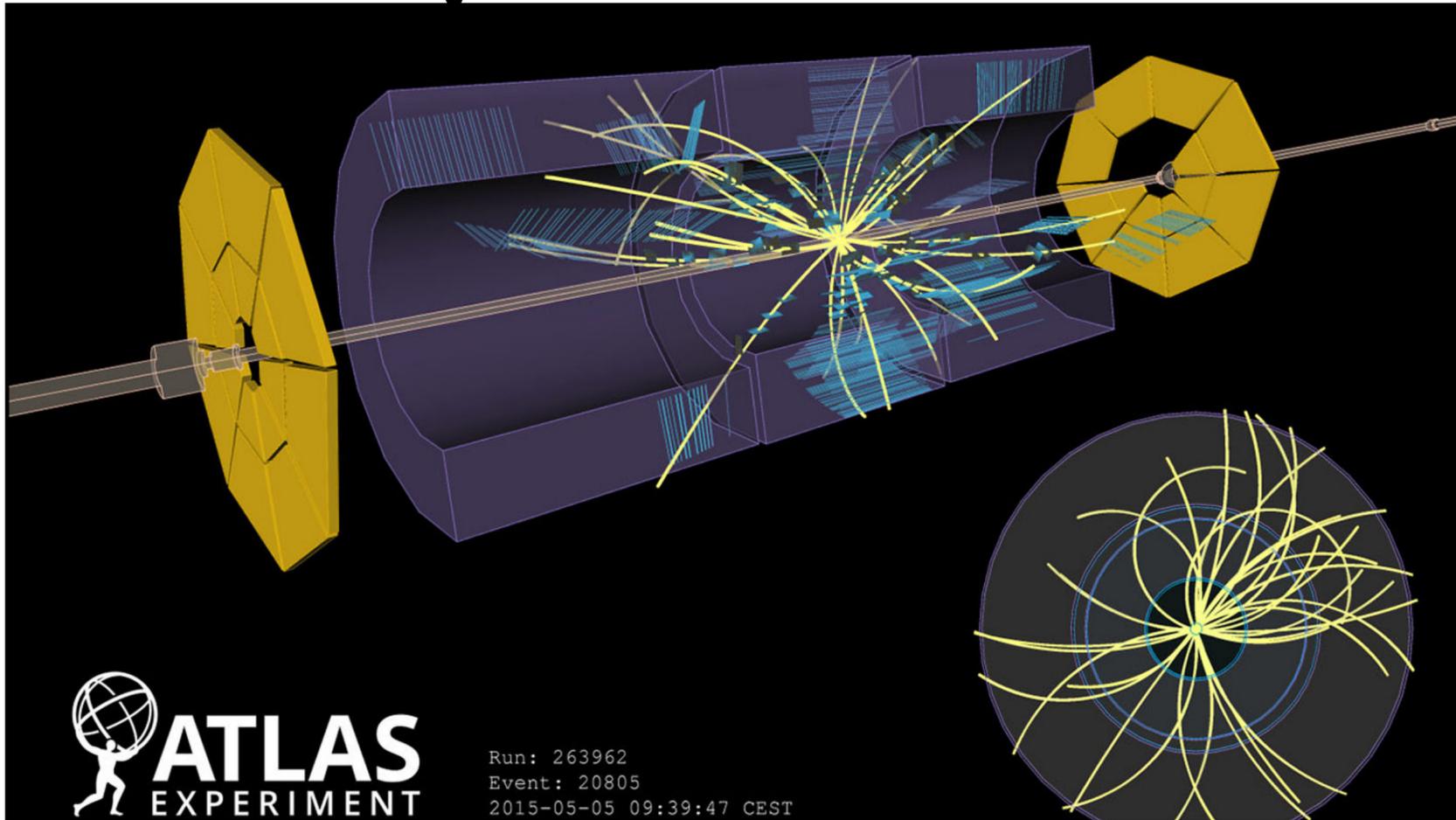


How do we make sense of it all?

Shut up and Calculate!

???

Experiment



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(Shut up and Calculate) Using Quantum Field Theory



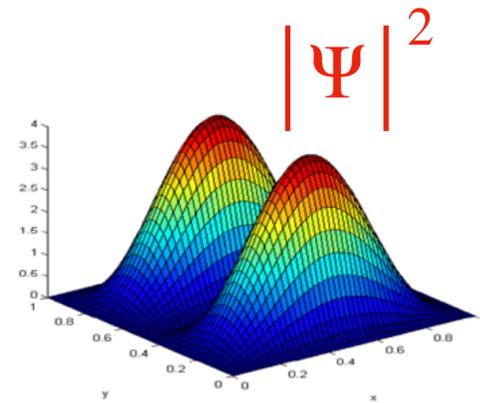
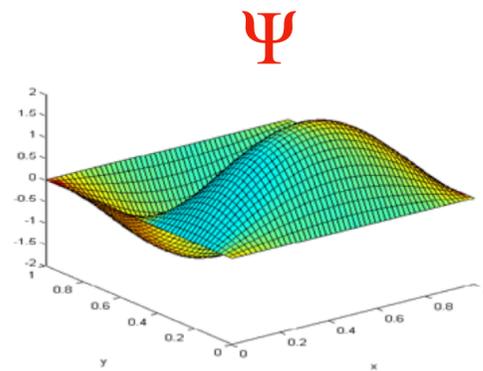
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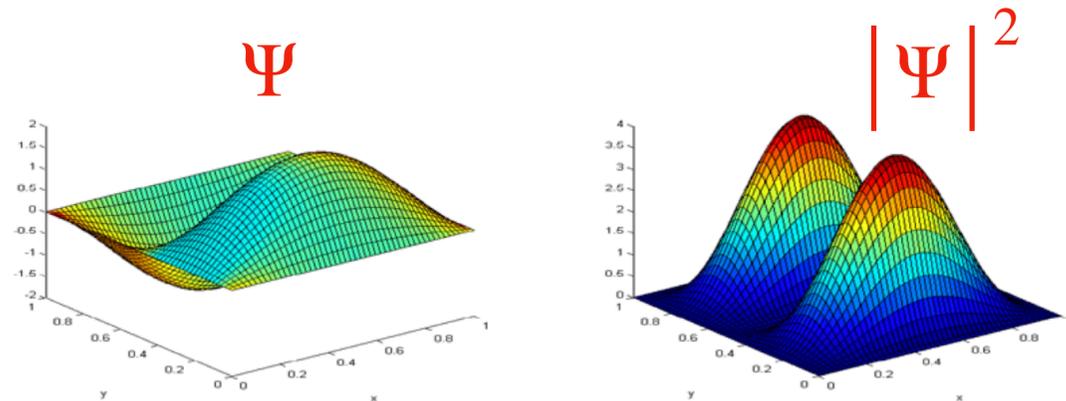
⇒ Probability to find particle somewhere $\propto |\Psi|^2$

[Ψ is called the “wave function”]



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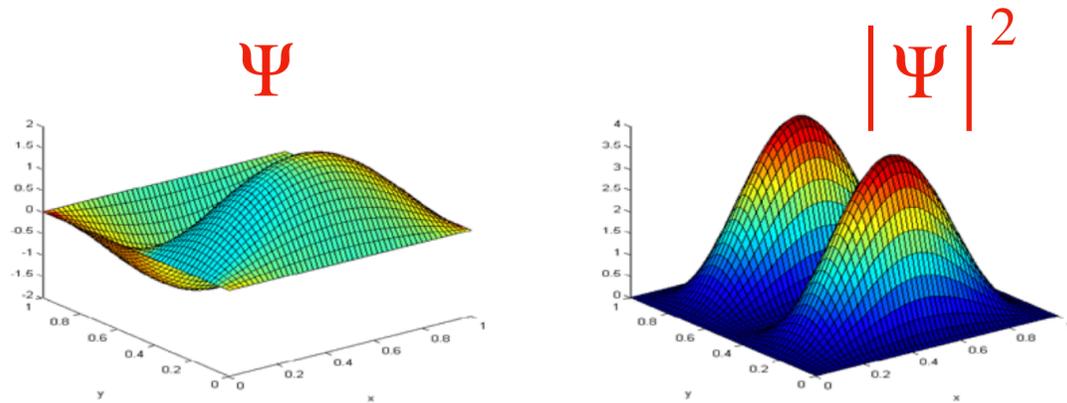
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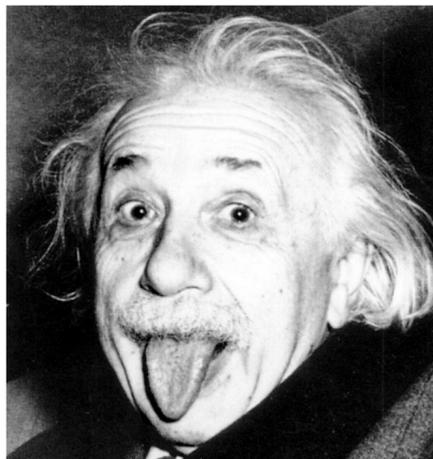


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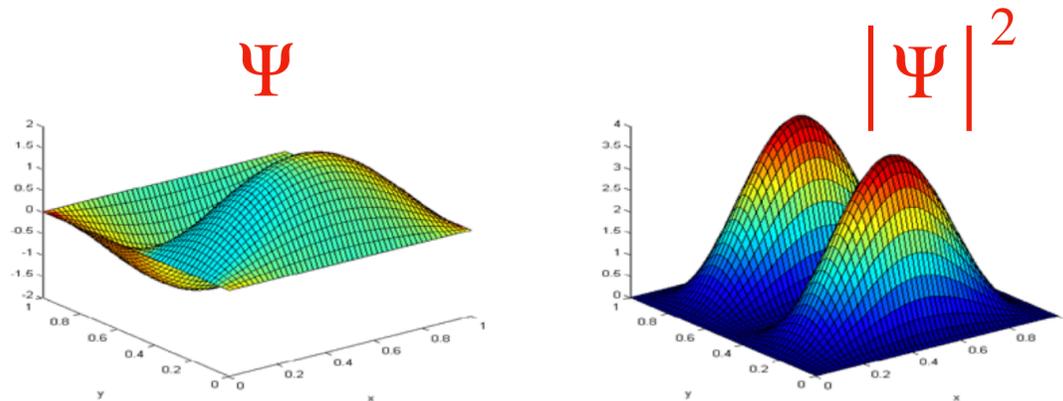


$$E = mc^2$$



(Shut up and Calculate) Using Quantum Field Theory

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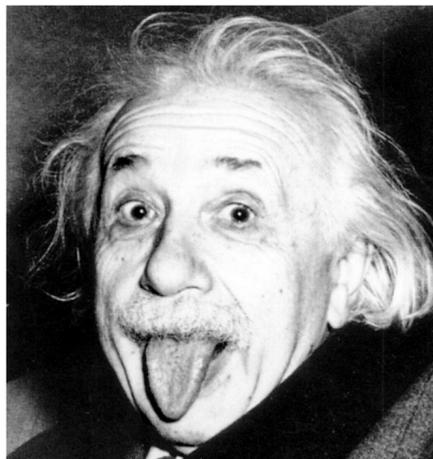


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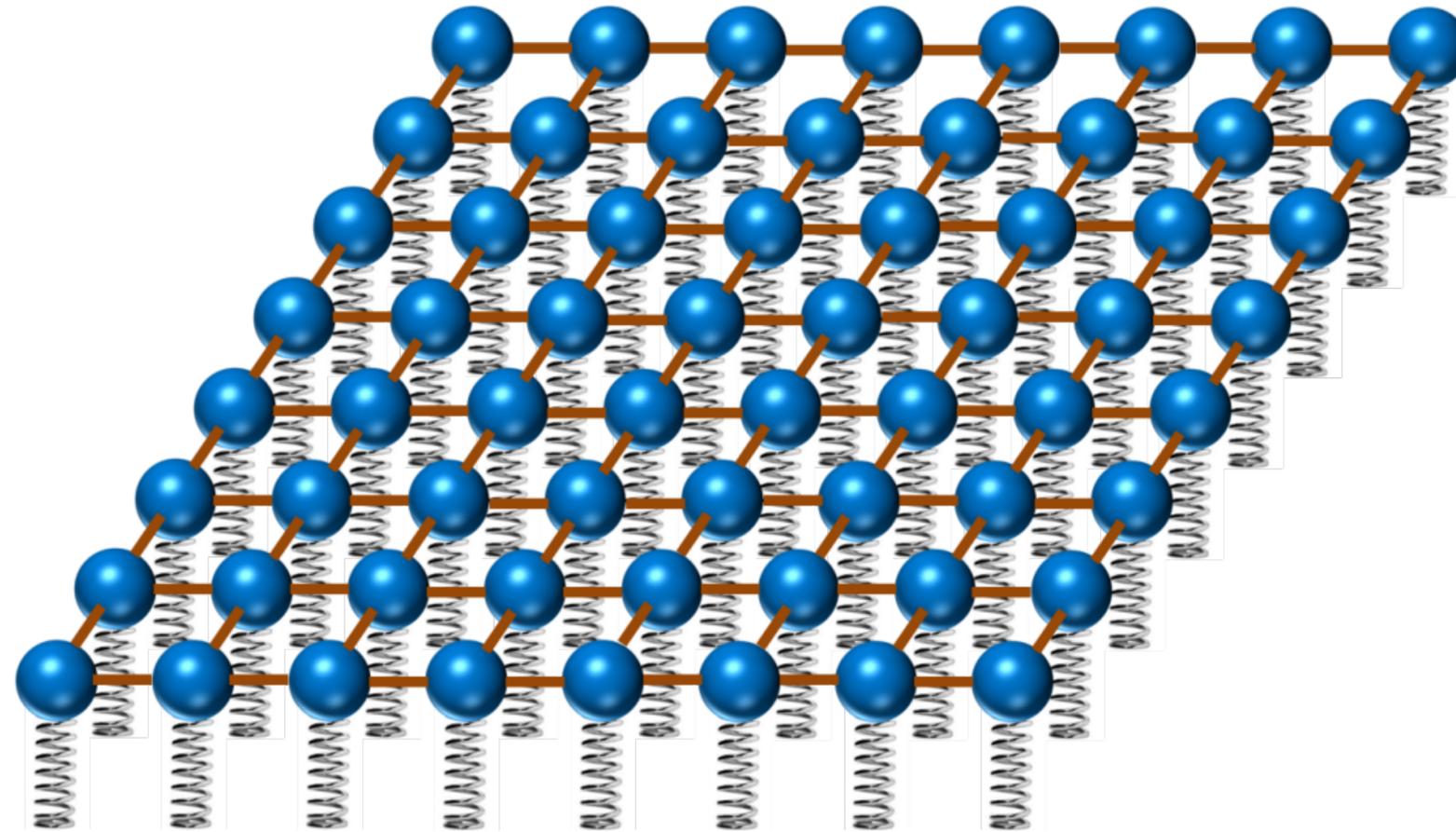


$$E = mc^2$$

- Mass and Energy are equivalent!
- ⇒ You can exchange one for the other!
- ⇒ Particle creation from Energy!



Quantum Mechanics + Special Relativity \approx **Quantum Field Theory**



From Theory to Experiment and Back Again



~~THEORY~~

~~EXPERIMENT~~



From Theory to Experiment and Back Again

$e^- \bullet$

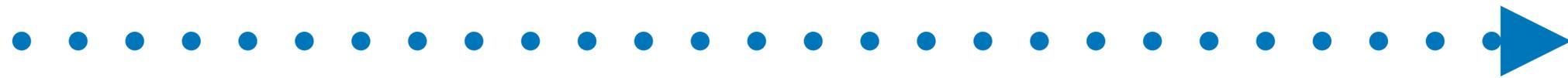
$e^- \bullet$



From Theory to Experiment and Back Again



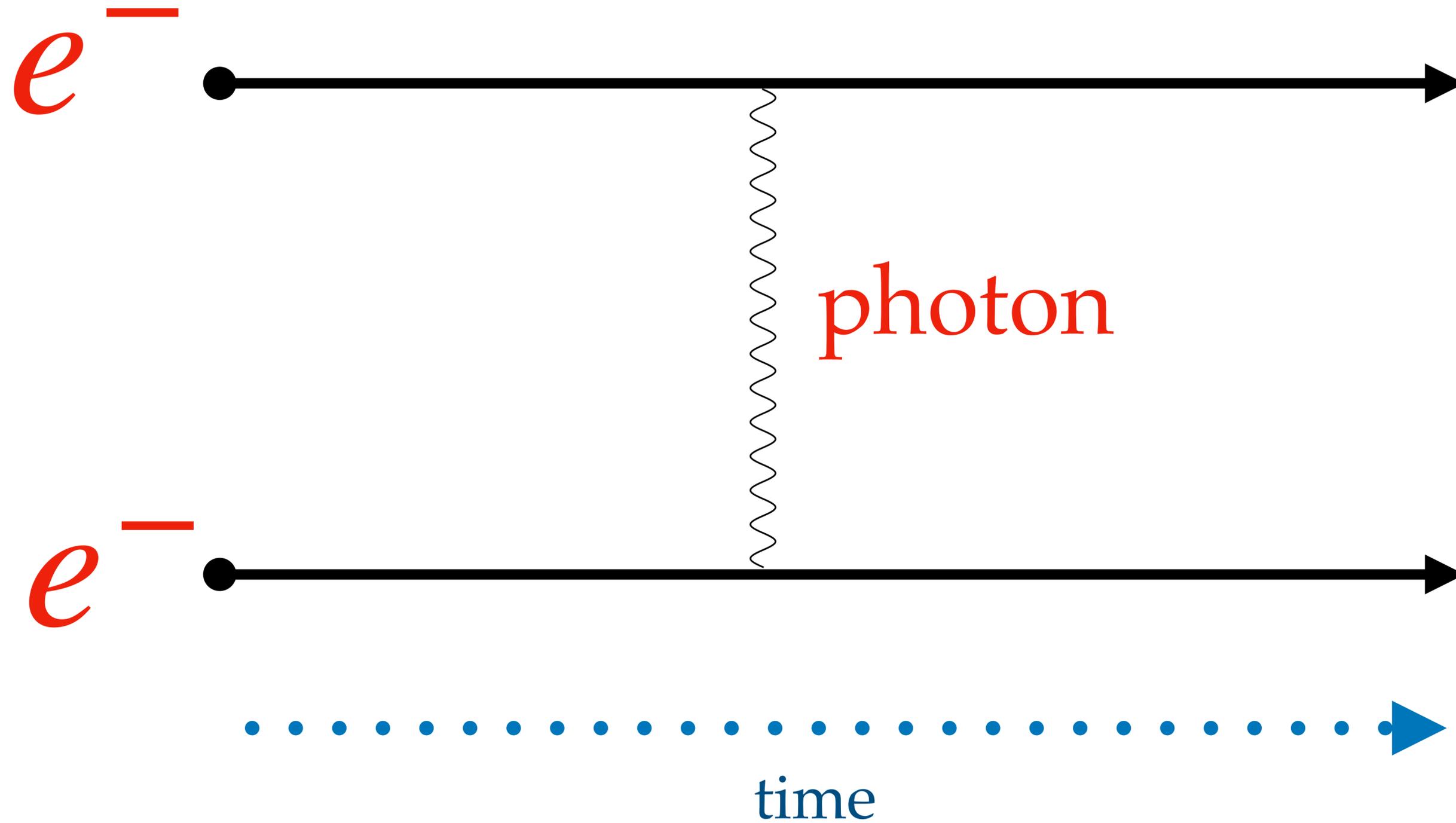
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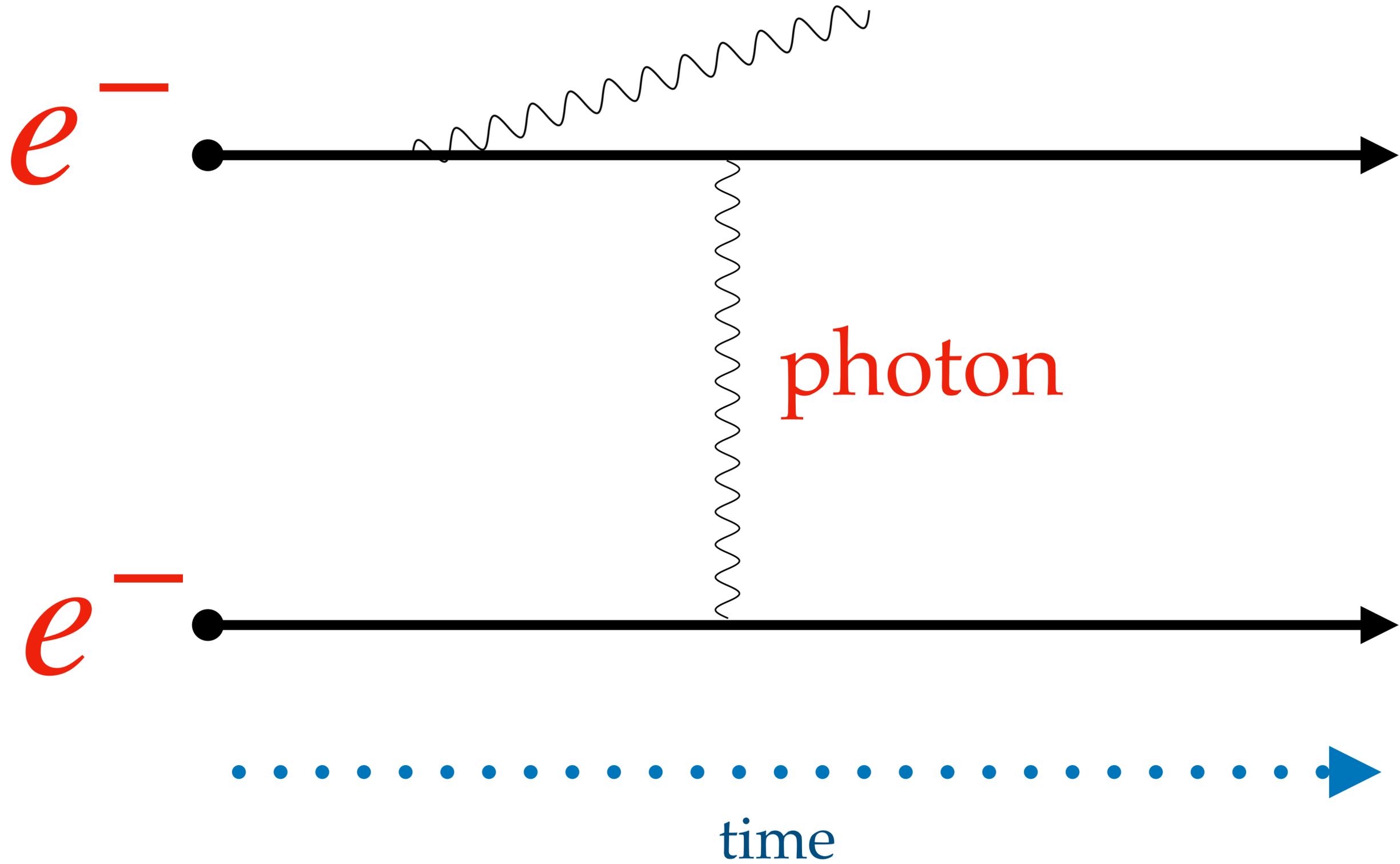
time



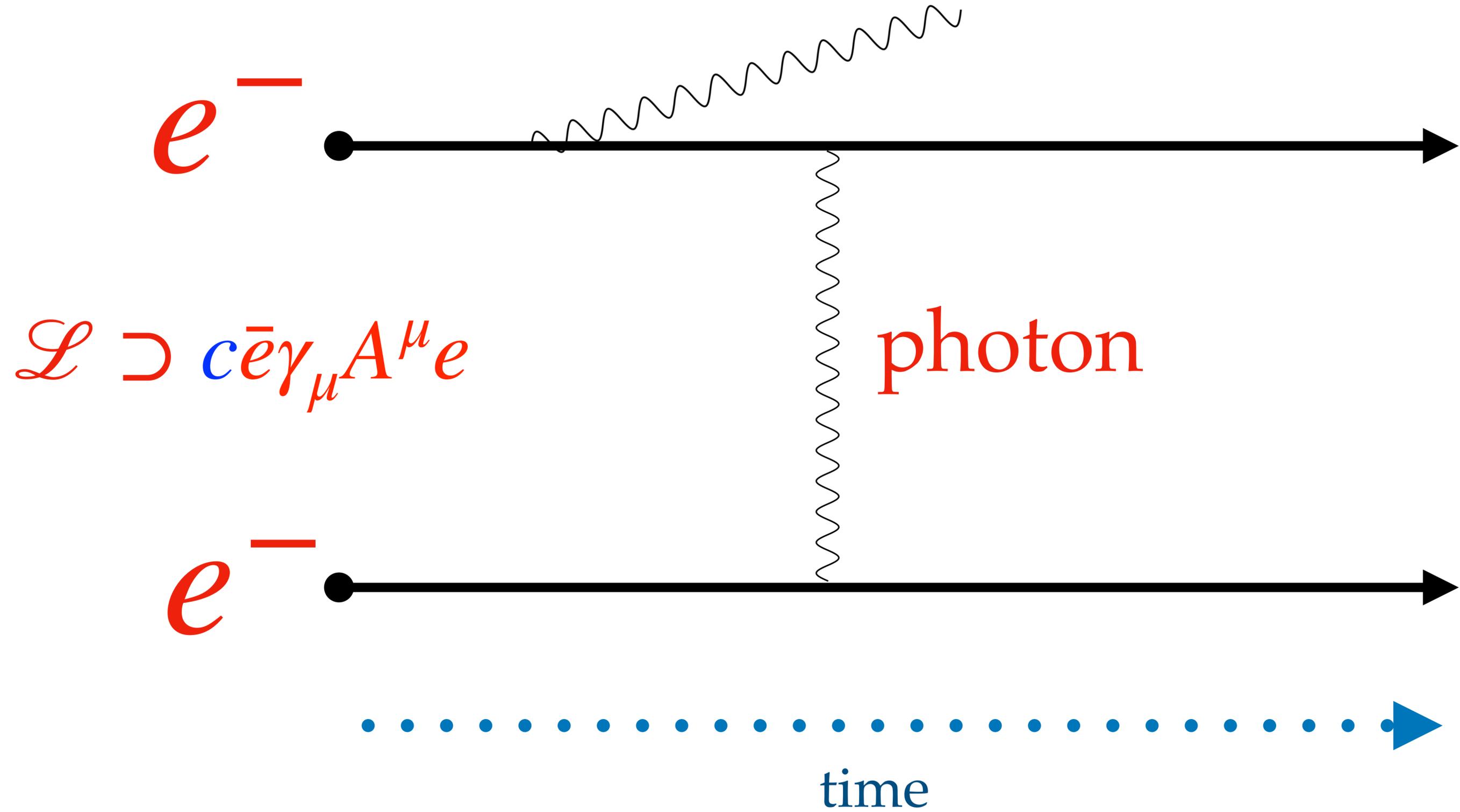
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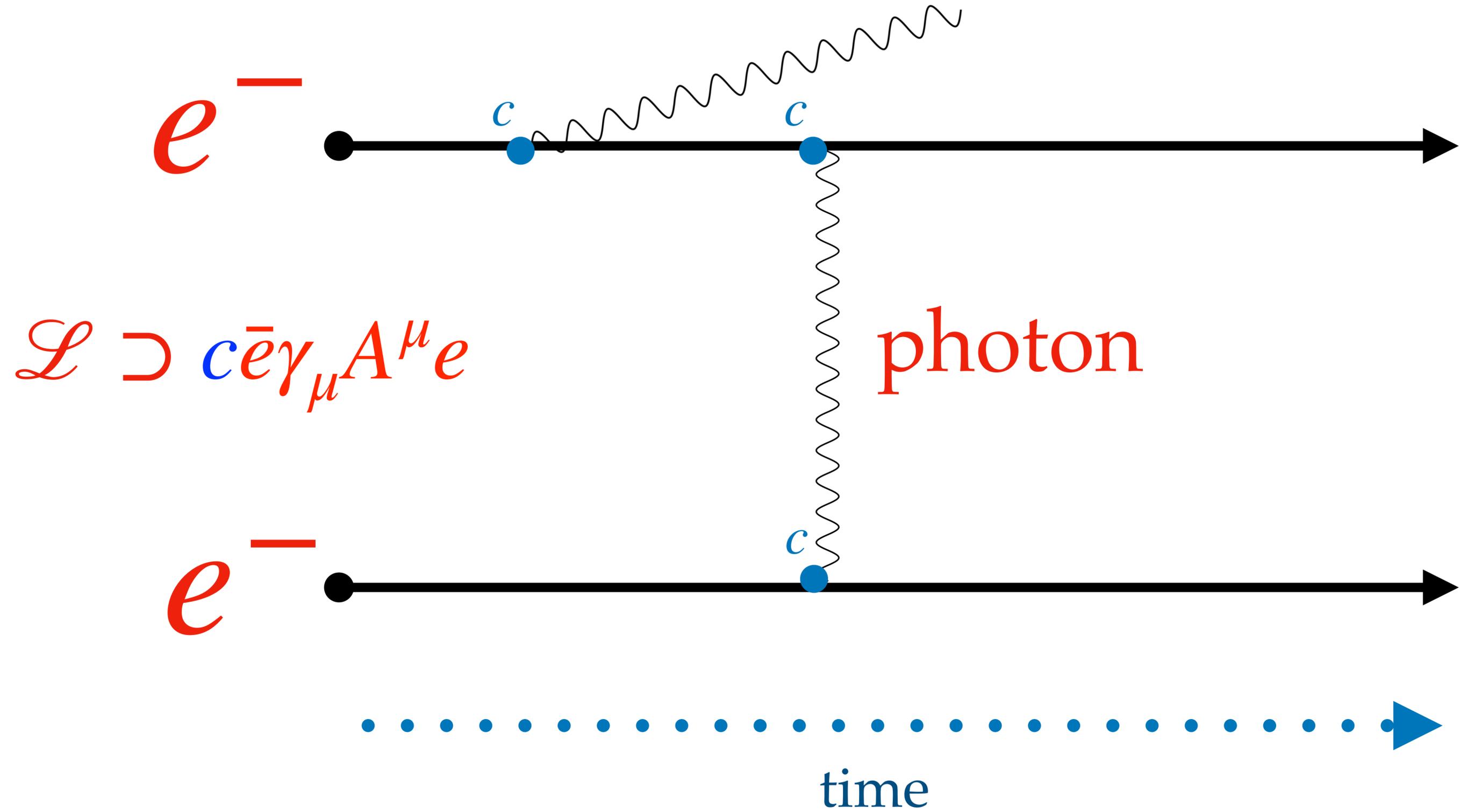
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From Theory to Experiment and Back Again



$$\mathcal{L} \supset c \bar{e} \gamma_{\mu} A^{\mu} e$$

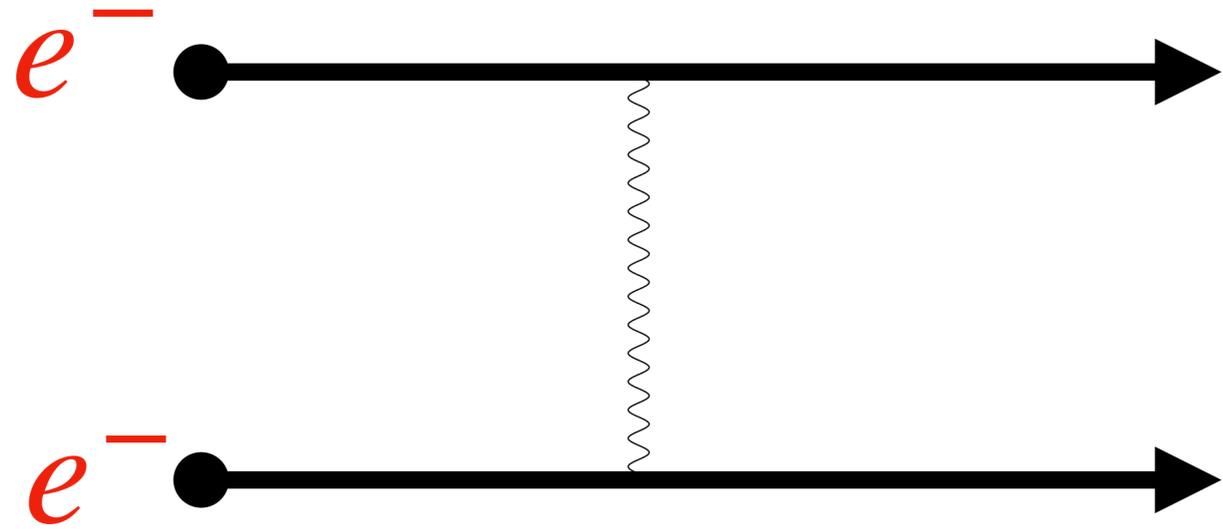
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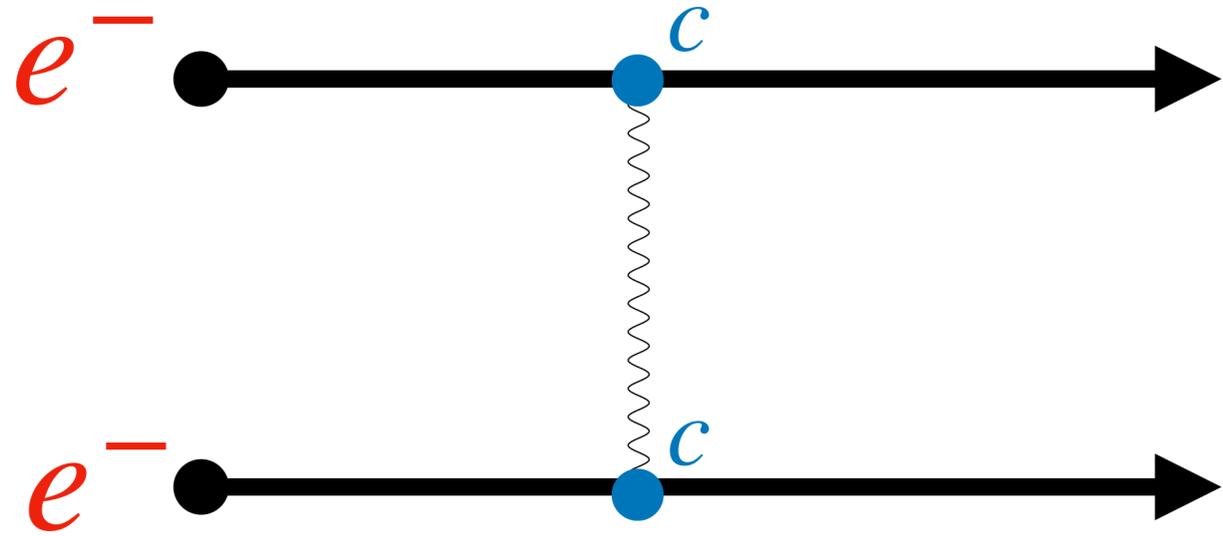
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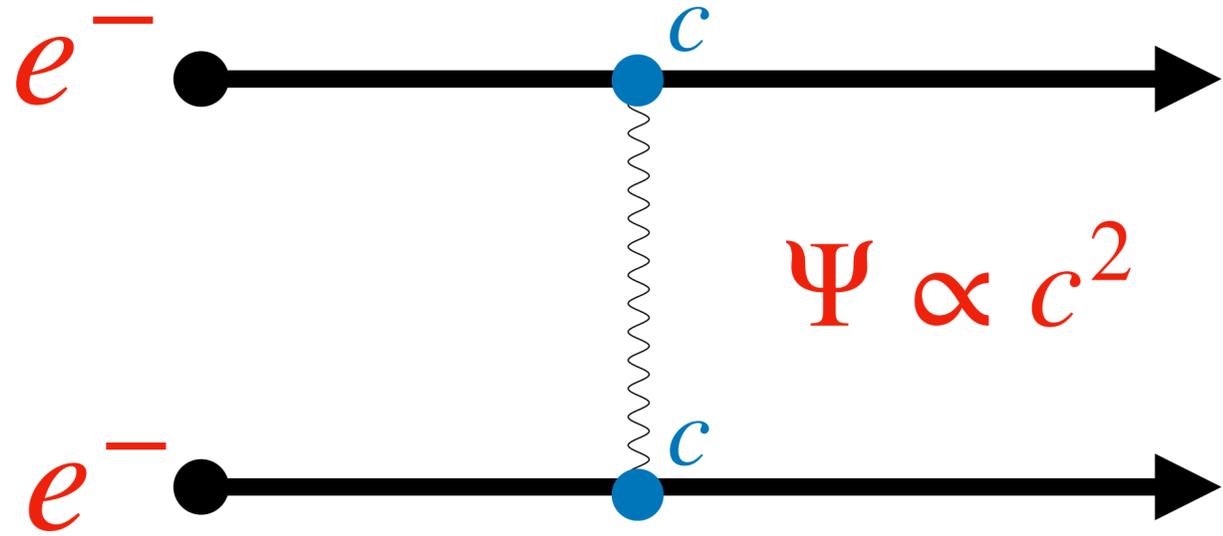
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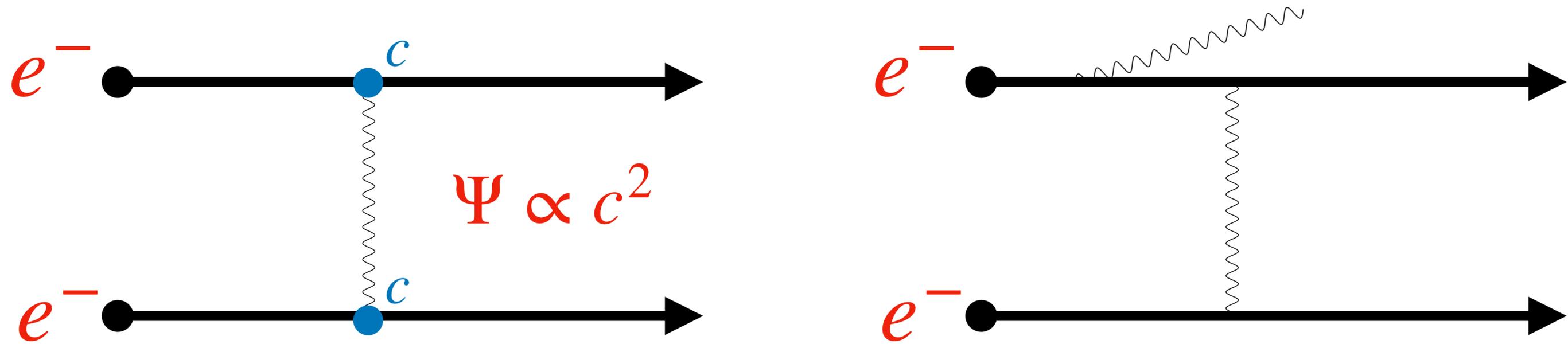
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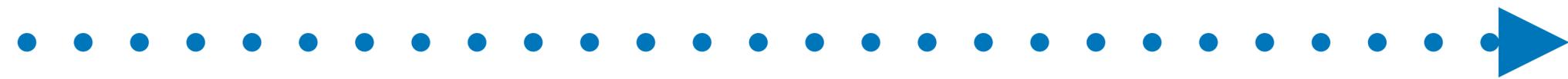
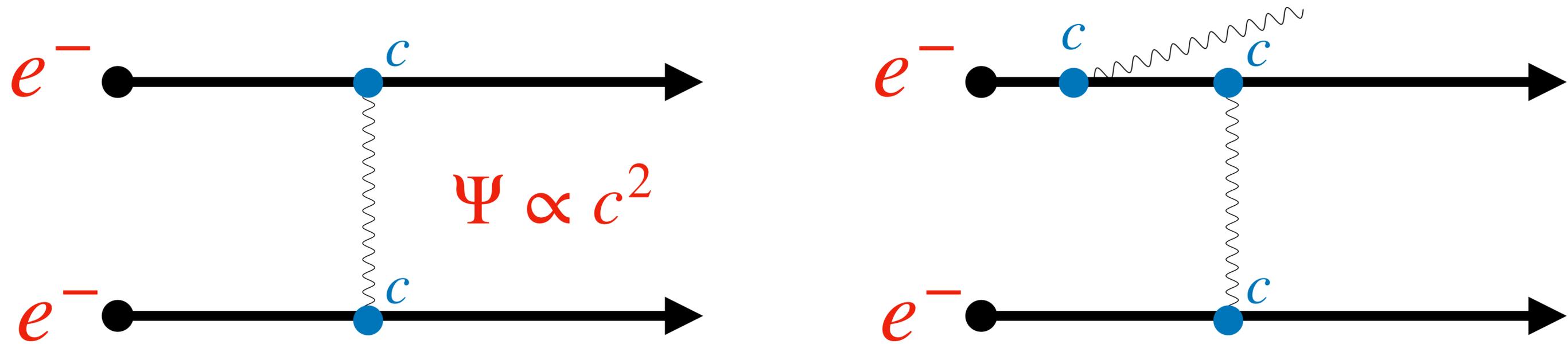
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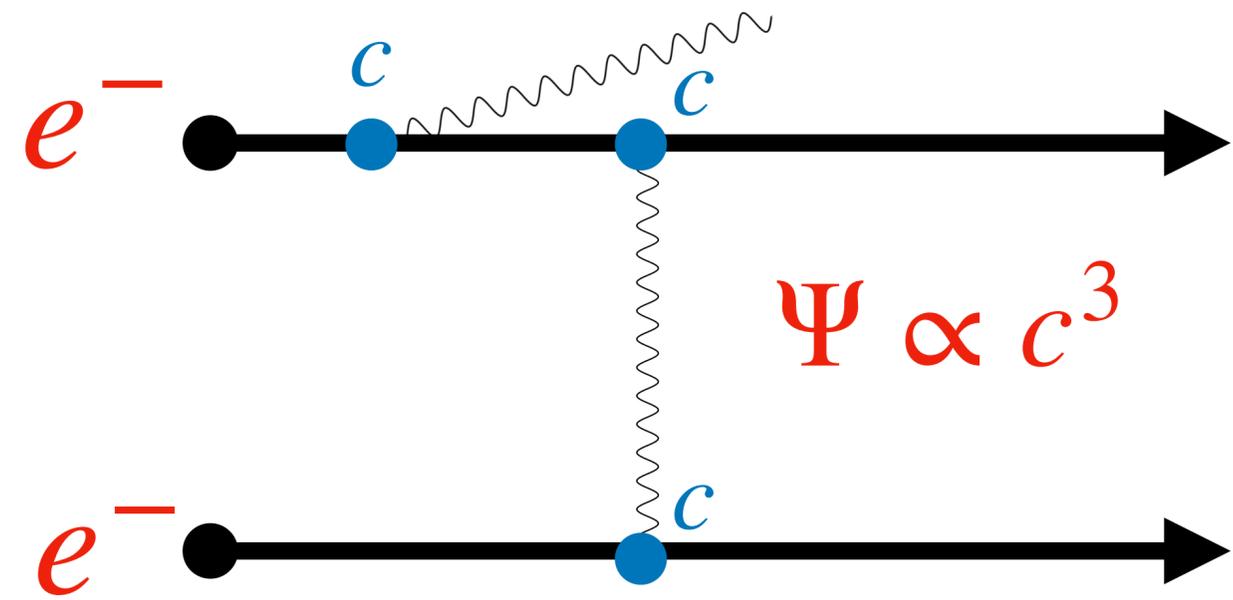
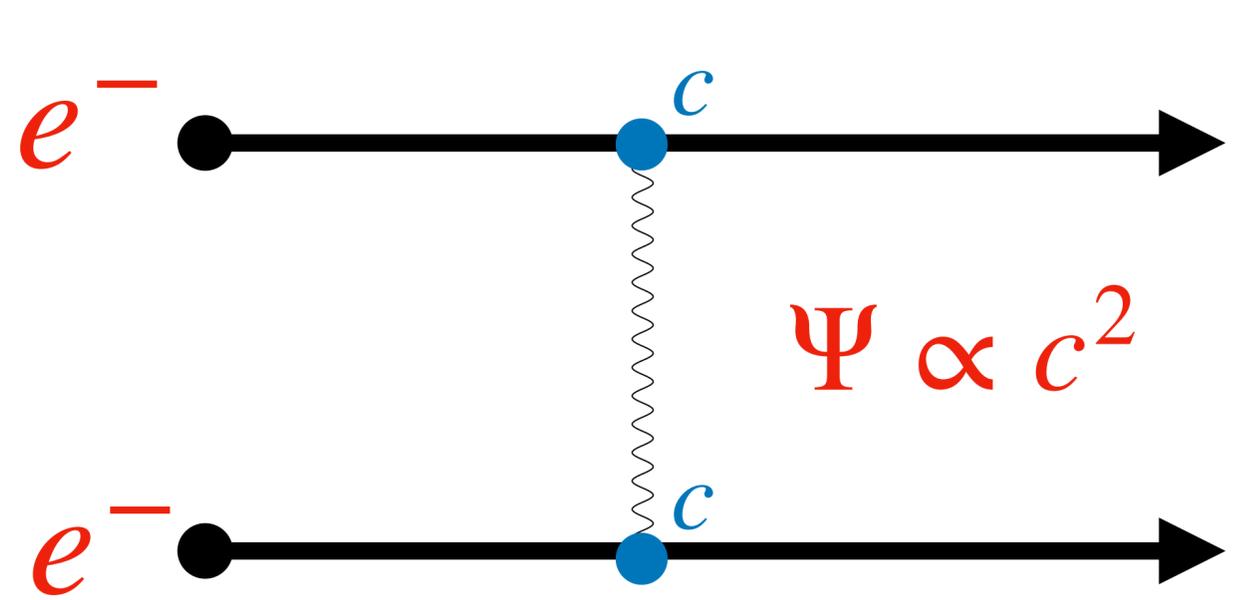
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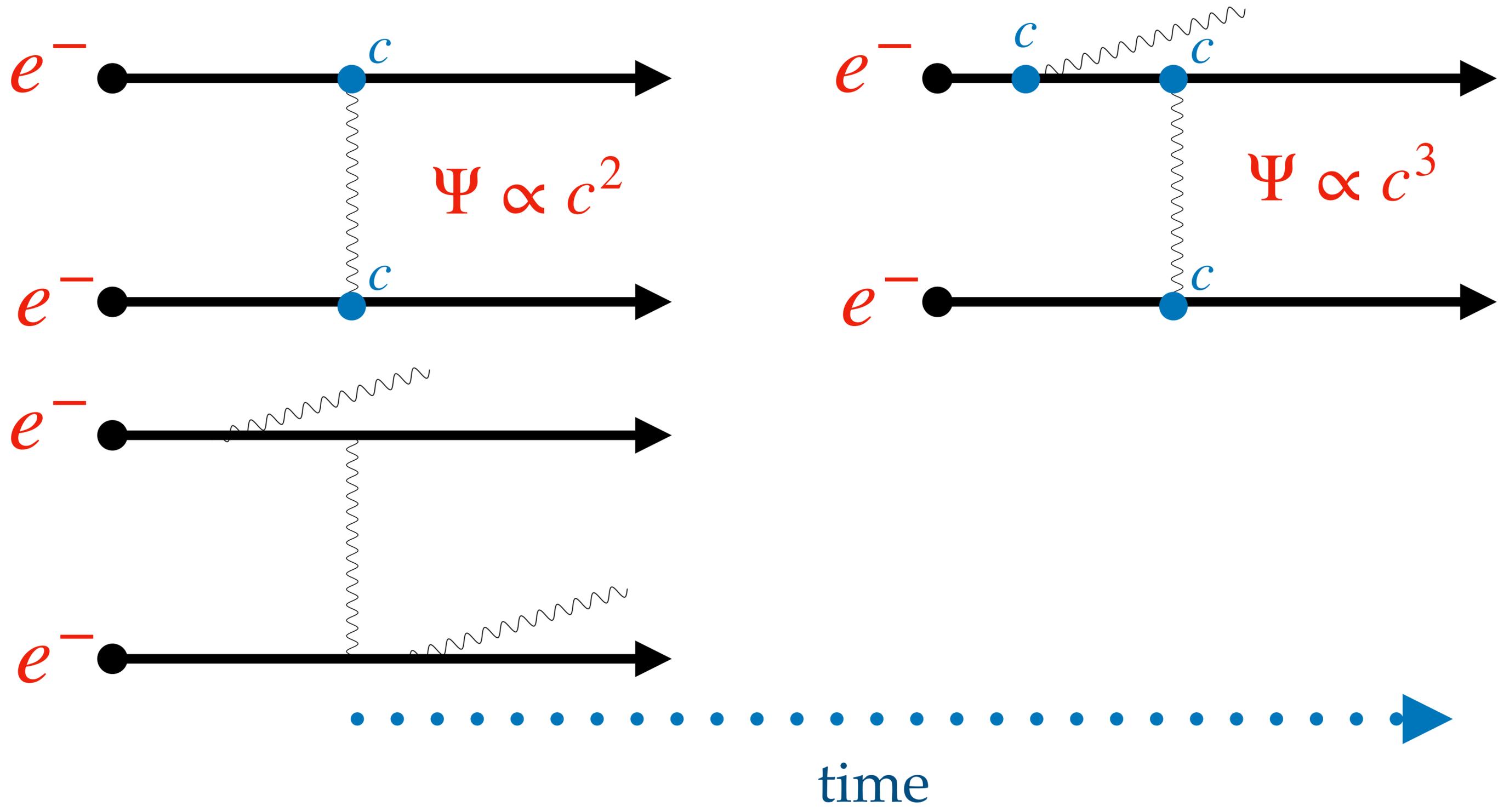
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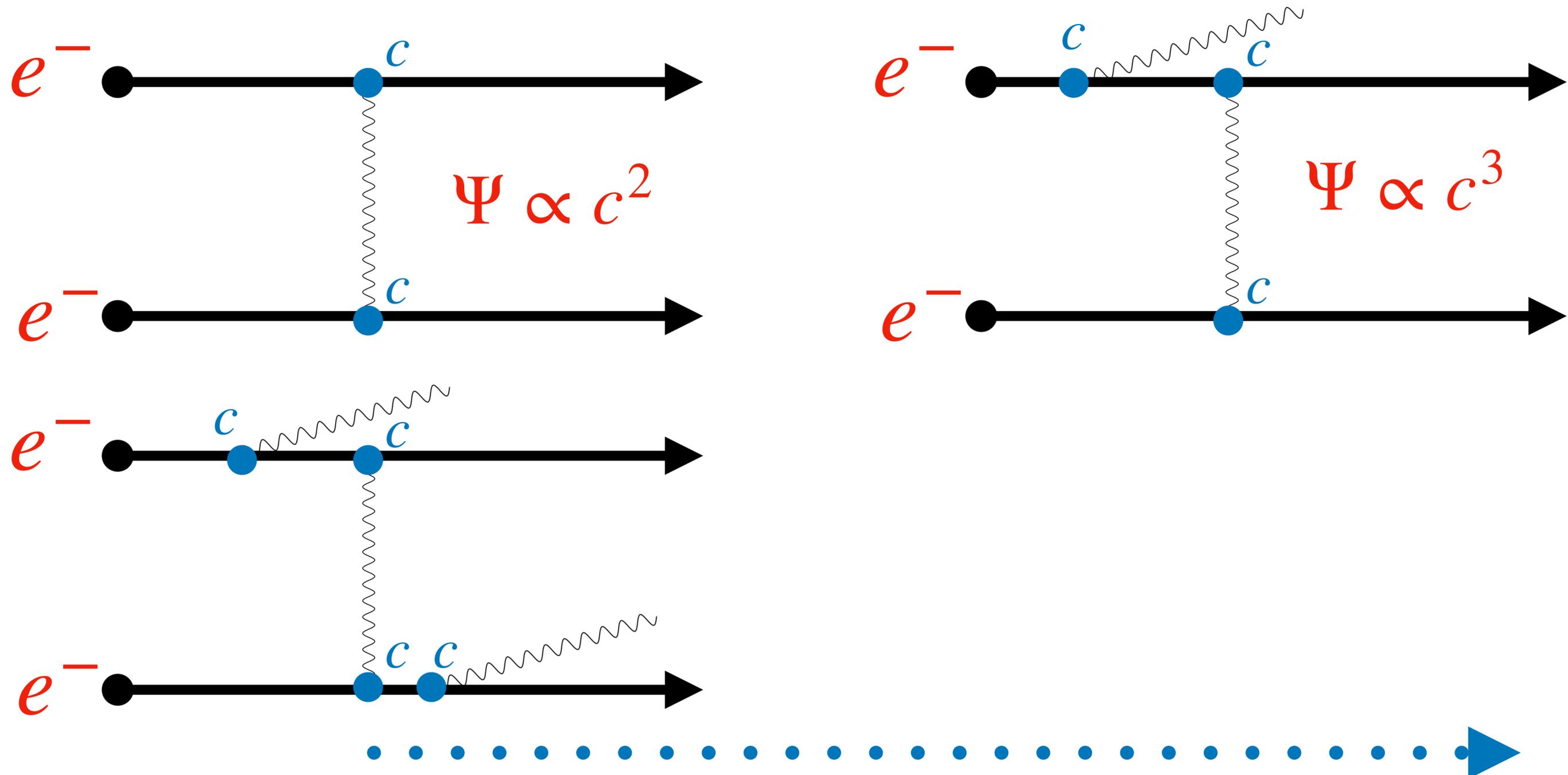
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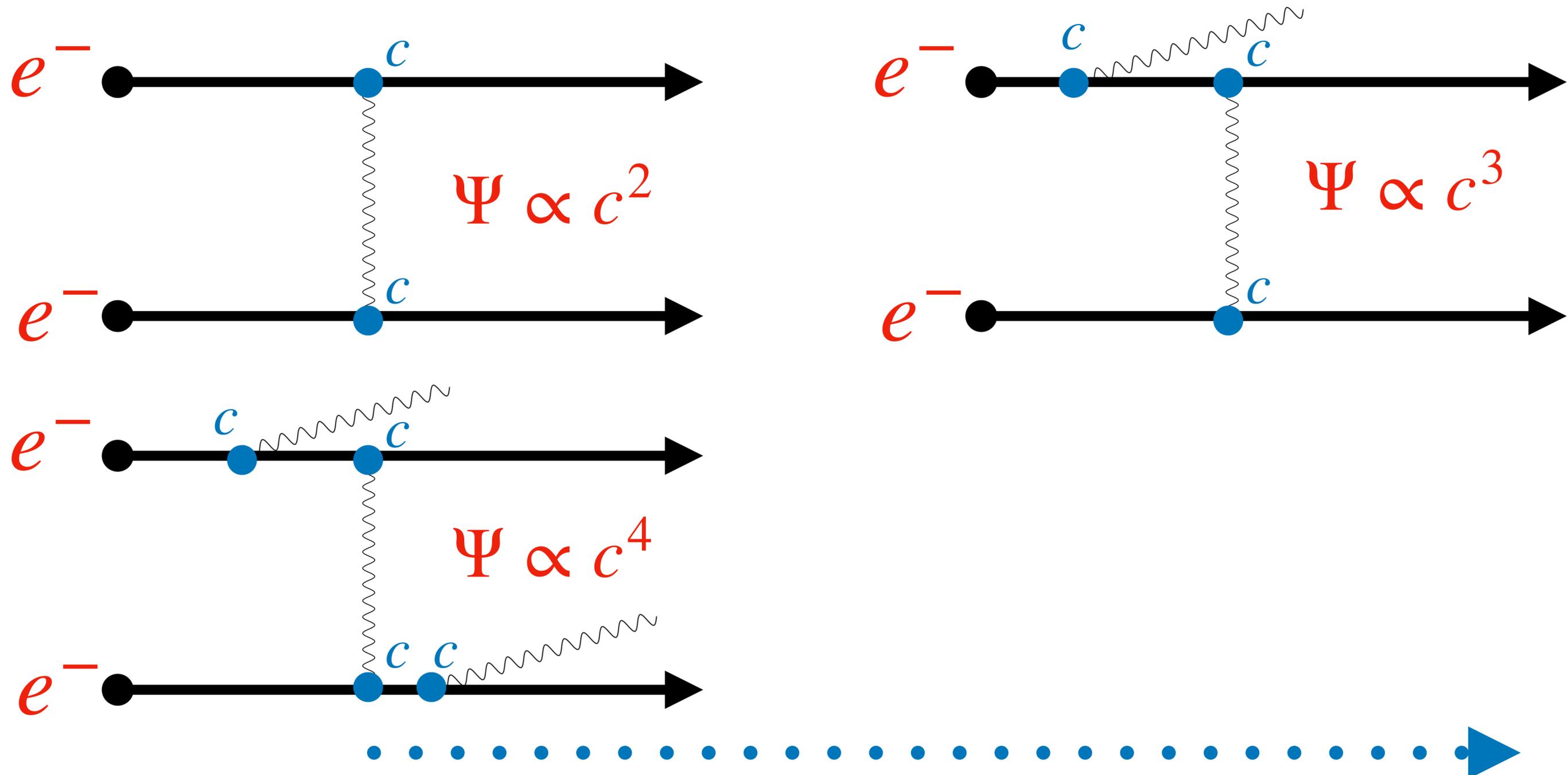


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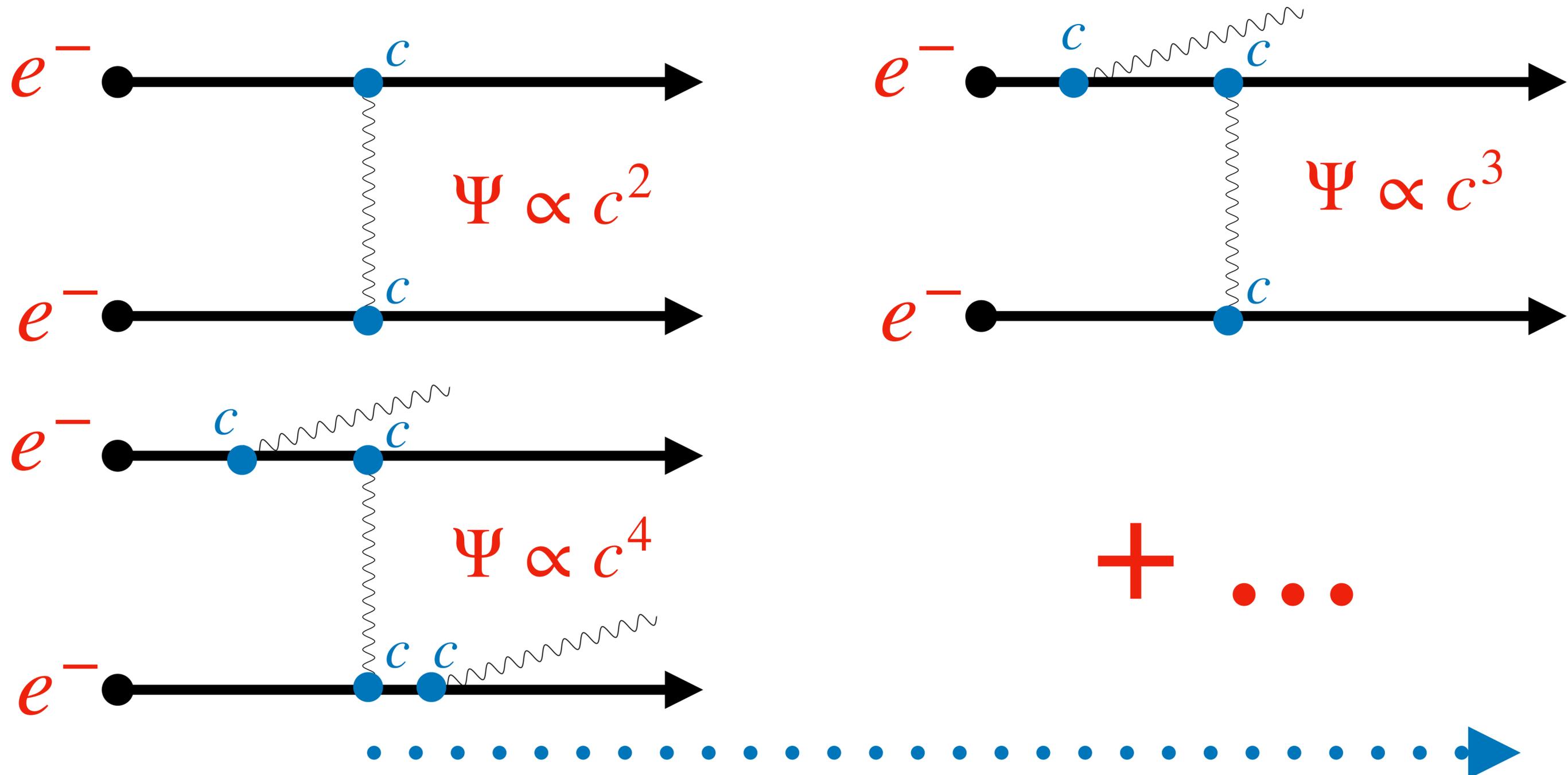


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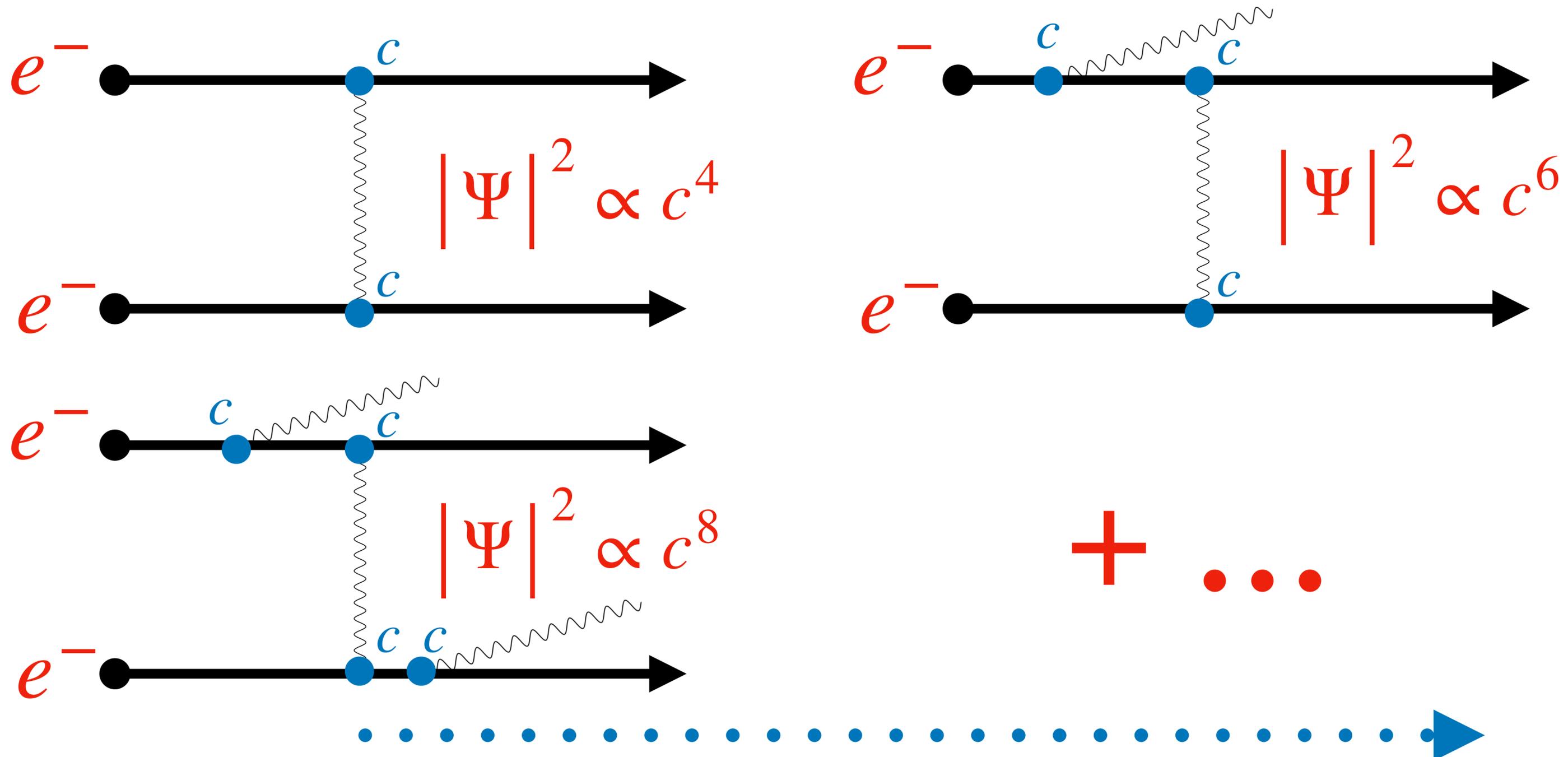
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From Theory to Experiment and Back Again

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$$|\Psi|^2 \propto \alpha c^4 + \beta c^6 + \gamma c^8 + \dots$$

If $c^2 < 1$ then we can simply add terms for increased precision!



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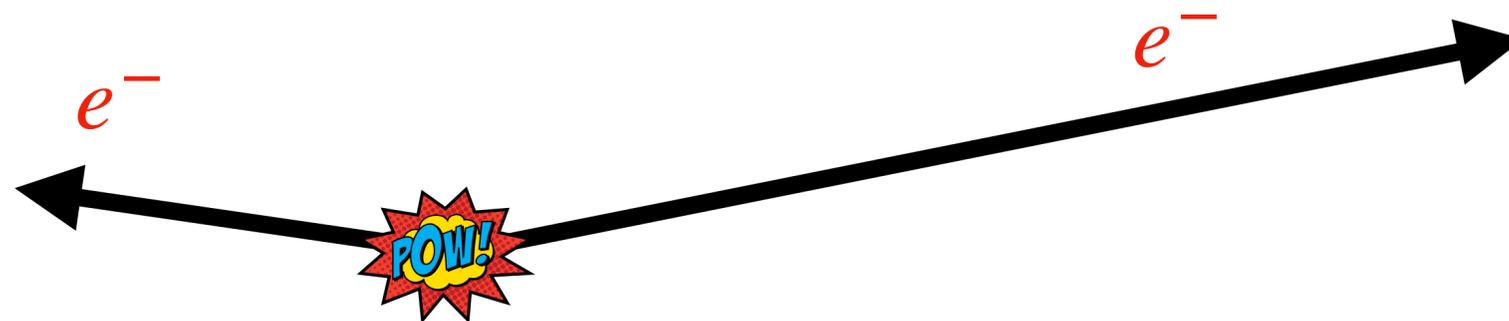
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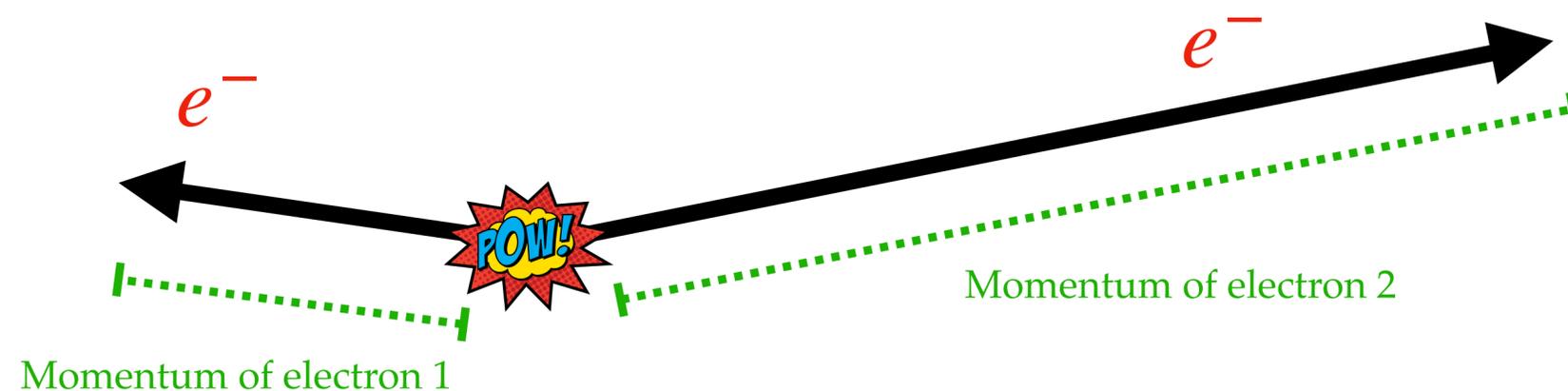
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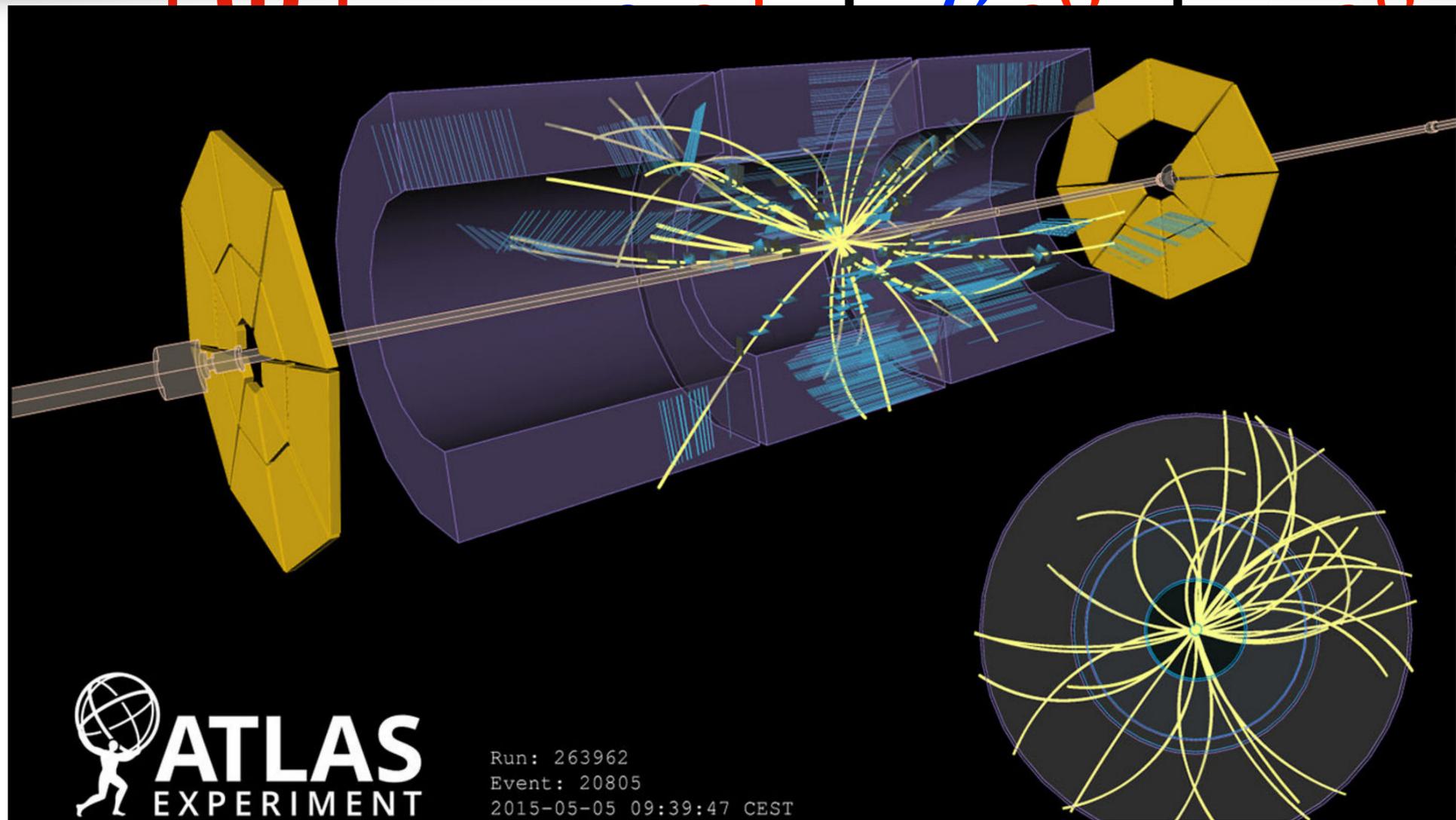


F

We have a way to describe what we observe at experiments!

If $c^2 < 1$

In reality:



+ ...

recision!

the particles involved!



Monte Carlo Simulations

- So, can we calculate everything using pen and paper? **Not really...**

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- **Solutions**: 1. Use approximate c^n for n large & 2. model situations where $c^2 \geq 1$!

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- **Solutions**: 1. Use **approximate** c^n for n large & 2. **model** situations where $c^2 \geq 1$!

- **Monte Carlo Simulations accomplish this!**



Based on **randomness** and **probability**.
Just like Quantum Mechanics!

Pictured: Casino Monte Carlo, Monaco.



How do we make sense of it all?



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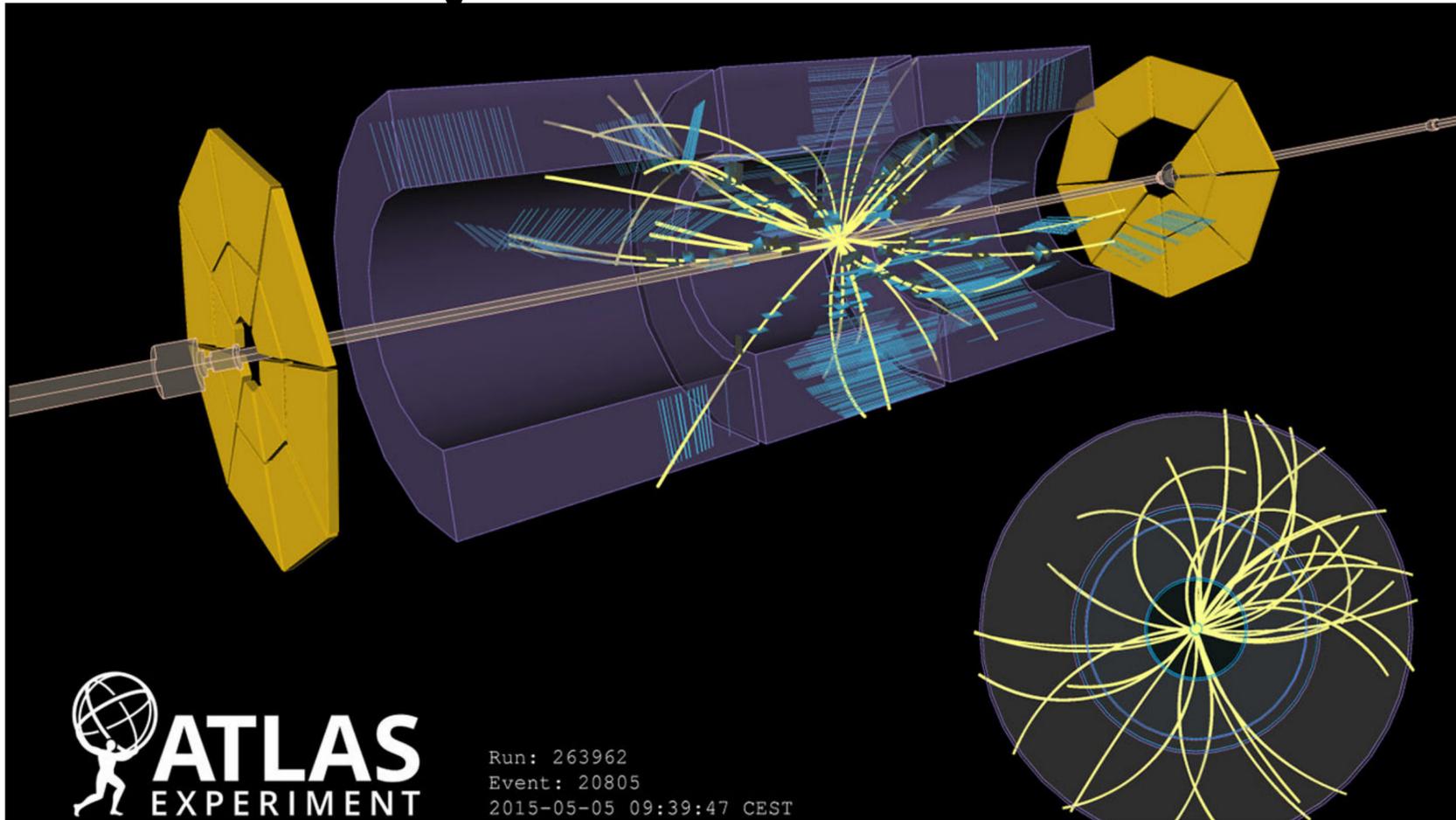
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Experiment

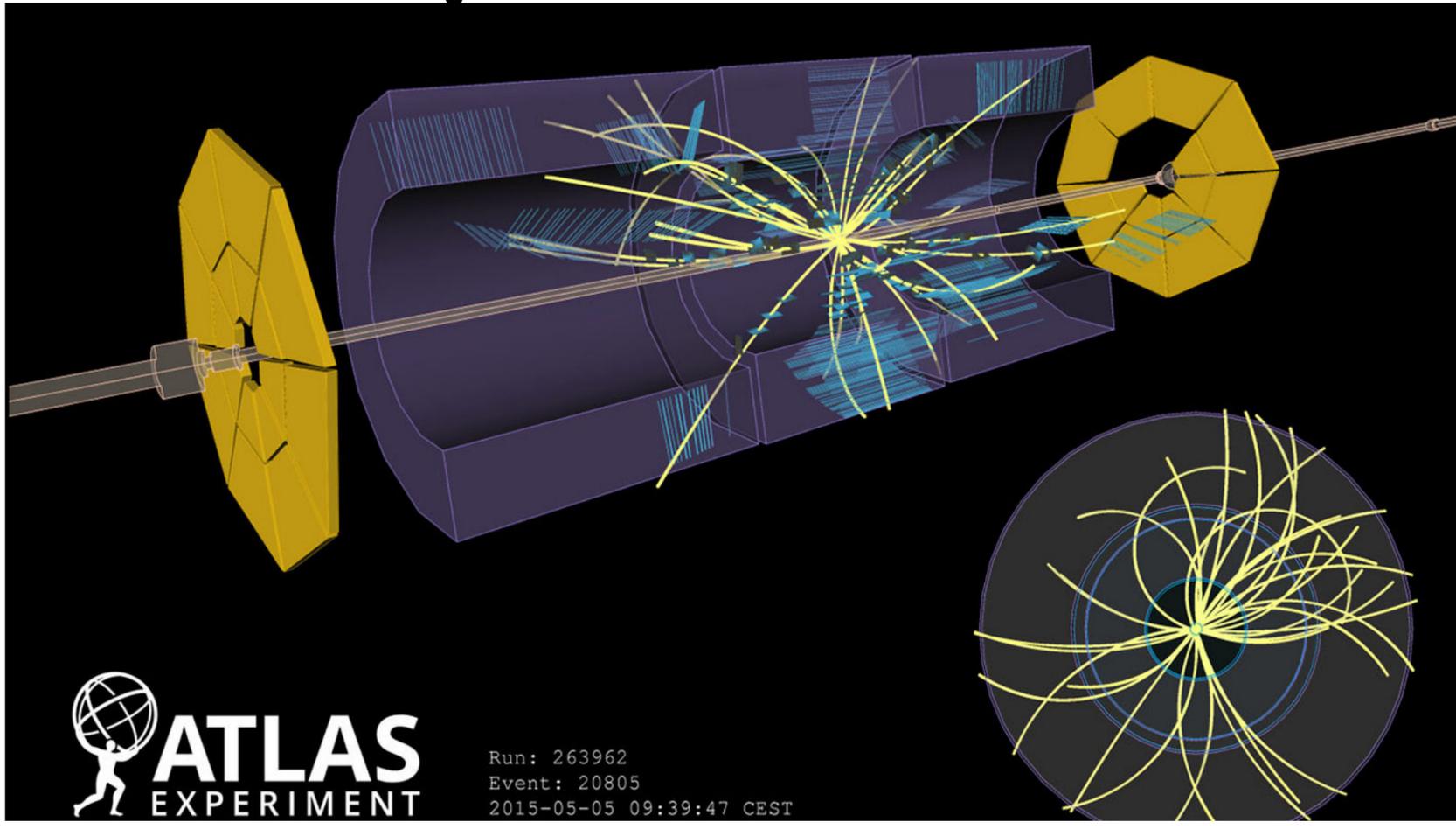


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Simulations

Experiment



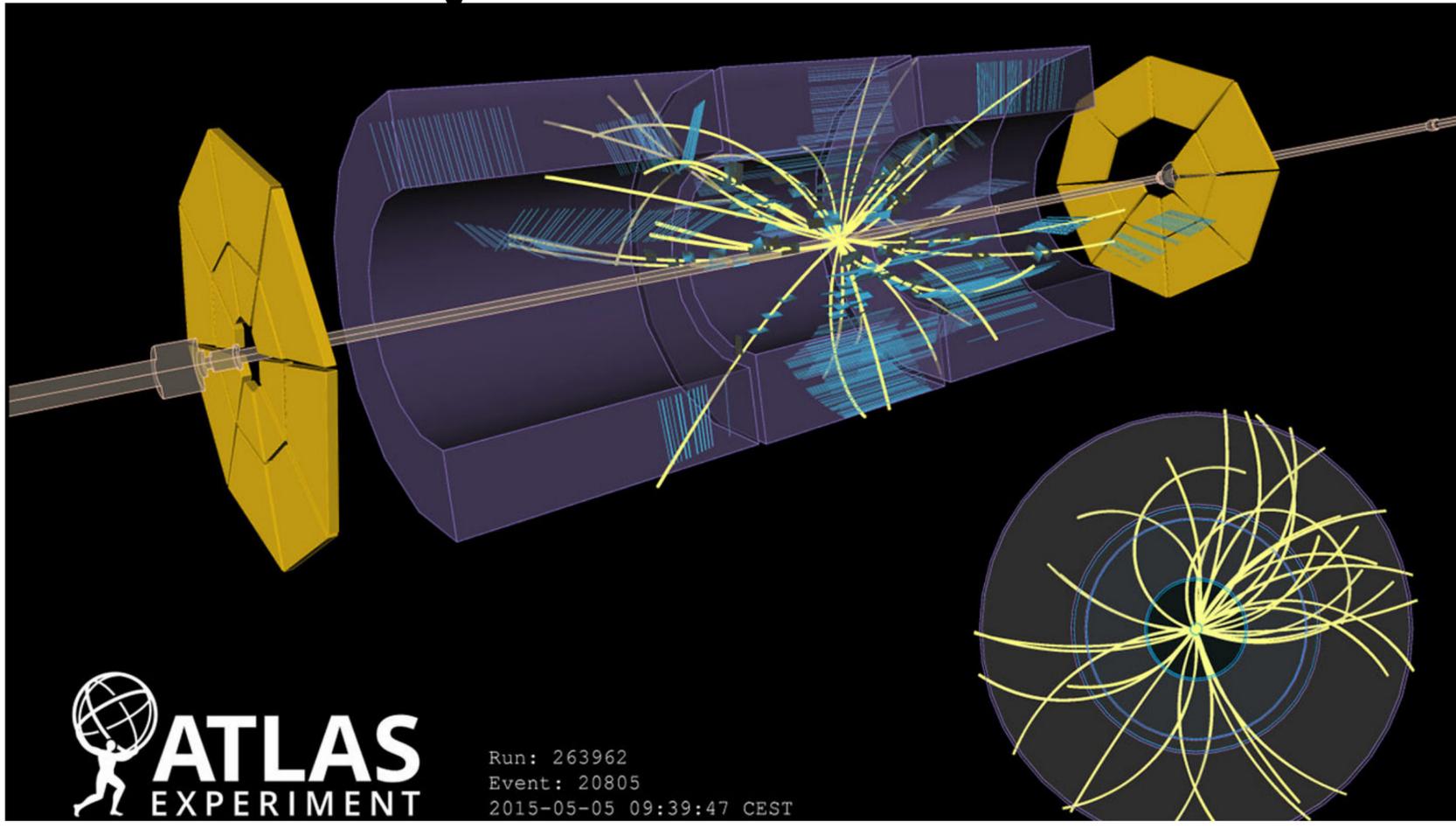
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Shut up and Simulate!

Simulations



Experiment

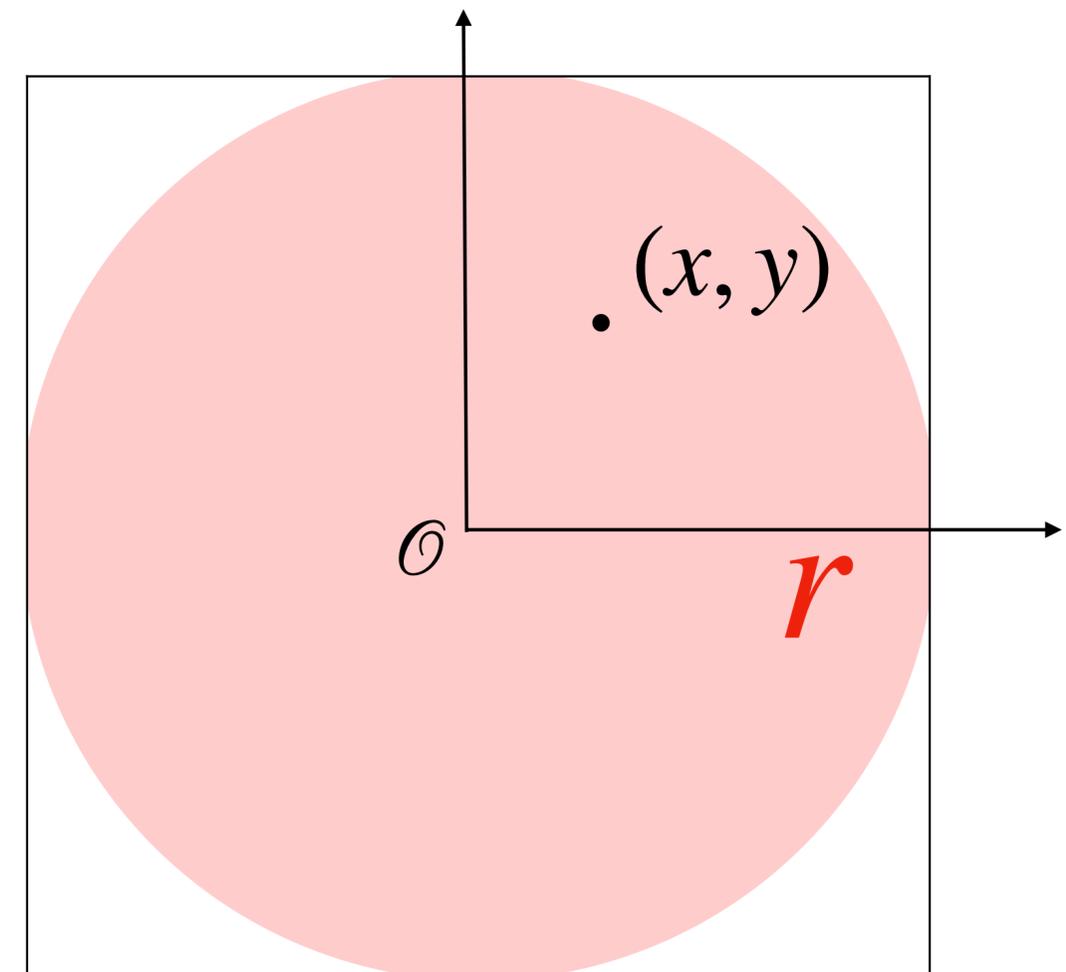


Simulations via the Monte Carlo method: Calculate π !



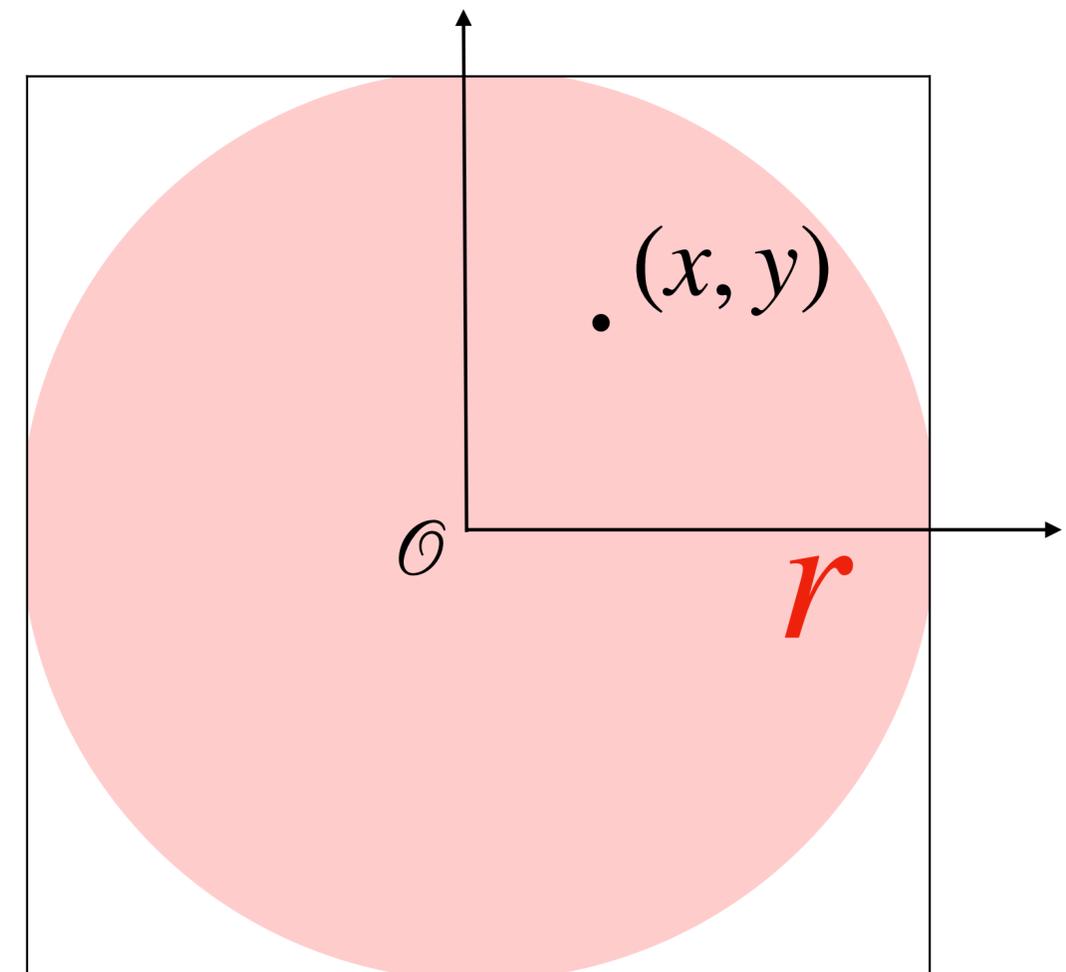
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 $\mathcal{P} = \pi r^2 / (4r^2) = \pi/4$.



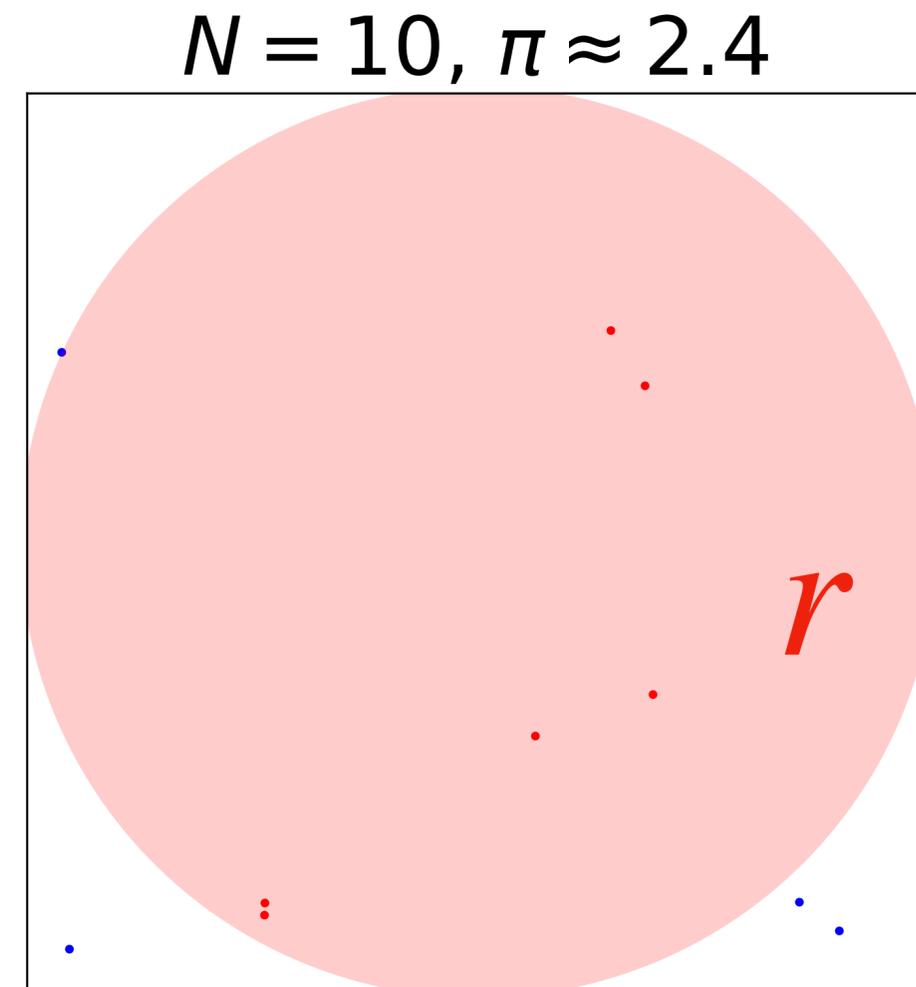
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 - Pick N random, uniform points inside the square.
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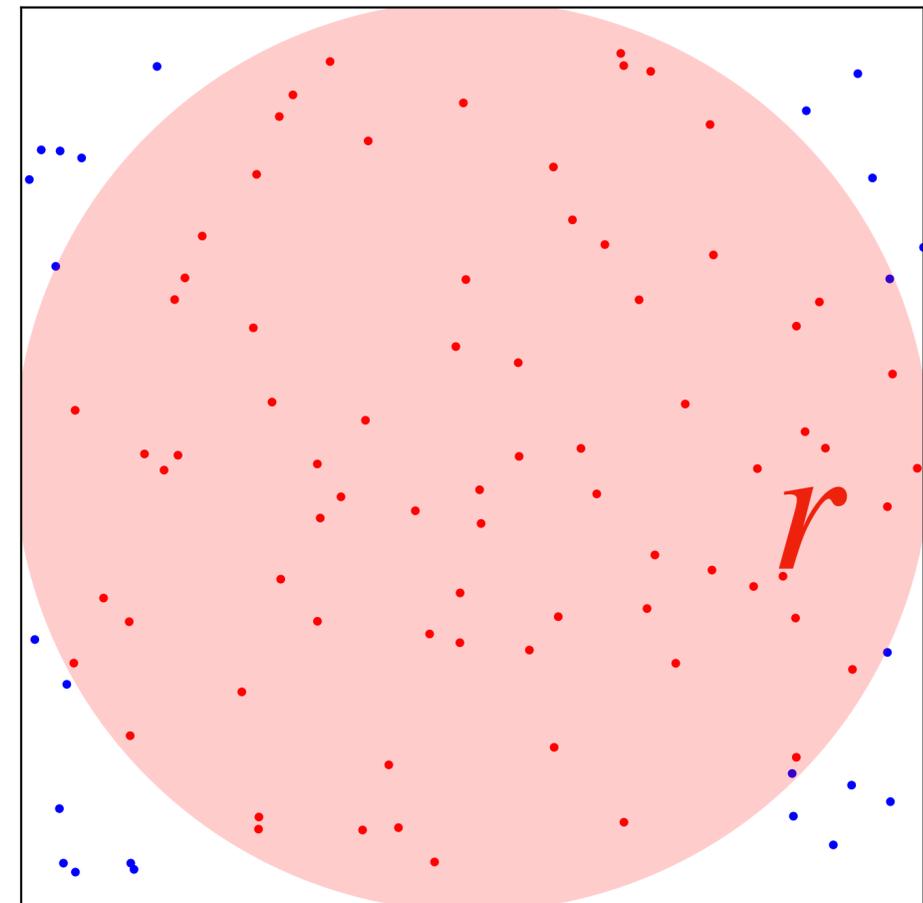
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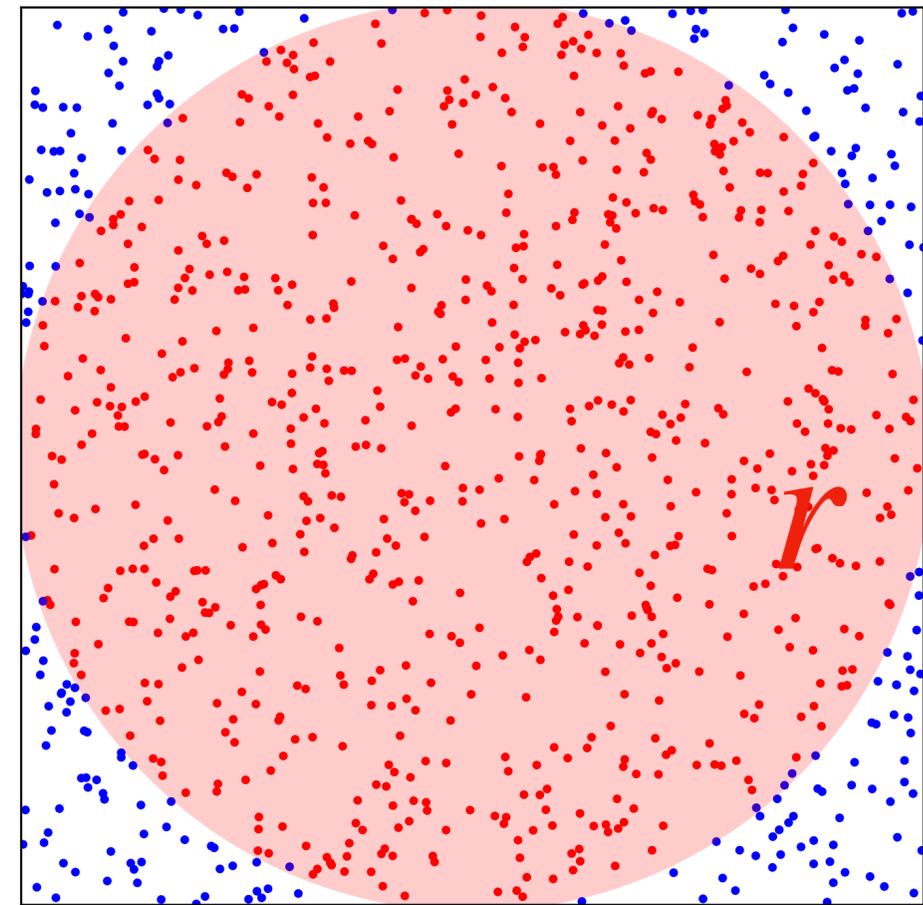
$N = 10^2, \pi \approx 3.04 \pm 0.35$



Simulations via the Monte Carlo method: Calculate π !

- “Theory”:
 - Area of a circle: $A = \pi r^2$.
 - Distance of any point (x, y) from origin \mathcal{O} , by Pythagorean theorem: $d = \sqrt{x^2 + y^2}$.
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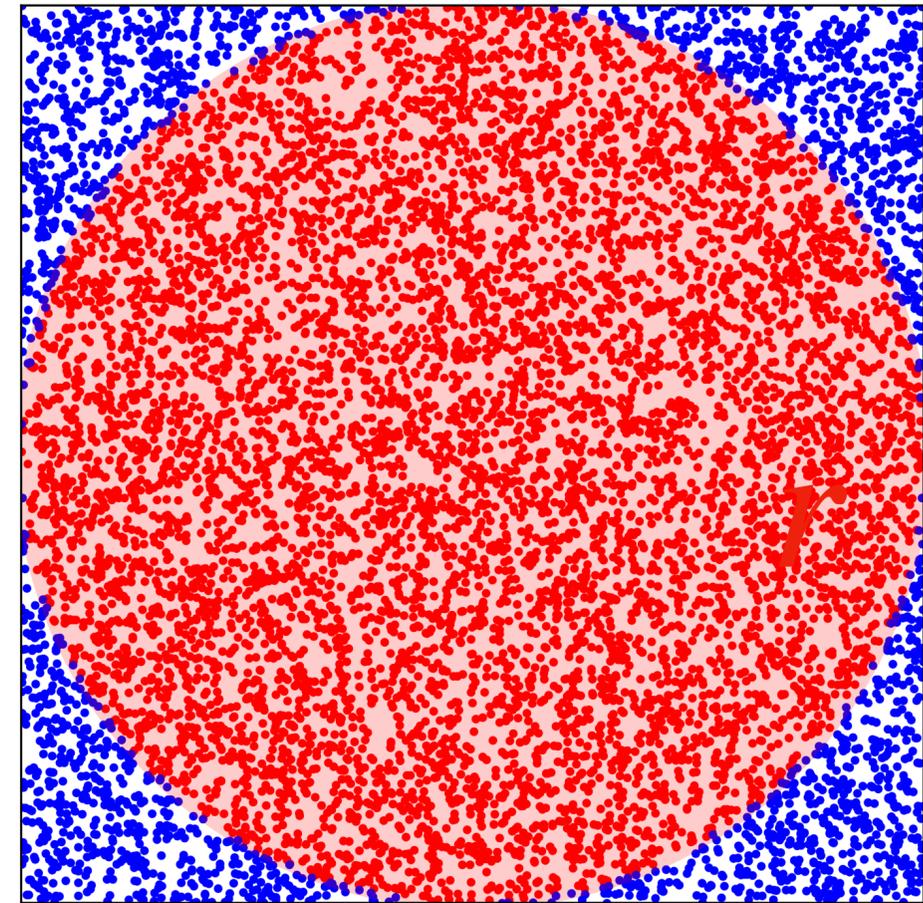
$$N = 10^3, \pi \approx 3.08 \pm 0.11$$



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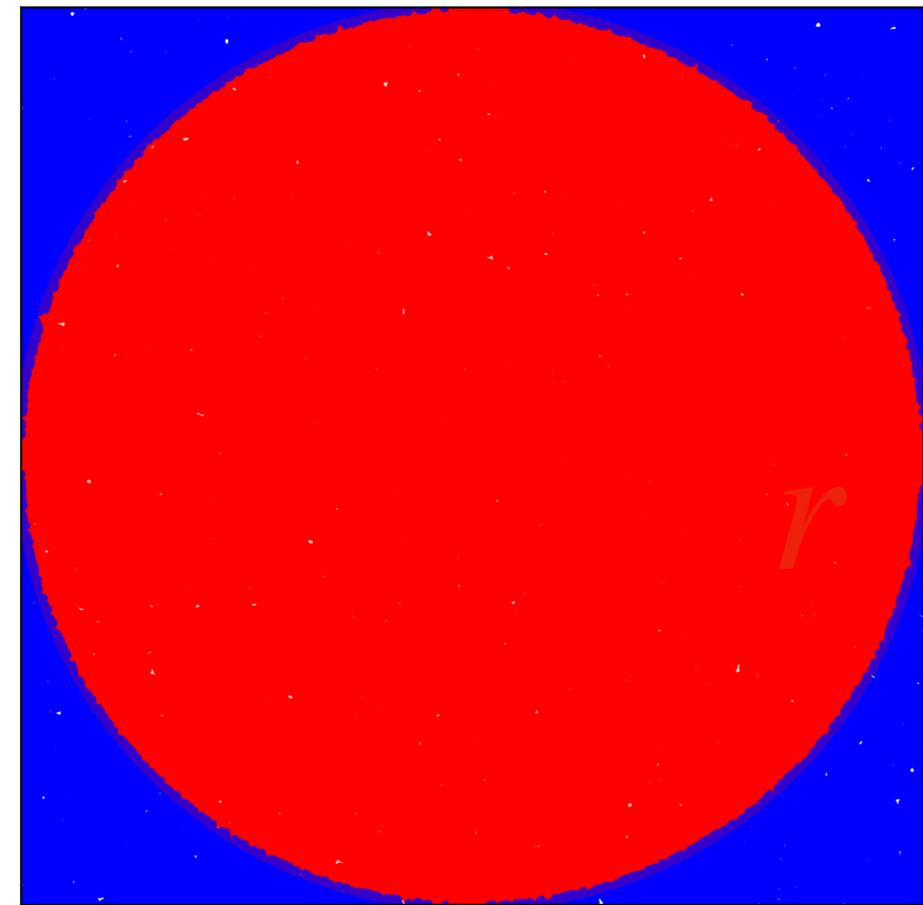
$$N = 10^4, \pi \approx 3.13 \pm 0.04$$



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$$N = 10^5, \pi \approx 3.14 \pm 0.01$$

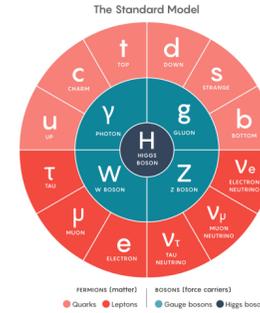


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- A set of particles and their interactions! e.g.:

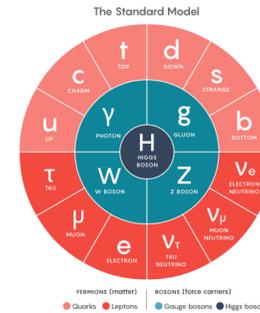


- Calculate probabilities via Quantum Field Theory → actually $|\mathcal{M}|^2$ (the matrix element).

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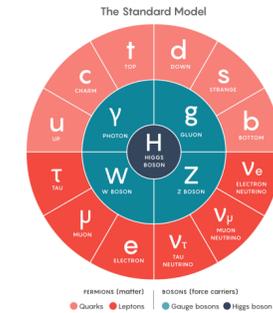
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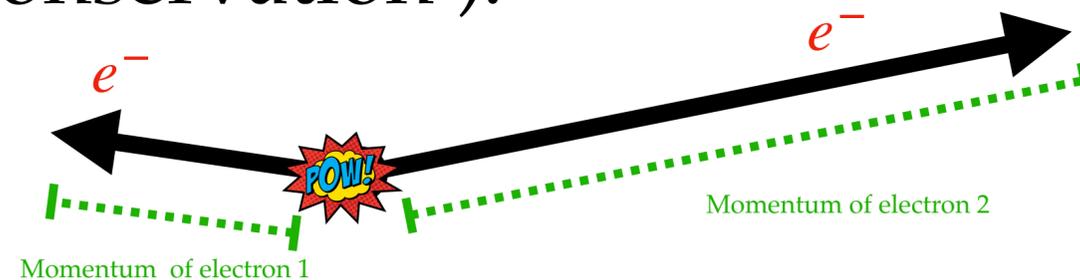


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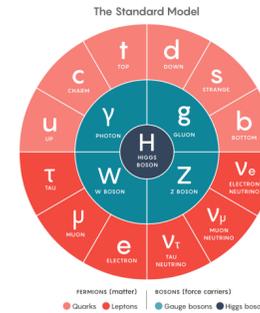


probability $\sim |\mathcal{M}|^2(p_1, p_2, \dots)$

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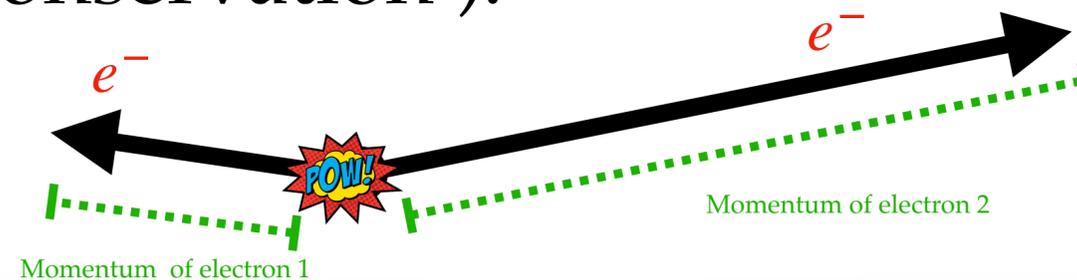
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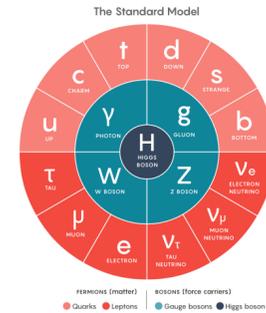
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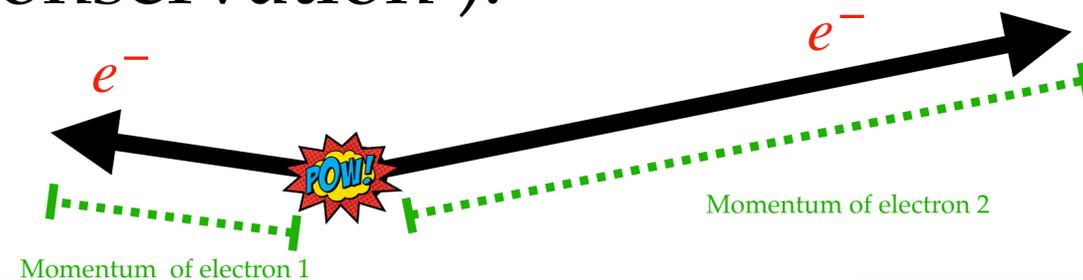
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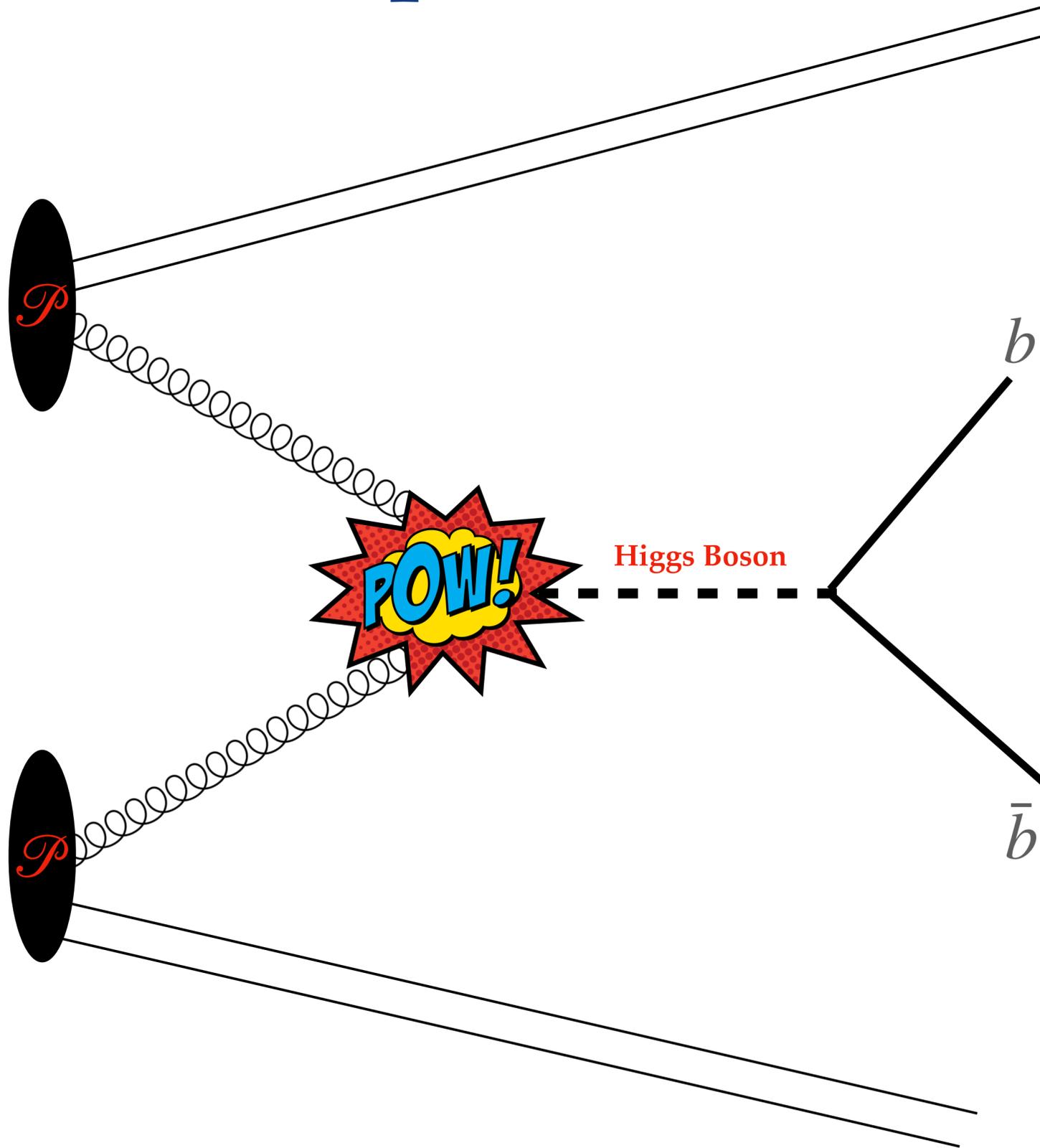


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[e.g. “How-to: write a parton-level Monte Carlo particle physics event generator”, AP, [arXiv:1412.4677](https://arxiv.org/abs/1412.4677)]

An Example of a Monte Carlo Event

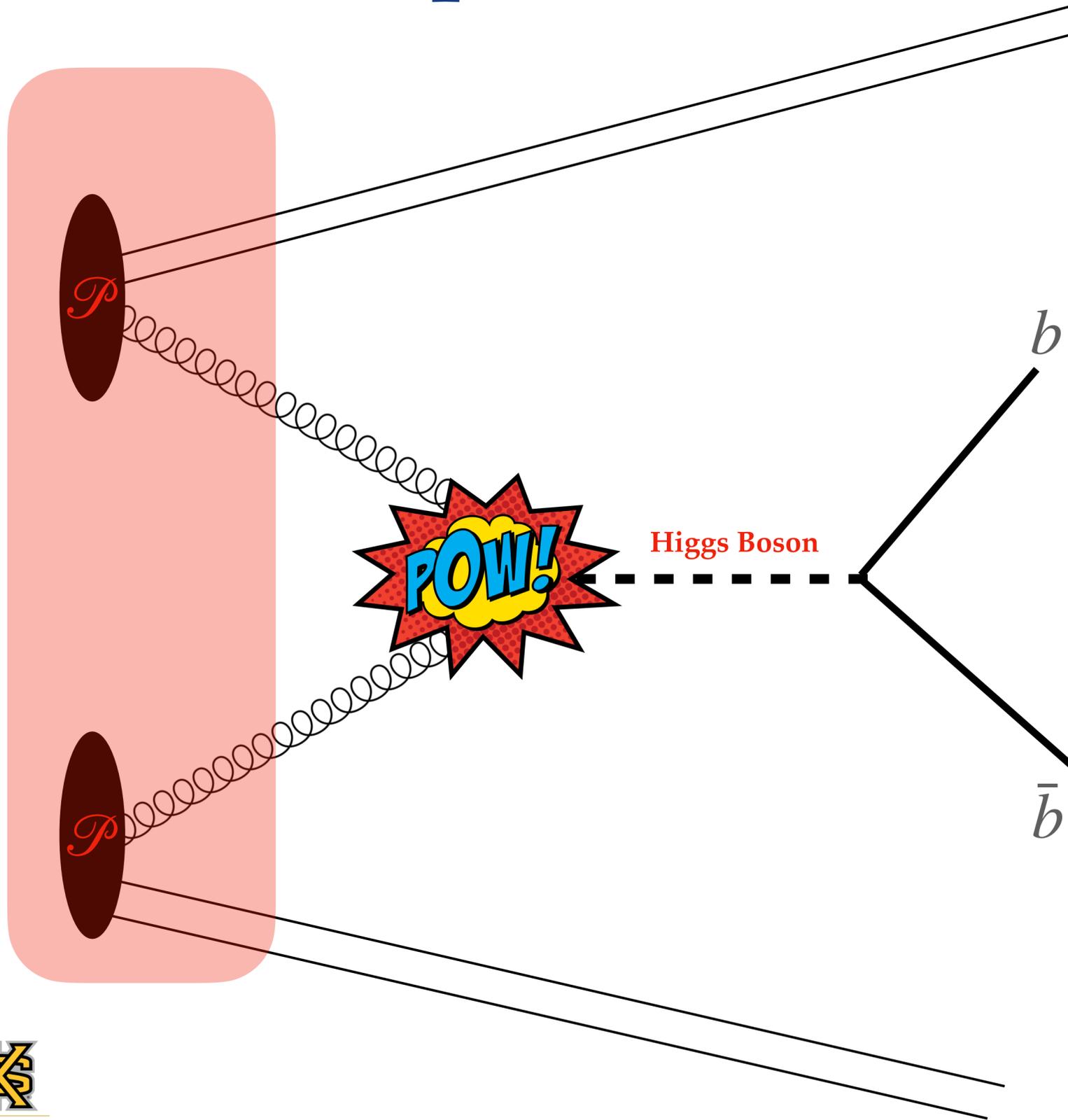


```
Event number 2 (id: LHC) performed by EventHandler
Weight: 1
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--- Colliding particles:
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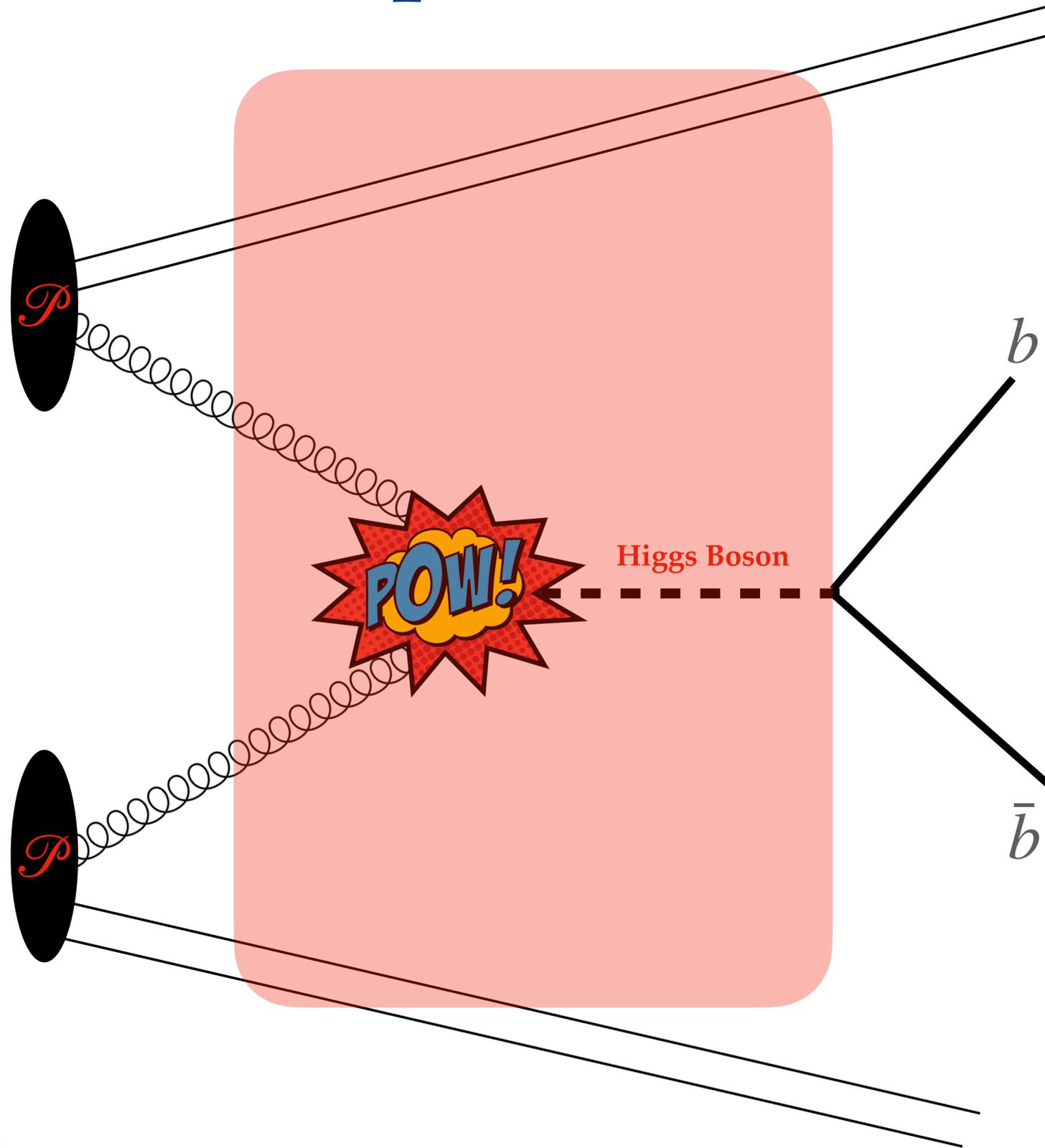


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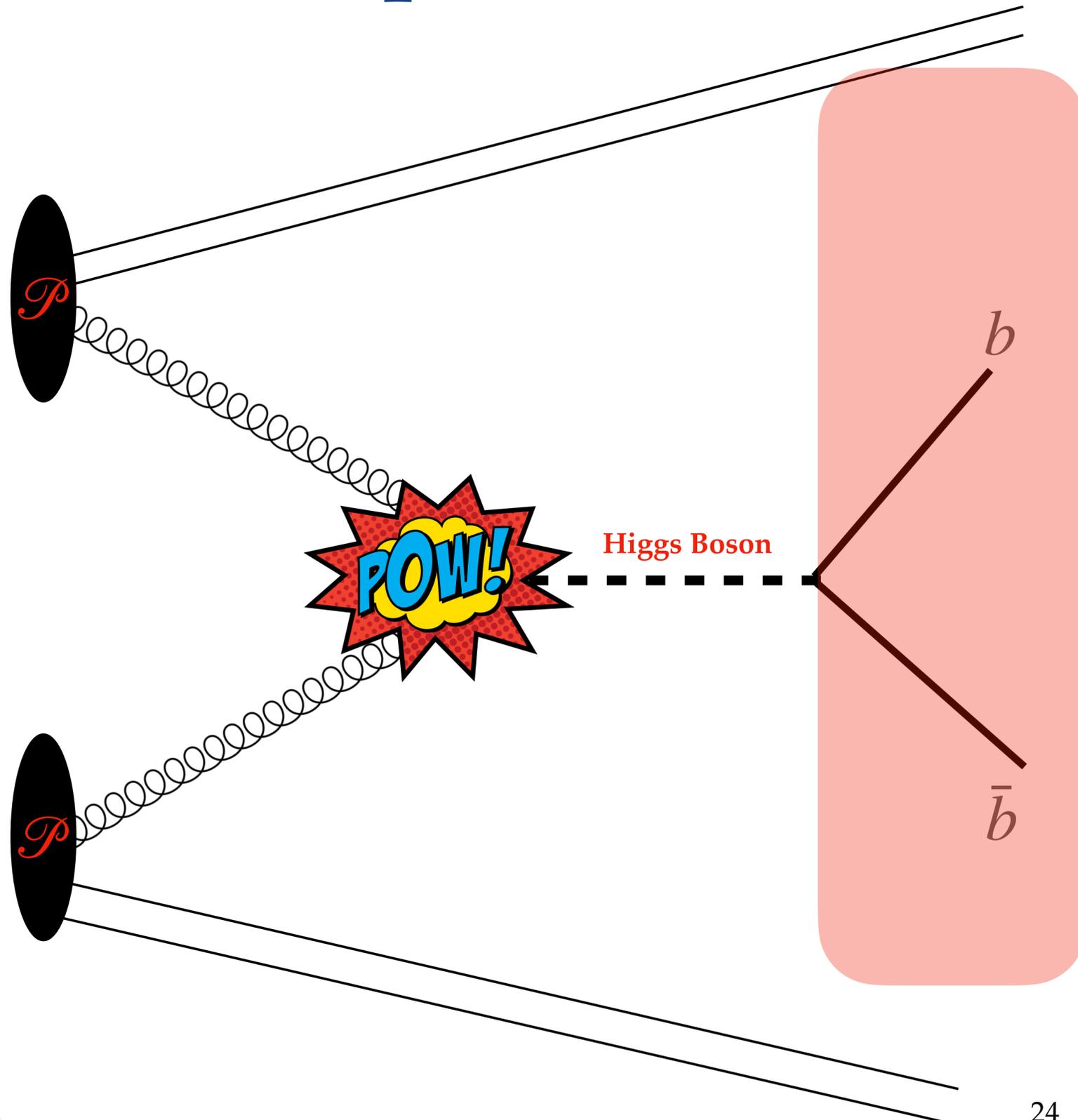


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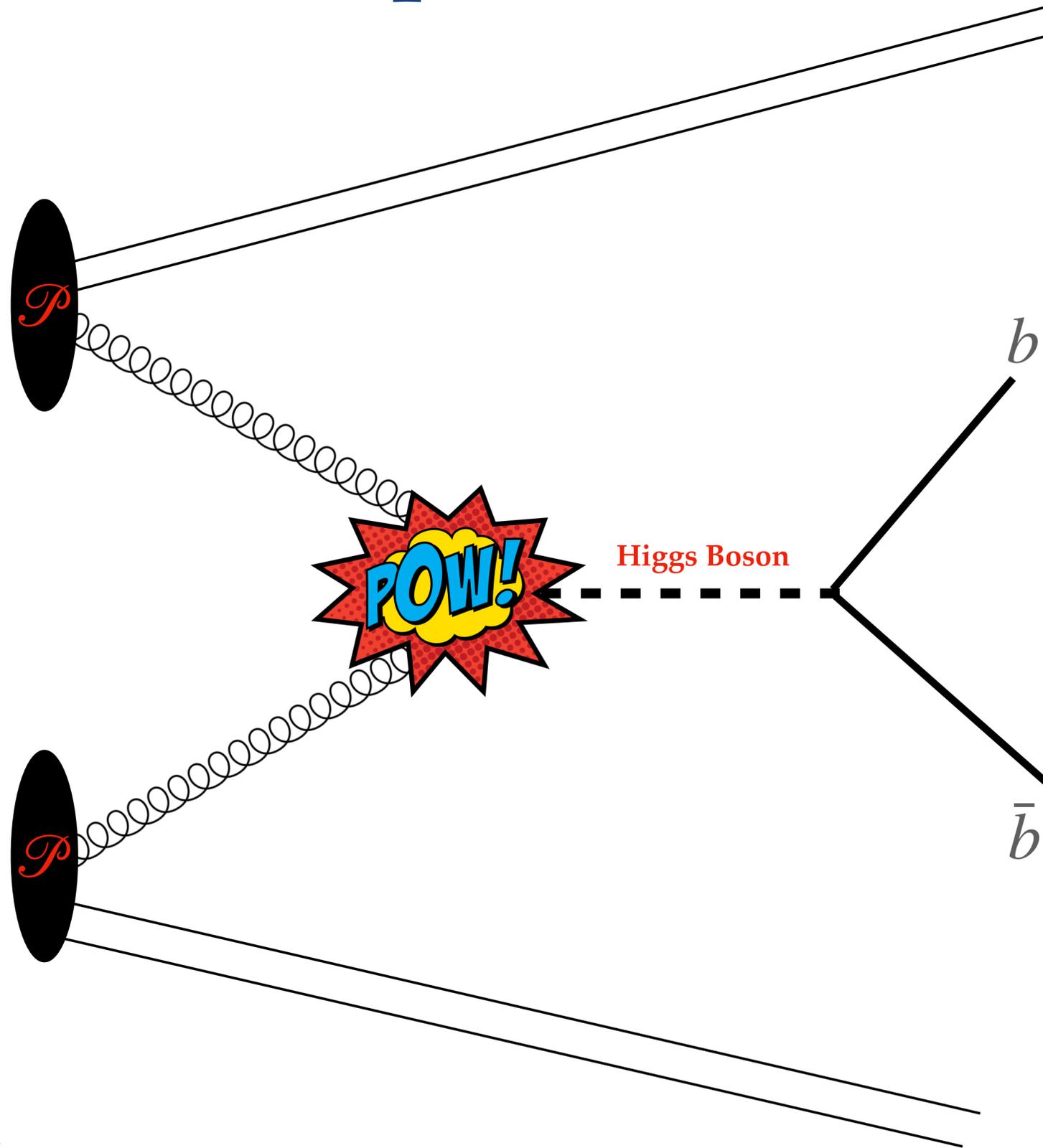


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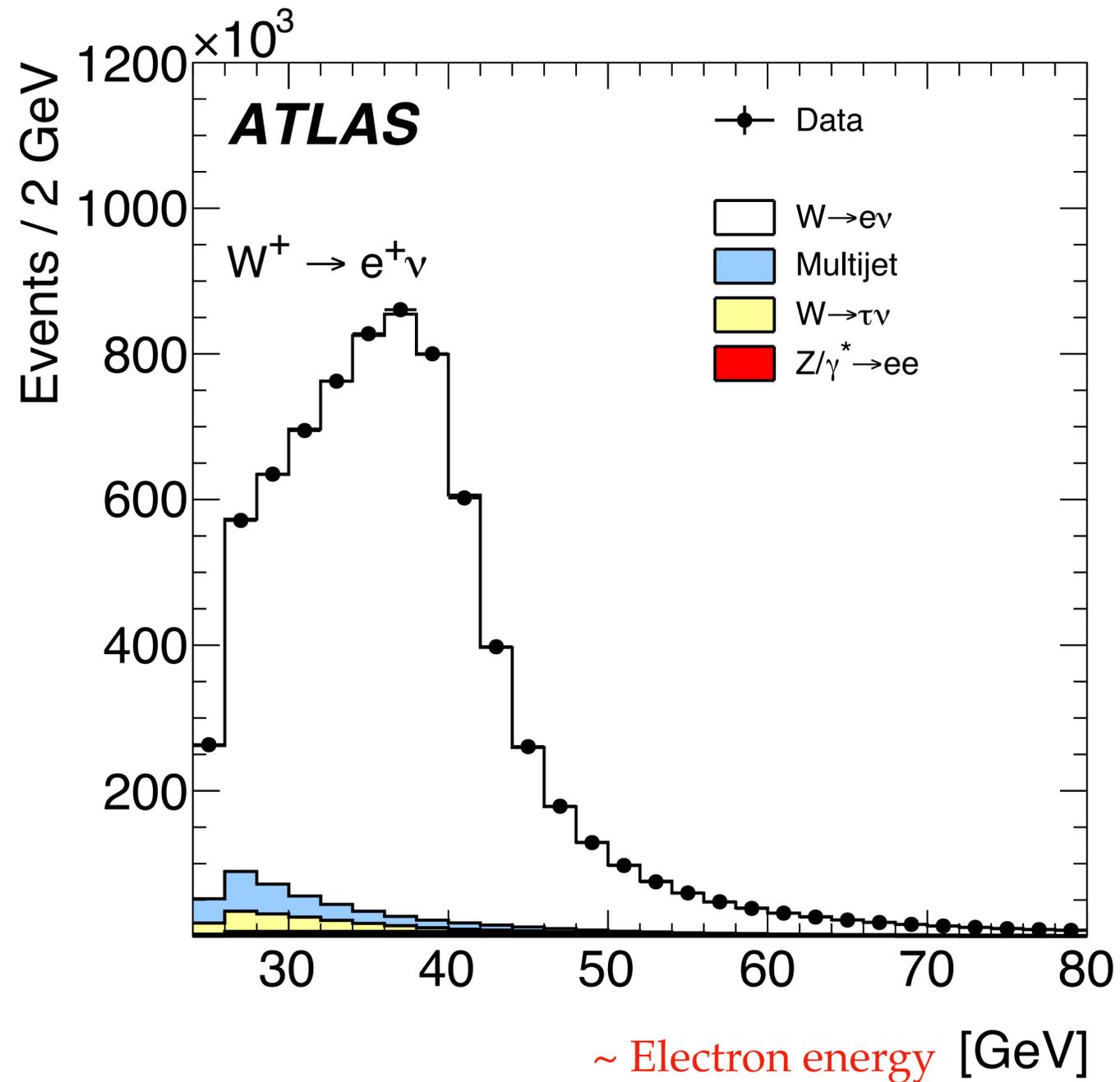
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Theory **VS** Experiments

⇒ Construct histograms & Compare to real data!

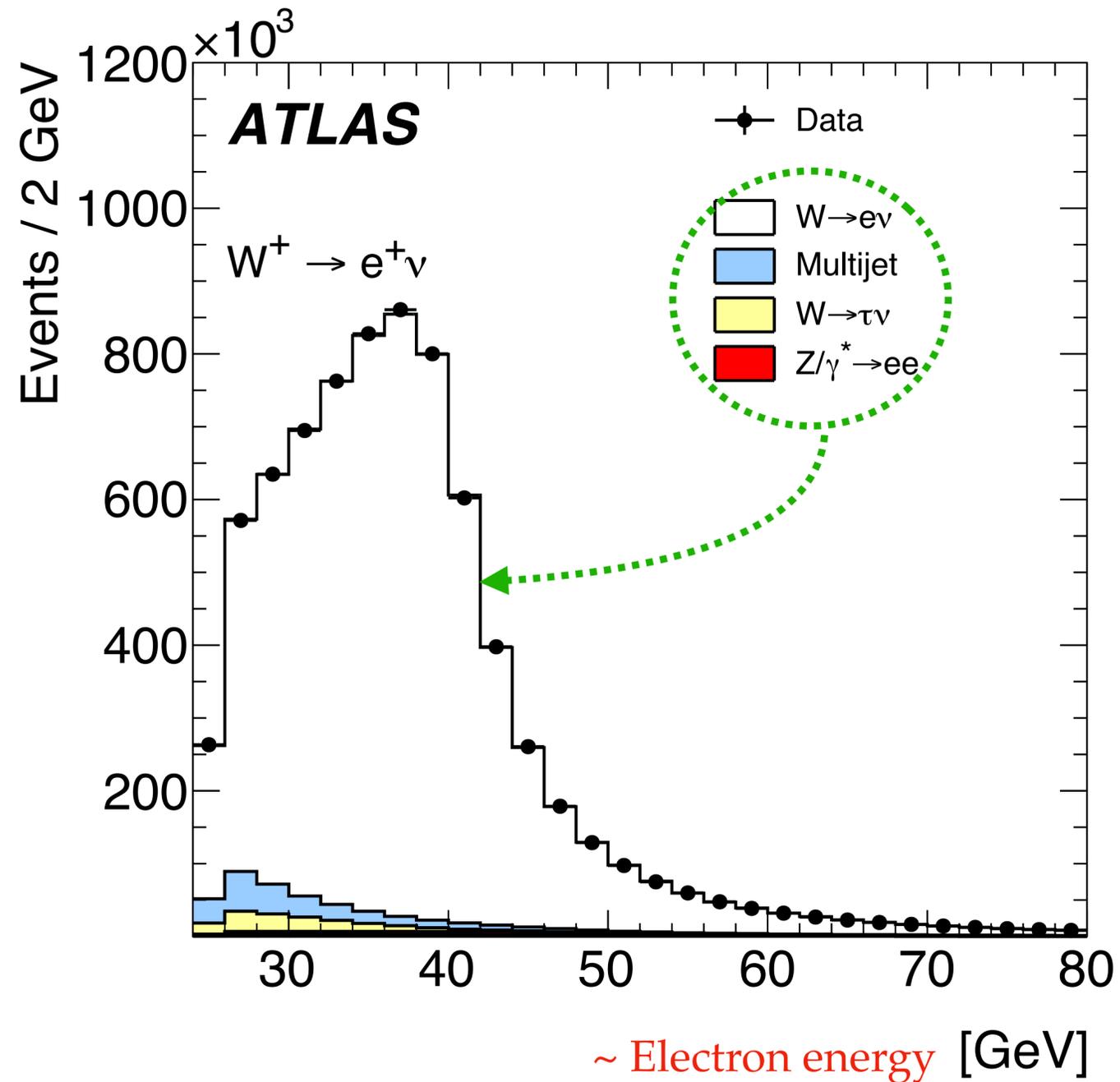
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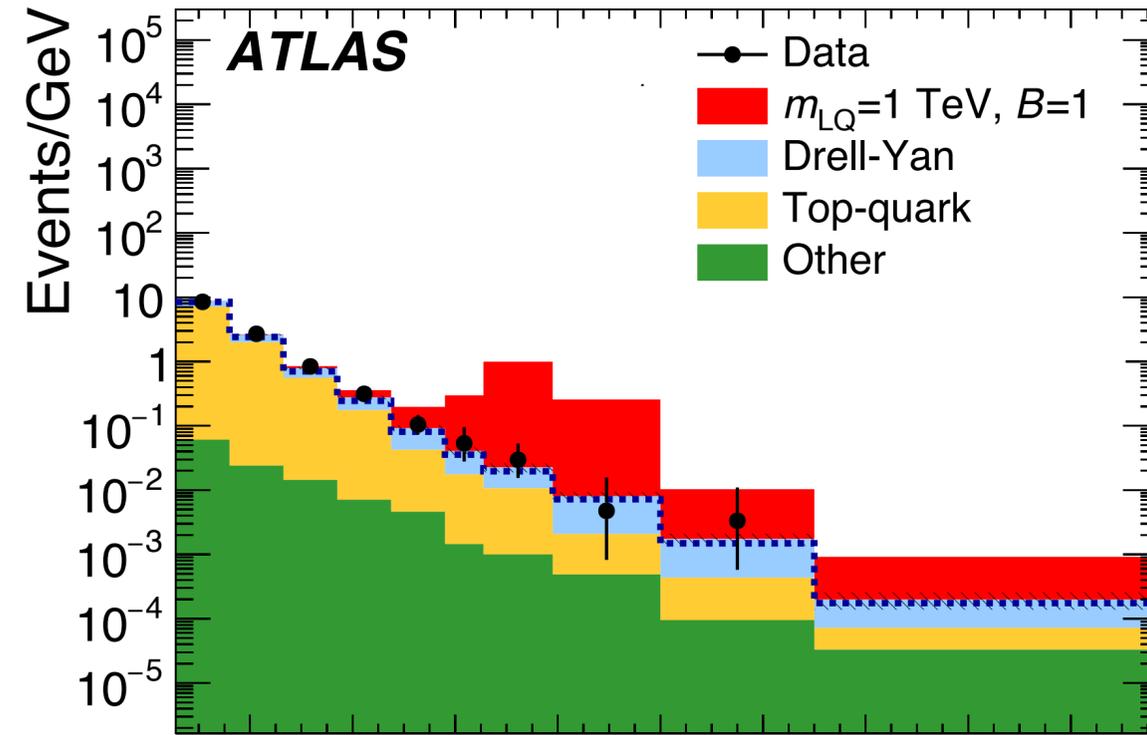
e.g.:



Searching for New Phenomena



⇒ Construct histograms &
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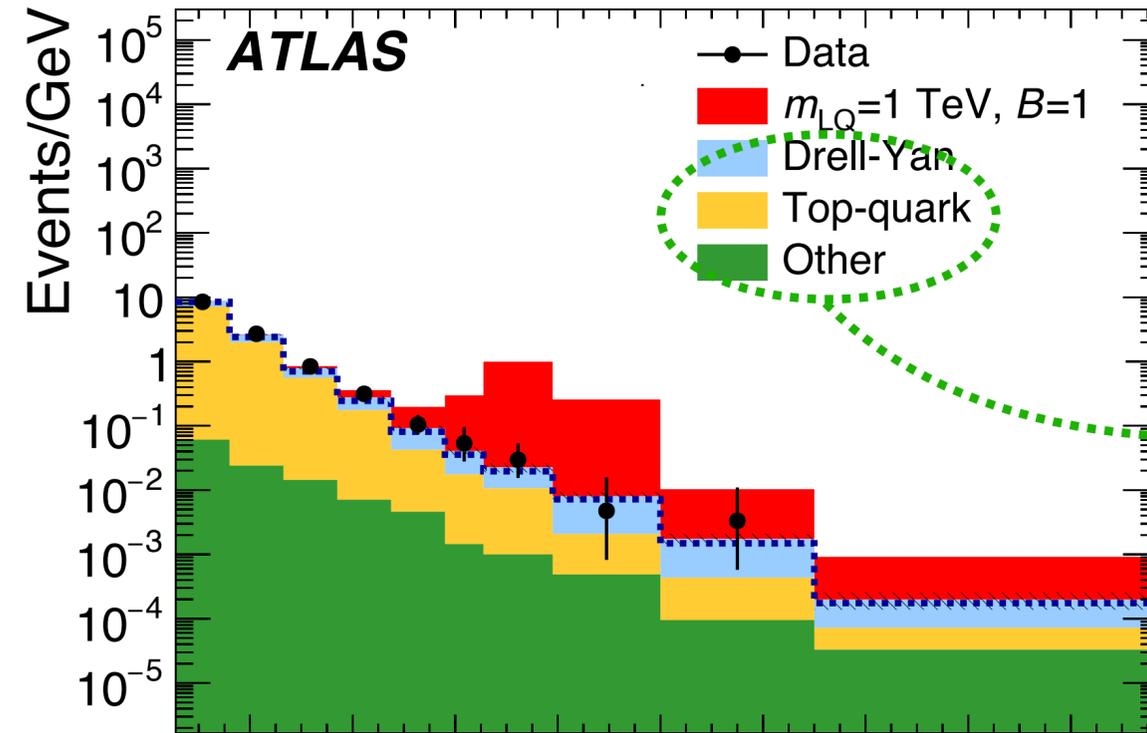
Some Observable Quantity [GeV]

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Simulation of
stuff that we
know exist!

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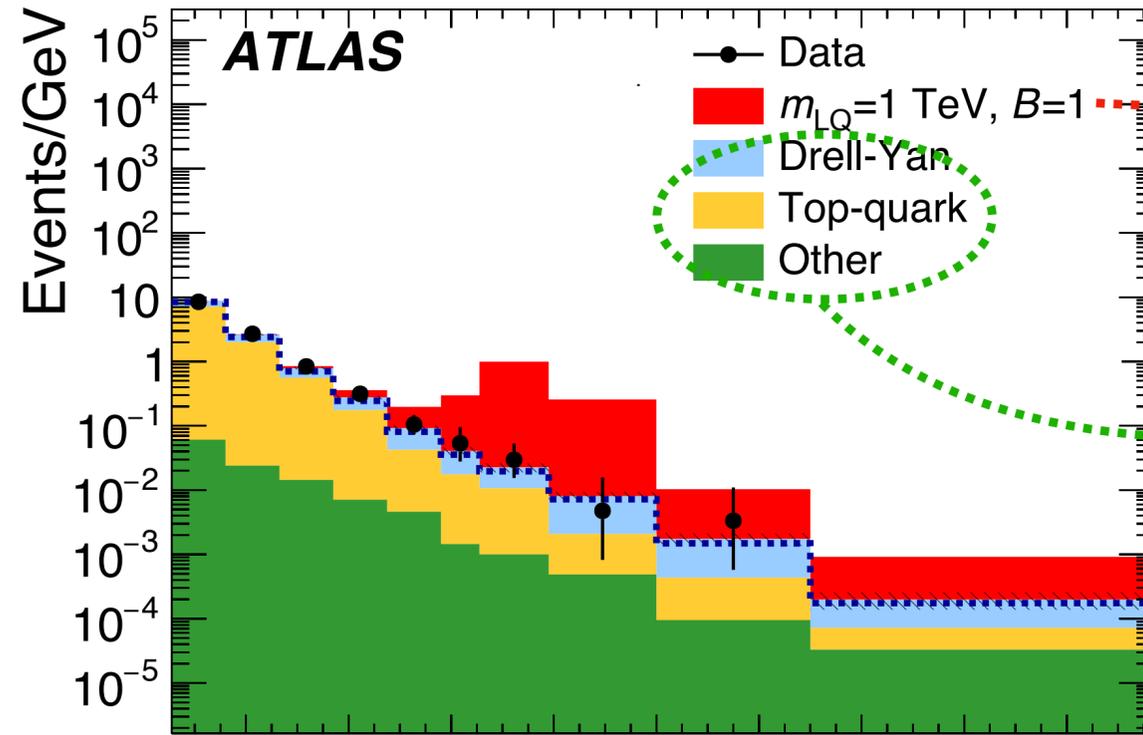


NEW!

Hypothetical particle

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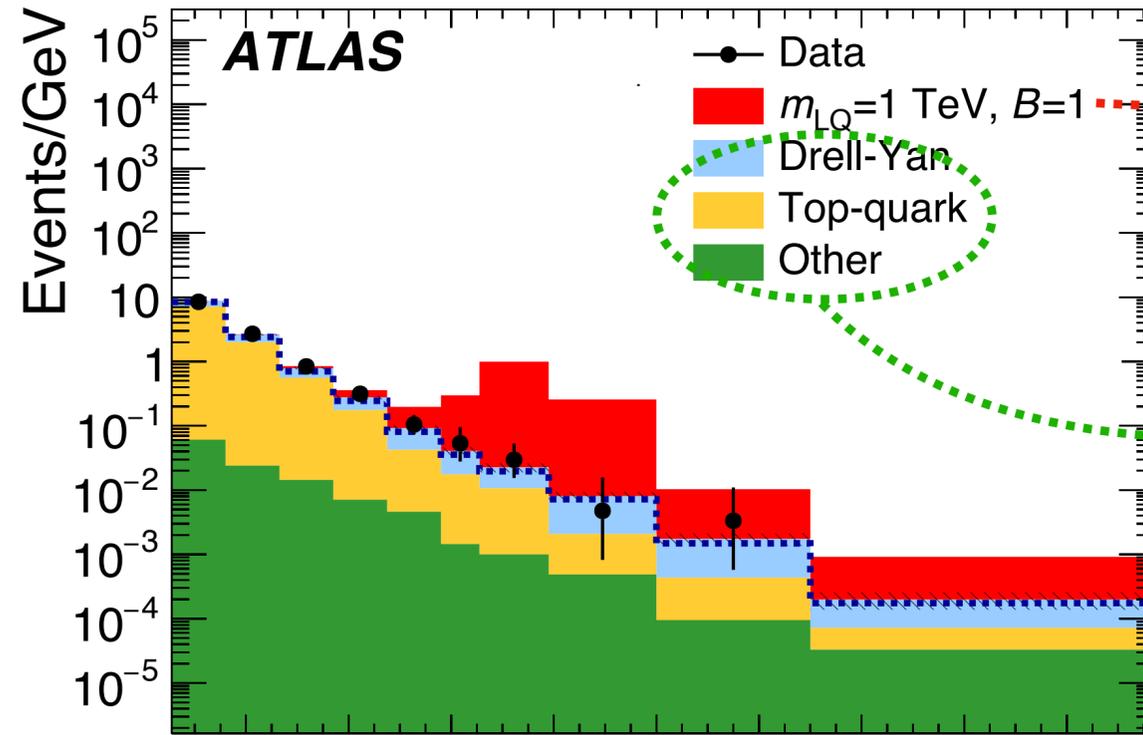


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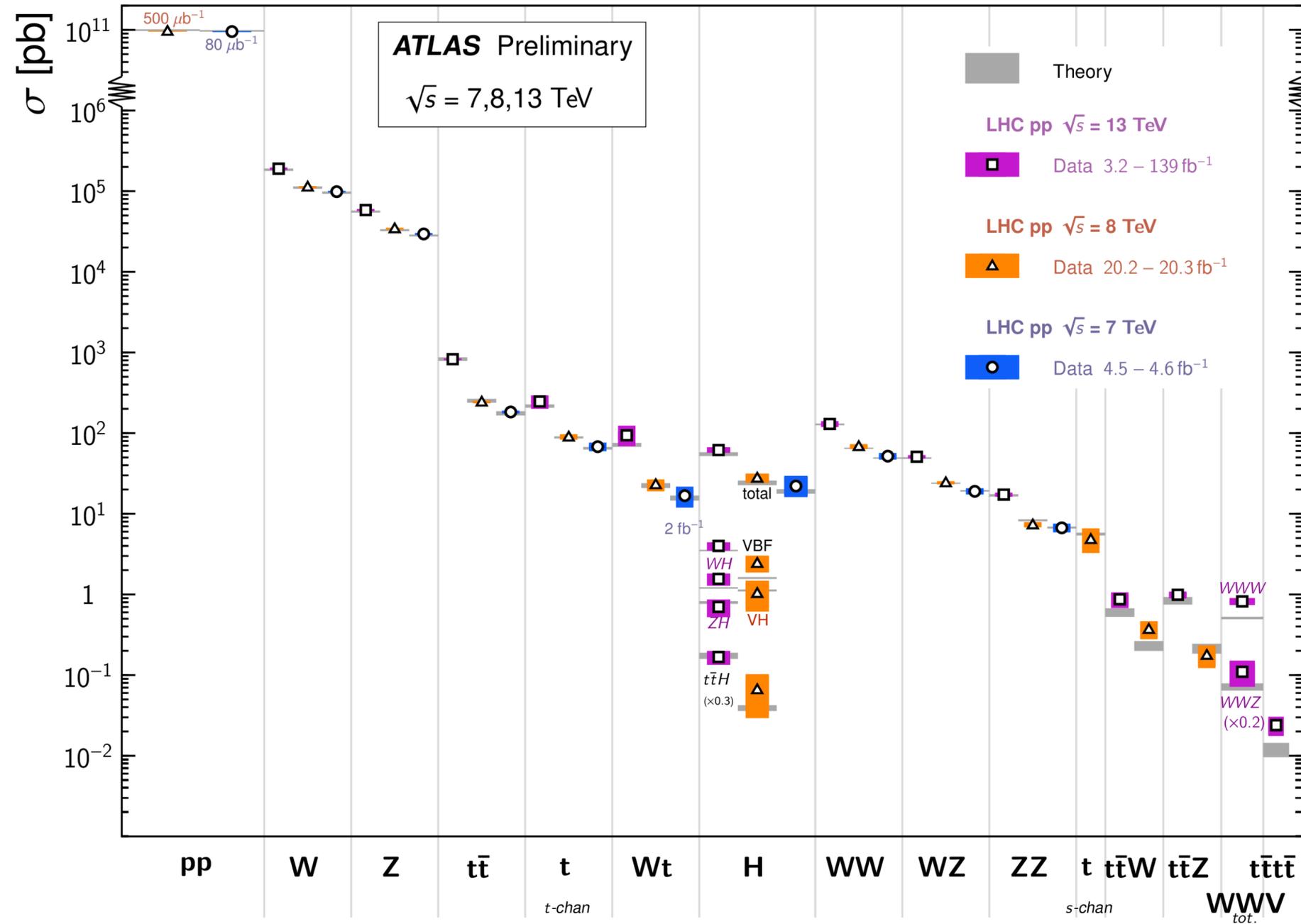
Some Observable Quantity [GeV]

⇒ Determine whether a hypothetical particle is compatible with data!

Experiment *VS* Theory

Standard Model Total Production Cross Section Measurements

Status: February 2022



Cross Section [pb]

~ Frequency of
type of event

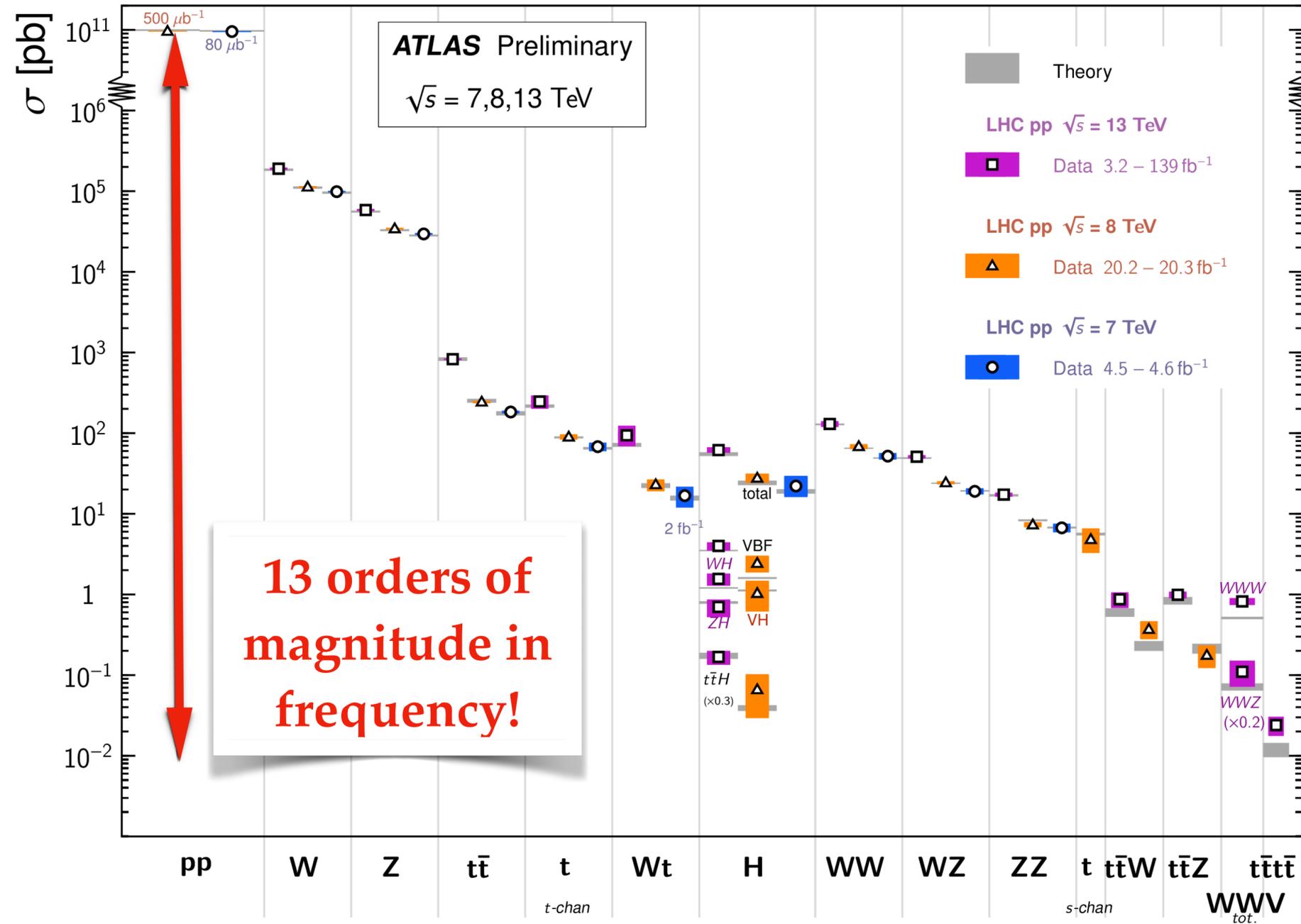
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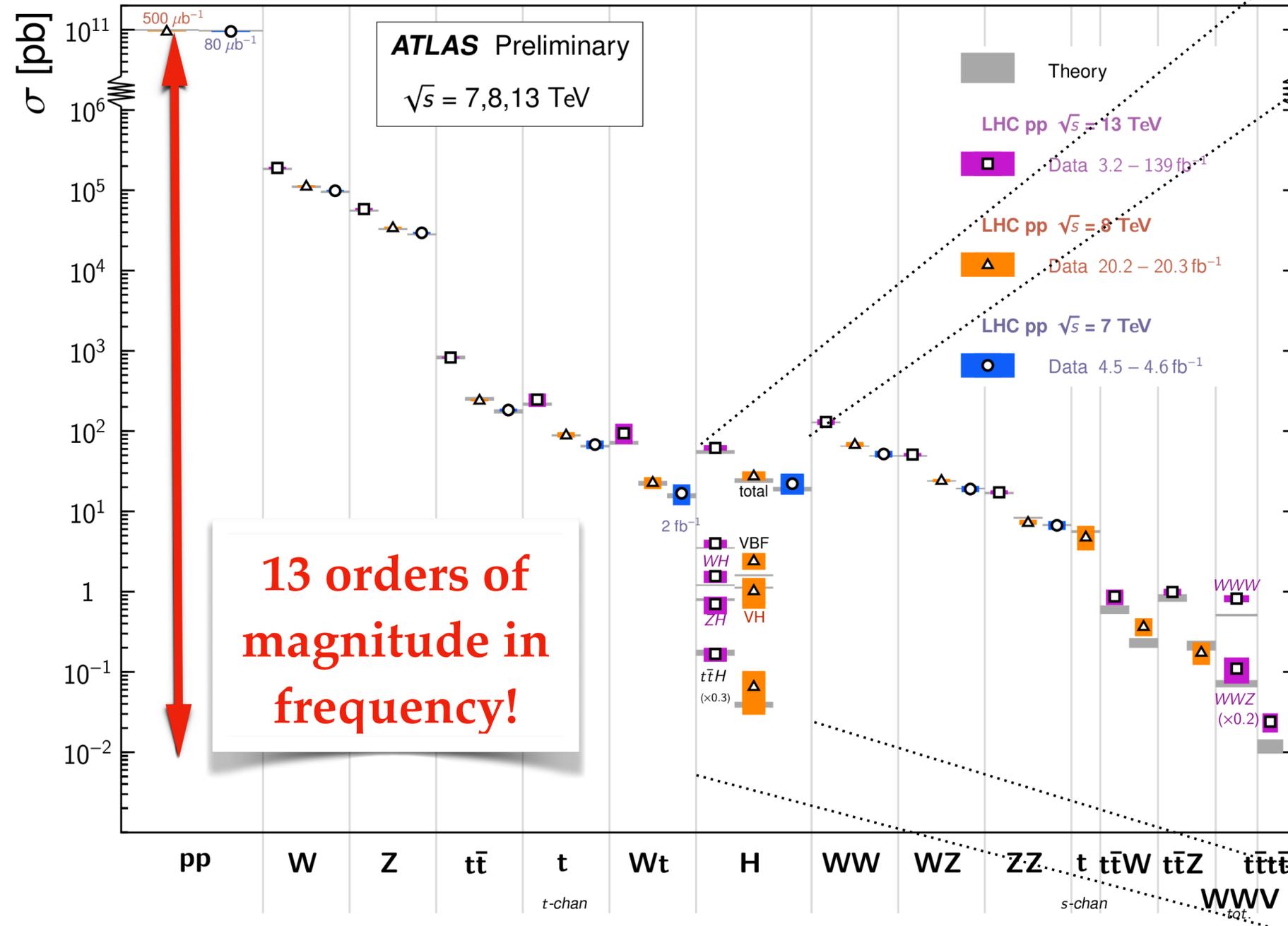
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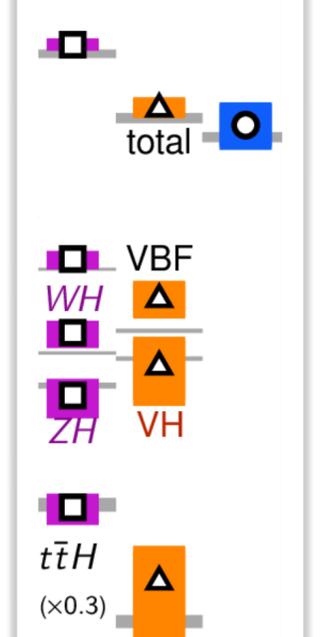
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Higgs Boson Measurements



H

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Experiment VS New Phenomena

ATLAS Heavy Particle Searches* - 95% CL Upper Exclusion Limits

Status: March 2023

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.6 - 139) \text{ fb}^{-1}$$

$$\sqrt{s} = 13 \text{ TeV}$$

Model	ℓ, γ	Jets [†]	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Limit	Reference	
Extra dimen.	ADD $G_{KK} + g/q$	$0 e, \mu, \tau, \gamma$	$1-4 j$	Yes	139	M_D 11.2 TeV $n=2$	2102.10874
	ADD non-resonant $\gamma\gamma$	2γ	-	-	36.7	M_S 8.6 TeV $n=3 \text{ HLZ NLO}$	1707.04147
	ADD QBH	-	$2 j$	-	139	M_{th} 9.4 TeV $n=6$	1910.08447
	ADD BH multijet	-	$\geq 3 j$	-	3.6	M_{th} 9.55 TeV $n=6, M_D = 3 \text{ TeV, rot BH}$	1512.02586
	RS1 $G_{KK} \rightarrow \gamma\gamma$	2γ	-	-	139	$G_{KK} \text{ mass}$ 4.5 TeV $k/\overline{M}_{Pl} = 0.1$	2102.13405
	Bulk RS $G_{KK} \rightarrow WW/ZZ$	multi-channel	-	-	36.1	$G_{KK} \text{ mass}$ 2.3 TeV $k/\overline{M}_{Pl} = 1.0$	1808.02380
	Bulk RS $g_{KK} \rightarrow tt$	$1 e, \mu$	$\geq 1 b, \geq 1 J/2 j$	Yes	36.1	$g_{KK} \text{ mass}$ 3.8 TeV $\Gamma/m = 15\%$	1804.10823
2UED / RPP	$1 e, \mu$	$\geq 2 b, \geq 3 j$	Yes	36.1	KK mass 1.8 TeV Tier (1,1), $\mathcal{B}(A^{(1,1)} \rightarrow tt) = 1$	1803.09678	
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	$2 e, \mu$	-	-	139	Z' mass 5.1 TeV	1903.06248
	SSM $Z' \rightarrow \tau\tau$	2τ	-	-	36.1	Z' mass 2.42 TeV	1709.07242
	Leptophobic $Z' \rightarrow bb$	-	$2 b$	-	36.1	Z' mass 2.1 TeV	1805.09299
	Leptophobic $Z' \rightarrow tt$	$0 e, \mu$	$\geq 1 b, \geq 2 J$	Yes	139	Z' mass 4.1 TeV $\Gamma/m = 1.2\%$	2005.05138
	SSM $W' \rightarrow \ell\nu$	$1 e, \mu$	-	Yes	139	W' mass 6.0 TeV	1906.05609
	SSM $W' \rightarrow \tau\nu$	1τ	-	Yes	139	W' mass 5.0 TeV	ATLAS-CONF-2021-025
	SSM $W' \rightarrow tb$	-	$\geq 1 b, \geq 1 J$	-	139	W' mass 4.4 TeV	ATLAS-CONF-2021-043
	HVT $W' \rightarrow WZ$ model B	$0-2 e, \mu$	$2 j / 1 J$	Yes	139	W' mass 4.3 TeV	2004.14636
	HVT $W' \rightarrow WZ \rightarrow \ell\nu \ell' \ell'$ model C	$3 e, \mu$	$2 j$ (VBF)	Yes	139	W' mass 340 GeV $g_V = 3, g_V c_H = 1, g_r = 0$	2207.03925
	HVT $Z' \rightarrow WW$ model B	$1 e, \mu$	$2 j / 1 J$	Yes	139	Z' mass 3.9 TeV $g_V = 3$	2004.14636
LRSM $W_R \rightarrow \mu N_R$	2μ	$1 J$	-	80	W_R mass 5.0 TeV $m(N_R) = 0.5 \text{ TeV, } g_L = g_R$	1904.12679	
CI	CI $qqqq$	-	$2 j$	-	37.0	Λ 21.8 TeV η_{LL}^-	1703.09127
	CI $\ell\ell qq$	$2 e, \mu$	-	-	139	Λ 35.8 TeV η_{LL}^-	2006.12946
	CI $e e b s$	$2 e$	$1 b$	-	139	Λ 1.8 TeV $g_s = 1$	2105.13847
	CI $\mu\mu b s$	2μ	$1 b$	-	139	Λ 2.0 TeV $g_s = 1$	2105.13847
	CI $tttt$	$\geq 1 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	Λ 2.57 TeV $ C_{4t} = 4\pi$	1811.02305
DM	Axial-vector med. (Dirac DM)	-	$2 j$	-	139	m_{med} 3.8 TeV $g_q = 0.25, g_\tau = 1, m(\chi) = 10 \text{ TeV}$	ATL-PHYS-PUB-2022-036
	Pseudo-scalar med. (Dirac DM)	$0 e, \mu, \tau, \gamma$	$1-4 j$	Yes	139	m_{med} 376 GeV $g_q = 1, g_\tau = 1, m(\chi) = 1 \text{ GeV}$	2102.10874
	Vector med. Z' -2HDM (Dirac DM)	$0 e, \mu$	$2 b$	Yes	139	$m_{Z'}$ 3.0 TeV $\tan\beta = 1, g_Z = 0.8, m(\chi) = 100 \text{ GeV}$	2108.13391
	Pseudo-scalar med. 2HDM+a	multi-channel	-	-	139	m_a 800 GeV $\tan\beta = 1, g_\tau = 1, m(\chi) = 10 \text{ GeV}$	ATLAS-CONF-2021-036
LQ	Scalar LQ 1 st gen	$2 e$	$\geq 2 j$	Yes	139	LQ mass 1.8 TeV $\beta = 1$	2006.05872
	Scalar LQ 2 nd gen	2μ	$\geq 2 j$	Yes	139	LQ mass 1.7 TeV $\beta = 1$	2006.05872
	Scalar LQ 3 rd gen	1τ	$2 b$	Yes	139	LQ_3^u mass 1.49 TeV $\mathcal{B}(LQ_3^u \rightarrow b\tau) = 1$	2303.01294
	Scalar LQ 3 rd gen	$0 e, \mu$	$\geq 2 j, \geq 2 b$	Yes	139	LQ_3^d mass 1.24 TeV $\mathcal{B}(LQ_3^d \rightarrow \nu\tau) = 1$	2004.14060
	Scalar LQ 3 rd gen	$\geq 2 e, \mu, \geq 1 \tau, \geq 1 j, \geq 1 b$	-	-	139	LQ_3^d mass 1.43 TeV $\mathcal{B}(LQ_3^d \rightarrow \tau\nu) = 1$	2101.11582
	Scalar LQ 3 rd gen	$0 e, \mu, \geq 1 \tau, 0-2 j, 2 b$	Yes	139	LQ_3^d mass 1.26 TeV $\mathcal{B}(LQ_3^d \rightarrow b\nu) = 1$	2101.12527	
	Vector LQ mix gen	multi-channel	$\geq 1 j, \geq 1 b$	Yes	139	LQ_3^u mass 2.0 TeV $\mathcal{B}(U_1 \rightarrow t\mu) = 1, \text{Y-M coupl.}$	ATLAS-CONF-2022-052
	Vector LQ 3 rd gen	$2 e, \mu, \tau$	$\geq 1 b$	Yes	139	LQ_3^d mass 1.96 TeV $\mathcal{B}(LQ_3^d \rightarrow b\tau) = 1, \text{Y-M coupl.}$	2303.01294
Vector-like fermions	VLQ $TT \rightarrow Zt + X$	$2e/2\mu \geq 3e, \mu$	$\geq 1 b, \geq 1 j$	-	139	T mass 1.46 TeV SU(2) doublet	2210.15413
	VLQ $BB \rightarrow Wt/Zb + X$	multi-channel	-	-	36.1	B mass 1.34 TeV SU(2) doublet	1808.02343
	VLQ $T_{5/3} T_{5/3} / T_{5/3} \rightarrow Wt + X$	$2(SS) \geq 3 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	$T_{5/3}$ mass 1.64 TeV $\mathcal{B}(T_{5/3} \rightarrow Wt) = 1, c(T_{5/3} Wt) = 1$	1807.11883
	VLQ $T \rightarrow Ht/Zt$	$1 e, \mu$	$\geq 1 b, \geq 3 j$	Yes	139	T mass 1.8 TeV SU(2) singlet, $\kappa_T = 0.5$	ATLAS-CONF-2021-040
	VLQ $Y \rightarrow Wb$	$1 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	Y mass 1.85 TeV $\mathcal{B}(Y \rightarrow Wb) = 1, c_R(Wb) = 1$	1812.07343
	VLQ $B \rightarrow Hb$	$0 e, \mu$	$\geq 2b, \geq 1j, \geq 1J$	-	139	B mass 2.0 TeV SU(2) doublet, $\kappa_B = 0.3$	ATLAS-CONF-2021-018
VLL $\tau' \rightarrow Z\tau/H\tau$	multi-channel	$\geq 1 j$	Yes	139	τ' mass 898 GeV SU(2) doublet	2303.05441	
Excited ferm.	Excited quark $q^* \rightarrow qg$	-	$2 j$	-	139	q^* mass 6.7 TeV only u^* and $d^*, \Lambda = m(q^*)$	1910.08447
	Excited quark $q^* \rightarrow q\gamma$	1γ	$1 j$	-	36.7	q^* mass 5.3 TeV only u^* and $d^*, \Lambda = m(q^*)$	1709.10440
	Excited quark $b^* \rightarrow bg$	-	$1 b, 1 j$	-	139	b^* mass 3.2 TeV	1910.08447
	Excited lepton τ^*	2τ	$\geq 2 j$	-	139	τ^* mass 4.6 TeV $\Lambda = 4.6 \text{ TeV}$	2303.09444
Other	Type III Seesaw	$2,3,4 e, \mu$	$\geq 2 j$	Yes	139	N^0 mass 910 GeV	2202.02039
	LRSM Majorana ν	2μ	$2 j$	-	36.1	N_R mass 3.2 TeV $m(W_R) = 4.1 \text{ TeV, } g_L = g_R$	1809.11105
	Higgs triplet $H^{\pm\pm} \rightarrow W^\pm W^\pm$	$2,3,4 e, \mu$ (SS)	various	Yes	139	$H^{\pm\pm}$ mass 350 GeV DY production	2101.11961
	Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$	$2,3,4 e, \mu$ (SS)	-	-	139	$H^{\pm\pm}$ mass 1.08 TeV DY production	2211.07505
	Multi-charged particles	-	-	-	139	multi-charged particle mass 1.59 TeV DY production, $ q = 5e$	ATLAS-CONF-2022-034
	Magnetic monopoles	-	-	-	34.4	monopole mass 2.37 TeV DY production, $ g = 1g_D, \text{spin } 1/2$	1905.10130

$\sqrt{s} = 13 \text{ TeV}$
partial data

$\sqrt{s} = 13 \text{ TeV}$
full data

10^{-1}

1

10

Mass scale [TeV]



*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).



Experiment VS New Phenomena

ATLAS Heavy Particle Searches* - 95% CL Upper Exclusion Limits

Status: March 2023

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.6 - 139) \text{ fb}^{-1}$$

$$\sqrt{s} = 13 \text{ TeV}$$

Model	ℓ, γ	Jets [†]	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Limit	Reference	
Extra dimen.	ADD $G_{KK} + g/q$	$0 e, \mu, \tau, \gamma$	$1-4 j$	Yes	139	M_D 11.2 TeV $n=2$	2102.10874
	ADD non-resonant $\gamma\gamma$	2γ	-	-	36.7	M_S 8.6 TeV $n=3 \text{ HLZ NLO}$	1707.04147
	ADD QBH	-	$2 j$	-	139	M_{th} 9.4 TeV $n=6$	1910.08447
	ADD BH multijet	-	$\geq 3 j$	-	3.6	M_{th} 9.55 TeV $n=6, M_D = 3 \text{ TeV, rot BH}$	1512.02586
	RS1 $G_{KK} \rightarrow \gamma\gamma$	2γ	-	-	139	$G_{KK} \text{ mass}$ 4.5 TeV $k/\overline{M}_{Pl} = 0.1$	2102.13405
	Bulk RS $G_{KK} \rightarrow WW/ZZ$	multi-channel	-	-	36.1	$G_{KK} \text{ mass}$ 2.3 TeV $k/\overline{M}_{Pl} = 1.0$	1808.02380
	Bulk RS $g_{KK} \rightarrow tt$	$1 e, \mu$	$\geq 1 b, \geq 1 J/2 j$	Yes	36.1	$g_{KK} \text{ mass}$ 3.8 TeV $\Gamma/m = 15\%$	1804.10823
2UED / RPP	$1 e, \mu$	$\geq 2 b, \geq 3 j$	Yes	36.1	KK mass 1.8 TeV Tier (1,1), $\mathcal{B}(A^{(1,1)} \rightarrow tt) = 1$	1803.09678	
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	$2 e, \mu$	-	-	139	Z' mass 5.1 TeV	1903.06248
	SSM $Z' \rightarrow \tau\tau$	2τ	-	-	36.1	Z' mass 2.42 TeV	1709.07242
	Leptophobic $Z' \rightarrow bb$	-	$2 b$	-	36.1	Z' mass 2.1 TeV	1805.09299
	Leptophobic $Z' \rightarrow tt$	$0 e, \mu$	$\geq 1 b, \geq 2 J$	Yes	139	Z' mass 4.1 TeV $\Gamma/m = 1.2\%$	2005.05138
	SSM $W' \rightarrow \ell\nu$	$1 e, \mu$	-	Yes	139	W' mass 6.0 TeV	1906.05609
	SSM $W' \rightarrow \tau\nu$	1τ	-	Yes	139	W' mass 5.0 TeV	ATLAS-CONF-2021-025
	SSM $W' \rightarrow tb$	-	$\geq 1 b, \geq 1 J$	-	139	W' mass 4.4 TeV	ATLAS-CONF-2021-043
	HVT $W' \rightarrow WZ$ model B	$0-2 e, \mu$	$2 j / 1 J$	Yes	139	W' mass 4.3 TeV	2004.14636
	HVT $W' \rightarrow WZ \rightarrow \ell\nu \ell' \ell'$ model C	$3 e, \mu$	$2 j$ (VBF)	Yes	139	W' mass 340 GeV $g_V = 3, g_V^{cH} = 1, g_r = 0$	2207.03925
	HVT $Z' \rightarrow WW$ model B	$1 e, \mu$	$2 j / 1 J$	Yes	139	Z' mass 3.9 TeV $g_V = 3$	2004.14636
LRSB $W_R \rightarrow \mu N_R$	2μ	$1 J$	-	80	W_R mass 5.0 TeV $m(N_R) = 0.5 \text{ TeV, } g_L = g_R$	1904.12679	
CI	CI $qqqq$	-	$2 j$	-	37.0	Λ 21.8 TeV η_{LL}	1703.09127
	CI $\ell\ell qq$	$2 e, \mu$	-	-	139	Λ 35.8 TeV η_{LL}	2006.12946
	CI $eebs$	$2 e$	$1 b$	-	139	Λ 1.8 TeV $g_s = 1$	2105.13847
	CI $\mu\mu bs$	2μ	$1 b$	-	139	Λ 2.0 TeV $g_s = 1$	2105.13847
	CI $tttt$	$\geq 1 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	Λ 2.57 TeV $ C_{4t} = 4\pi$	1811.02305
DM	Axial-vector med. (Dirac DM)	-	$2 j$	-	139	m_{med} 3.8 TeV $g_a = 0.25, g_s = 1, m(\chi) = 10 \text{ TeV}$	ATL-PHYS-PUB-2022-036
	Pseudo-scalar med. (Dirac DM)	$0 e, \mu, \tau, \gamma$	$1-4 j$	Yes	139	m_{med} 376 GeV $g_a = 1, g_s = 1, m(\chi) = 1 \text{ GeV}$	2102.10874
	Vector med. Z' -2HDM (Dirac DM)	$0 e, \mu$	$2 b$	Yes	139	$m_{Z'}$ 3.0 TeV $\tan\beta = 1, g_z = 0.8, m(\chi) = 100 \text{ GeV}$	2108.13391
	Pseudo-scalar med. 2HDM+a	multi-channel	-	-	139	m_a 800 GeV $\tan\beta = 1, g_s = 1, m(\chi) = 10 \text{ GeV}$	ATLAS-CONF-2021-036
LQ	Scalar LQ 1 st gen	$2 e$	$\geq 2 j$	Yes	139	LQ mass 1.8 TeV $\beta = 1$	2006.05872
	Scalar LQ 2 nd gen	2μ	$\geq 2 j$	Yes	139	LQ mass 1.7 TeV $\beta = 1$	2006.05872
	Scalar LQ 3 rd gen	1τ	$2 b$	Yes	139	LQ_3^u mass 1.49 TeV $\mathcal{B}(LQ_3^u \rightarrow b\tau) = 1$	2303.01294
	Scalar LQ 3 rd gen	$0 e, \mu$	$\geq 2 j, \geq 2 b$	Yes	139	LQ_3^d mass 1.24 TeV $\mathcal{B}(LQ_3^d \rightarrow \nu\tau) = 1$	2004.14060
	Scalar LQ 3 rd gen	$\geq 2 e, \mu, \geq 1 \tau, \geq 1 j, \geq 1 b$	-	-	139	LQ_3^u mass 1.43 TeV $\mathcal{B}(LQ_3^u \rightarrow t\tau) = 1$	2101.11582
	Scalar LQ 3 rd gen	$0 e, \mu, \geq 1 \tau, 0-2 j, 2 b$	Yes	139	LQ_3^d mass 1.26 TeV $\mathcal{B}(LQ_3^d \rightarrow b\nu) = 1$	2101.12527	
	Vector LQ mix gen	multi-channel	$\geq 1 j, \geq 1 b$	Yes	139	LQ_3^u mass 2.0 TeV $\mathcal{B}(\tilde{U}_1 \rightarrow t\mu) = 1, \text{Y-M coupl.}$	ATLAS-CONF-2022-052
	Vector LQ 3 rd gen	$2 e, \mu, \tau$	$\geq 1 b$	Yes	139	LQ_3^d mass 1.96 TeV $\mathcal{B}(LQ_3^d \rightarrow b\tau) = 1, \text{Y-M coupl.}$	2303.01294
Vector-like fermions	VLQ $TT \rightarrow Zt + X$	$2e/2\mu \geq 3e, \mu$	$\geq 1 b, \geq 1 j$	-	139	T mass 1.46 TeV SU(2) doublet	2210.15413
	VLQ $BB \rightarrow Wt/Zb + X$	multi-channel	-	-	36.1	B mass 1.34 TeV SU(2) doublet	1808.02343
	VLQ $T_{5/3} T_{5/3} / T_{5/3} \rightarrow Wt + X$	$2(SS) \geq 3 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	$T_{5/3}$ mass 1.64 TeV $\mathcal{B}(T_{5/3} \rightarrow Wt) = 1, c(T_{5/3} Wt) = 1$	1807.11883
	VLQ $T \rightarrow Ht/Zt$	$1 e, \mu$	$\geq 1 b, \geq 3 j$	Yes	139	T mass 1.8 TeV SU(2) singlet, $\kappa_T = 0.5$	ATLAS-CONF-2021-040
	VLQ $Y \rightarrow Wb$	$1 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	Y mass 1.85 TeV $\mathcal{B}(Y \rightarrow Wb) = 1, c_R(Wb) = 1$	1812.07343
	VLQ $B \rightarrow Hb$	$0 e, \mu$	$\geq 2b, \geq 1j, \geq 1J$	-	139	B mass 2.0 TeV SU(2) doublet, $\kappa_B = 0.3$	ATLAS-CONF-2021-018
VLL $\tau' \rightarrow Z\tau/H\tau$	multi-channel	$\geq 1 j$	Yes	139	τ' mass 898 GeV SU(2) doublet	2303.05441	
Excited ferm.	Excited quark $q^* \rightarrow qg$	-	$2 j$	-	139	q^* mass 6.7 TeV only u^* and $d^*, \Lambda = m(q^*)$	1910.08447
	Excited quark $q^* \rightarrow q\gamma$	1γ	$1 j$	-	36.7	q^* mass 5.3 TeV only u^* and $d^*, \Lambda = m(q^*)$	1709.10440
	Excited quark $b^* \rightarrow bg$	-	$1 b, 1 j$	-	139	b^* mass 3.2 TeV	1910.08447
	Excited lepton τ^*	2τ	$\geq 2 j$	-	139	τ^* mass 4.6 TeV $\Lambda = 4.6 \text{ TeV}$	2303.09444
Other	Type III Seesaw	$2,3,4 e, \mu$	$\geq 2 j$	Yes	139	N^0 mass 910 GeV	2202.02039
	LRSB Majorana ν	2μ	$2 j$	-	36.1	N_R mass 3.2 TeV $m(W_R) = 4.1 \text{ TeV, } g_L = g_R$	1809.11105
	Higgs triplet $H^{\pm\pm} \rightarrow W^\pm W^\pm$	$2,3,4 e, \mu$ (SS)	various	Yes	139	$H^{\pm\pm}$ mass 350 GeV DY production	2101.11961
	Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$	$2,3,4 e, \mu$ (SS)	-	-	139	$H^{\pm\pm}$ mass 1.08 TeV DY production	2211.07505
	Multi-charged particles	-	-	-	139	multi-charged particle mass 1.59 TeV DY production, $ q = 5e$	ATLAS-CONF-2022-034
Magnetic monopoles	-	-	-	34.4	monopole mass 2.37 TeV DY production, $ g = 1g_D, \text{spin } 1/2$	1905.10130	

$\sqrt{s} = 13 \text{ TeV}$
partial data

$\sqrt{s} = 13 \text{ TeV}$
full data

10^{-1}

1

10

Mass scale [TeV]



Mass Scale limits: 1-10 TeV

*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).



STANDARD MODEL WINS!



⇒ No new phenomena!?



[credit: Alvaro Rujula,
Found in a CERN TH office]



STANDARD MODEL WINS!



⇒ No new phenomena!?



[credit: Alvaro Rujula, Found in a CERN TH office]





Is that all folks?

Questions



Questions



Answers



Questions



Answers



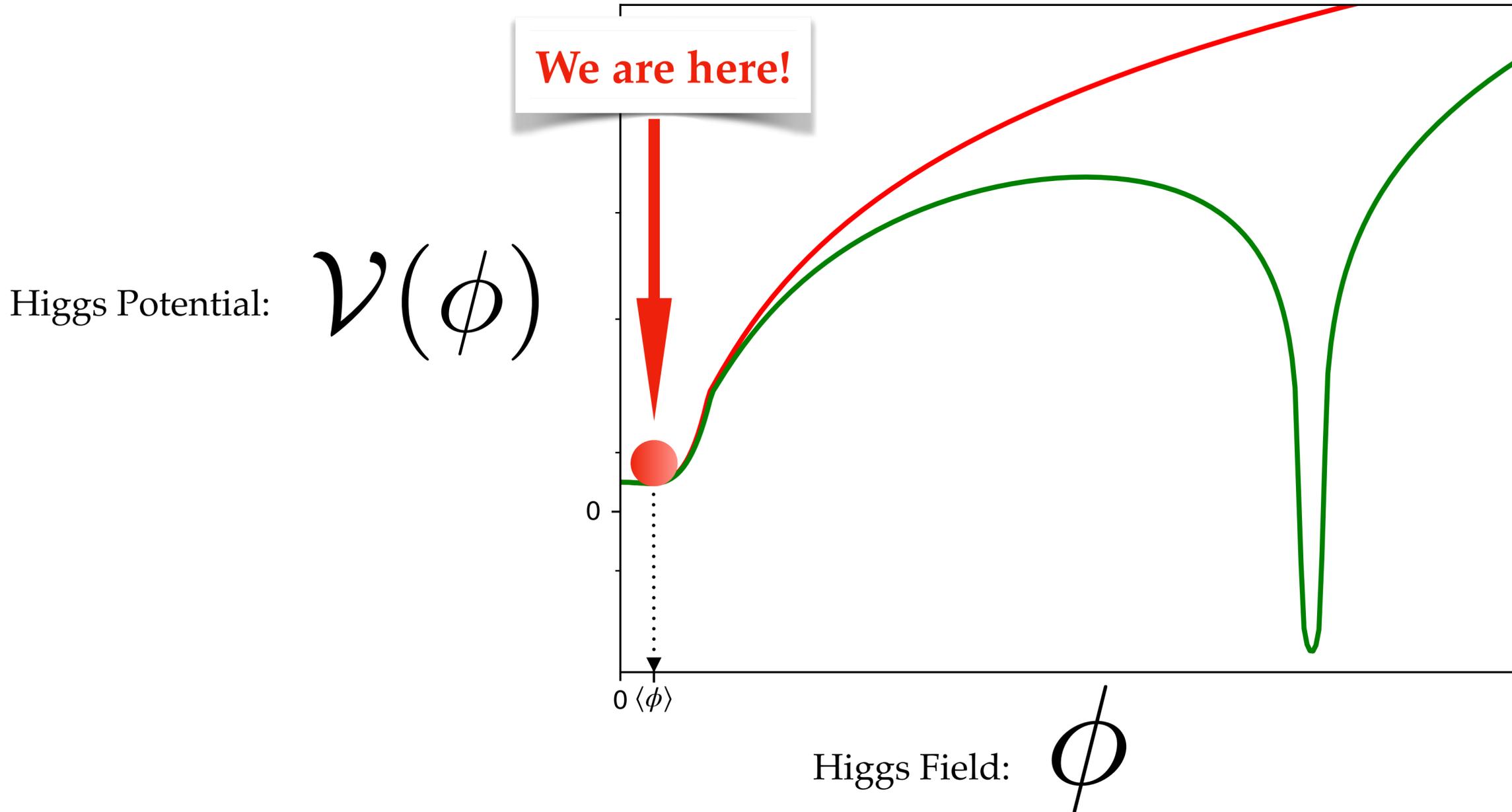
Questions



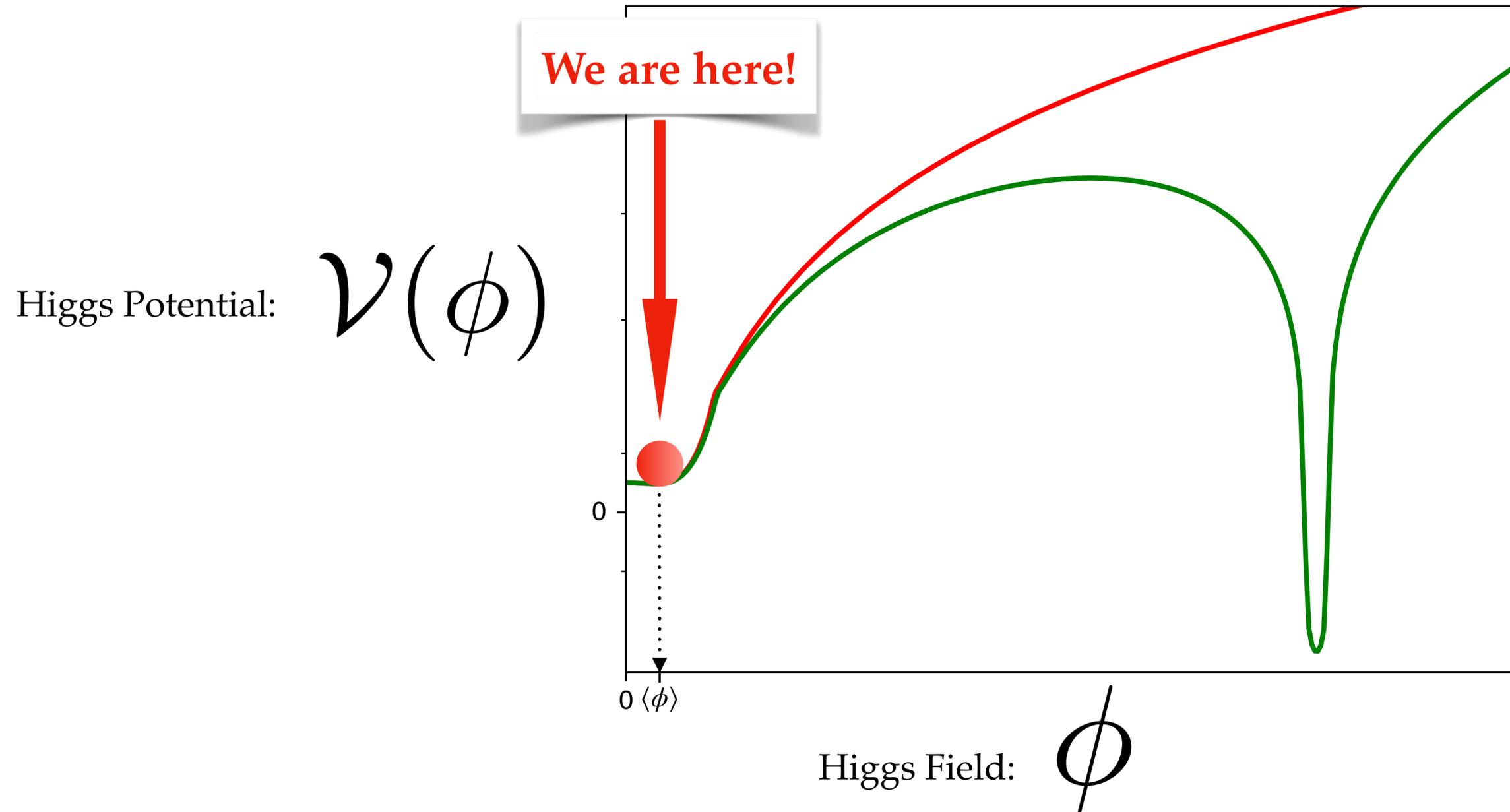
Answers



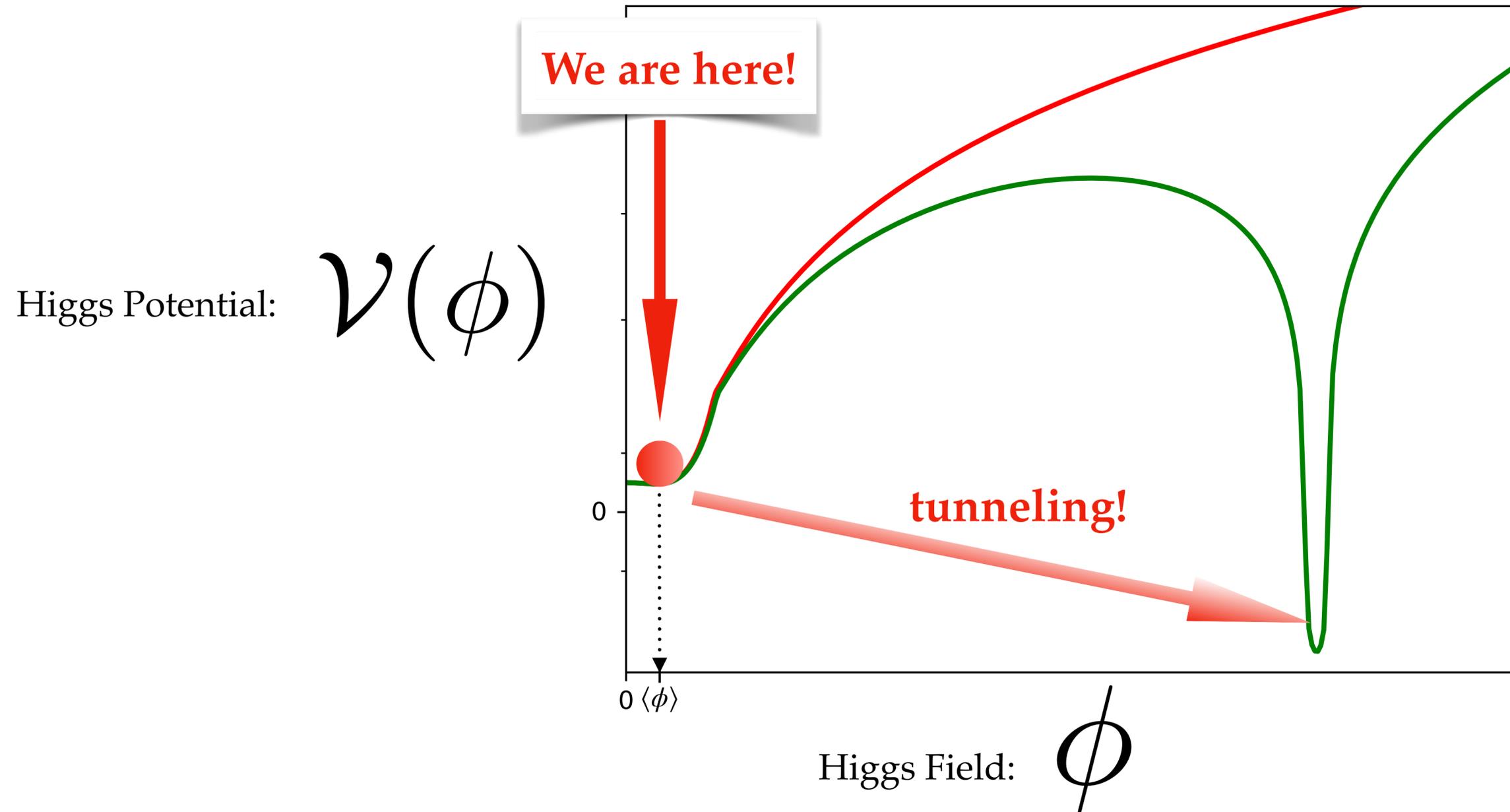
2 Q: What is the ultimate fate of our Universe?



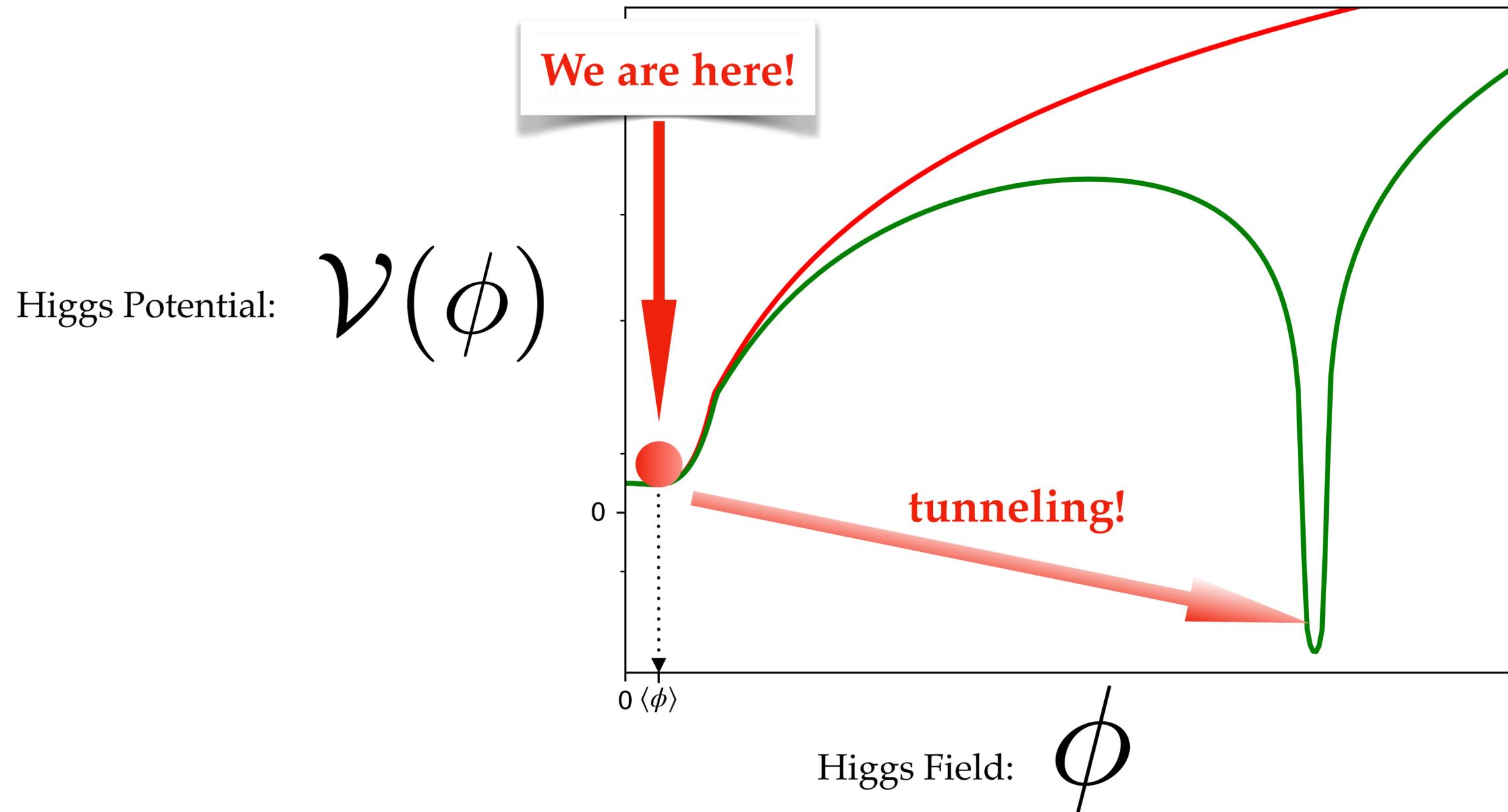
Q: What is the ultimate fate of our Universe?



Q: What is the ultimate fate of our Universe?



Q: What is the ultimate fate of our Universe?



→ the **“Vacuum Stability Problem”**.



Q: Why does the matter we are made of exist?

[or: Why is there so much more matter than anti-matter?]



[credit: QuantaMagazine]



Andreas Papaefstathiou

Q: Why does the matter we are made of exist?

[or: Why is there so much more matter than anti-matter?]



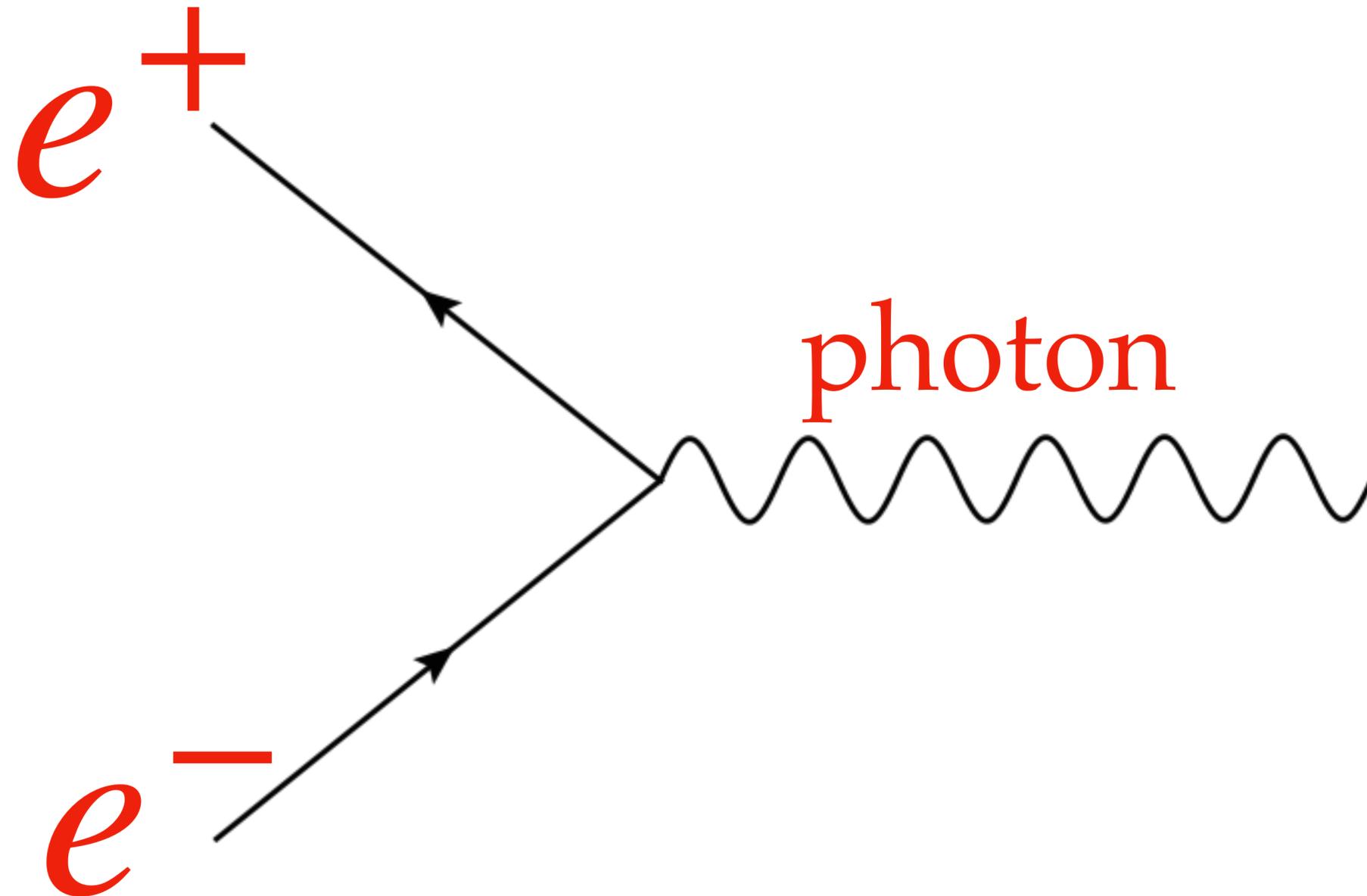
[credit: QuantaMagazine]



Andreas Papaefstathiou

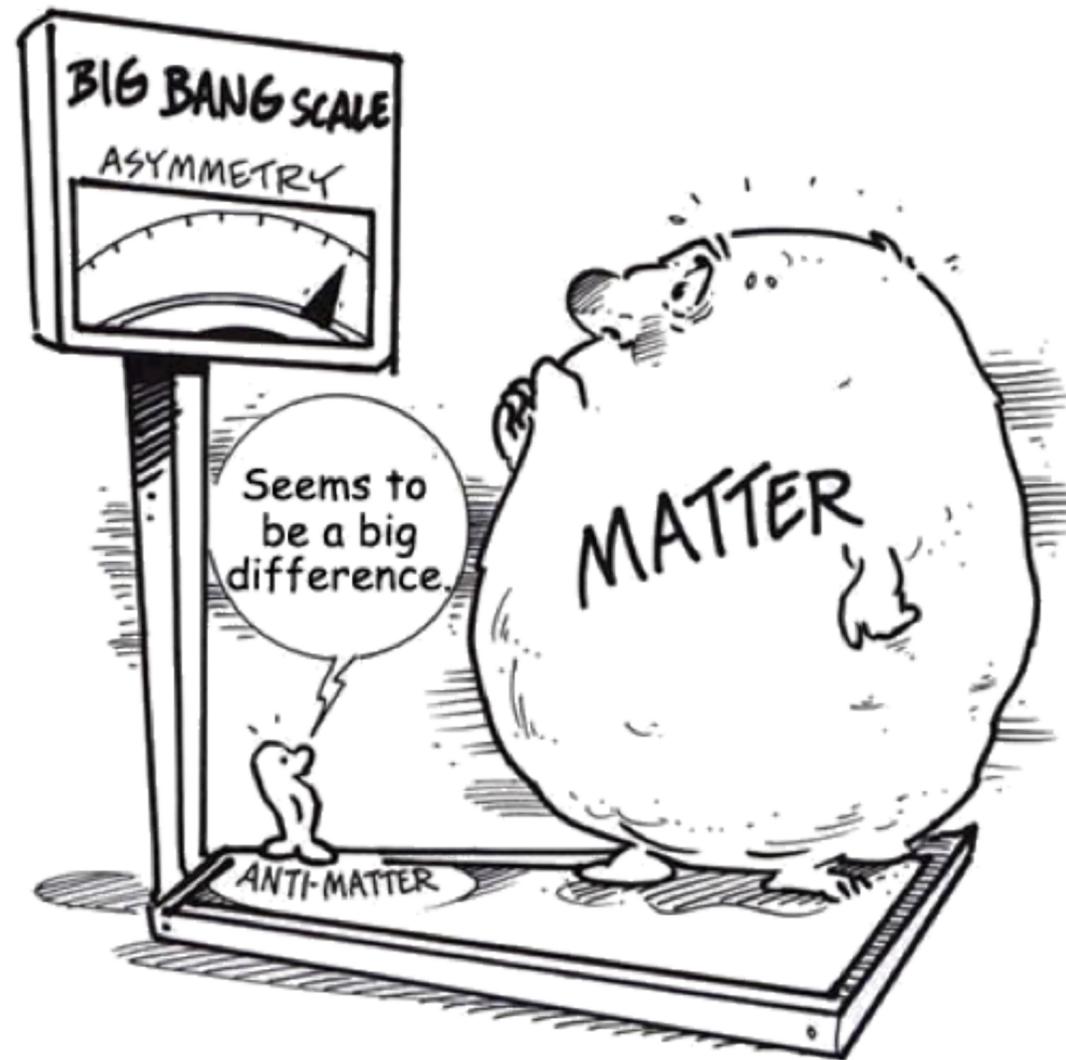
Q: Why does the matter we are made of exist?

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Q: Why does the matter we are made of exist?

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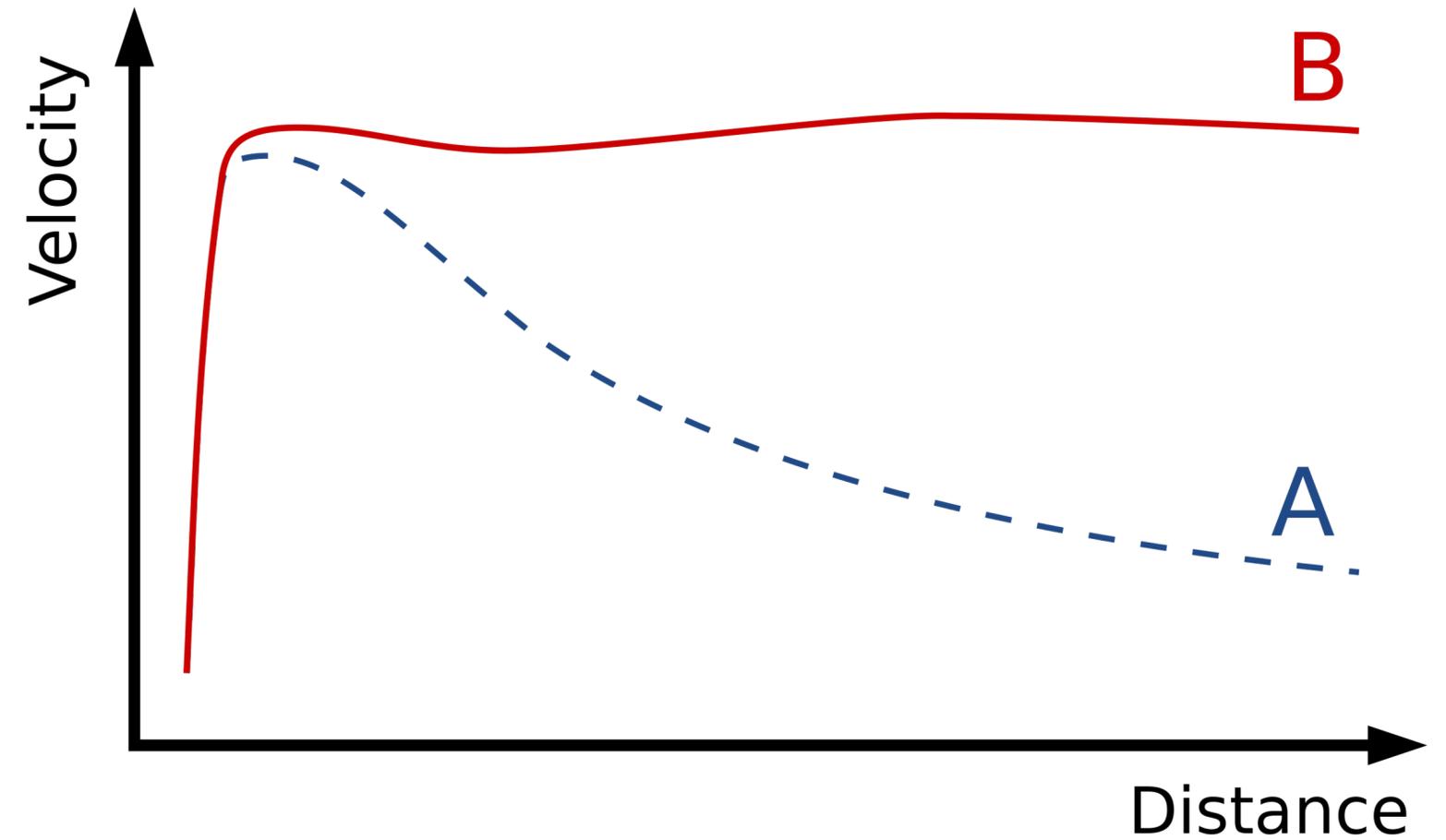
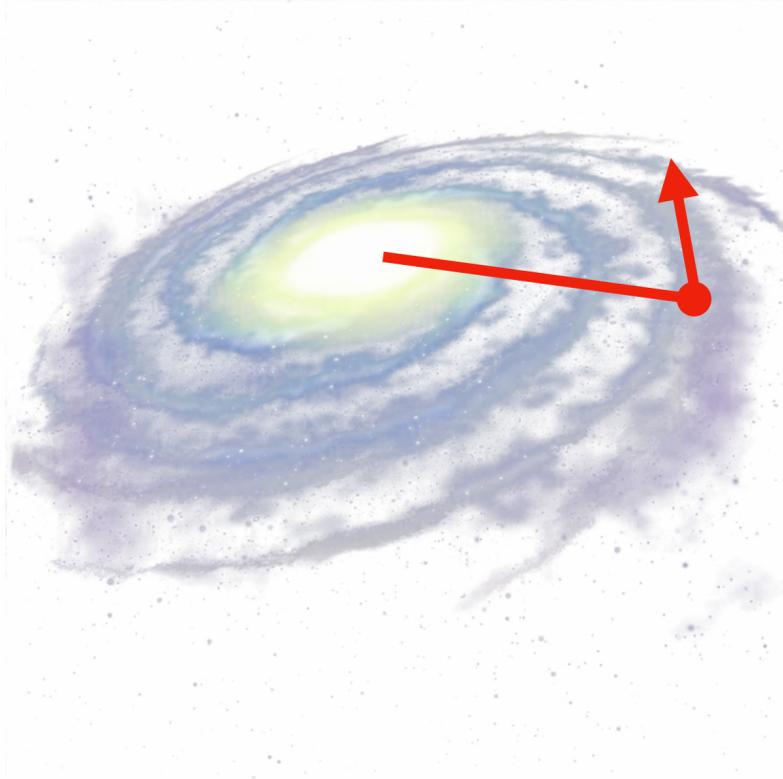


→ the **“Matter-Anti-Matter Asymmetry”**.



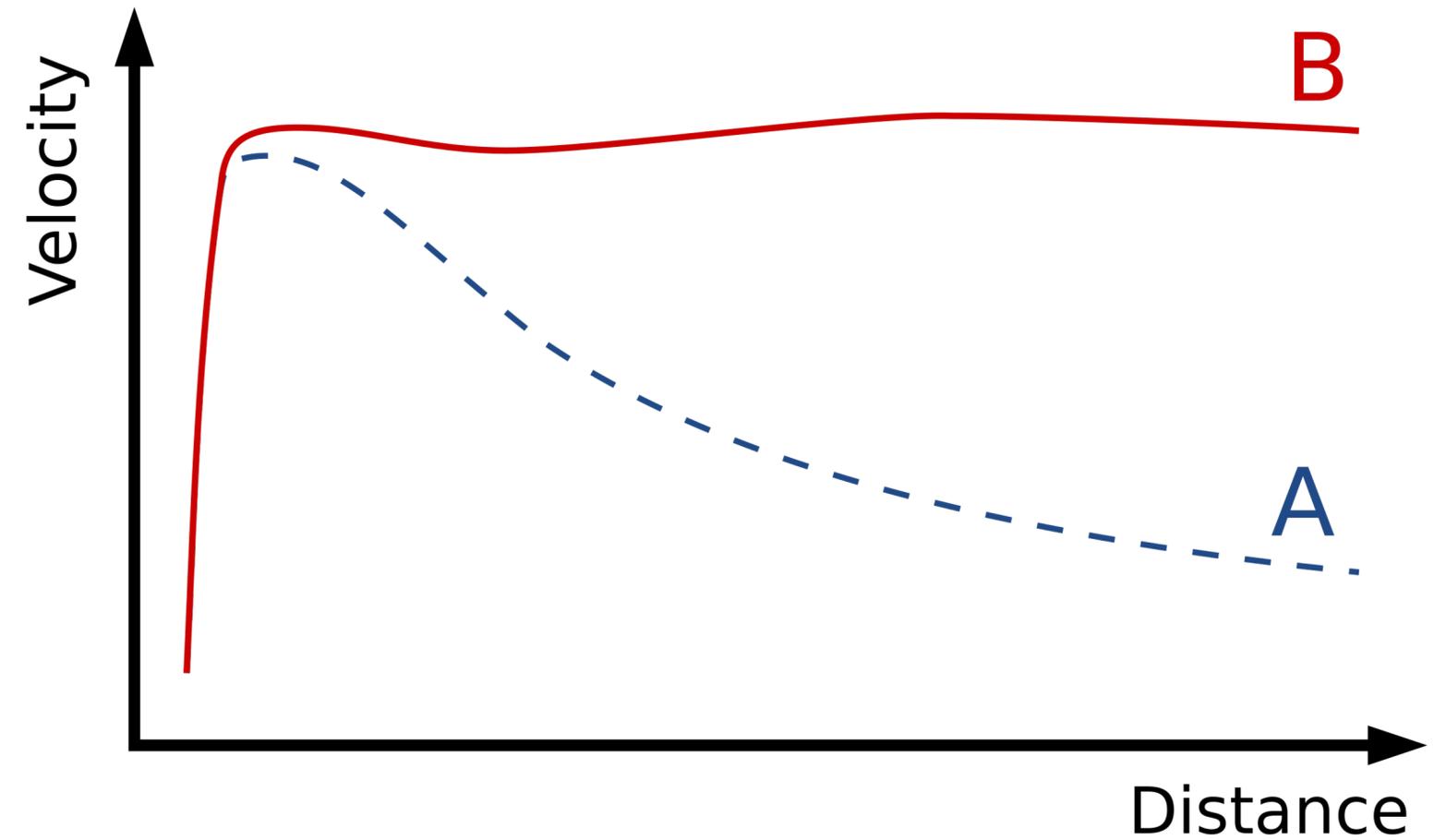
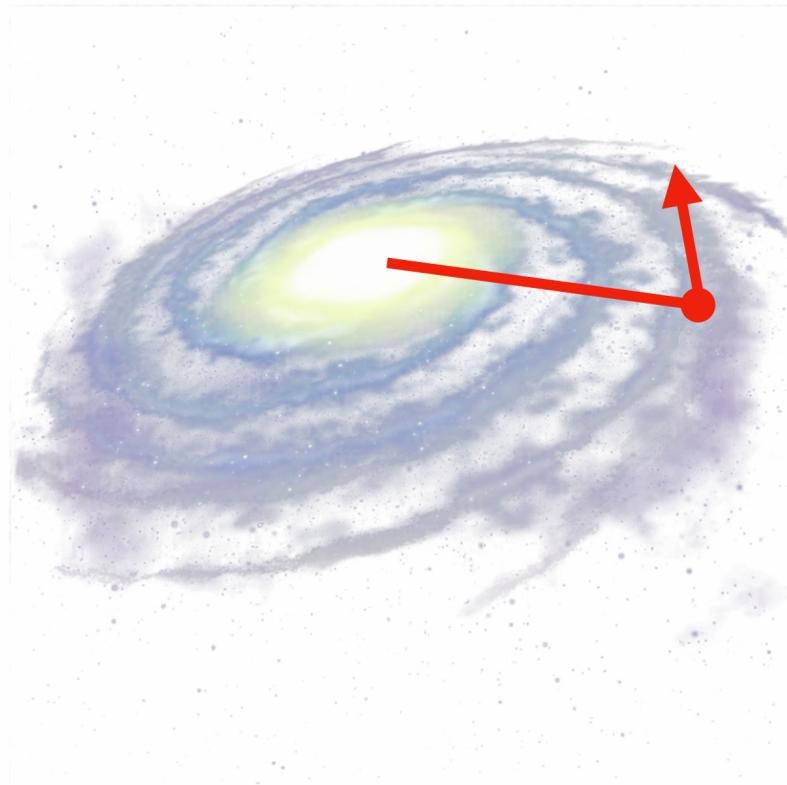
Q: What is the nature of **Dark Matter**?

Rotation curve of a typical spiral galaxy: predicted (A) and observed (B).



Q: What is the nature of **Dark Matter**?

Rotation curve of a typical spiral galaxy: predicted (A) and observed (B).



Q: ...





Dark Matter

Vacuum Stability

**Matter-anti-matter
Asymmetry**



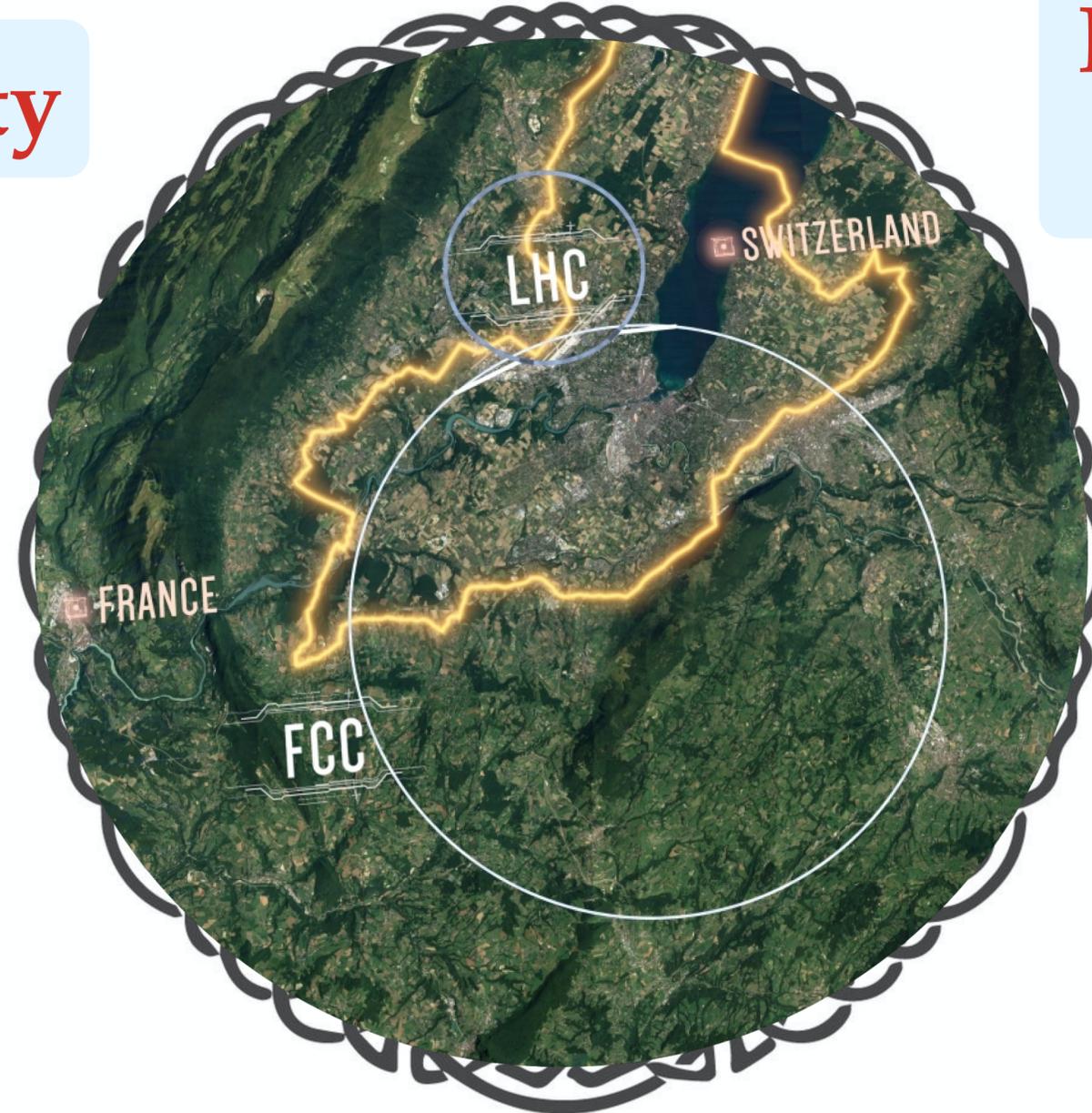
The Higgs Field & Symmetry Breaking



Dark Matter

Vacuum Stability

**Matter-anti-matter
Asymmetry**



e.g. Future Circular Collider:
 $pp@100 \text{ TeV}, e^+e^-$.

e.g. "High-Energy" LHC:
 $pp@27 \text{ TeV}$.

e.g. Muon Collider.

The Higgs Field & Symmetry Breaking



3

Breaking the Symmetry



Breaking the Symmetry

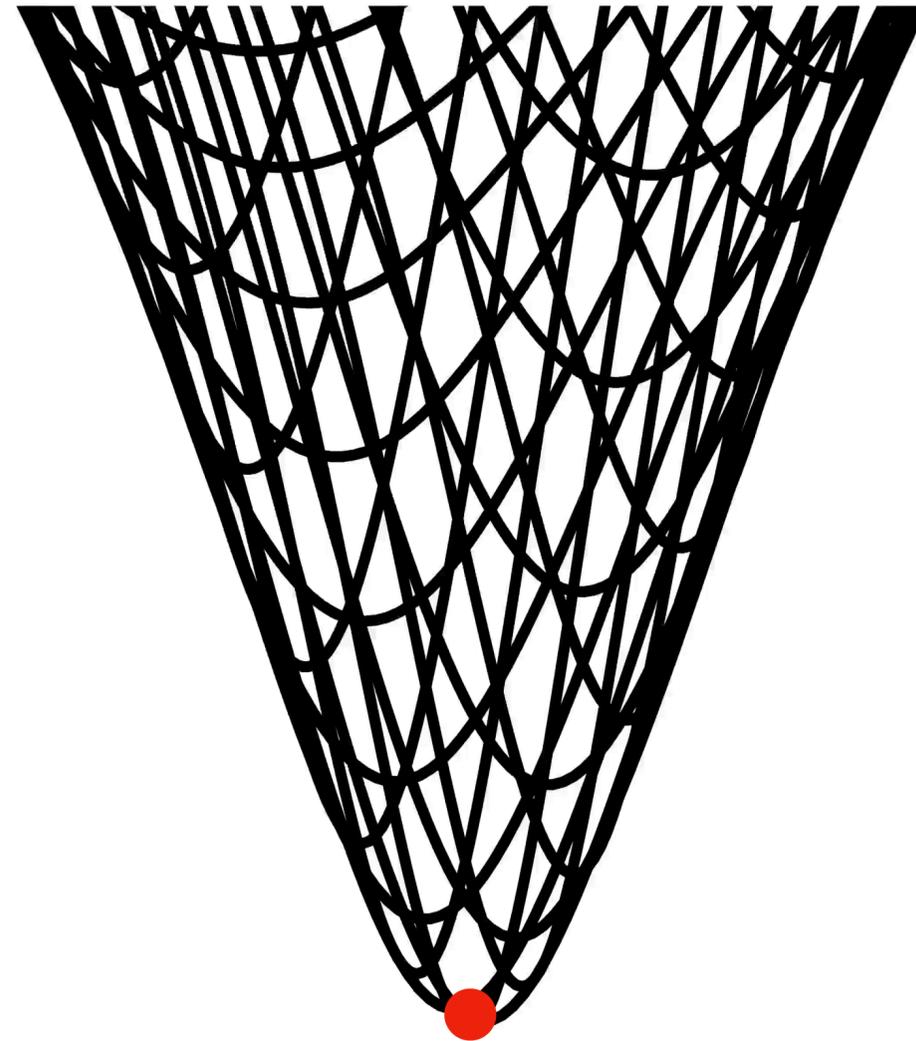


Breaking the Symmetry in the Standard Model

The potential of the Higgs field (ϕ), a complex doublet:

$$\phi = \begin{pmatrix} \phi_1 + i\phi_3 \\ \phi_2 + i\phi_4 \end{pmatrix}$$

(Arbitrarily) Set $\phi_3 = \phi_4 = 0$
to illustrate potential in
 (ϕ_1, ϕ_2) plane.



$\mathcal{V}(\phi)$

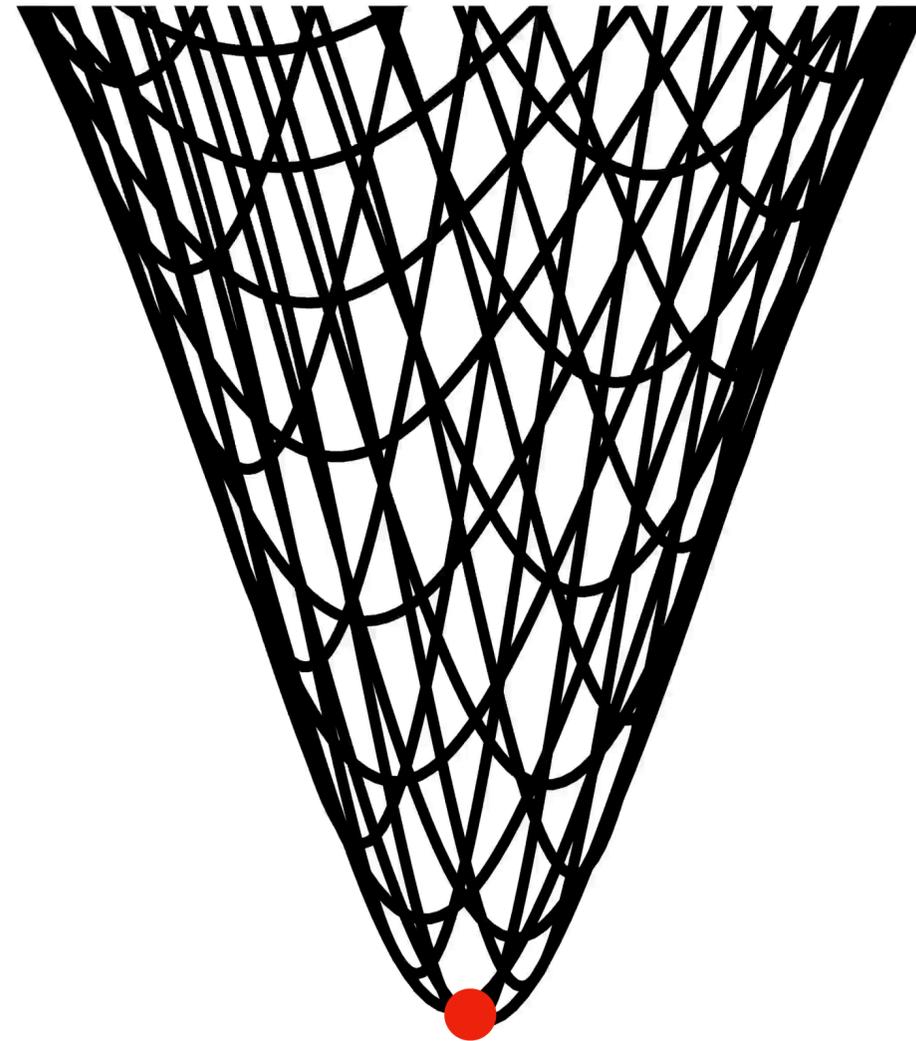


Breaking the Symmetry in the Standard Model

The potential of the Higgs field (ϕ), a complex doublet:

$$\phi = \begin{pmatrix} \phi_1 + i\phi_3 \\ \phi_2 + i\phi_4 \end{pmatrix}$$

(Arbitrarily) Set $\phi_3 = \phi_4 = 0$
to illustrate potential in
 (ϕ_1, ϕ_2) plane.



$\mathcal{V}(\phi)$

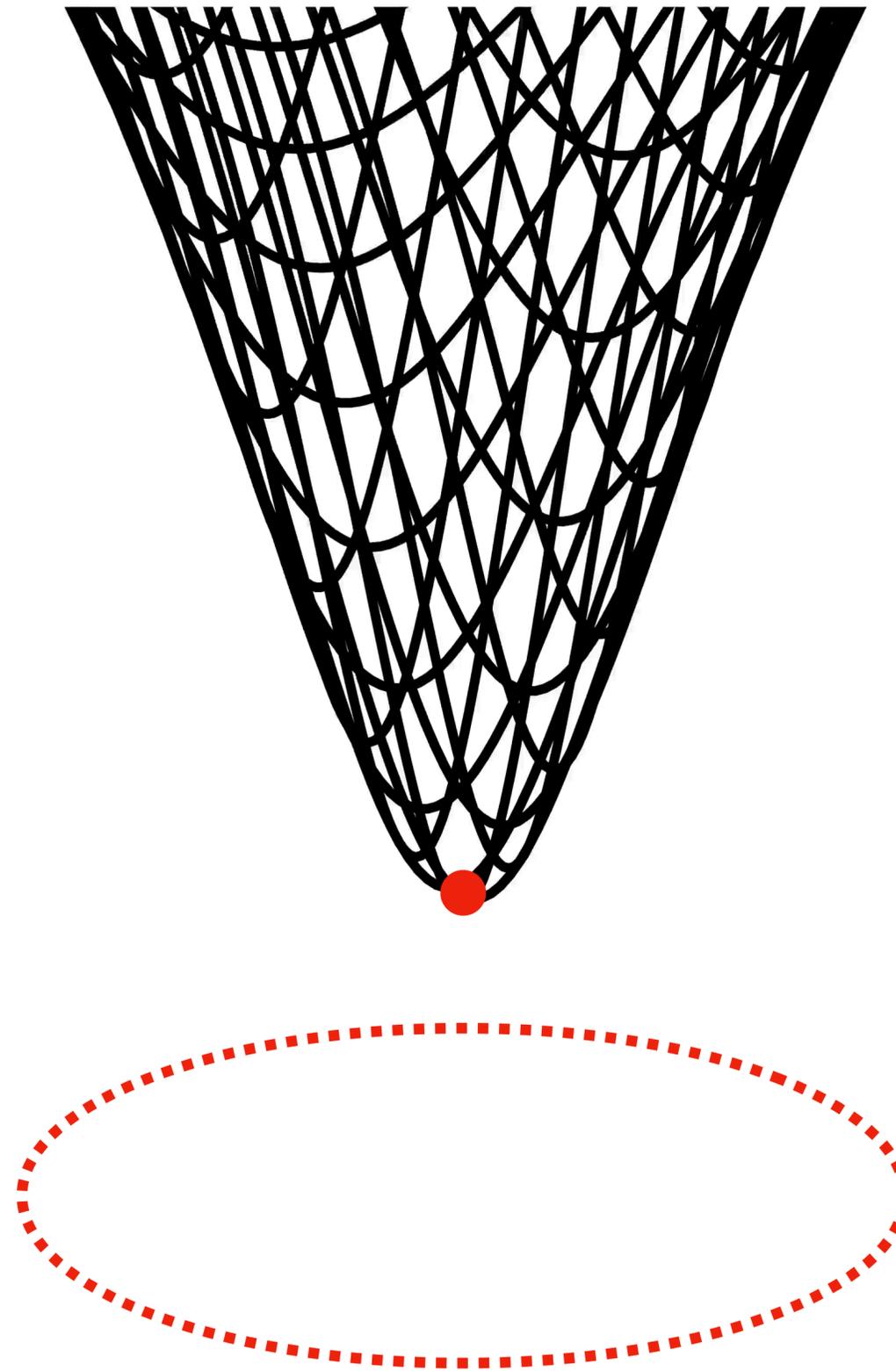


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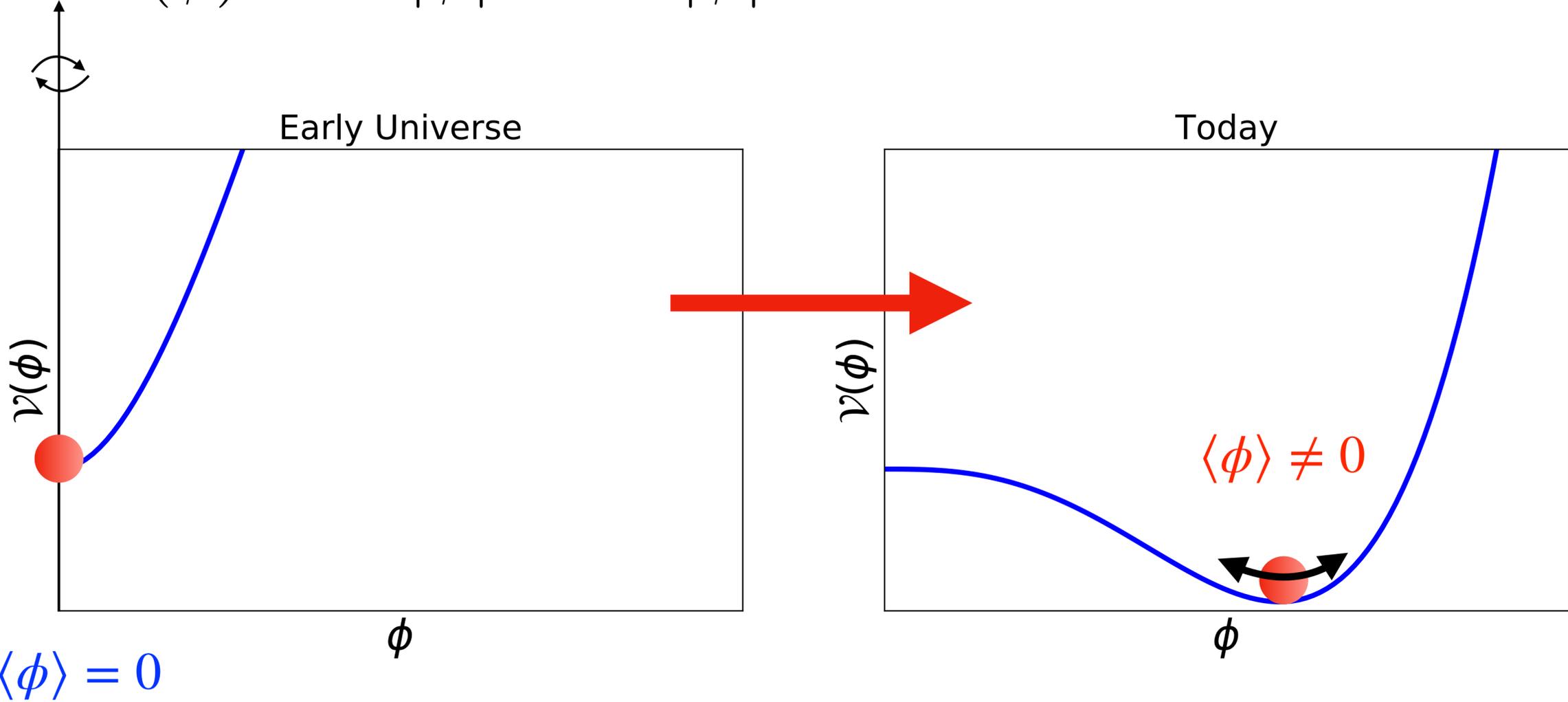
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$$\mathcal{V}(\phi) = \text{●} |\phi|^2 + \text{■} |\phi|^4$$

Higgs field (ϕ) potential



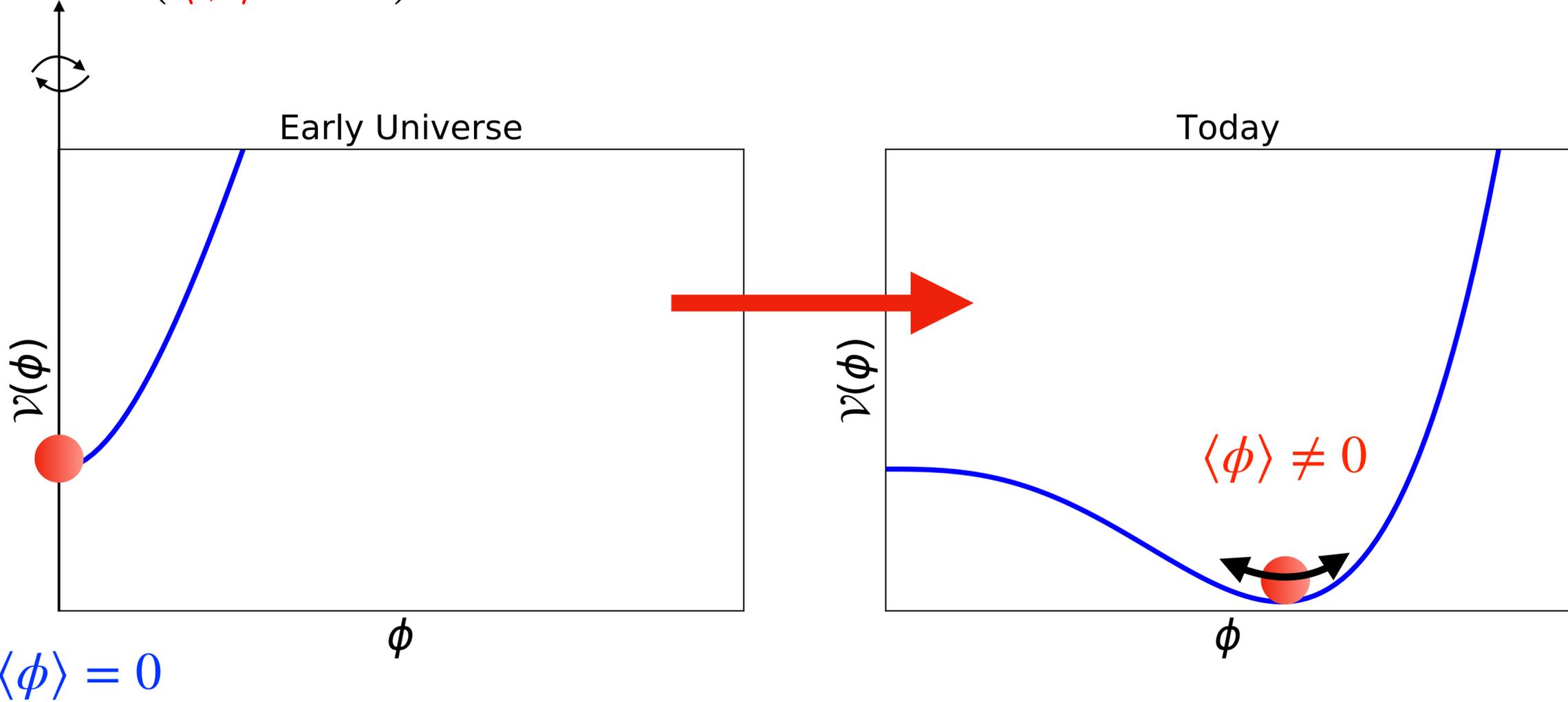
More Symmetry \rightarrow Less (obvious) Symmetry



Breaking the Symmetry in the Standard Model

$$\mathcal{V}(\langle\phi\rangle + h) = \bullet h^2 + \blacktriangle h^3 + \blacksquare h^4 \rightarrow h \text{ is the Higgs boson! (LHC, 2012)}$$

Higgs field (ϕ) potential



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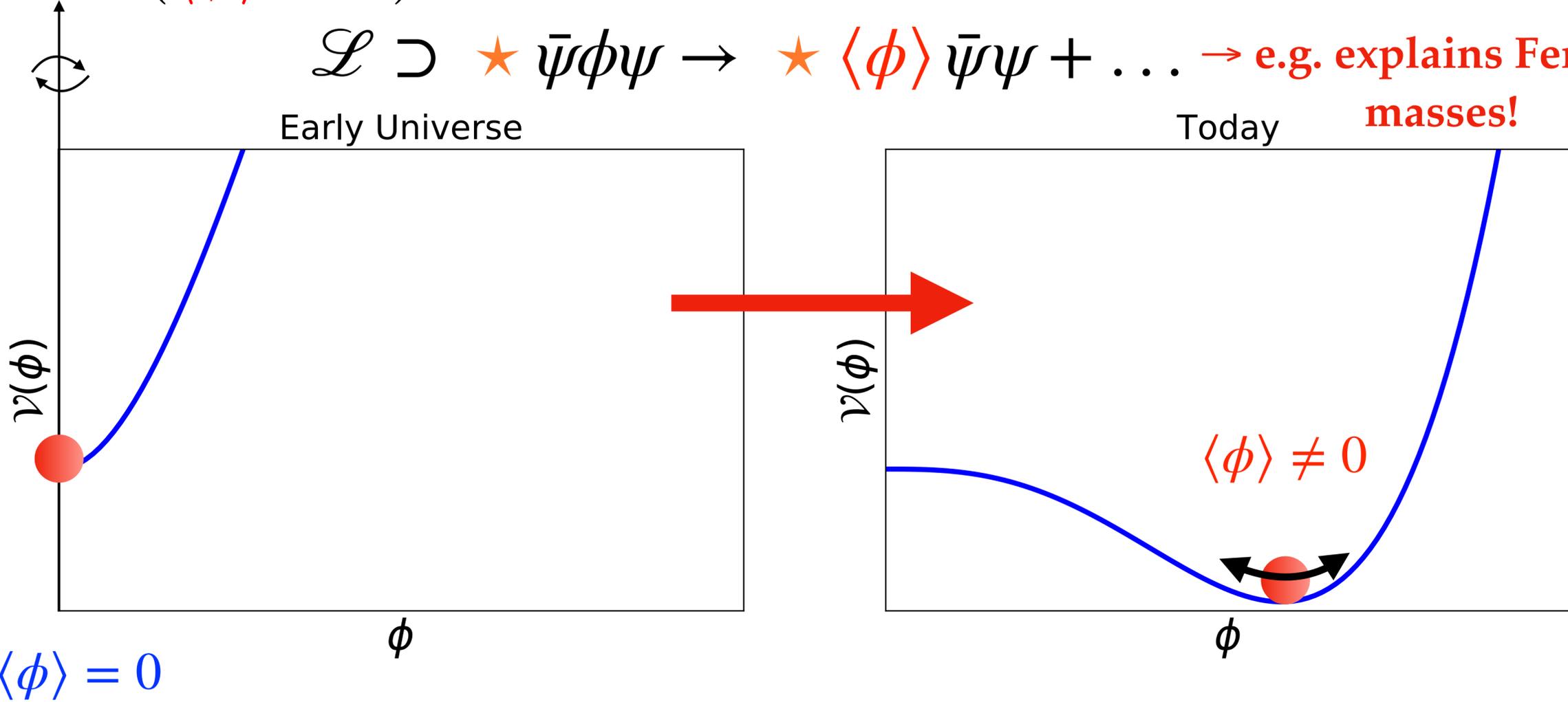


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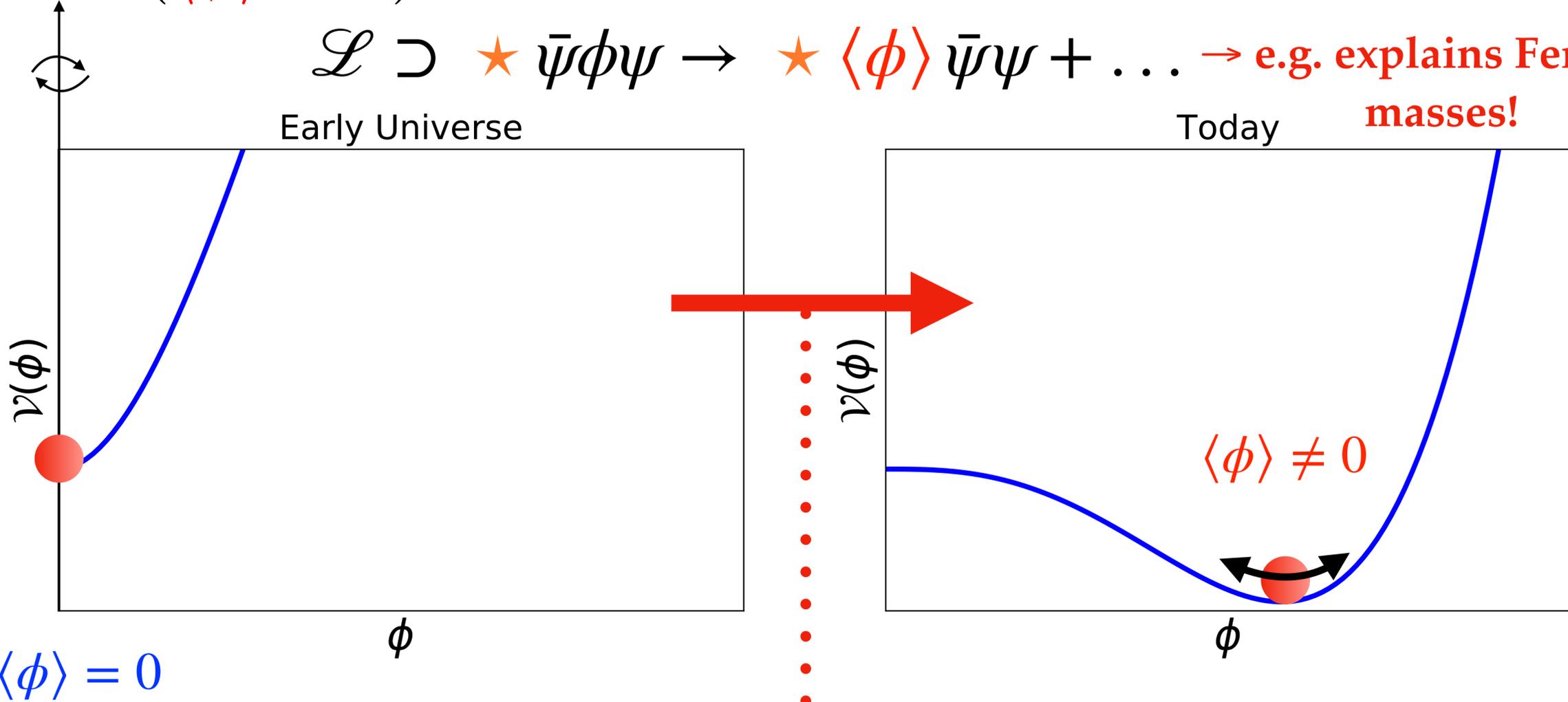


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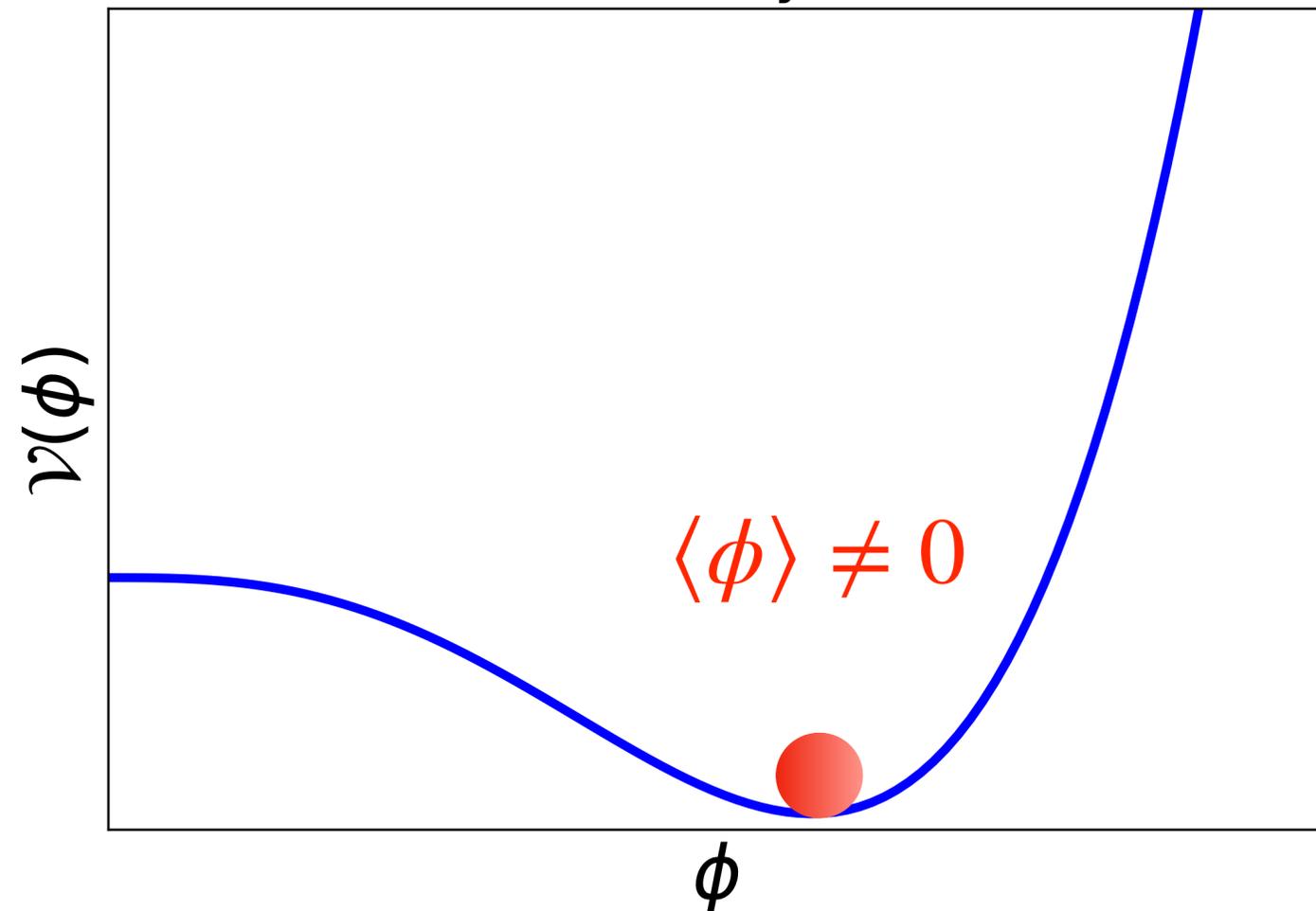
\equiv Electro-Weak Phase Transition (EWPT)



The Higgs Boson's Potential

$$\mathcal{V}(\langle\phi\rangle + h) = \bullet h^2 + \blacktriangle h^3 + \blacksquare h^4 \rightarrow \text{the Higgs boson's self-interactions.}$$

Today



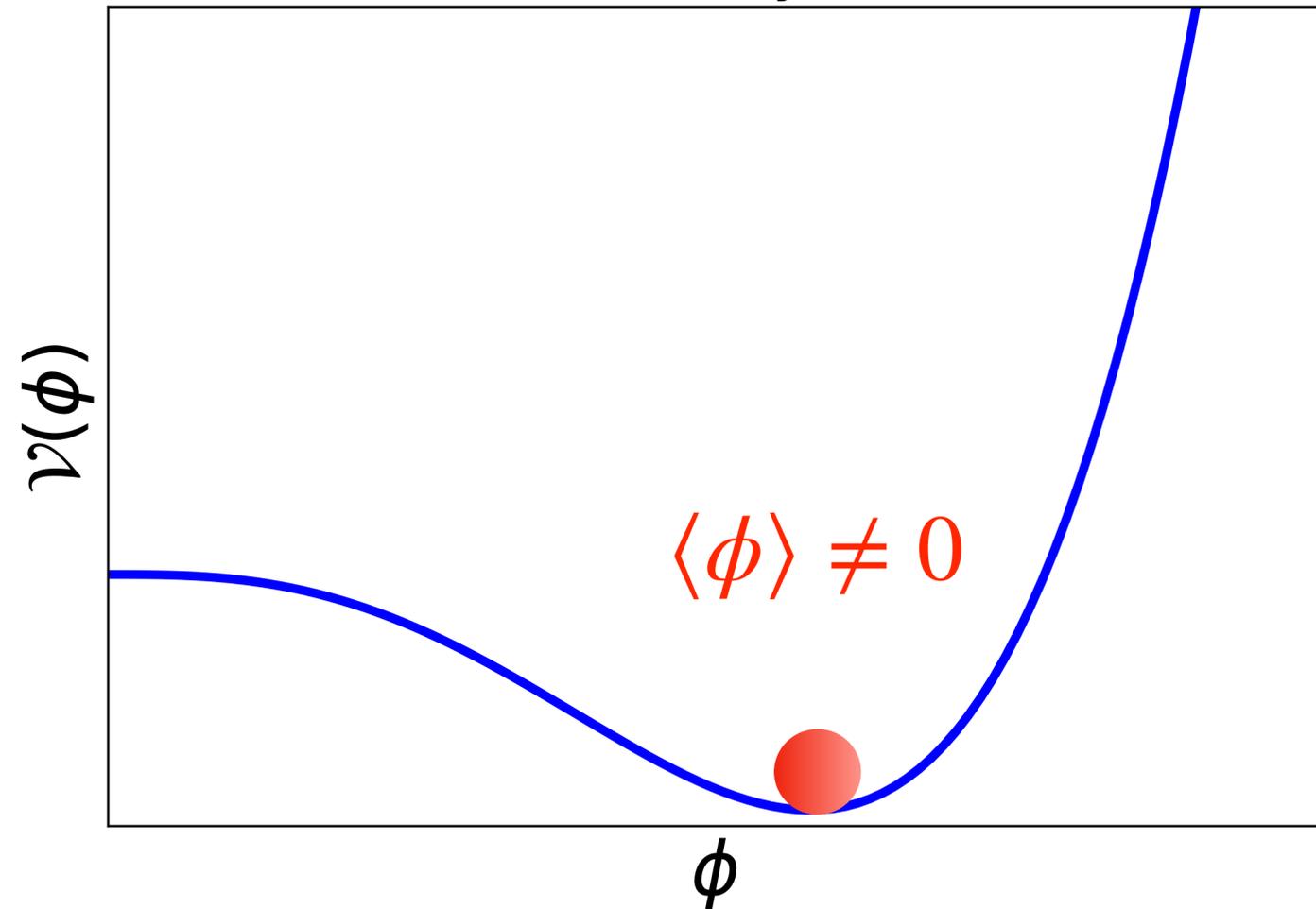
→ Determine shape of potential by measuring:

$$\{\bullet, \blacktriangle, \blacksquare\}$$

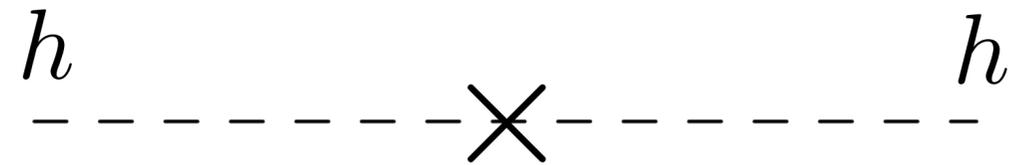
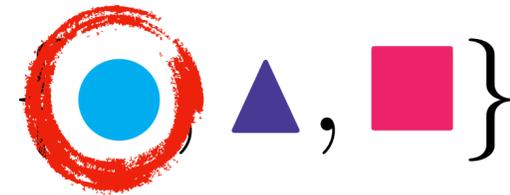
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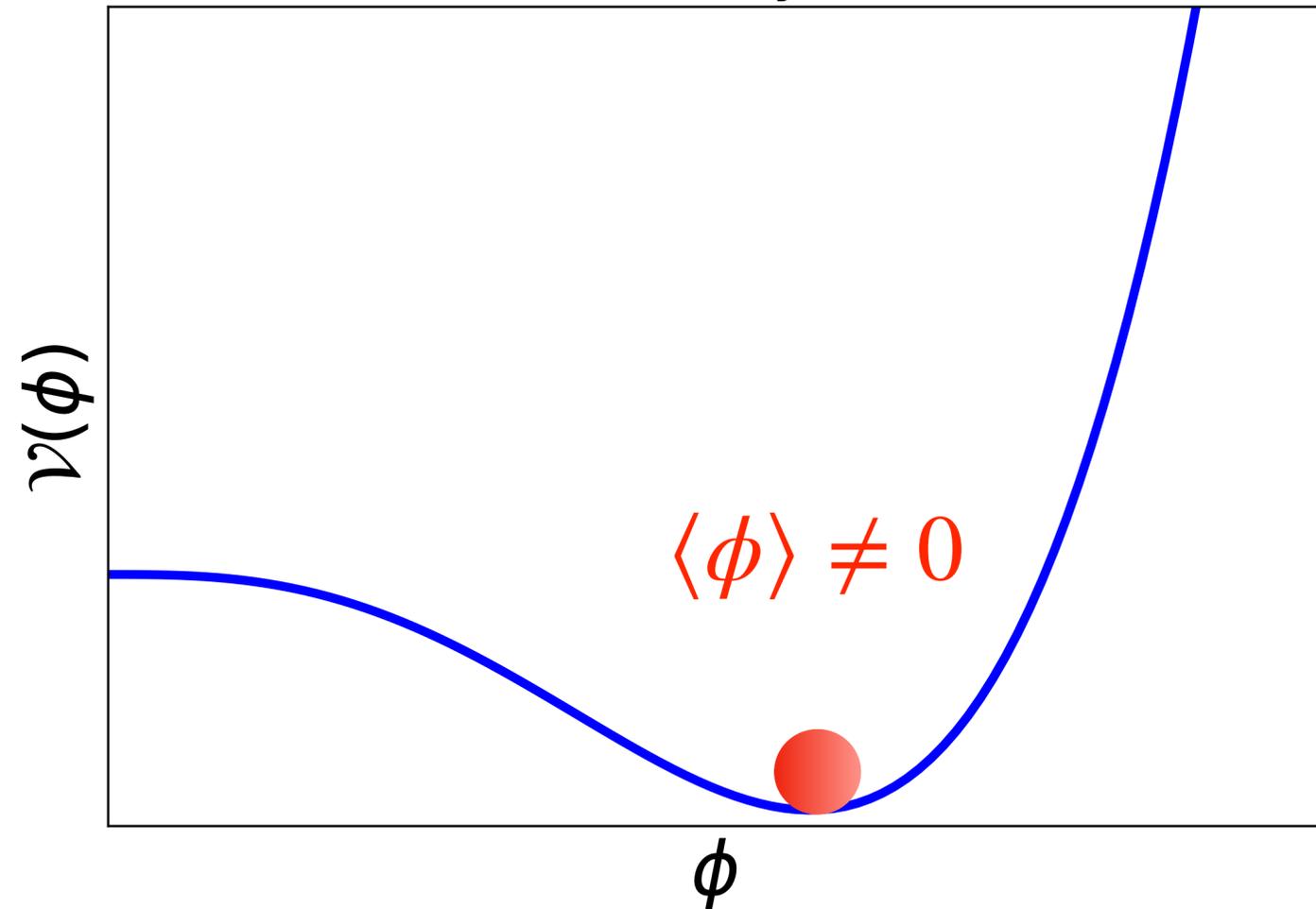
Higgs boson discovery @ LHC, 2012



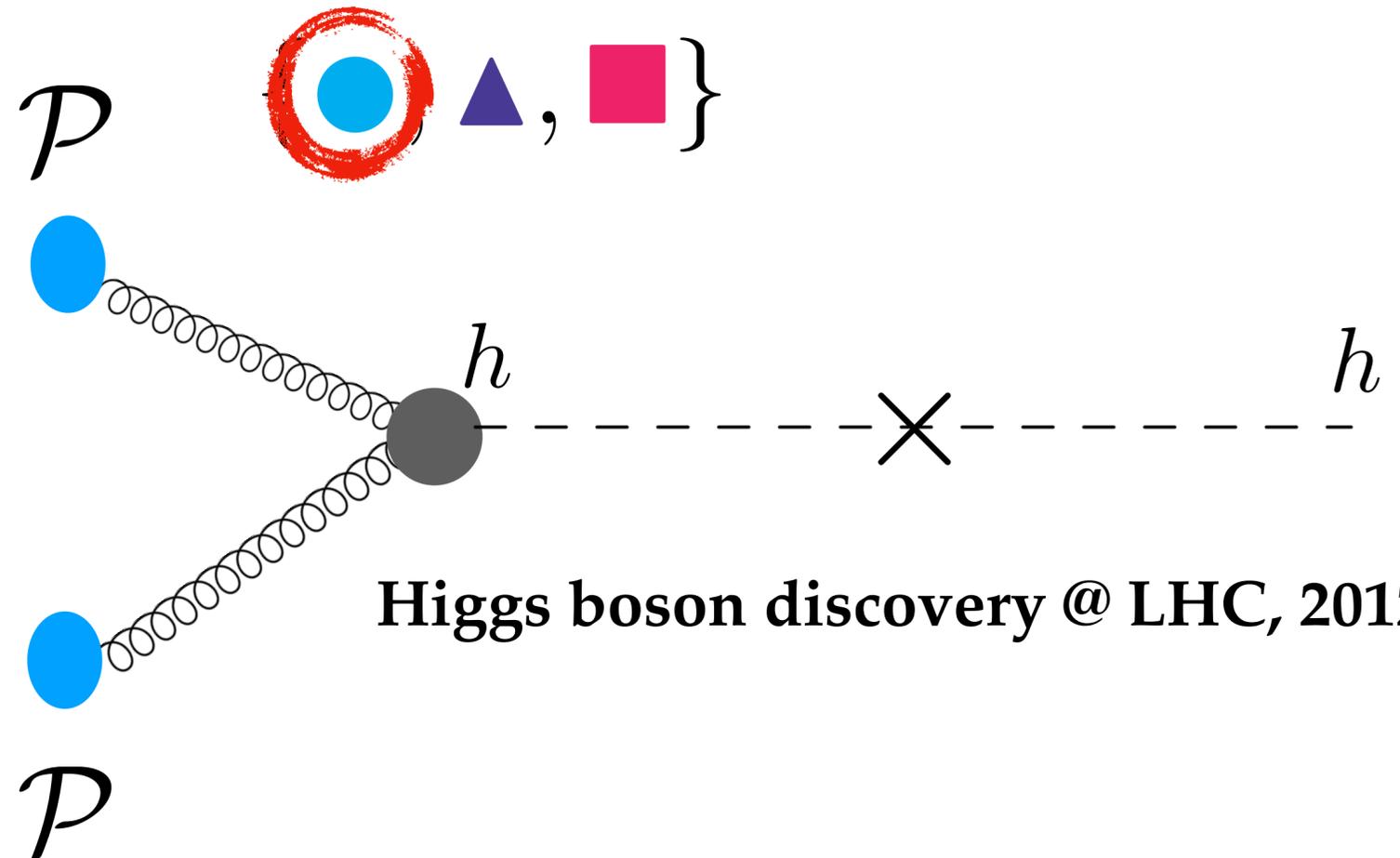
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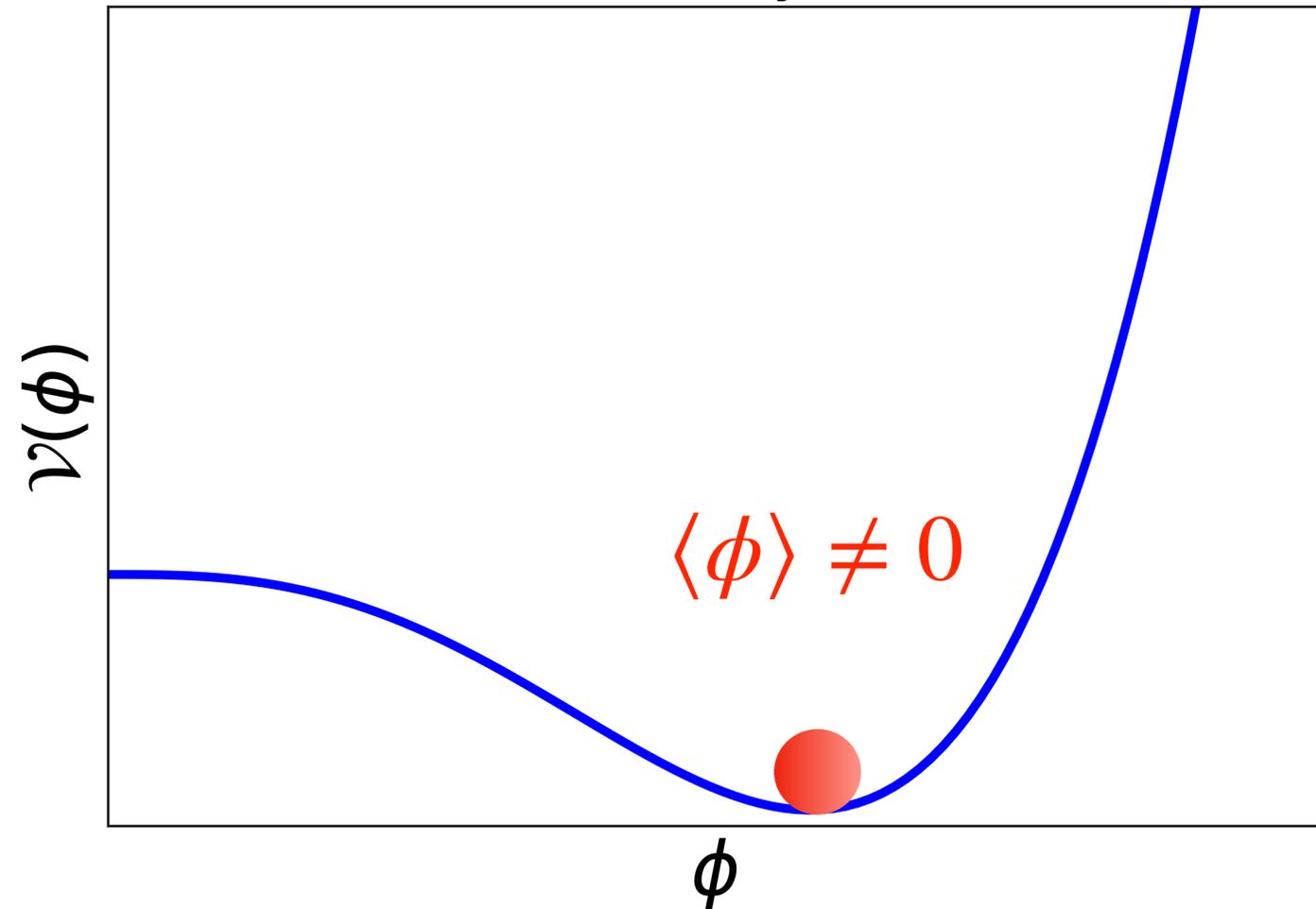
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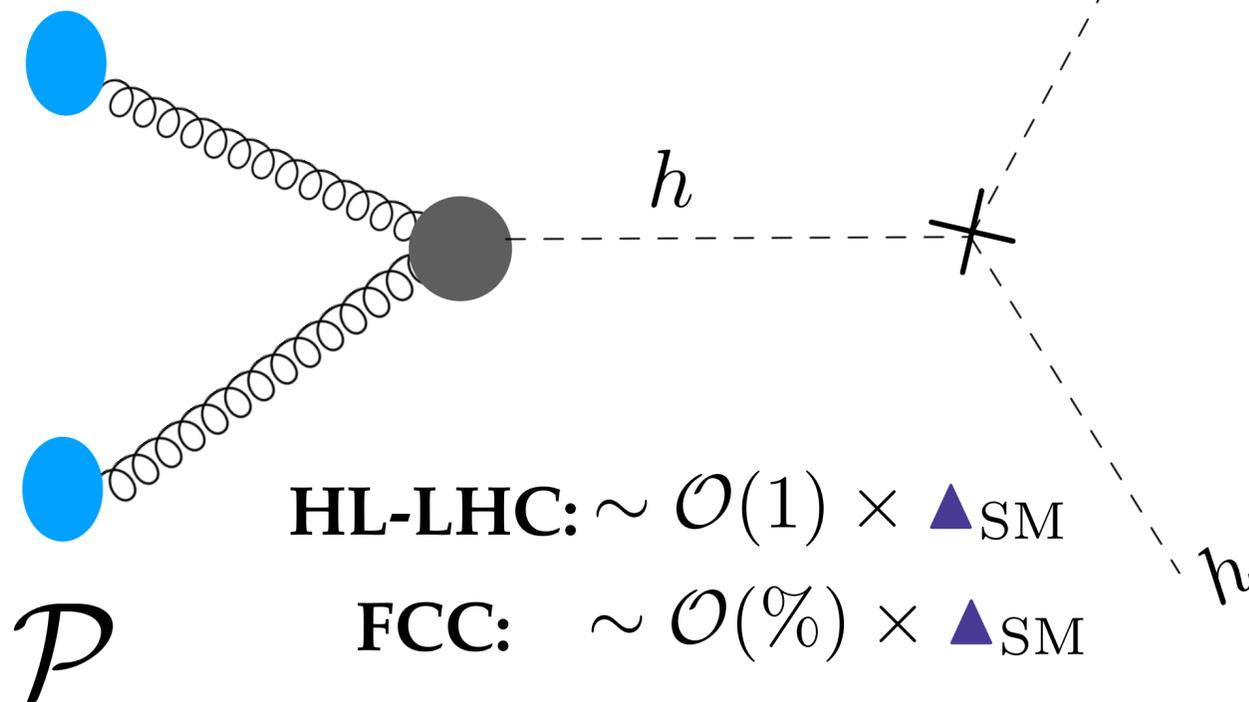
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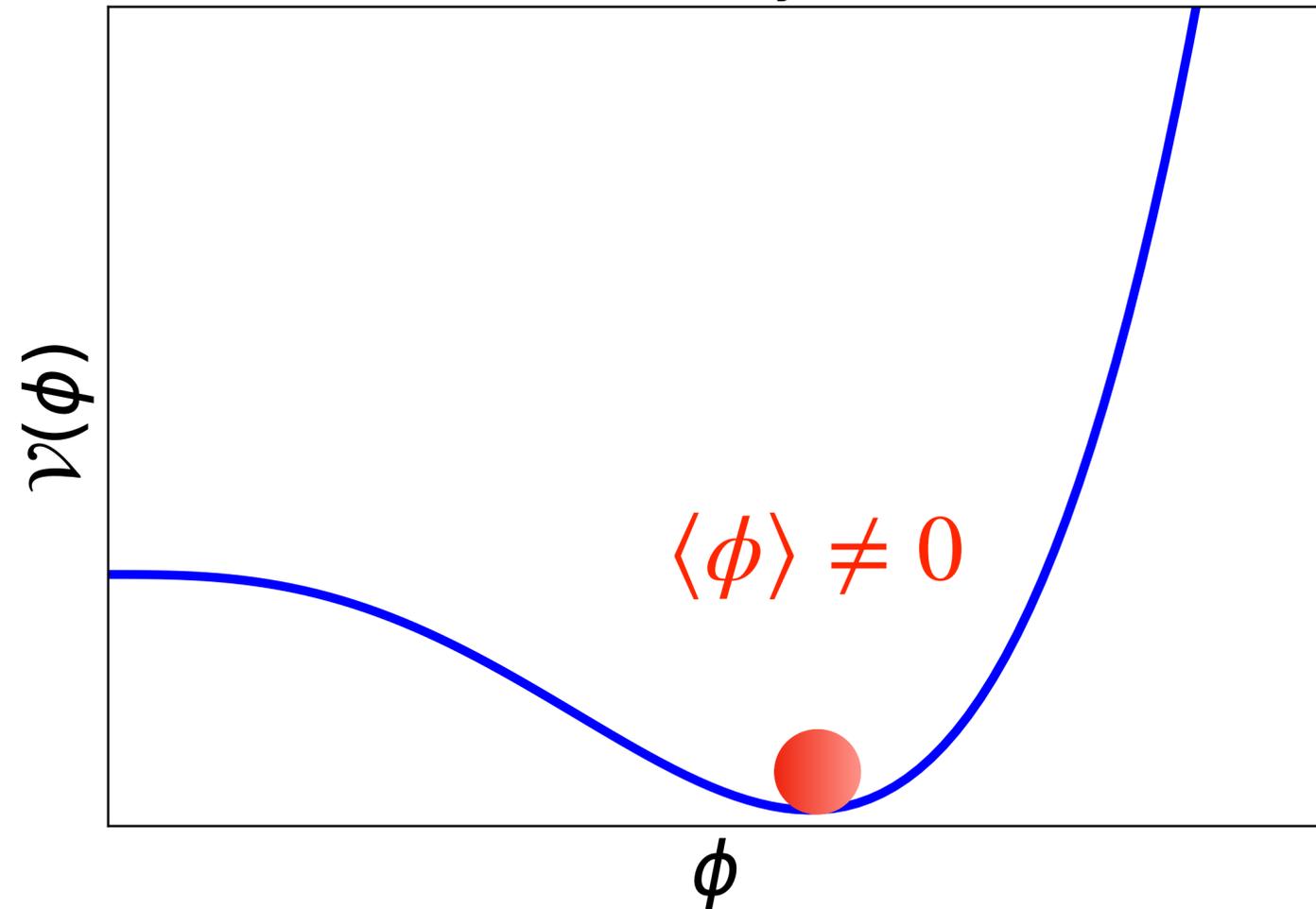
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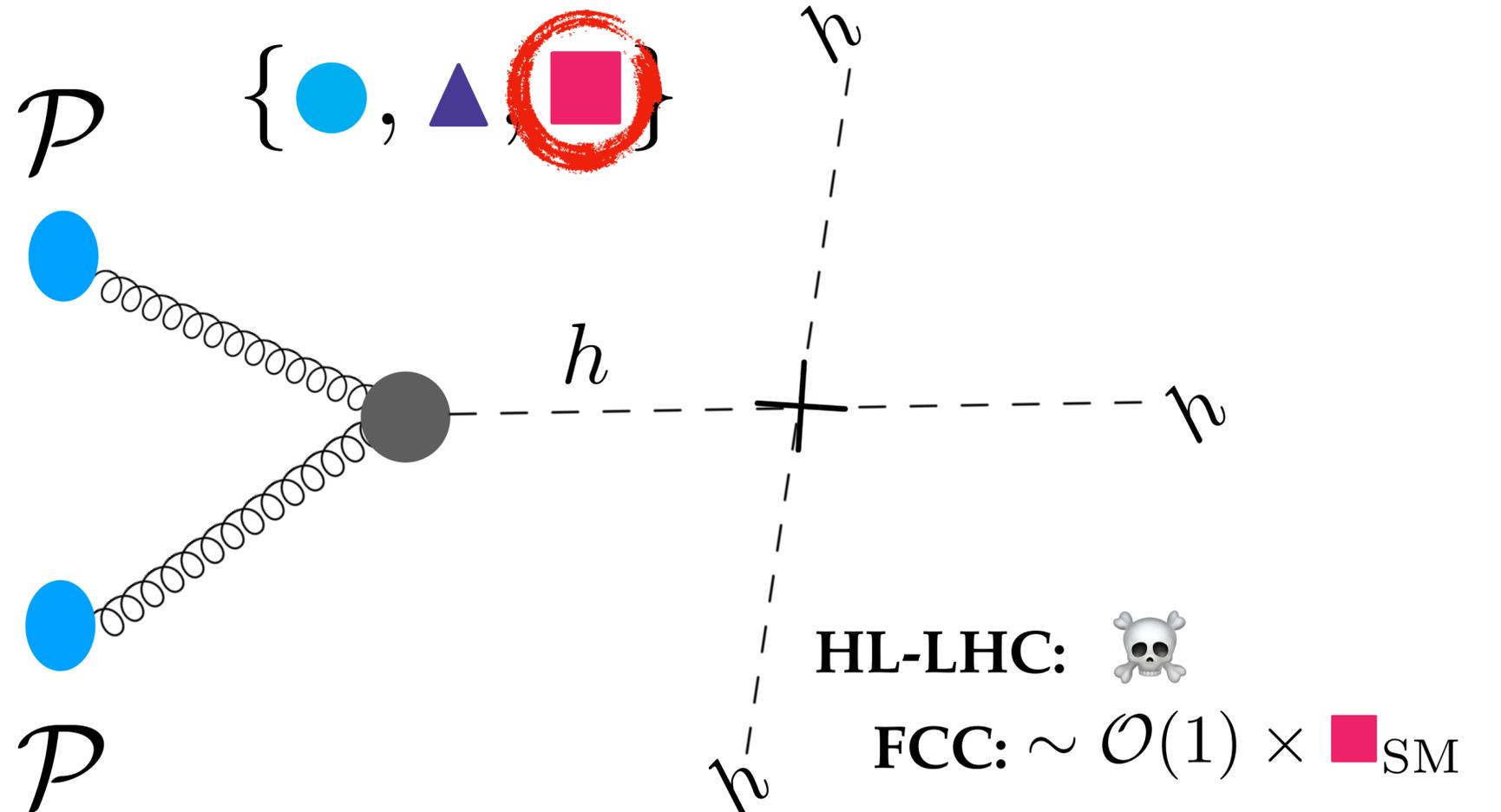
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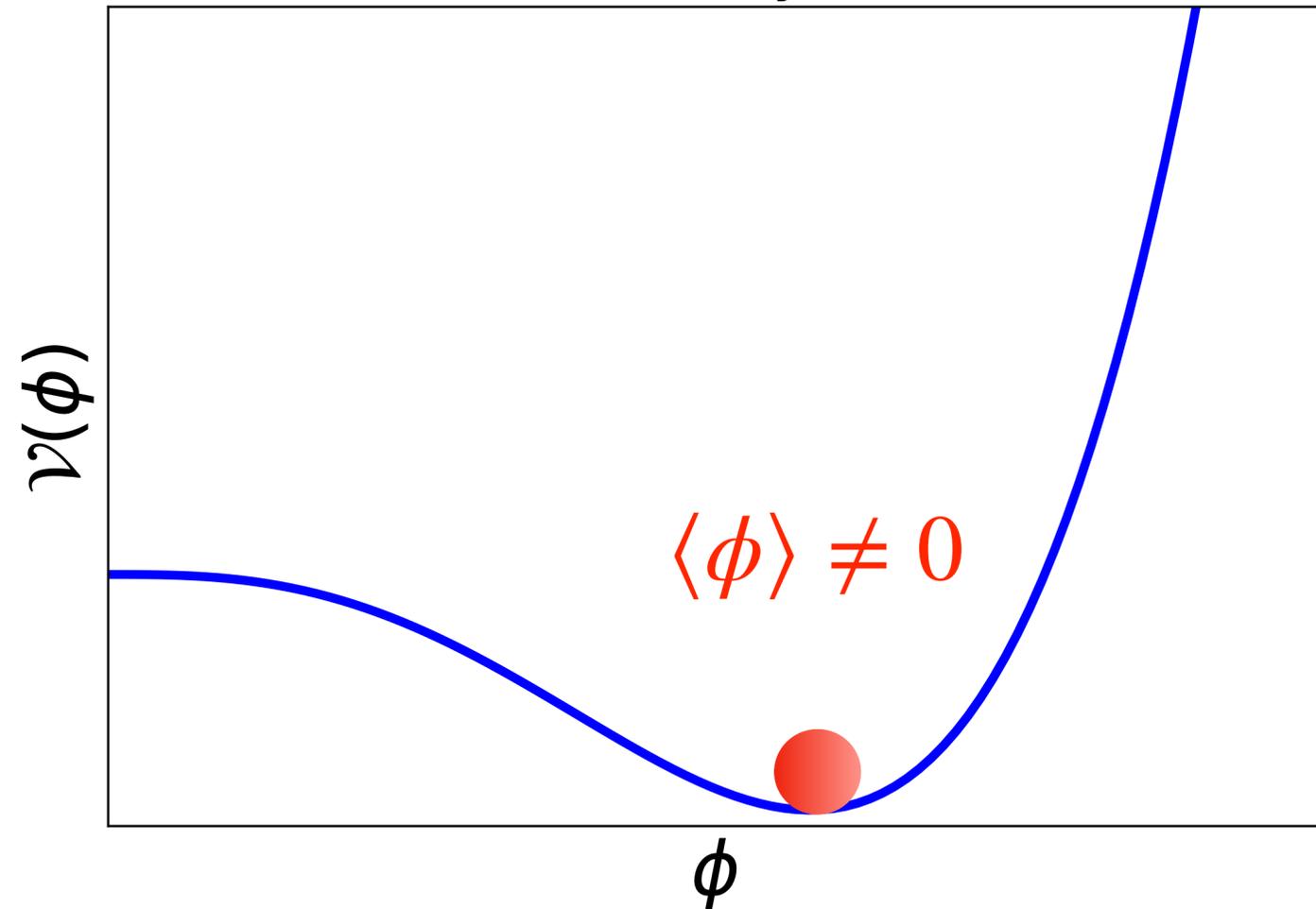


[e.g. [AP](#), Sakurai, arXiv:1508.06524, [AP](#),
Tetlalmatzi-Xolocotzi, Zaro,
arXiv:1909.09166]

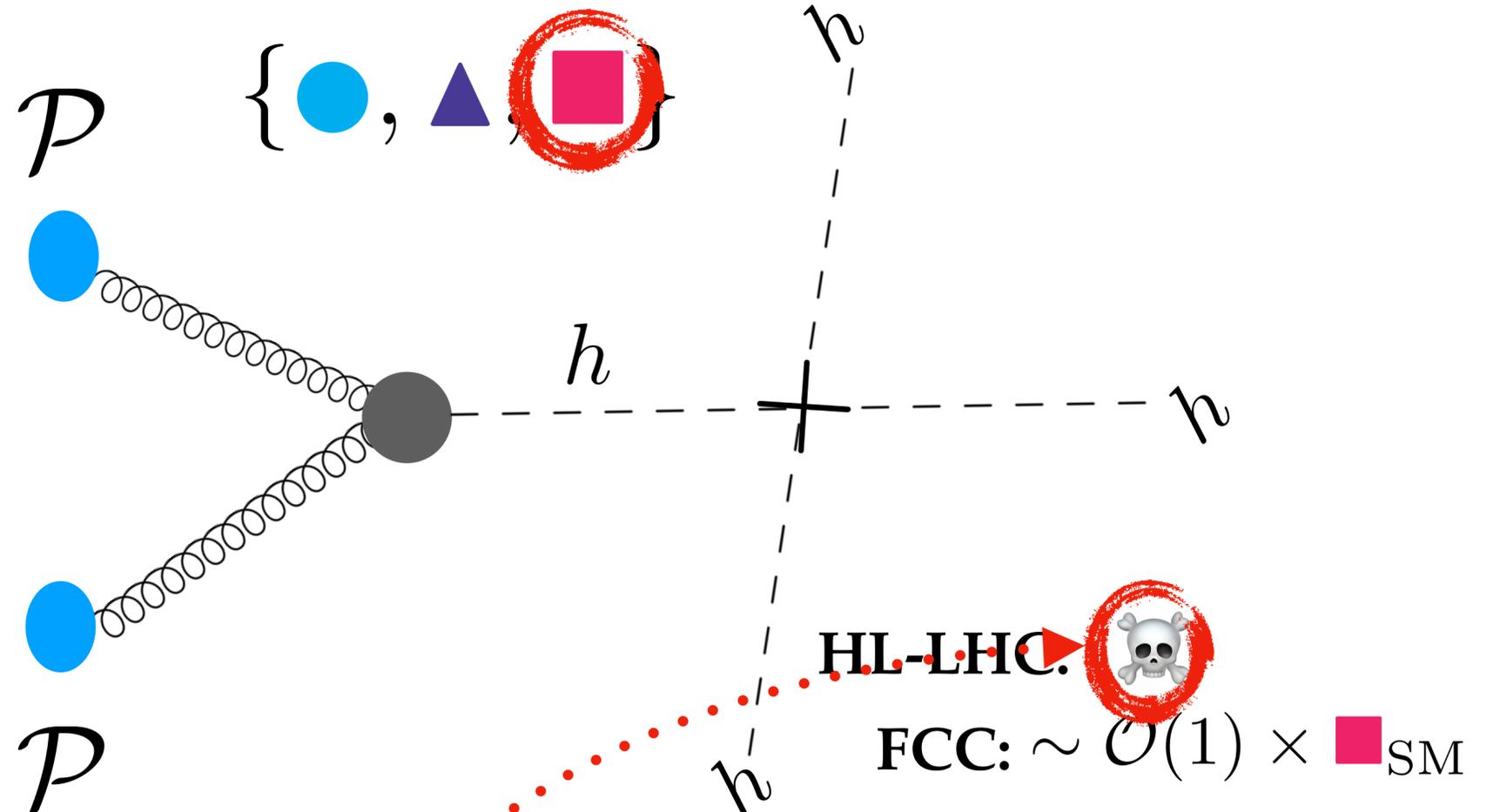
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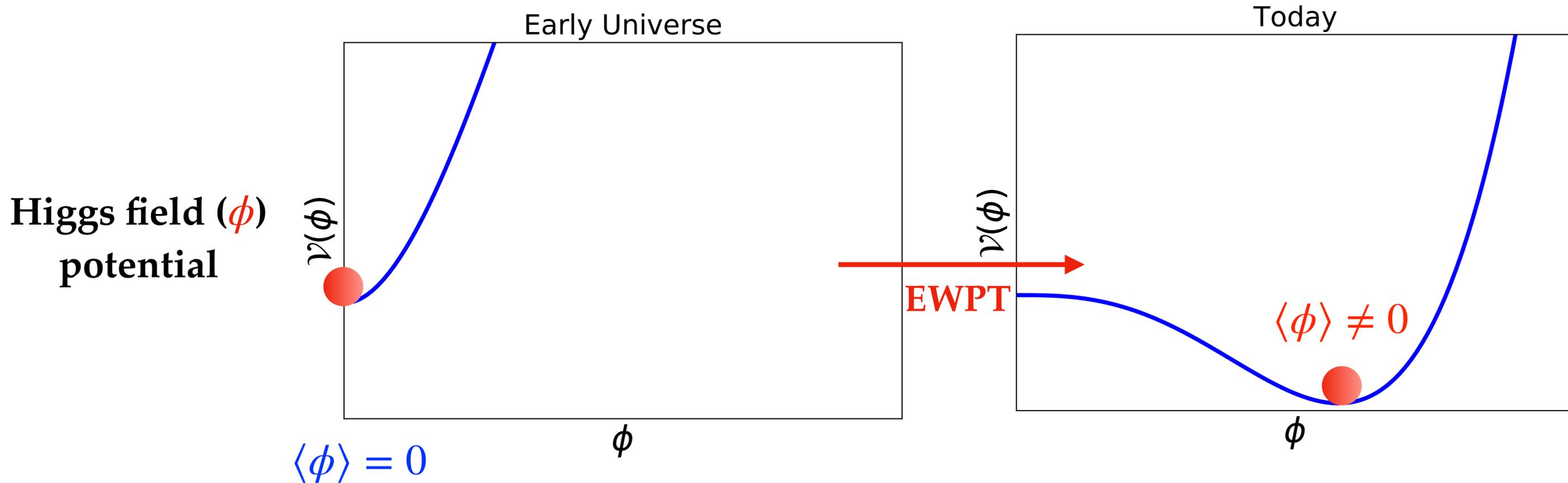
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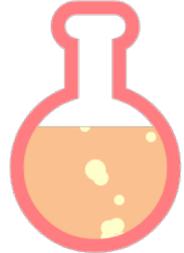


[[AP](#), Robens, Tetlalmatzi-Xolocotzi, arXiv:2101.00037 + Karkout, [AP](#), Postma, du Pree, Tetlalmatzi-Xolocotzi, van de Vis, arXiv:2404.12425, [AP](#), Tetlalmatzi-Xolocotzi, aXiv:2312.13562]

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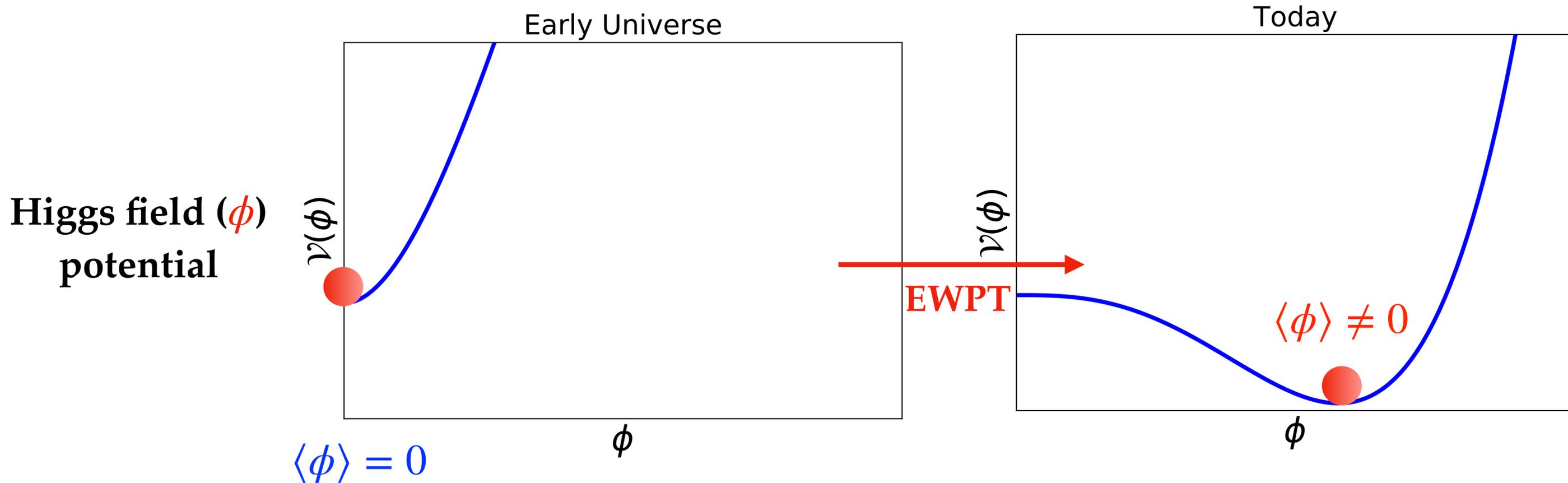
Breaking the Symmetry in the SM



- **Nature** of EWPT \rightarrow Important open question, e.g. **its order**:
 - A **First-Order transition** (e.g. the boiling of water)? 
 - or a **Second-Order transition** (e.g. the superfluid transition) or **cross-over**?



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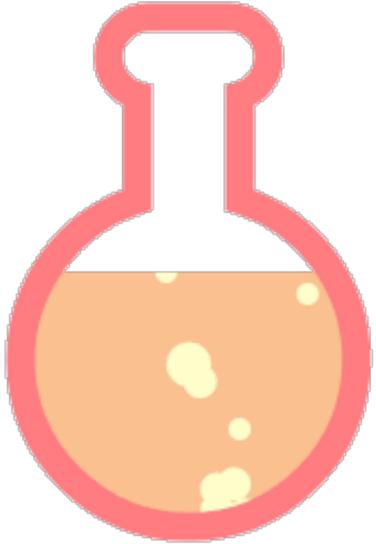
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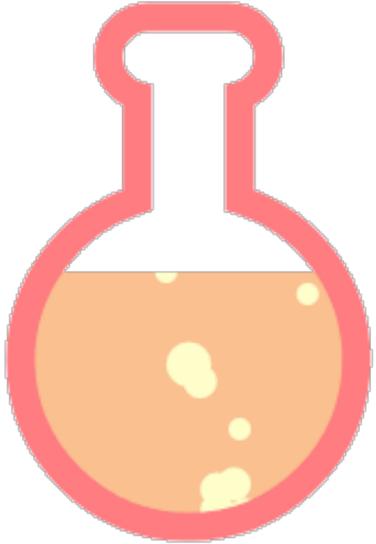
The Nature of the Phase Transition

- Clues to the origin of **matter-anti-matter asymmetry**.
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 - “Electro-Weak Baryogenesis” (**EWBG**).
- Pre-requisite: **a First-Order transition**. → 
- *Note:* This **does not** occur in the SM!

[Kajantie, Laine, Rummukainen, Shaposhnikov hep-ph/9605288]



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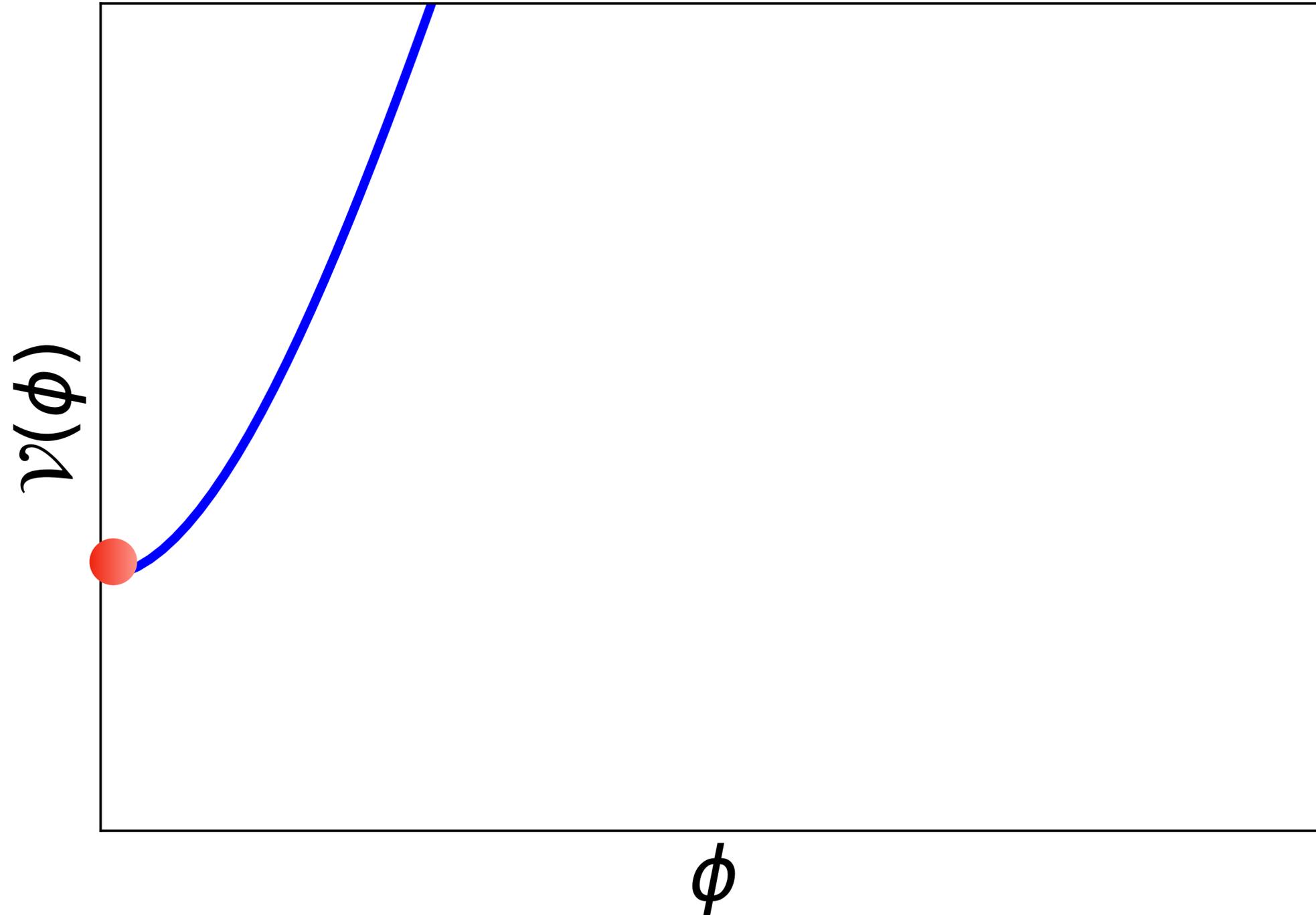
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A First-Order Electro-Weak Phase Transition: As the Universe Cools Down!

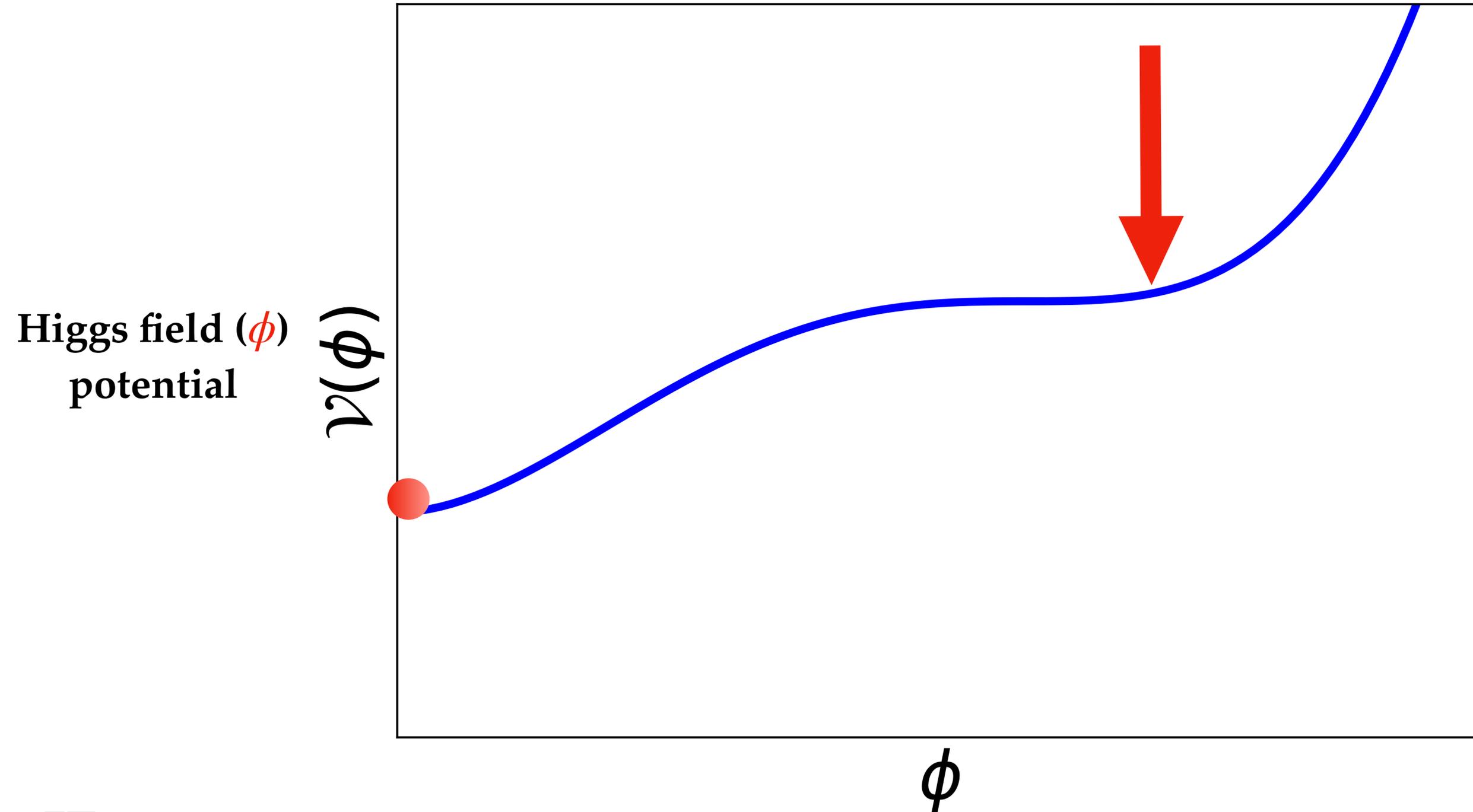
Early Universe

Higgs field (ϕ)
potential



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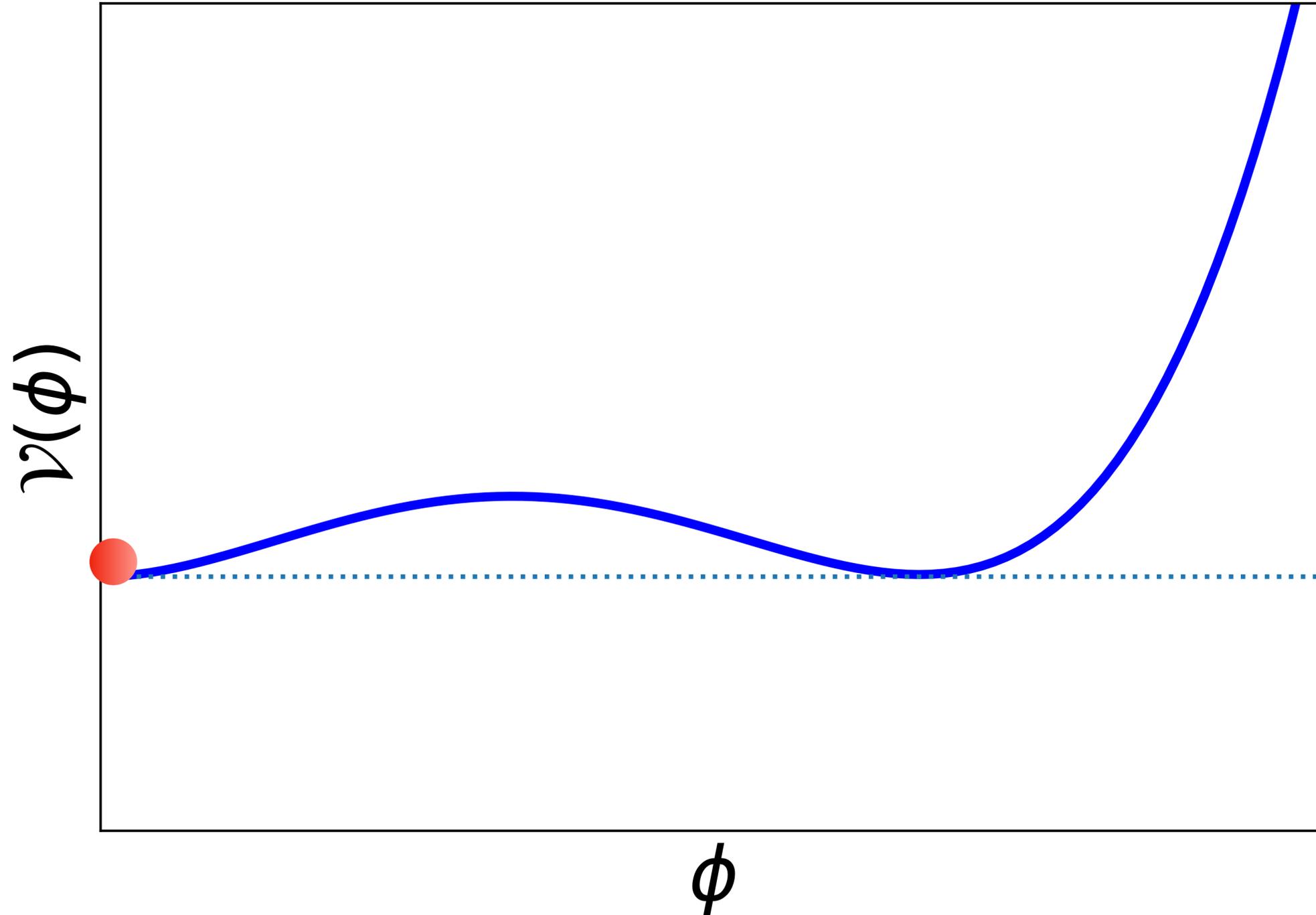
Second minimum appears



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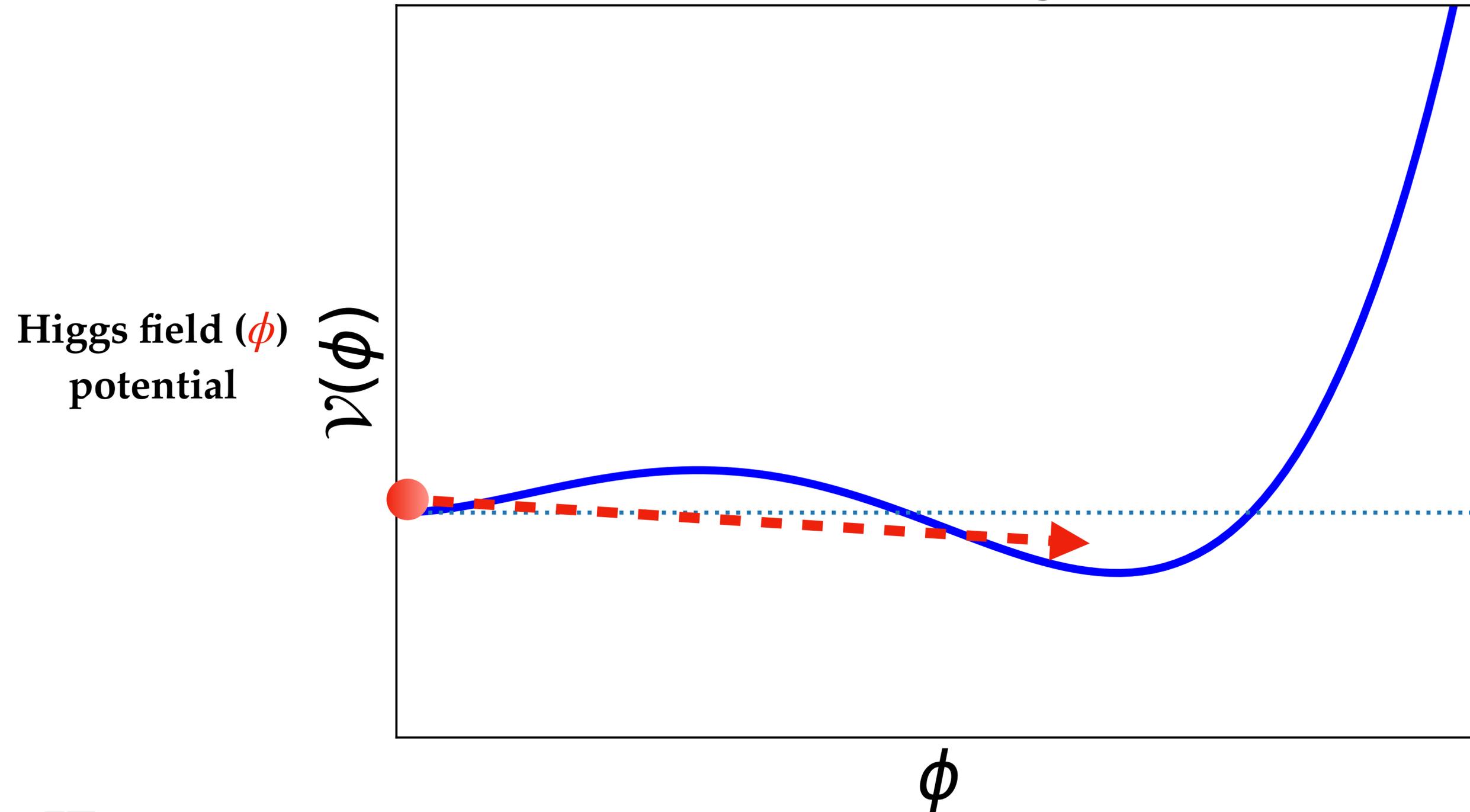
Critical temperature reached

Higgs field (ϕ)
potential



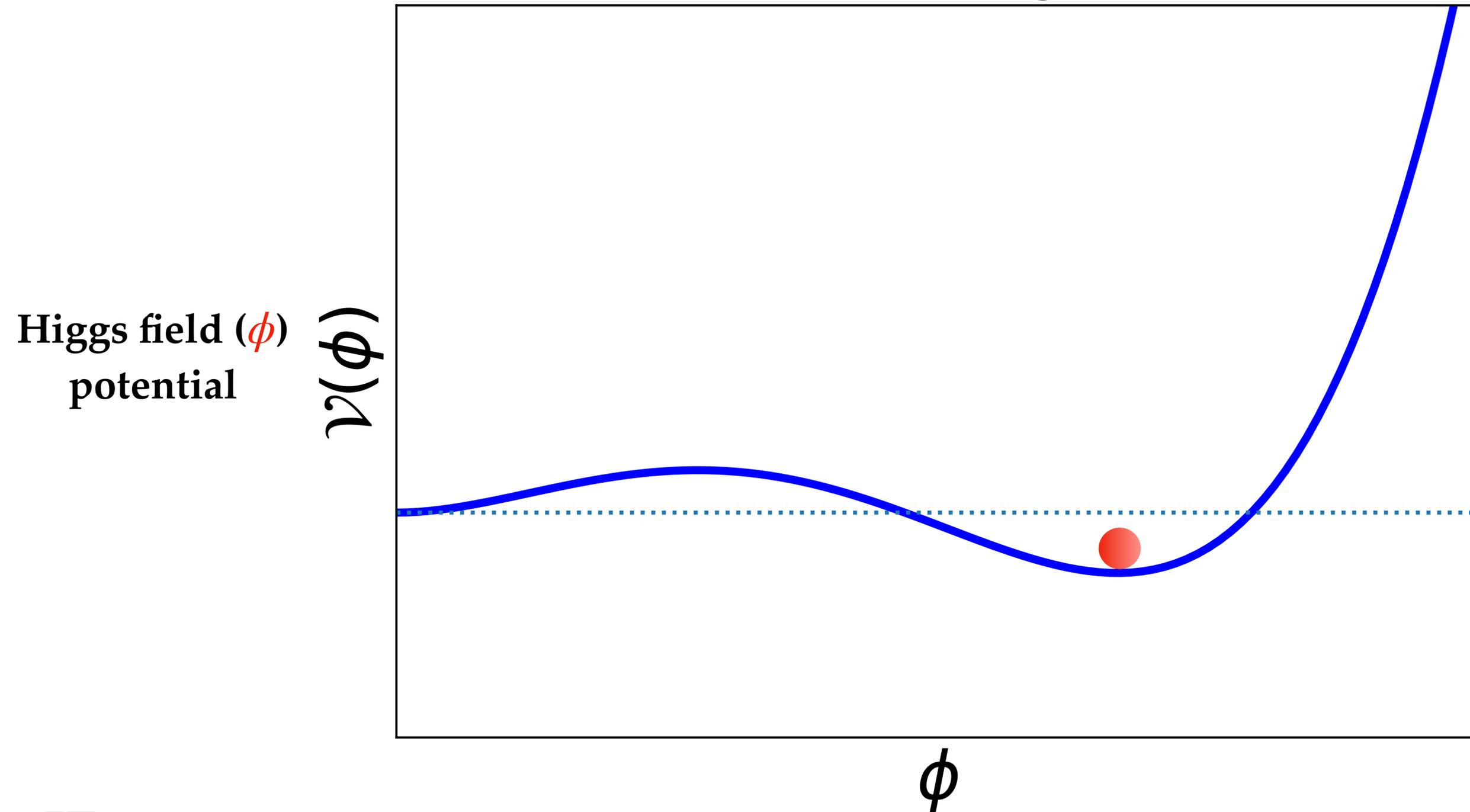
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Tunneling!



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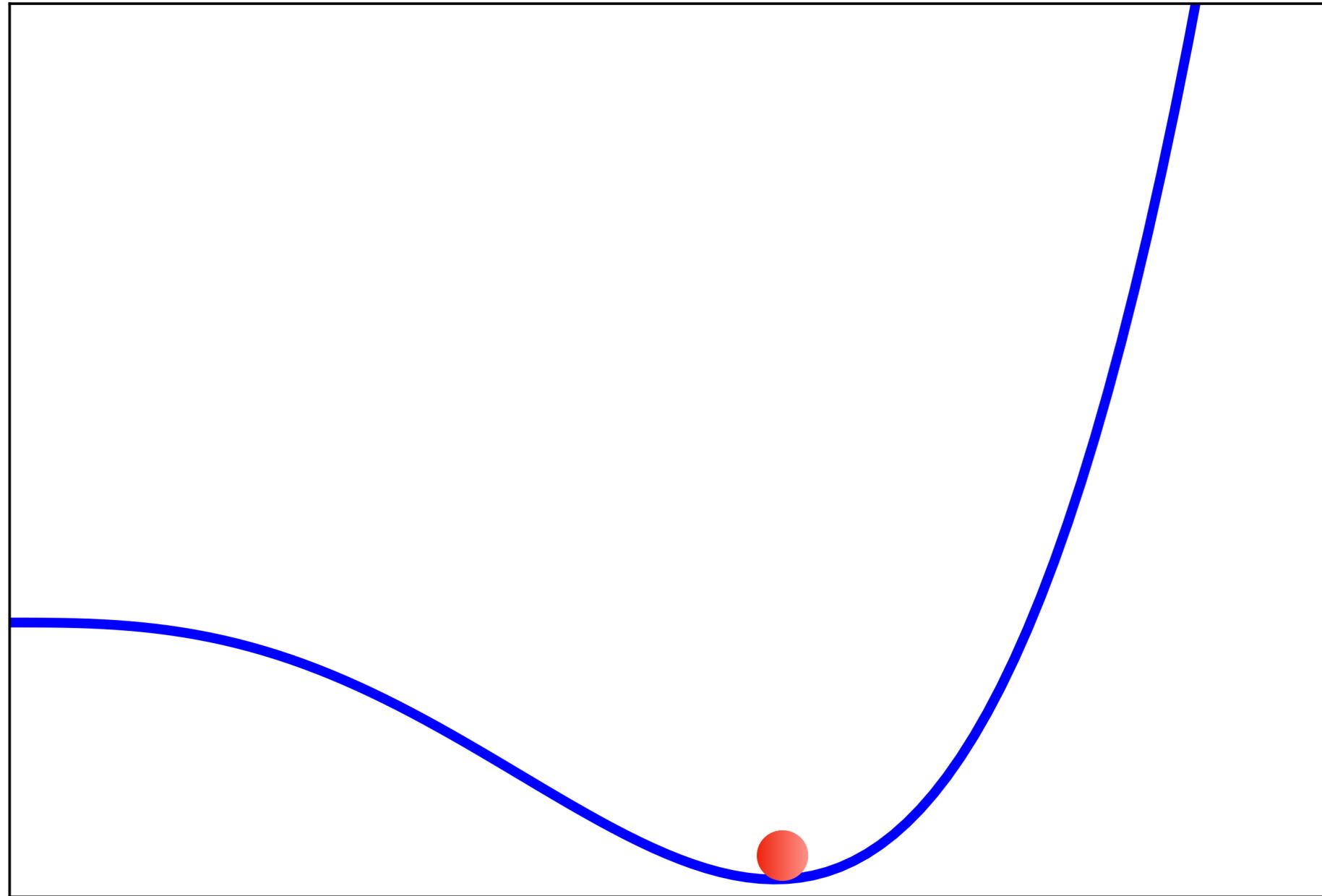


A First-Order Electro-Weak Phase Transition: As the Universe Cools Down!

Today

Higgs field (ϕ)
potential

$V(\phi)$

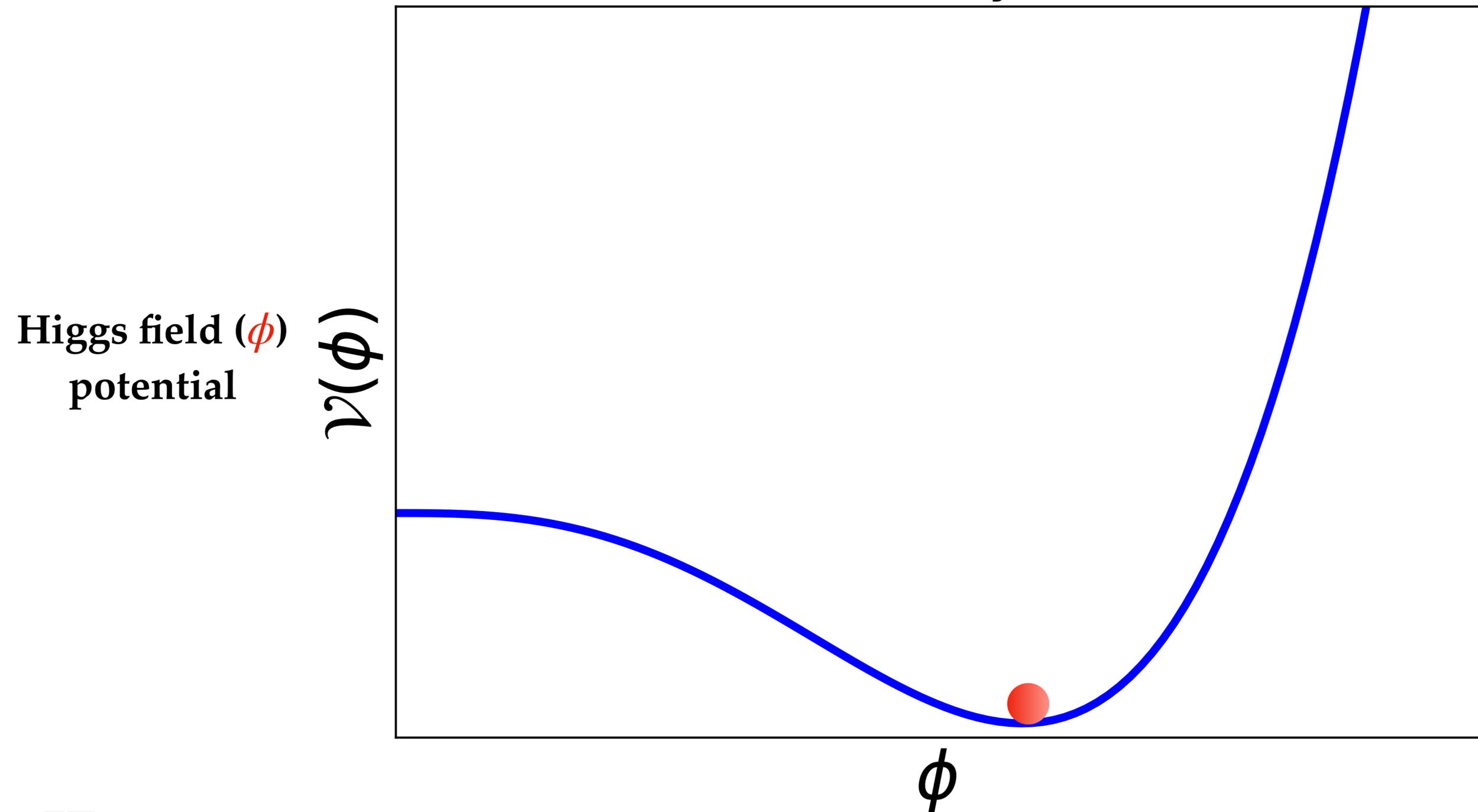


ϕ

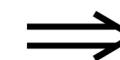


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Today



$$\langle \phi \rangle \neq 0$$



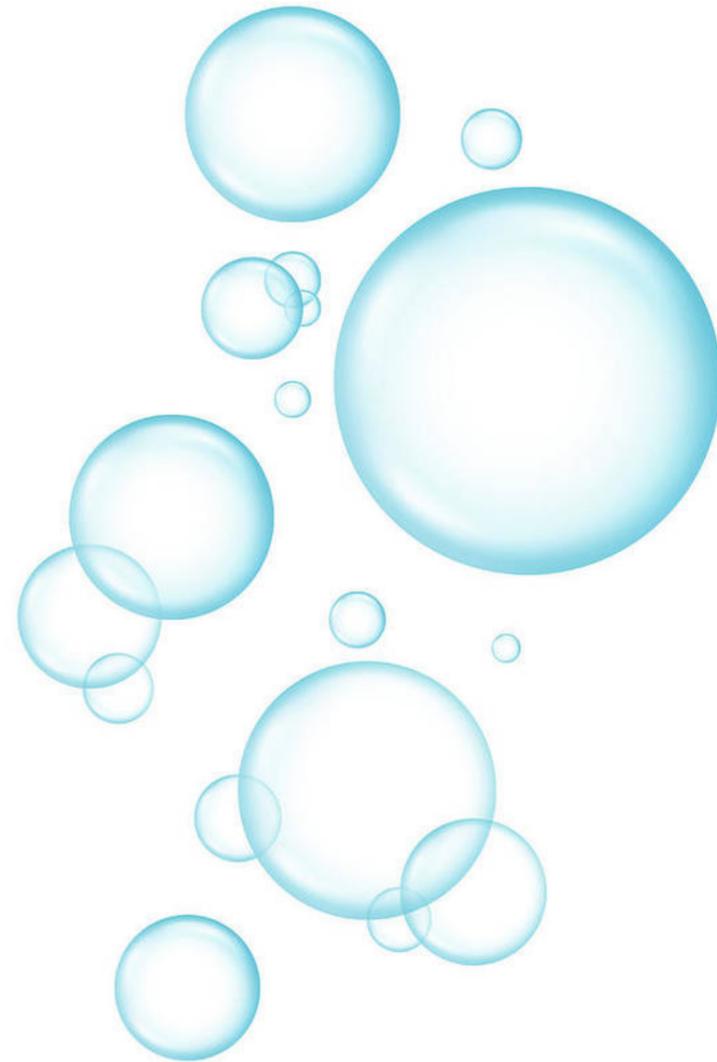
**Symmetry
Breaking.**



A First-Order EWPT

Some time after critical temperature is reached:

→ **Bubbles** of the broken phase nucleate and expand.*



**Note:* Gravitational Waves can form during this period!

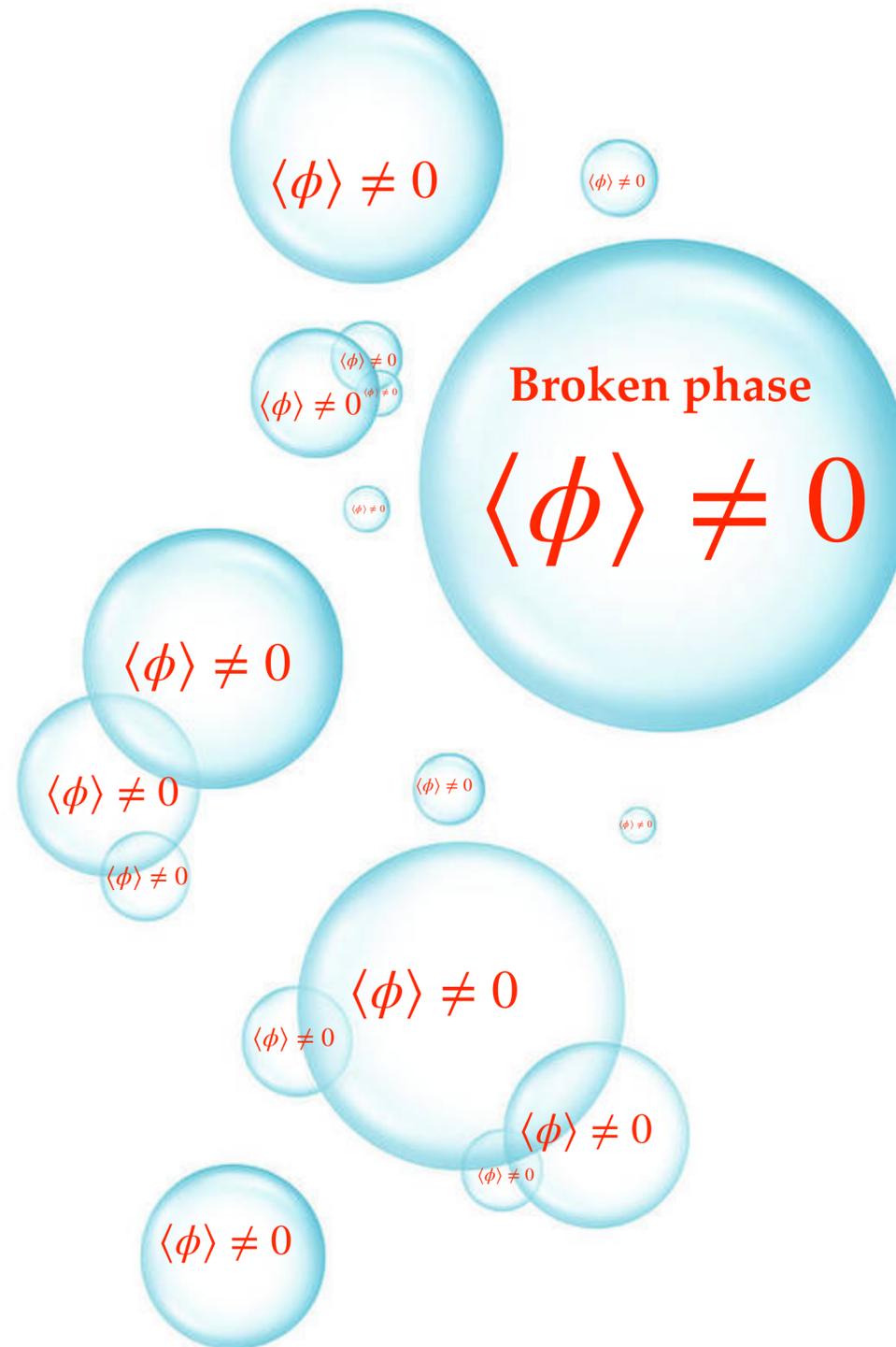


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$\langle \phi \rangle = 0$
Symmetric phase

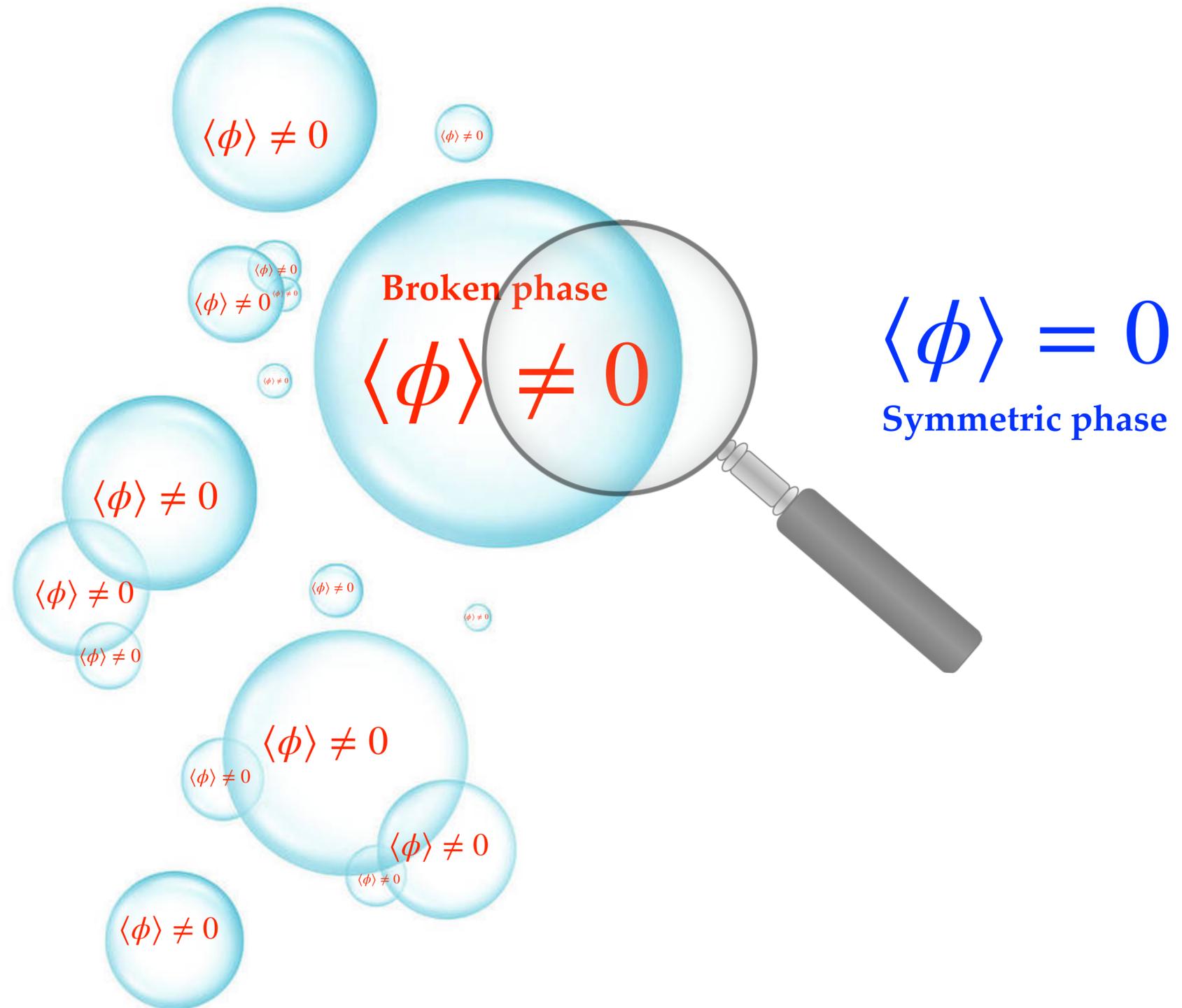


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Electro-Weak Baryogenesis

$\langle \phi \rangle \neq 0$
Broken phase

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Electro-Weak Baryogenesis

Left/Right-Handed Fermions

$$\psi_L + \psi_R$$

$$\langle \phi \rangle \neq 0$$

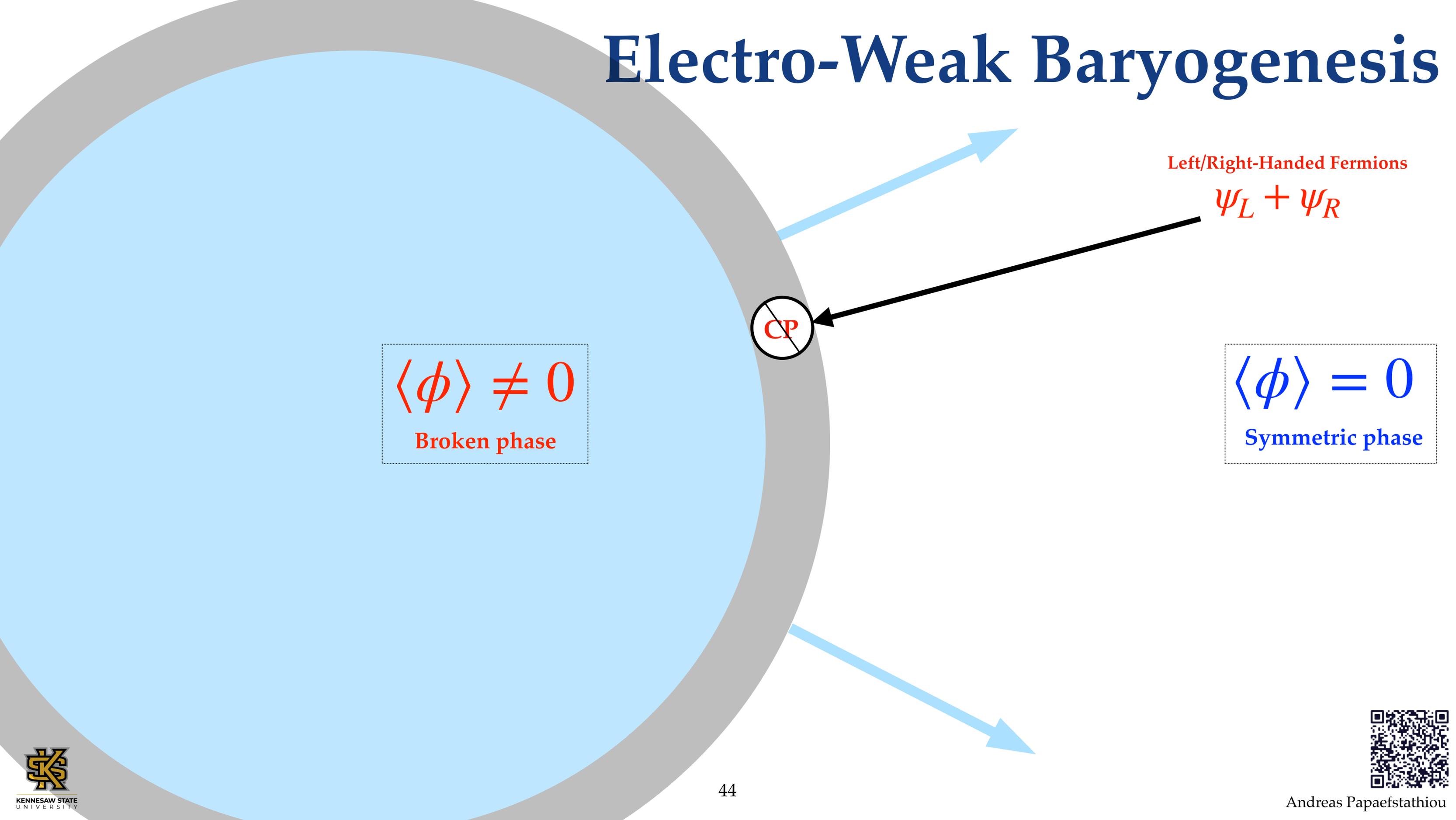
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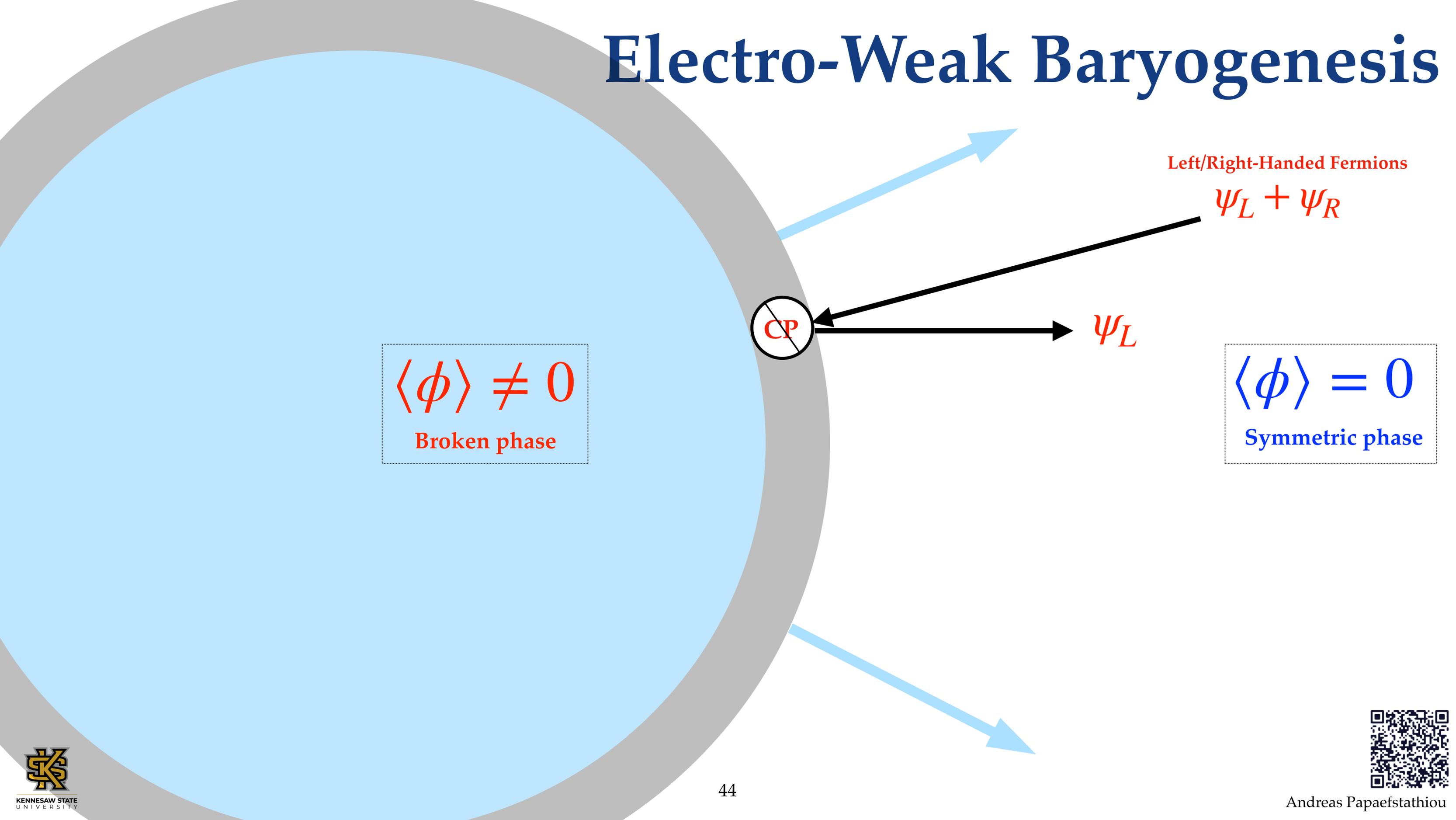
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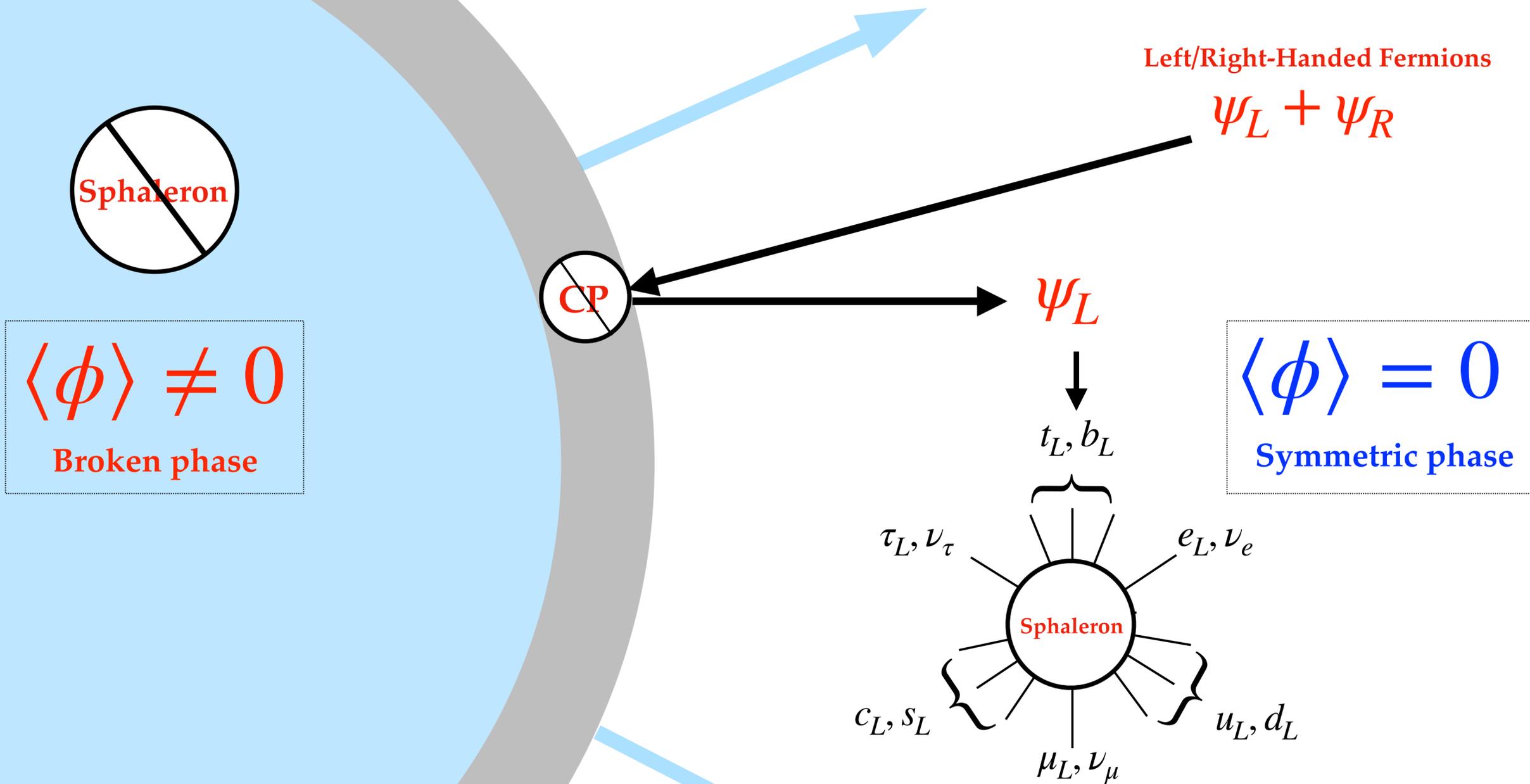
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ψ_L

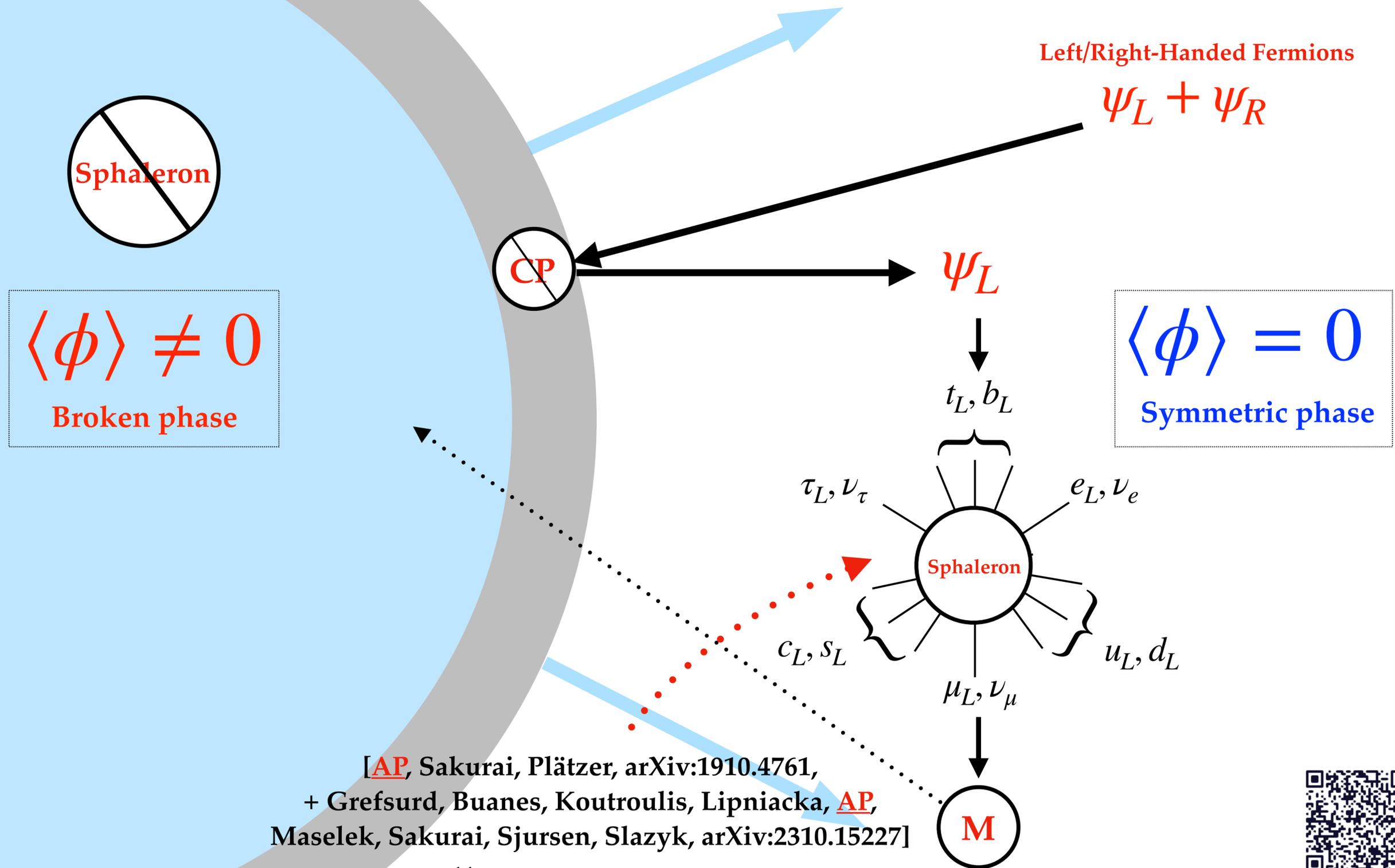
CP



Electro-Weak Baryogenesis



Electro-Weak Baryogenesis



[AP, Sakurai, Plätzer, arXiv:1910.4761,
 + Grefsurd, Buanes, Koutroulis, Lipniacka, AP,
 Maselek, Sakurai, Sjrursen, Slazyk, arXiv:2310.15227]



Electro-Weak Baryogenesis

~~Sphaleron~~

~~CP~~

Left/Right-Handed Fermions

$$\psi_L + \psi_R$$

$$\psi_L$$

$$\langle \phi \rangle = 0$$

Symmetric phase

$$t_L, b_L$$

$$\tau_L, \nu_\tau$$

$$e_L, \nu_e$$

Sphaleron

$$c_L, s_L$$

$$u_L, d_L$$

$$\mu_L, \nu_\mu$$

M

**A First-Order
Transition requires
New Phenomena
beyond the SM!**

[[AP](#), Sakurai, Plätzer, arXiv:1910.4761,
+ Grefsurd, Buanes, Koutroulis, Lipniacka, [AP](#),
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Electro-Weak Archaeology



We live here!



$$\langle \phi \rangle \neq 0$$

Broken phase



Electro-Weak Archaeology



We live here!

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Broken phase



→ What are the imprints of
Electro-Weak Baryogenesis
at Colliders?



4 Extending the Scalar Sector [AP, White, arXiv:2010.00597]

- A First-Order EWPT dictates new phenomena. [Kajantie, Laine, Rummukainen, Shaposhnikov hep-ph/9605288]
- Consider the **simplest possible extension to the SM!**

$$\mathcal{V}(\phi, S) = \bullet |\phi|^2 + \blacksquare |\phi|^4$$

Add: S , a new scalar field,
No SM “charges” \equiv Singlet.

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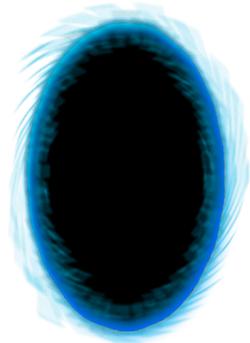
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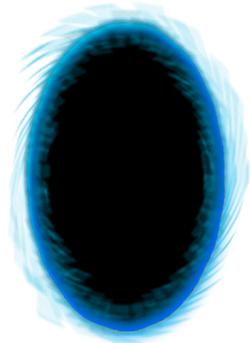
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$$+ S \times (\text{Hidden Sector}) + \dots \leftarrow \text{Dark Matter?}$$



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EWSB \leftrightarrow VEVs:

$$\phi \rightarrow \langle \phi \rangle + h$$

$$S \rightarrow \langle S \rangle + \chi$$



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\Rightarrow Mass (squared) matrix:

$$M^2 = \begin{pmatrix} \frac{\partial^2 \mathcal{V}}{\partial h^2} & \frac{\partial^2 \mathcal{V}}{\partial h \partial \chi} \\ \frac{\partial^2 \mathcal{V}}{\partial h \partial \chi} & \frac{\partial^2 \mathcal{V}}{\partial \chi^2} \end{pmatrix}$$



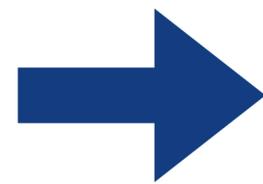
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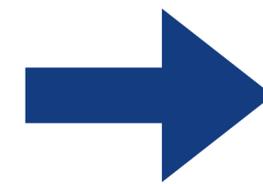
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Diagonalize!

\Rightarrow Mass (squared) matrix:

Mass Eigenstates

$$M^2 = \begin{pmatrix} \frac{\partial^2 \mathcal{V}}{\partial h^2} & \frac{\partial^2 \mathcal{V}}{\partial h \partial \chi} \\ \frac{\partial^2 \mathcal{V}}{\partial h \partial \chi} & \frac{\partial^2 \mathcal{V}}{\partial \chi^2} \end{pmatrix}$$

$$\begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} h \\ \chi \end{pmatrix}$$

θ : mixing angle



Extending the Scalar Sector [[AP](#), White, arXiv:2010.00597]

$$\mathcal{V}(\phi, S) = \color{green}{\bullet} |\phi|^2 + \color{blue}{\blacksquare} |\phi|^4 + \color{magenta}{\bullet} S^2 + \color{cyan}{\blacktriangle} S^3 + \color{maroon}{\blacksquare} S^4 + \color{red}{\blacktriangle} |\phi|^2 S + \color{purple}{\blacksquare} |\phi|^2 S^2$$

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$$\begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} h \\ \chi \end{pmatrix} \rightarrow \begin{matrix} h_1 \rightarrow \text{“SM-like” Higgs} \\ \text{boson.} \\ h_2 \rightarrow \text{new scalar} \\ \text{resonance.} \end{matrix}$$

θ : mixing angle

i.e. choose: $|\theta| \gtrsim 0$, and:

$$h_1 = h \cos \theta + \chi \sin \theta$$

$$h_2 = -h \sin \theta + \chi \cos \theta$$



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$h_1 \rightarrow$ "SM-like" Higgs boson.

$h_2 \rightarrow$ new scalar resonance.

i.e. choose: $|\theta| \gtrsim 0$, and:

$$h_1 = h \cos \theta + \chi \sin \theta$$

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Primary targets for collider studies!

$h_1 \Rightarrow$ **Reductions in Higgs boson rates.**

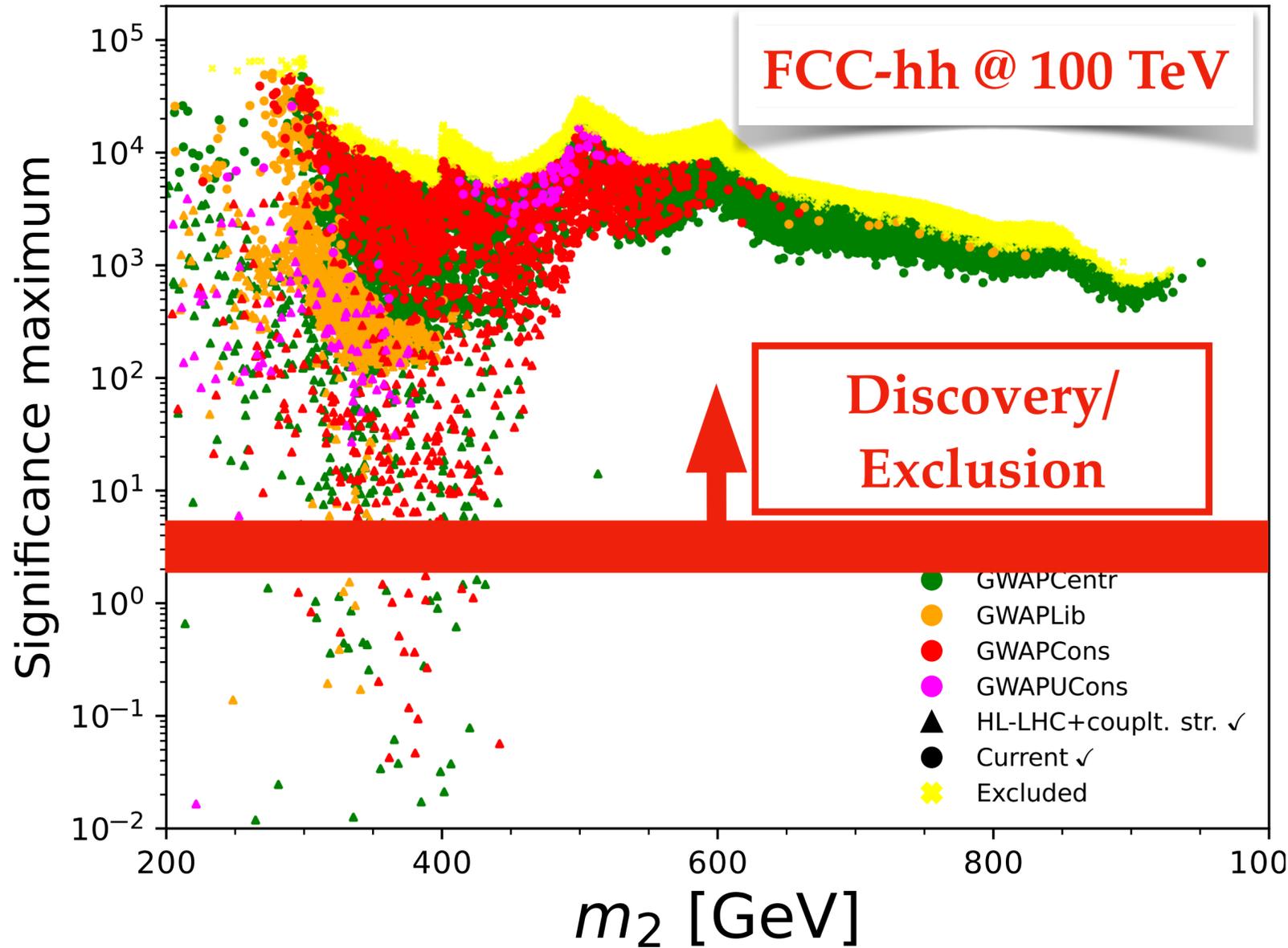
$h_2 \Rightarrow$ **New resonance searches.**



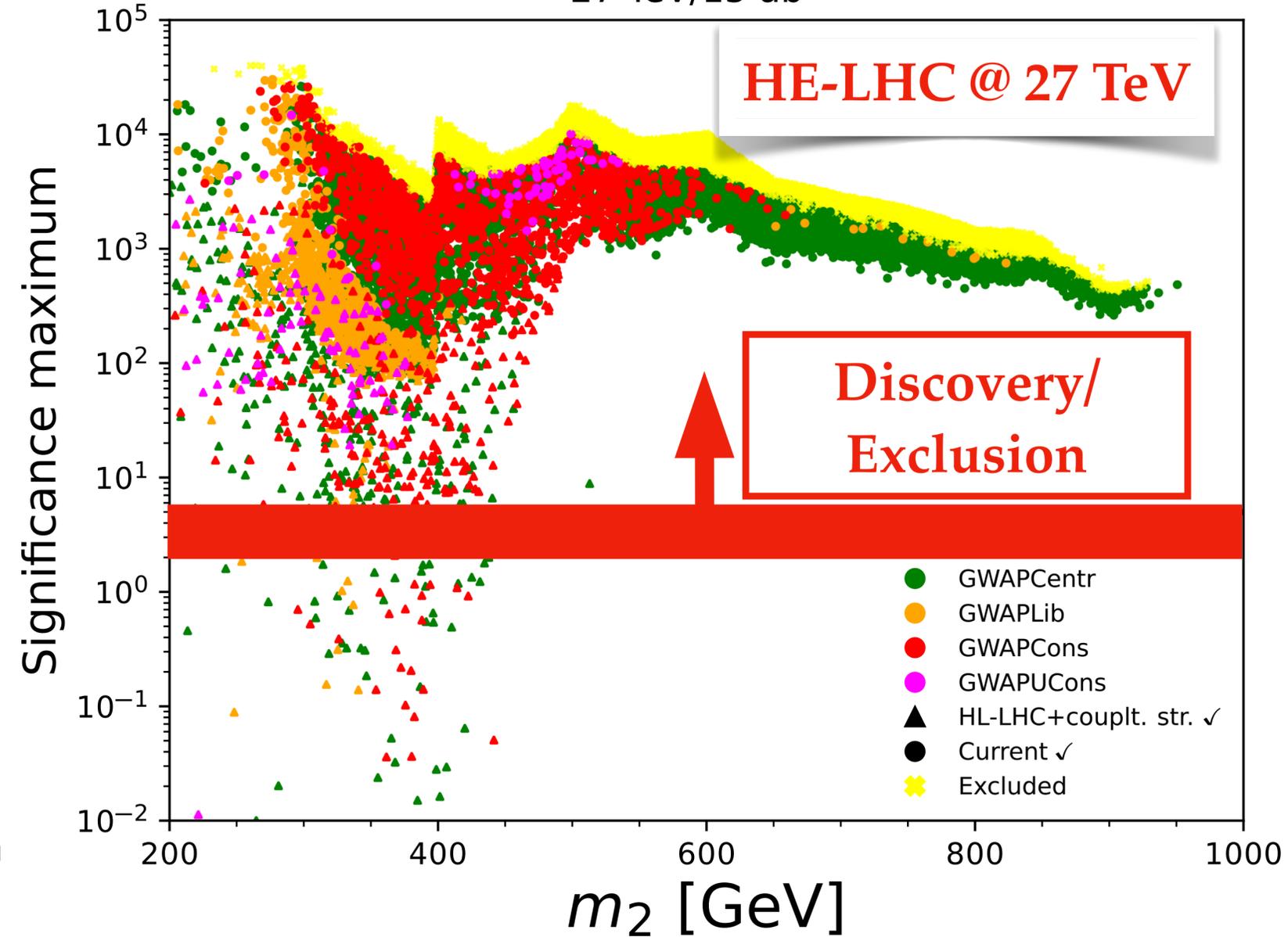
Significance @ Future Colliders

Color-coding of parameter-space points denotes **theoretical uncertainty**.

100 TeV/30 ab⁻¹



27 TeV/15 ab⁻¹

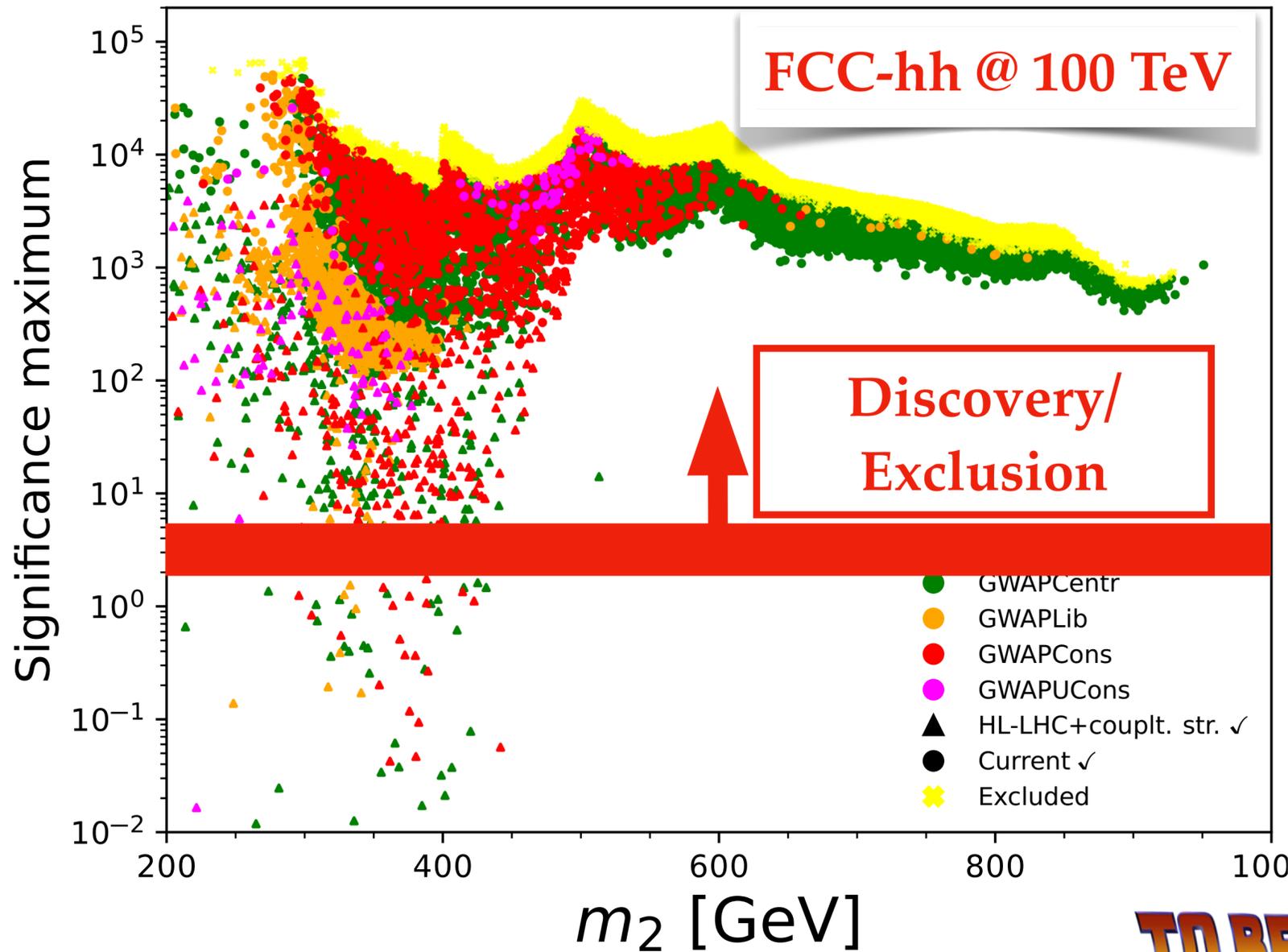


✖ = Point is excluded today.

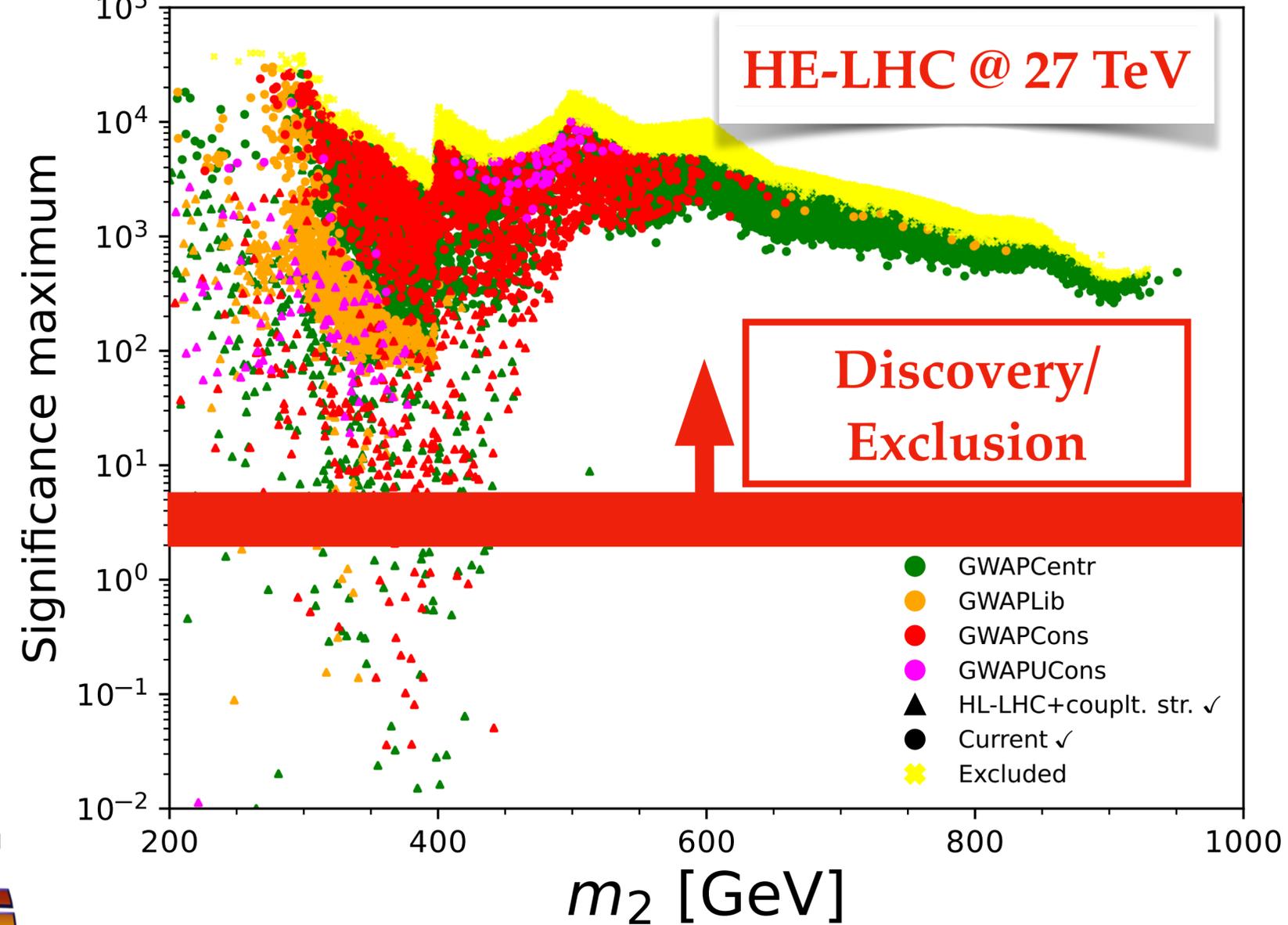
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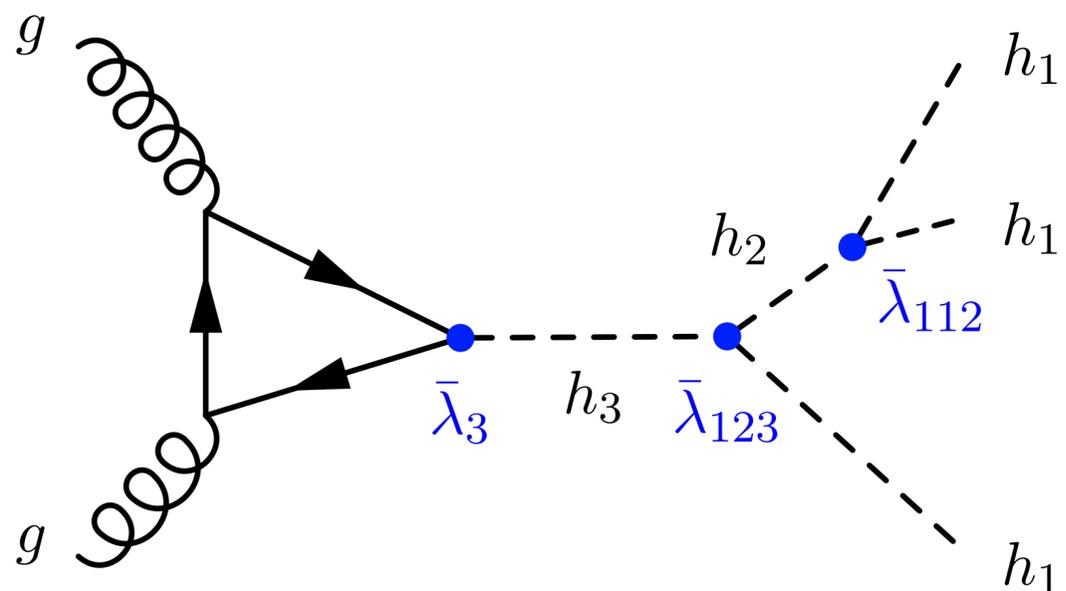


TO BE CONTINUED...

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Extending the Scalar Sector with Two Singlets

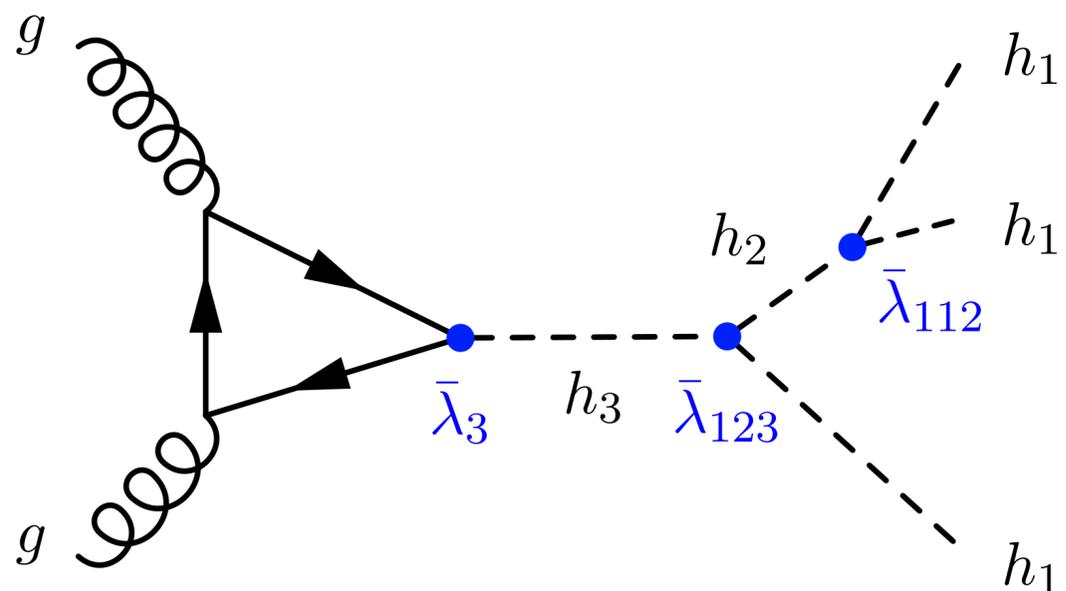
- More elaborate extensions? E.g. Two Singlet Scalar fields S_1 and S_2 .
- (+ two \mathcal{F}_2 symmetries forbidding some interactions).
- **Diagonalization** of the (3×3) mass matrix leads to **three** physical scalar states: $h_{1,2,3}$.
- **Enhanced triple Higgs boson production at the LHC!?** $> 100 \times$ **enhancement possible!** [Karkout, [AP](#), Postma, du Pree, Tetlalmatzi-Xolocotzi, van de Vis, arXiv:2404.12425]



- **Unfortunately:** Mutually exclusive with a first-order phase transition...
- **Removing the \mathcal{F}_2 restriction might help!**

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Discovery Post-Mortem

“With 4 parameters I can fit an elephant and with 5 I can make him wiggle his trunk.”
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Can we fit the shape of an elephant with 4 parameters?



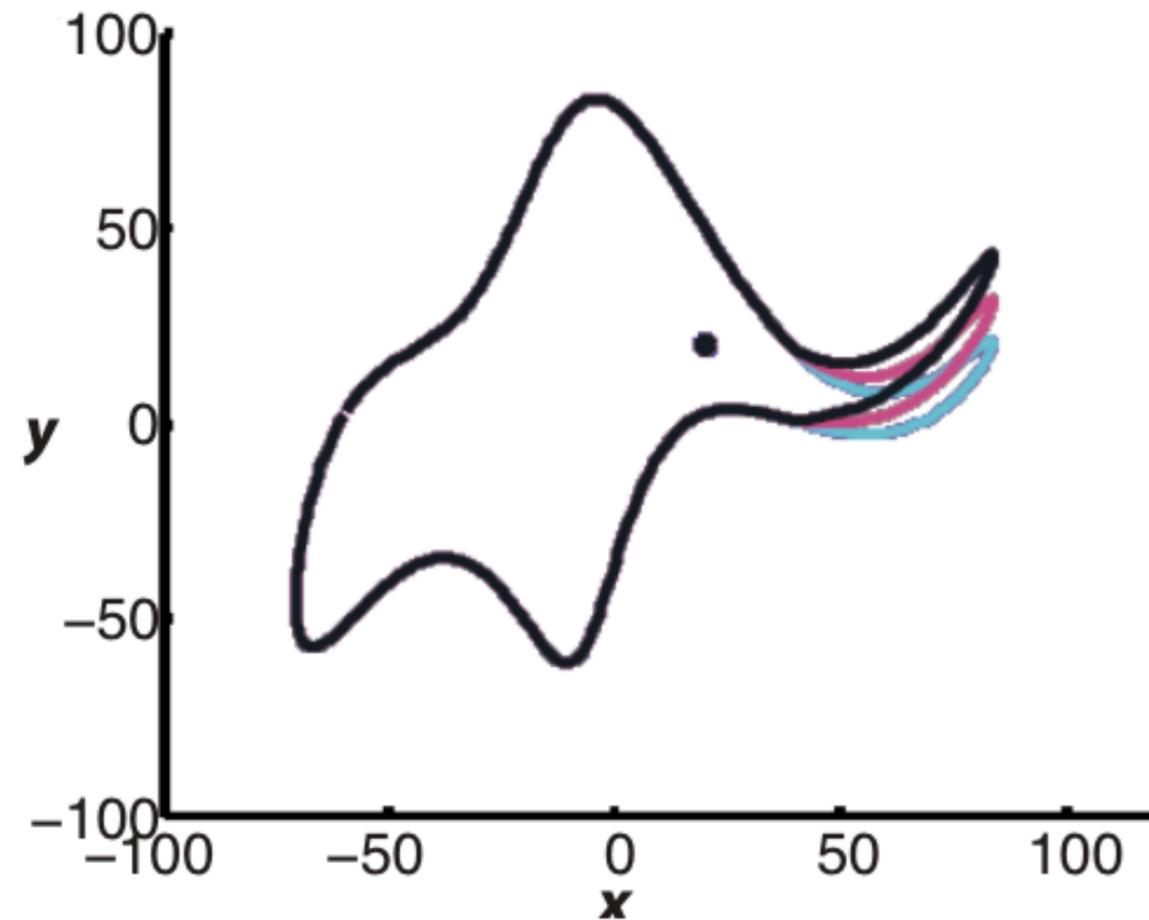
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Can we fit the shape of an elephant with 4 parameters?

→ **Yes!** With **four** complex parameters,
[and with **five** we can make it wiggle its trunk.]

[Mayer, Khairy, Howard, Am. J. Phys., Vol. 78, No. 6, June 2010]



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If we discover e.g. a new scalar particle at colliders,

→ Can we verify that it is indeed the remnant of a singlet field that generates a first-order electro-weak phase transition?



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Experimental observations → { ●, ▲, ■, ▲, ■ }

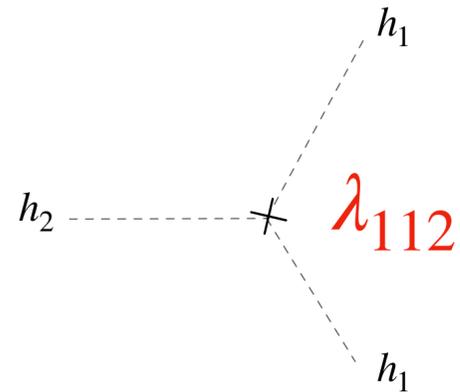


Discovery Post-Mortem Example [AP, White, arXiv:2108.11394]

Combine possible measurements:

$$pp \rightarrow h_2 \rightarrow ZZ$$

$$pp \rightarrow h_2 \rightarrow h_1 h_1$$

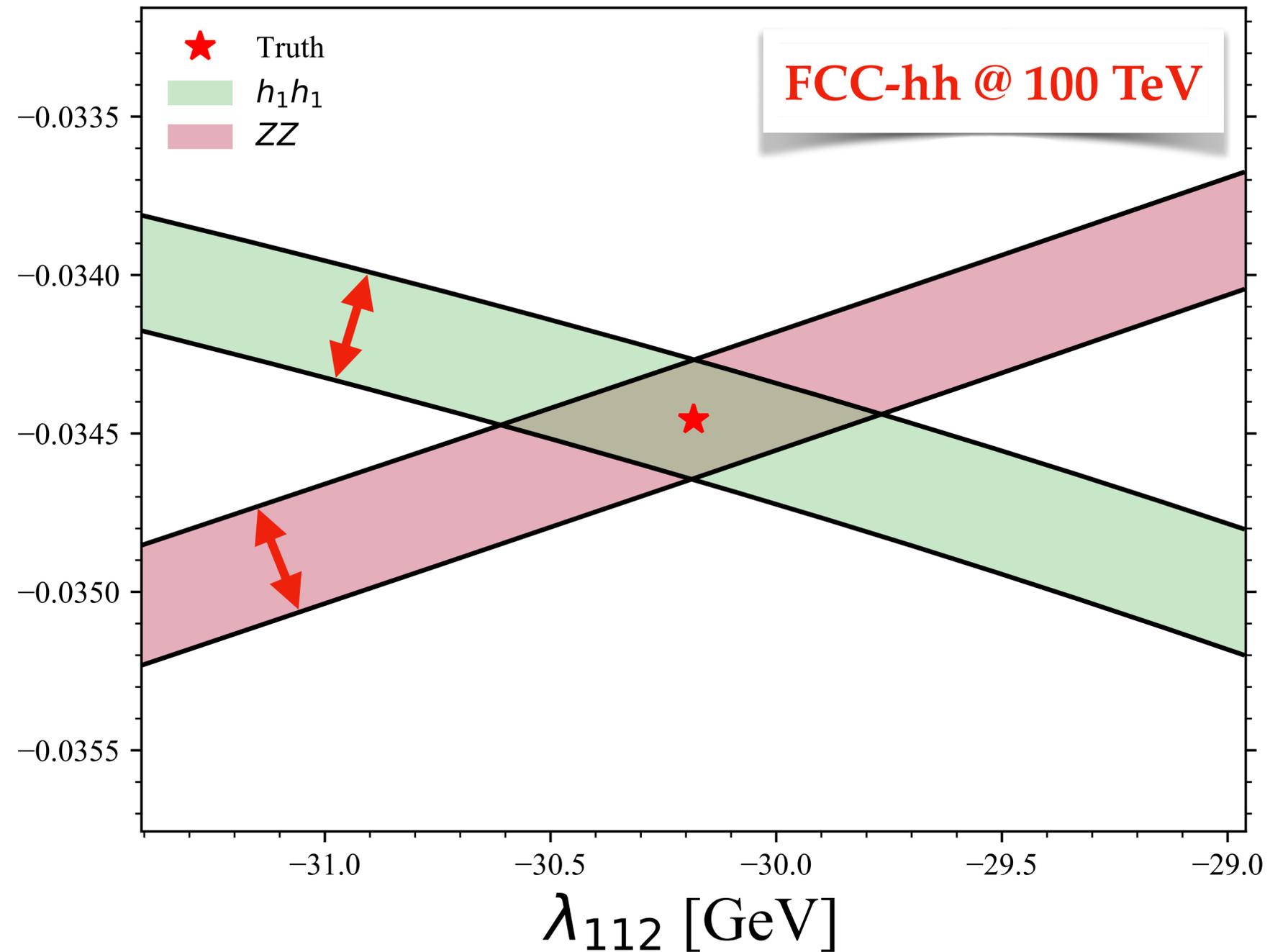


→ both functions of $\sin \theta$ & λ_{112} !

→ in turn f^{ns} of { , , , , }.

[Width of bands represents expected measurement uncertainty].

pp@100 TeV/30000 fb⁻¹, UCons1

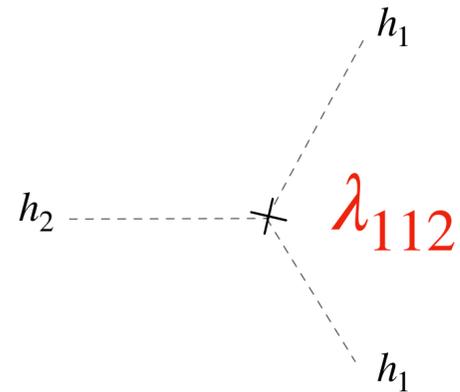


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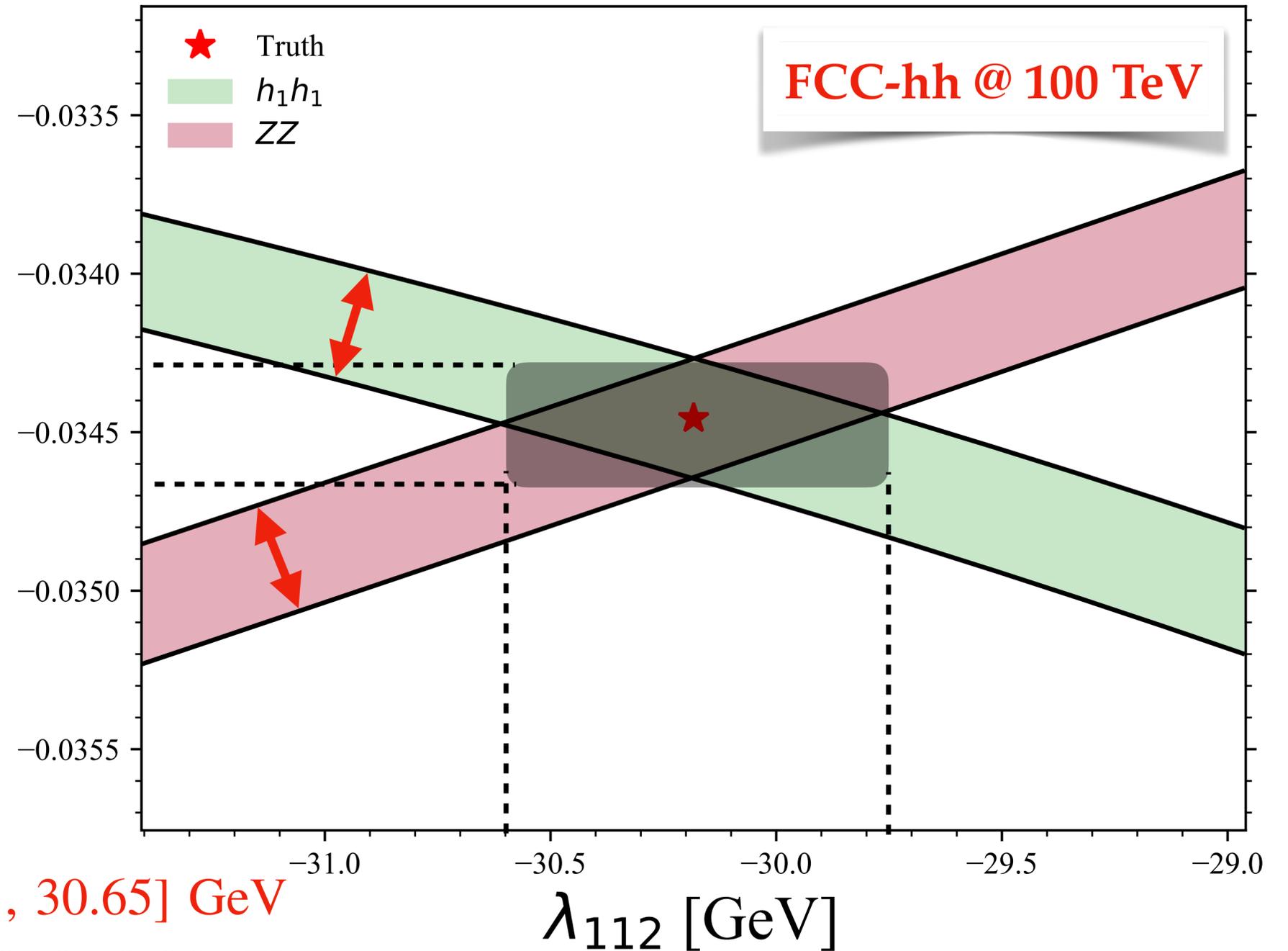


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$$\lambda_{112} \in [29.81, 30.65] \text{ GeV}$$

$$\sin \theta \in [0.0343, 0.0347]$$

The Inverse Problem in Extended Scalar Sectors:

Multi-scalar processes should play a **crucial rôle**:

$$pp \rightarrow h_1 h_2 \quad | \mathcal{M} |^2 \sim \lambda_{122}^2, \lambda_{112}^2 + \dots$$

$$pp \rightarrow h_2 h_2 \quad | \mathcal{M} |^2 \sim \lambda_{222}^2, \lambda_{122}^2 + \dots$$

$$pp \rightarrow h_1 h_1 h_1 \quad | \mathcal{M} |^2 \sim f[\lambda_{1111}, \lambda_{1112}, \lambda_{111}, \lambda_{112}]$$

[...]

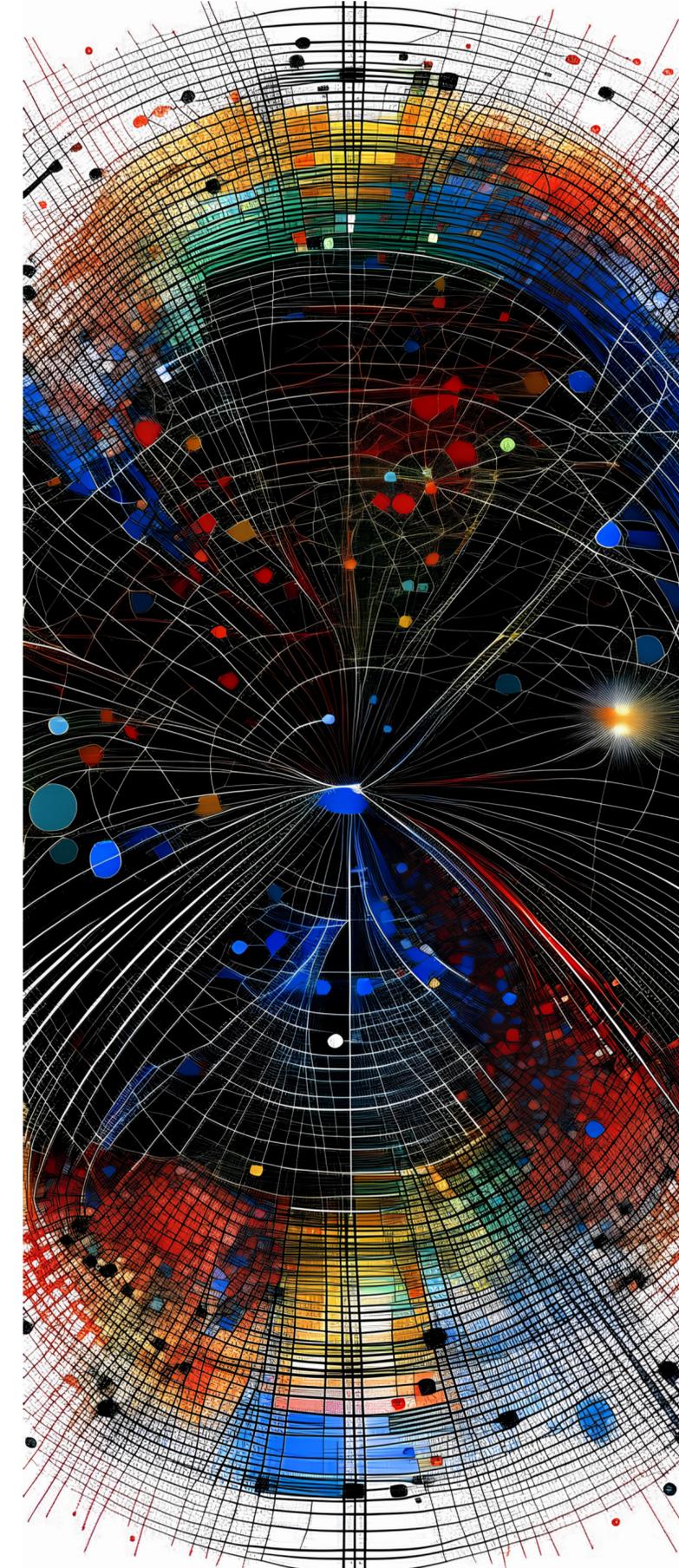


**TO BE
CONTINUED** 



Summary

- Within the realm of the Standard Model, we can describe a whole of range of particle interactions with extraordinary precision.
 - **Monte Carlo Event Generators** are invaluable tools of communication between theory and experiment.
- However, there remains a multitude of **open questions**, that may be linked via the **Electro-Weak Phase Transition**.
- The **Nature of the Electro-Weak Phase Transition** is an important scientific enquiry.
 - **(Strong) First-Order EWPT [not in SM!] → Matter-Anti-Matter asymmetry.**
 - **Extending the scalar sector of the SM can be the necessary catalyst!**
 - **Future particle colliders have the potential to probe this mechanism.**



Outlook

- Following **any** discovery, **solving the inverse problem** would be the **crucial** next step.
- I discussed possible **first steps in this exercise**, following the discovery of a new scalar.
- **Multi-scalar production processes** (e.g. triple Higgs boson production) will play a crucial role in this endeavor.
- [Could they also be **discovery channels** themselves?]
- Questions merit investigation both at the LHC and other future colliders (e.g. FCC, Muon Collider, ...).



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Thanks!
Questions?

Appendix



Theoretical Uncertainties in SM+Singlet

- **Uncertainties** → Can affect e.g. the strength of the transition, $\langle \phi(T) \rangle / T$.
- Due to:
 - **gauge dependence**, [Patel, Ramsey-Musolf, arXiv:1101.4665]
 - **scale dependence** → Linde's IR problem: expansion parameter is $gn_B \sim gT/m$, (n_B mode occupation), diverges as $m \rightarrow 0 \Rightarrow$ perturbativity breaks down.
[Linde, Phys. Lett. 96B (1980) 289.]
- \Rightarrow **To make reliable and sensible statements on colliders prospects:**
 - **Crucial to take uncertainties into account.**

Theoretical Uncertainty Bands

- Define “uncertainty band” by:

1. Deriving 1-loop effective potential in the covariant gauge,

[Arnold, Espinosa, hep-ph/9212235], [Andreassen, MSc, Norwegian U. Sci. Tech., 2013]

2. Run couplings $\lambda \rightarrow \lambda(\mu)$, μ is RGE scale, [SARAH, Staub, arXiv:0806.0538]

3. Scan parameter space of Lagrangian,

4. Vary $\mu \in [\frac{1}{2} \times m_Z, 5 \times m_Z]$ & gauge params. $\xi_i \in [0,3] \rightarrow$ band of 8 pts.

5. Use **PhaseTracer** for each point in band \rightarrow Get phase transitions, $\langle \phi(T_c) \rangle / T_c$.

[Athron, Balázs, Fowlie, Zhang, arXiv:2003.02859]

Parameter-space Categories

1. Define two conditions:
 - i. VEV at 1-loop: $\langle \phi(T = 0) \rangle = 246 \pm 30 \text{ GeV}$ & deepest minimum.
 - ii. $\langle \phi(T_c) \rangle / T_c > 1$ & no other transition with higher T_c .
2. Define four mutually-exclusive categories*:

*An alternative classification appears in our article: see Appendix.

SFO-EWPT more certain



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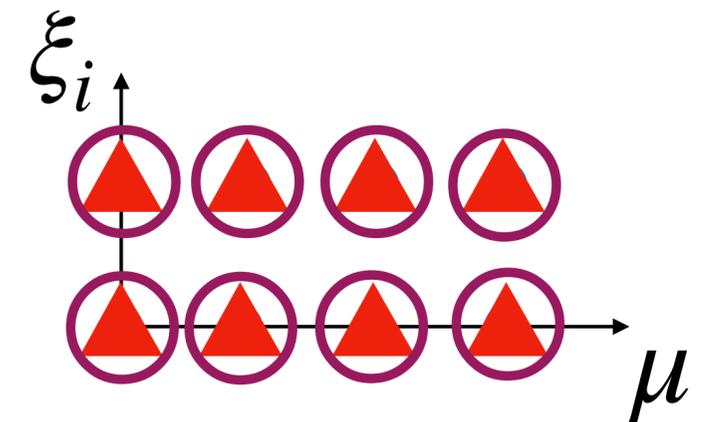
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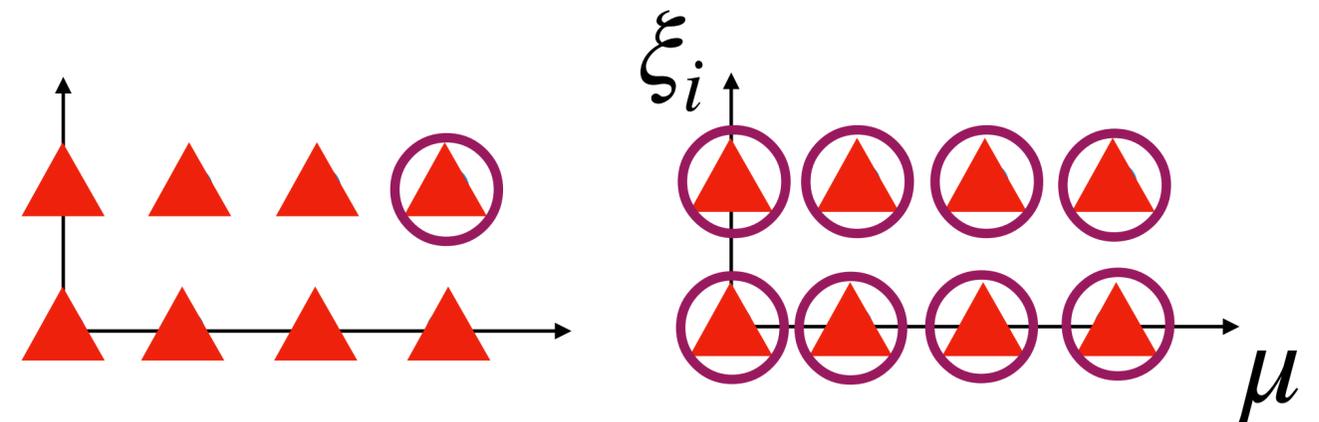
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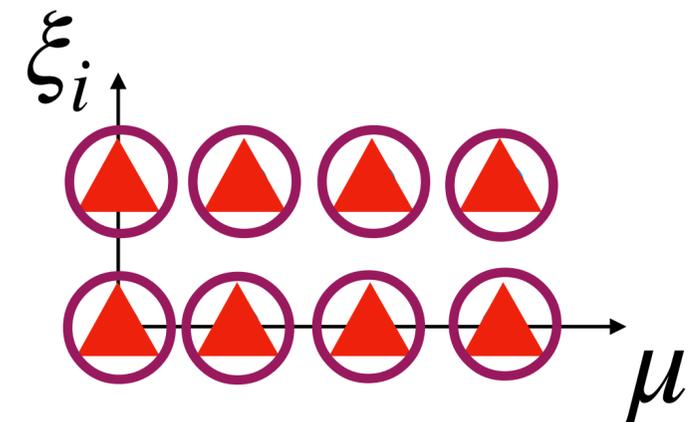
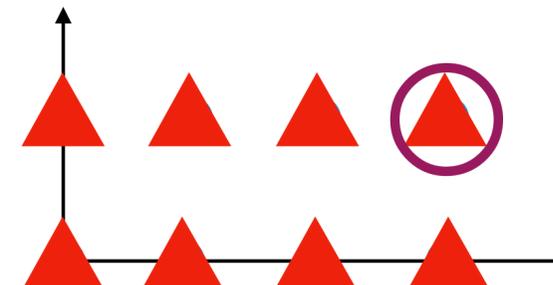
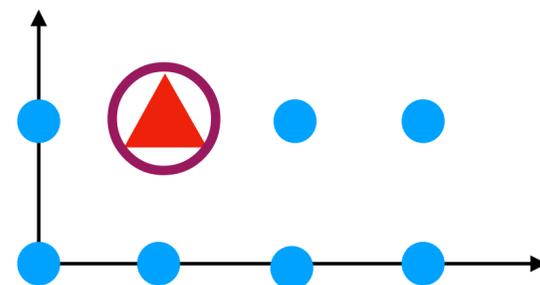
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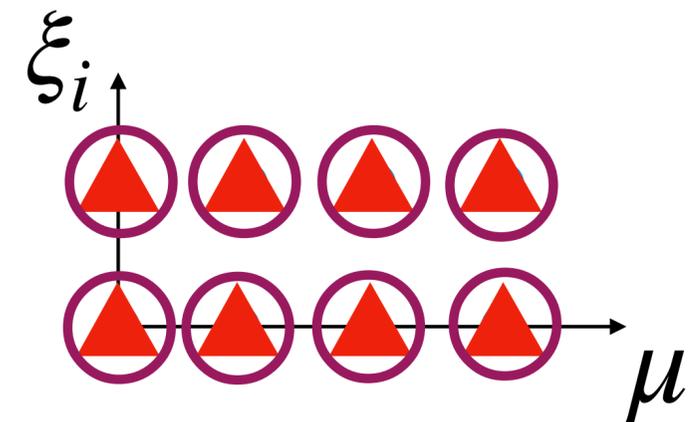
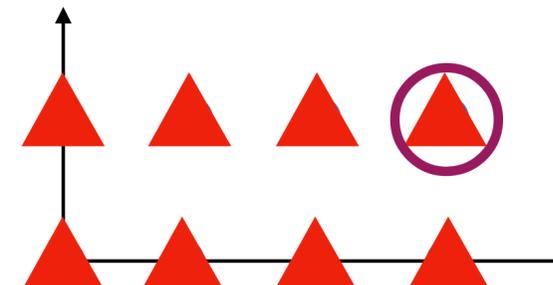
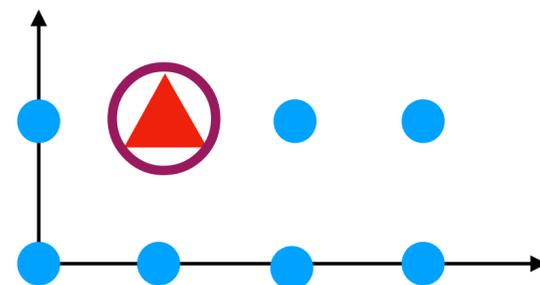
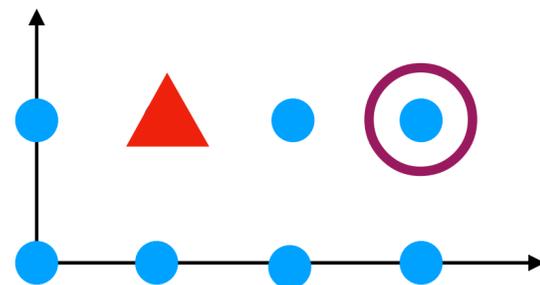
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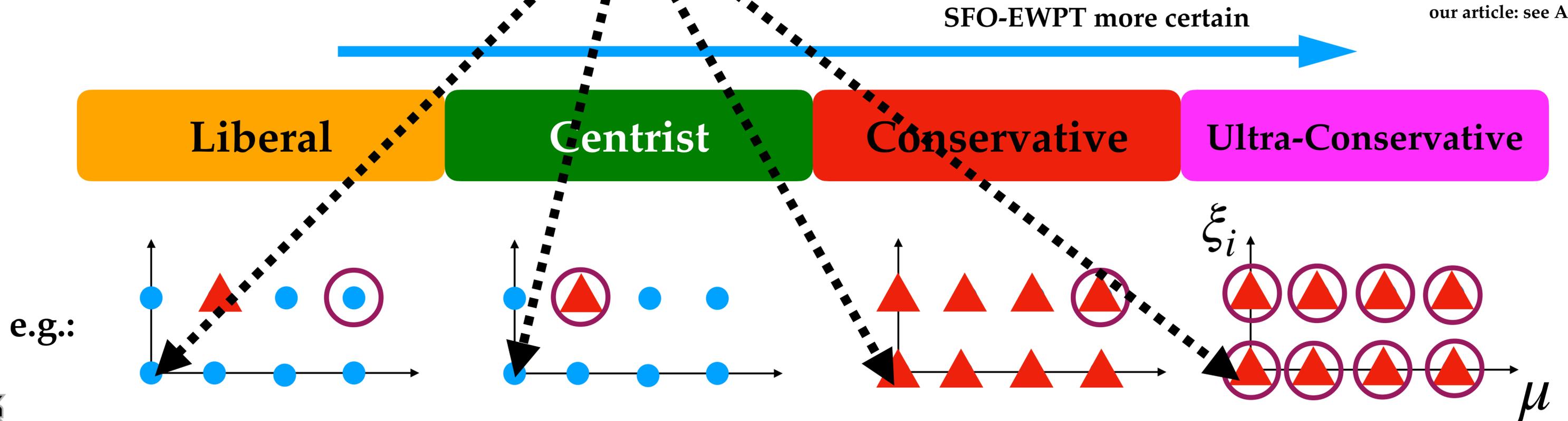
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1. Define two conditions:

Note: For phenomenological analyses, take “central” μ and ξ_i .

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Electro-Weak Precision Observables

- Real singlet scalar field
 - \rightarrow modifies Higgs contributions to diagonal weak gauge boson vacuum polarisation diagrams,
 - & introduces additional contributions.
- Quantify via S, T, U parameters. [Hagiwara, Matsumoto, Haidt, Kim, hep-ph/9409380]
- Change in EWPO \mathcal{O} ($= S, T, U$):

$$\Delta\mathcal{O} = (\mathcal{O}(m_2^2) - \mathcal{O}(m_1^2)) \times \sin^2 \theta$$

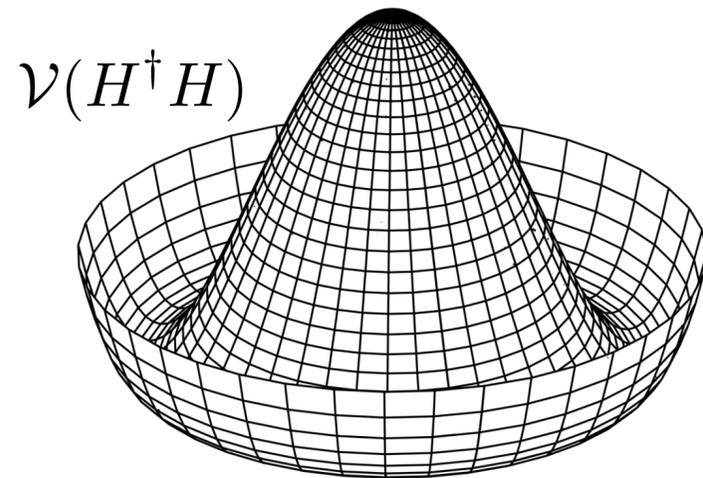
\Rightarrow calculate compatibility with experimental measurement $\Delta\mathcal{O}^{\text{EXP}}$.

The Higgs Potential & Vacuum Stability

The Importance of the Higgs sector

- the Higgs boson: the central protagonist of EWSB:

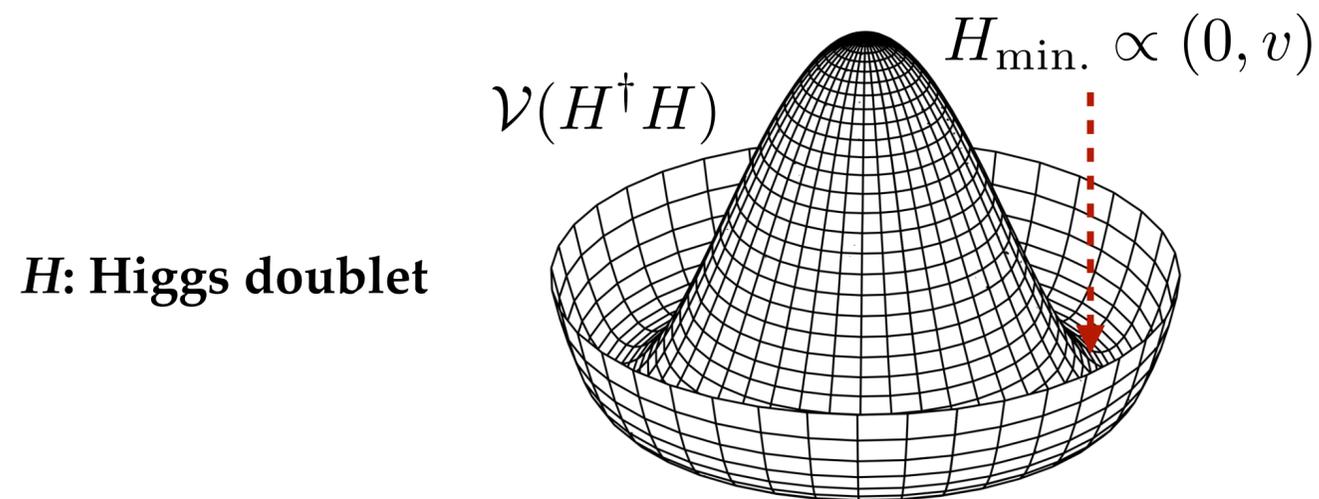
H: Higgs doublet



- an important characteristic of the Higgs boson is the way it couples to itself:

The Importance of the Higgs sector

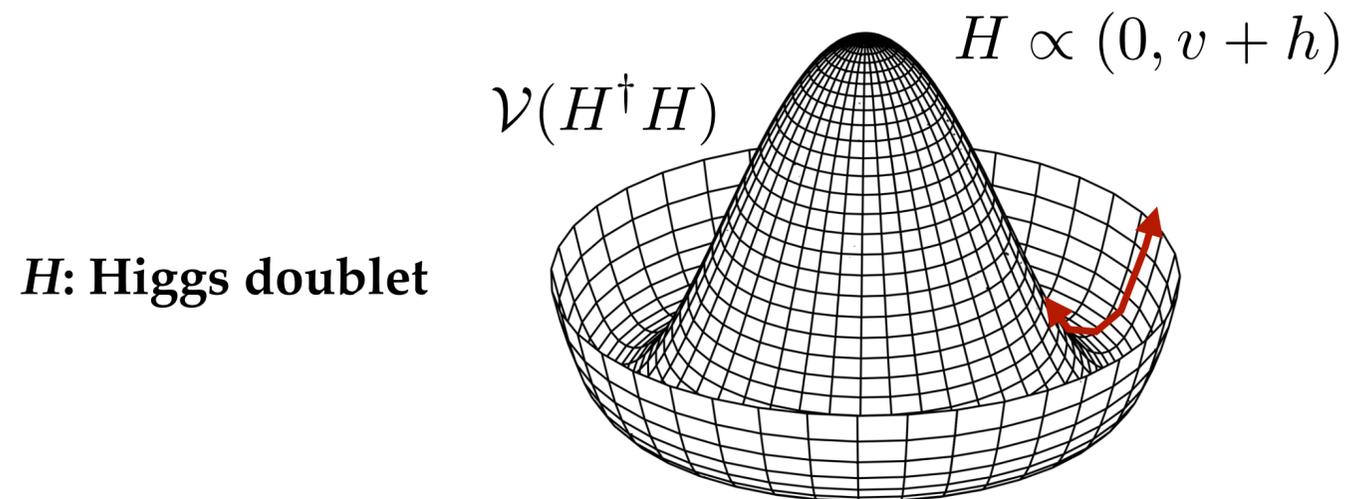
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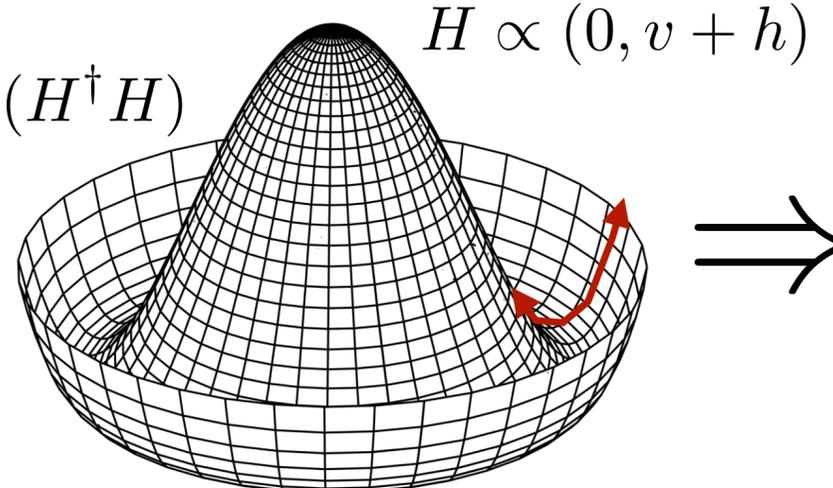
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$\mathcal{V}(H^\dagger H)$

$H \propto (0, v + h)$



e.g. fermion masses
& interactions:

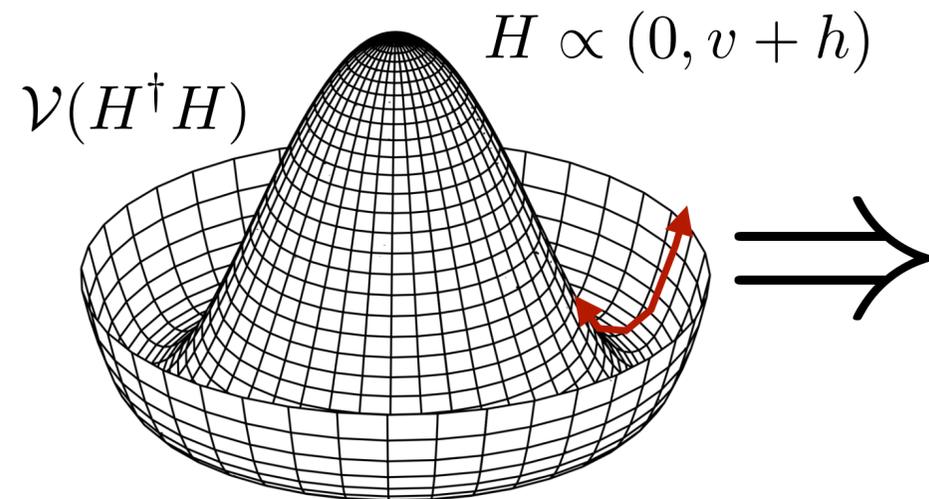
$$\mathcal{L} \supset -m_f \bar{f}_L f_R - \frac{m_f}{v} h \bar{f}_L f_R + \text{h.c.}$$

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The Importance of the Higgs sector

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e.g. gauge boson masses
& interactions:

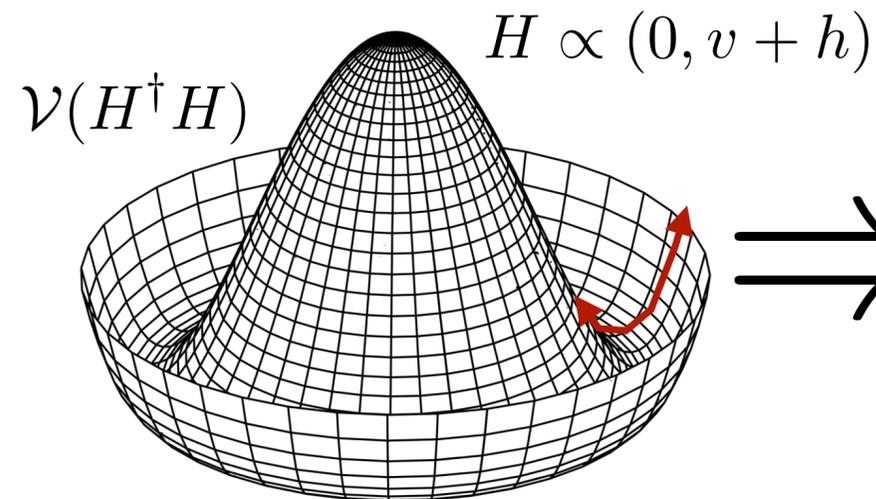
$$\mathcal{L} \supset [m_W^2 W^{\mu+} W_{\mu}^- + \frac{1}{2} m_Z^2 Z^{\mu} Z_{\mu}] \times \left(1 + \frac{h}{v}\right)^2$$

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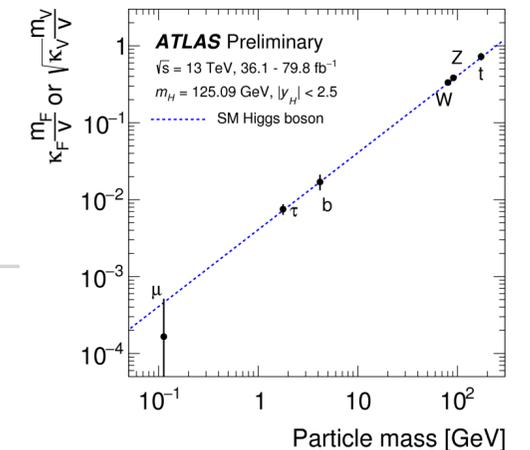
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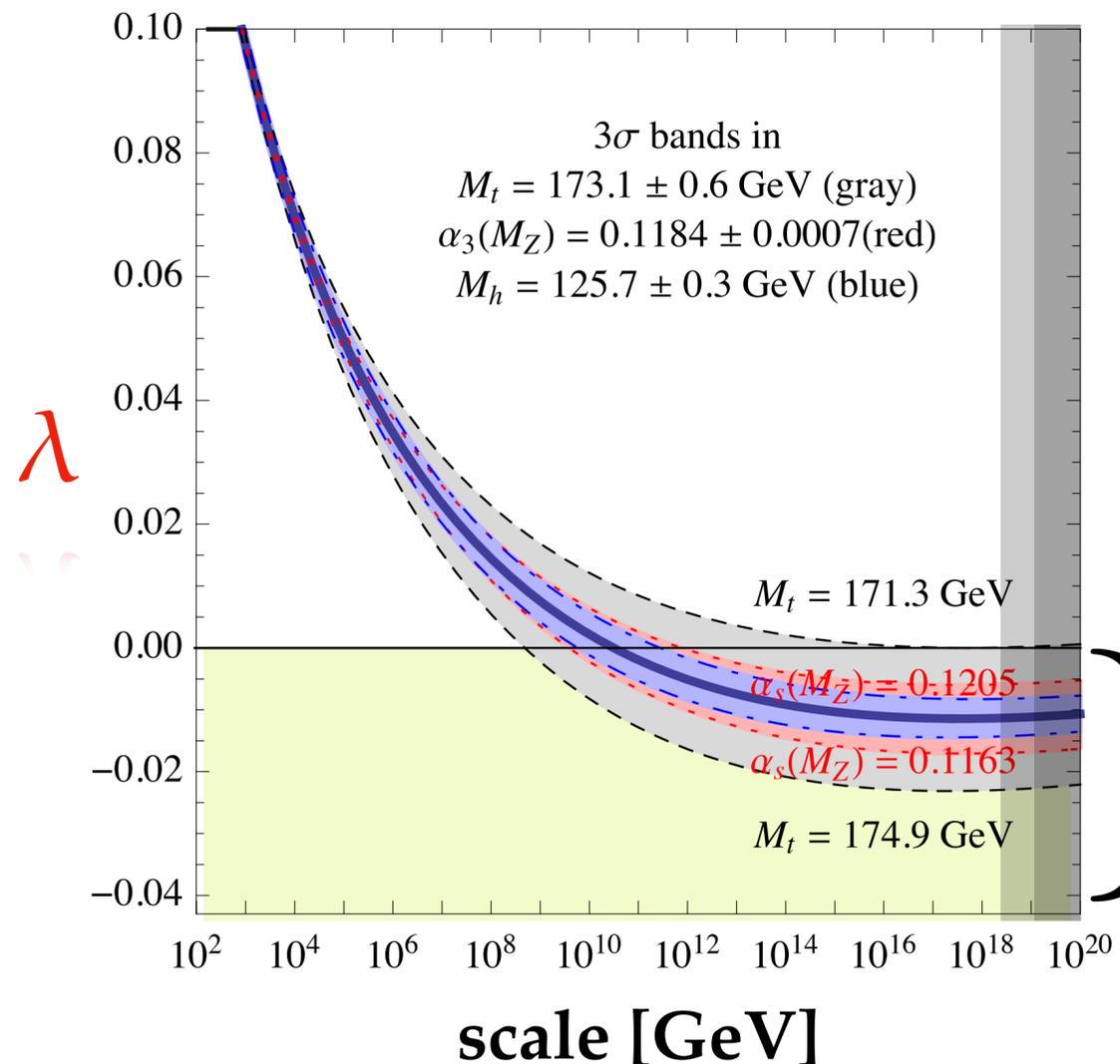
- an important characteristic of the Higgs boson is the way it couples to itself:

Vacuum Stability

- SM potential for the Higgs doublet:

$$\mathcal{V}(H^\dagger H) = -m^2(H^\dagger H) + \lambda(H^\dagger H)^2$$

- renormalisation group evolution of the coupling λ :

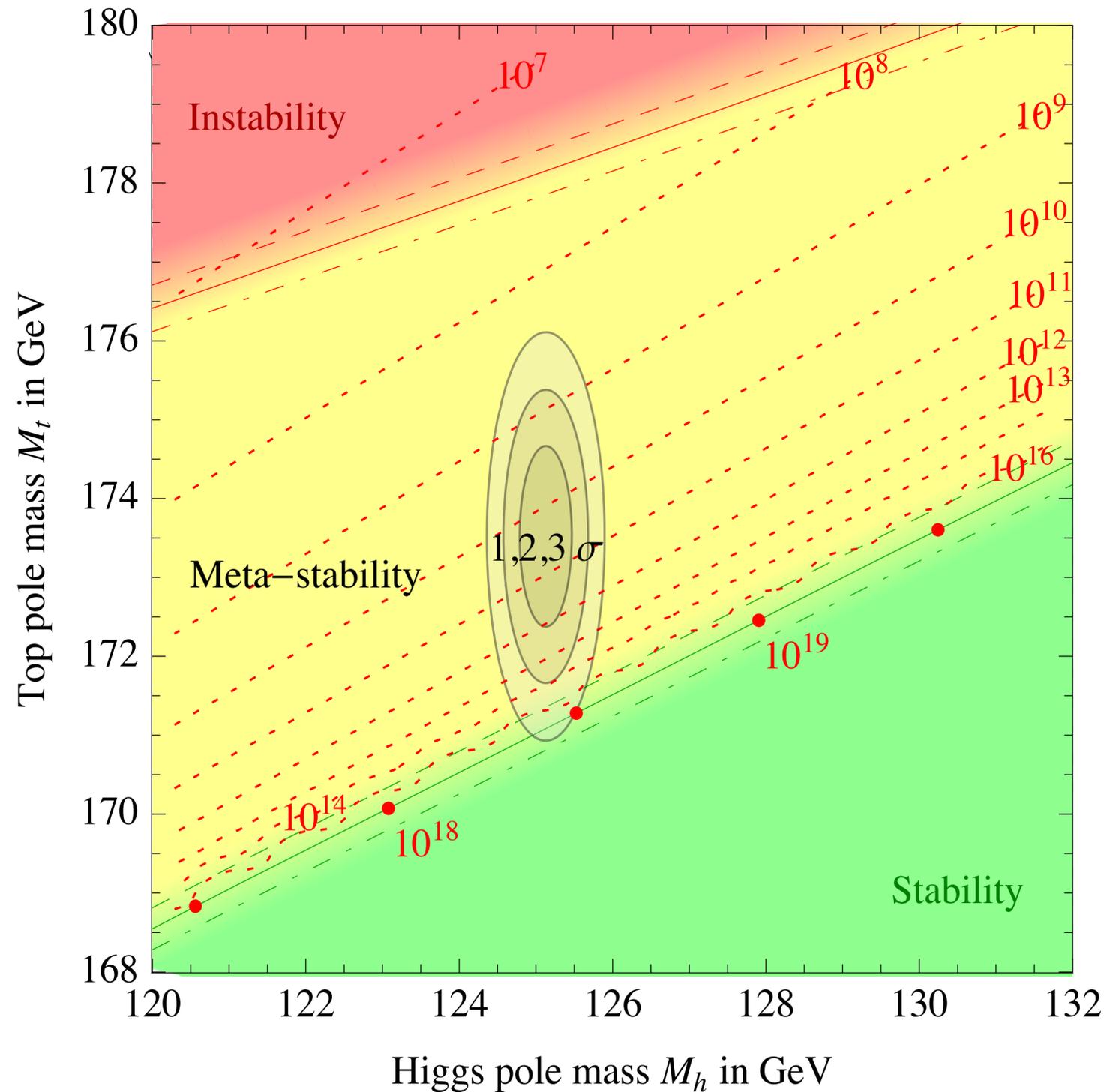


[Cabibbo, Maiani, Parisi, Petronzio, 1979, Hung, 1979, ..., Degrassi, Di Vita, Elias-Miró, Giudice, Isidori, Strumia, 1205.6497, Buttazzo, Degrassi, Giardino, Giudice, Sala, Salvio, Strumia, 1307.3536 ..., Espinosa, 1512.01222]

**potentially
unstable or
meta-stable
vacuum!**

vacuum stability

[Buttazzo, Degrassi, Giardino, Giudice, Sala, Salvio, Strumia, 1307.3536, Espinosa, 1512.01222]



in deriving this: **assumed Standard Model.**

➡ a hint for a non-standard Higgs sector?

➡ further investigation necessary.

Sphaleron/Instanton Processes

Instantons and Baryon-# Violation

- toy model:

(1+1)-dimensions, Abelian gauge field A^μ , complex scalar Φ , Dirac fermion of unit charge Ψ .

- Euclidean space action:

$$S = \int d^2x \left[\frac{1}{4} F_{\mu\nu}^2 + |(\partial_\mu - ieA_\mu)\Phi|^2 + V(\Phi) + i\bar{\Psi}(\partial_\mu - ieA_\mu)\gamma^\mu\Psi \right]$$

“Higgs potential”: $V(\Phi) = \lambda(\Phi^*\Phi - v^2)^2 \implies$ “EWSB” $\implies M_A, M_h$

Instantons and Baryon-# Violation

- consider the current:

$$K_\mu = \frac{e}{2\pi} \epsilon_{\mu\nu} A_\nu$$

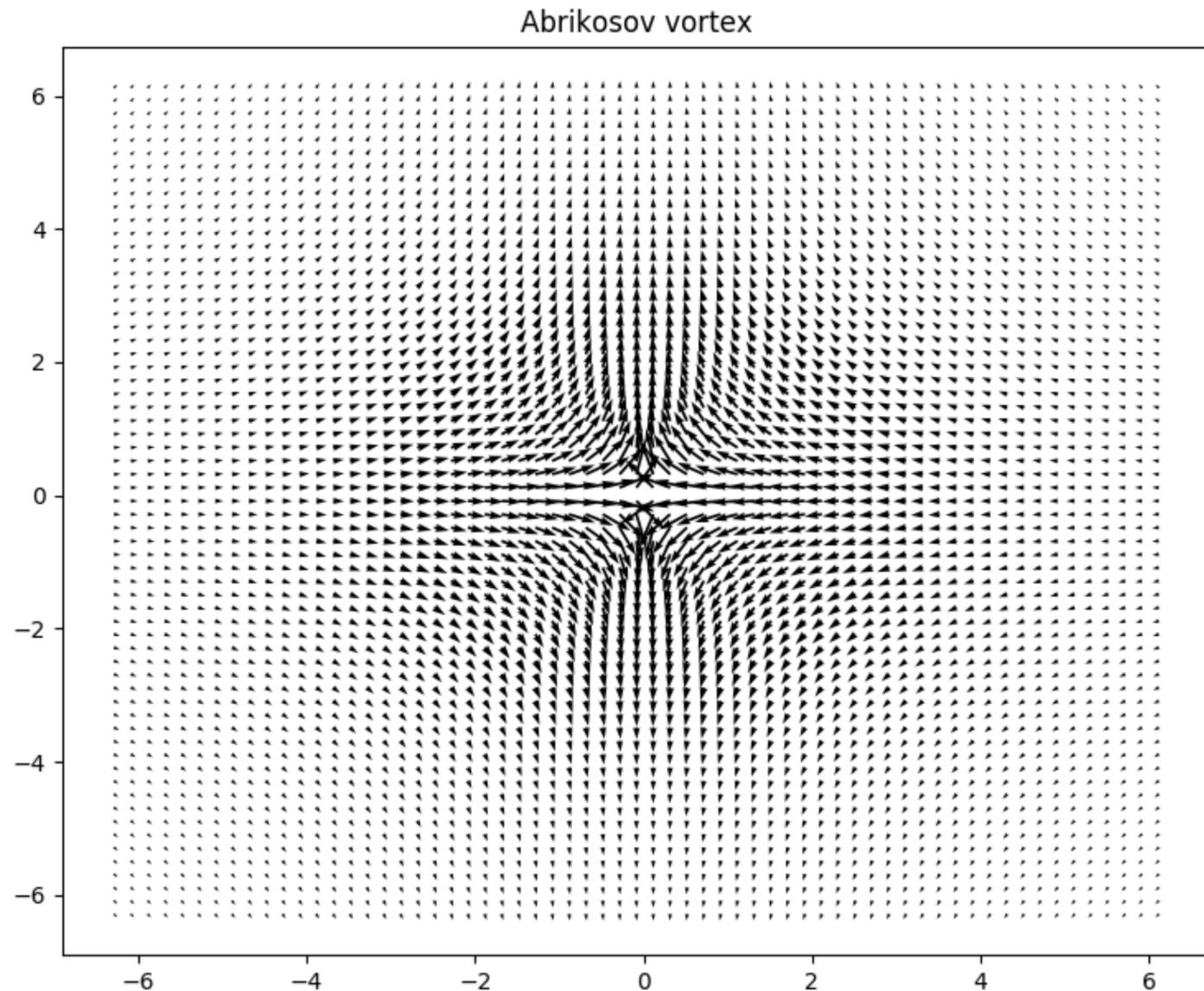
- corresponds to “charge density”:

$$N_{\text{CS}} = \int dx K_0 = \frac{e}{2\pi} \int dx A_1$$

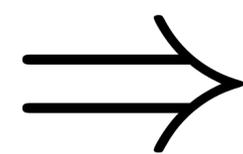
- known as the “winding” or “Chern-Simons” number.

Instantons and Baryon-# Violation

- a classical solution to equations of motion is the “Abrikosov vortex”:



$$A_r = 0, \quad A_\theta = \frac{1}{er} f(r),$$
$$f(0) = 0, \quad 1 - f(r) \sim e^{-M_A r}$$



changes the Chern-Simons number by one unit:

$$\Delta N_{\text{CS}} = \int d^2x \partial_\mu K^\mu = \frac{e}{4\pi} \int d^2x \epsilon_{\mu\nu} F^{\mu\nu} = 1$$

Instantons and Baryon-# Violation

- “instanton” transition necessarily accompanied by change of chirality of fermions by two units:

$$j_{\mu}^5 = \bar{\Psi} \gamma_{\mu} \gamma_5 \Psi$$

$$\frac{1}{2} \partial_{\mu} j^{5\mu} = \frac{e}{4\pi} \epsilon_{\mu\nu} F^{\mu\nu} \quad \text{anomalous divergence of the axial-vector current.}$$

$$\Delta N_{\text{CS}} = \int d^2x \partial_{\mu} K^{\mu} = \frac{e}{4\pi} \int d^2x \epsilon_{\mu\nu} F^{\mu\nu} = 1$$

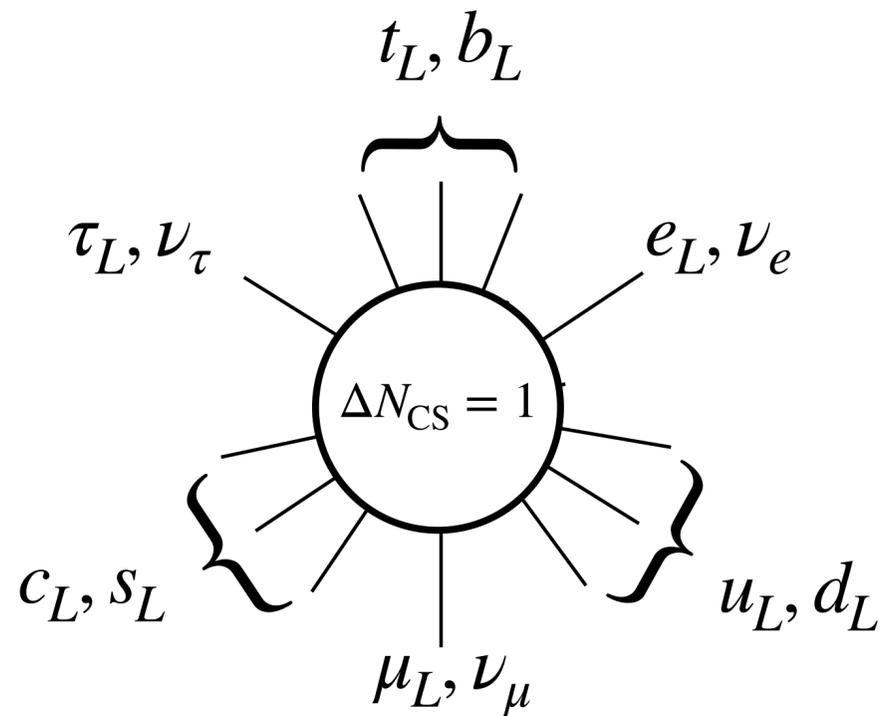
$$\Rightarrow \Delta Q_5 = \Delta \int dx j_0^5 = 2$$

EW Sphalerons at colliders?

- Rate and observability of sphaleron processes at colliders debated.

e.g. [Bezrukov, Levkov, Rebbi, Rybakov, Tinyakov, hep-ph/0304180] VS. [Tye, Wong, 1505.0360, 1710.07223].

- Ponder: **Sphaleron-induced interactions at hadron colliders:**

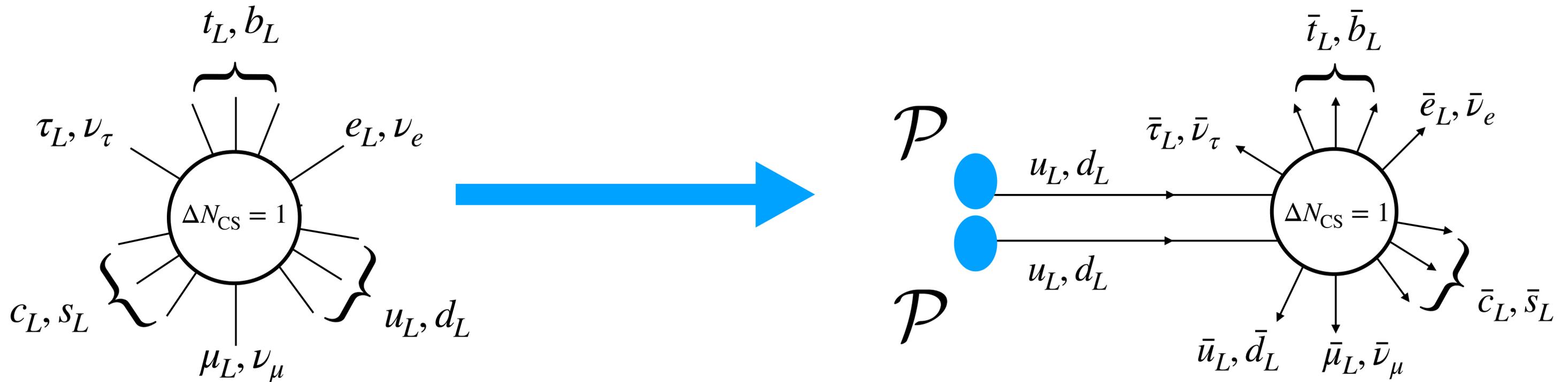


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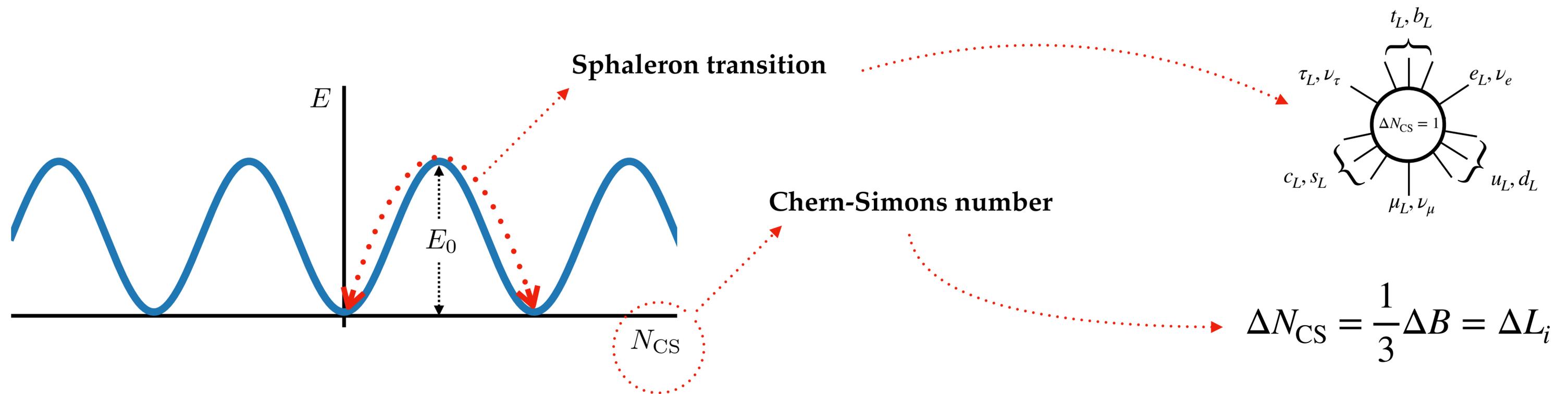
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What is the Sphaleron?



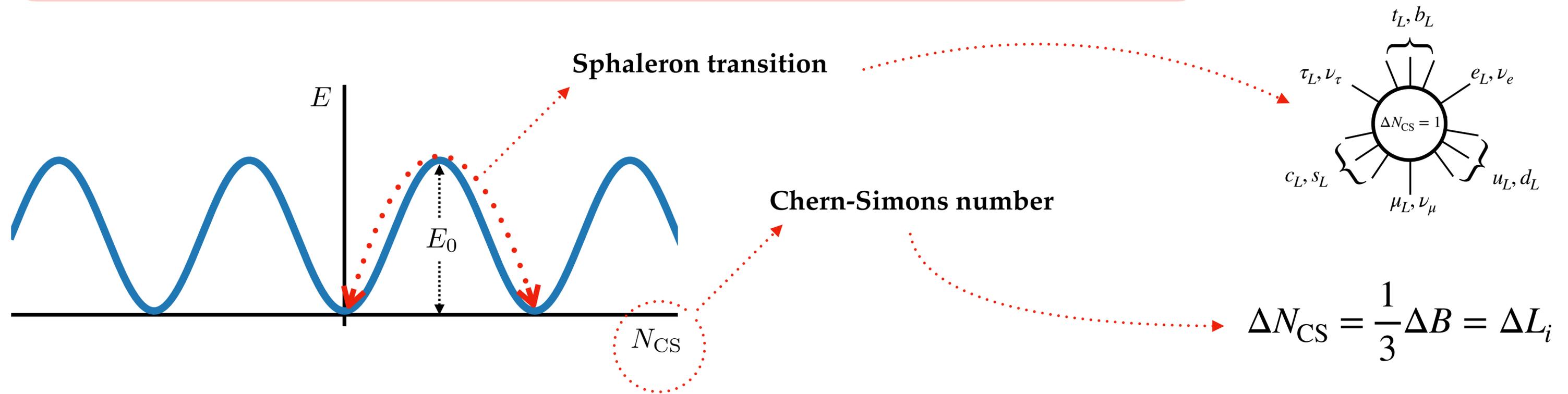
- $E_0 \sim \mathcal{O}(10)$ TeV, separates degenerate Electro-Weak vacua.



What is the Sphaleron?



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The Sphaleron energy depends crucially on the Higgs sector!

A Note on Sphaleron Suppression

- **Suppression of sphaleron rate inside bubble**

⇒ Baryon Asymmetry “swept in” broken phase and “frozen in”.

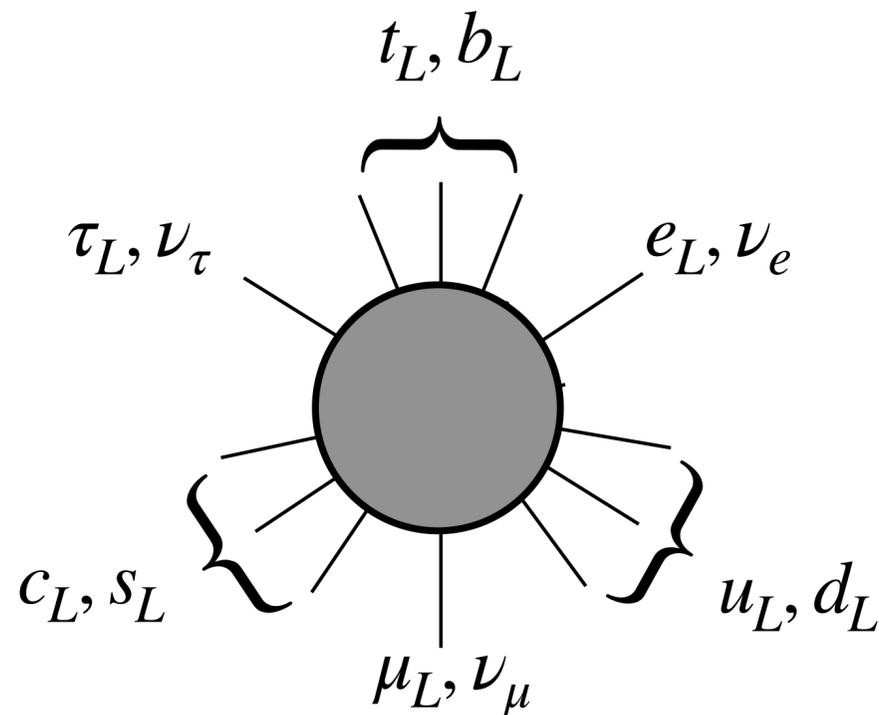
- **Rate** $\sim \exp[-\langle \phi(T_C) \rangle / T_C \times \dots]$,

[T_C : the critical temperature.]

- ⇒ Require: $\langle \phi(T_C) \rangle / T_C \geq 1 \Rightarrow$ a “**Strong**” First-Order EWPT (**SFO-EWPT**).

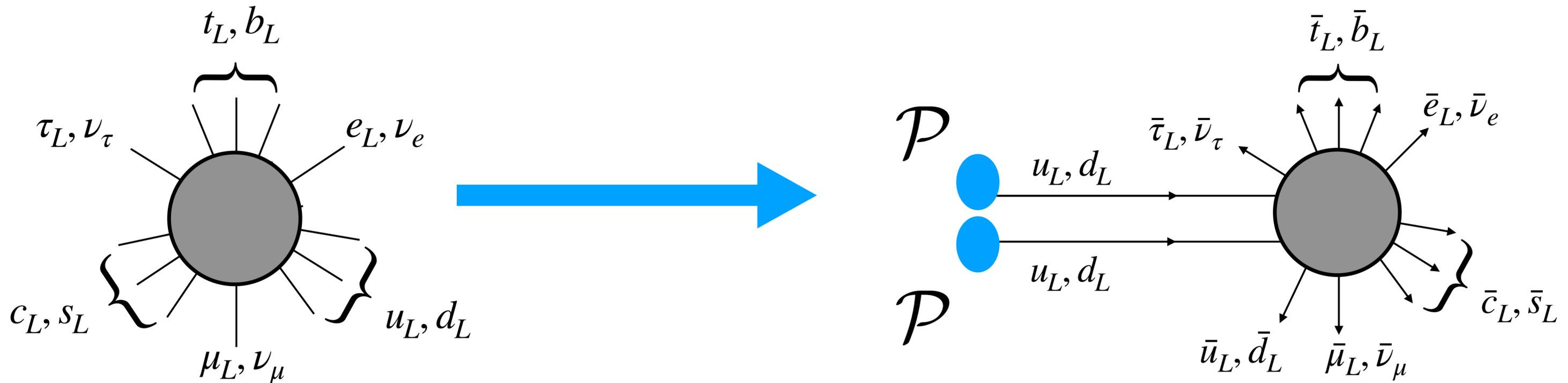
Sphaleron Suppression

- Inside the bubble: ~~Sphaleron~~
- Suppression requires “**Strong**” First-Order EWPT (**SFO-EWPT**).
- **Despite suppression:** Sphalerons @ colliders?

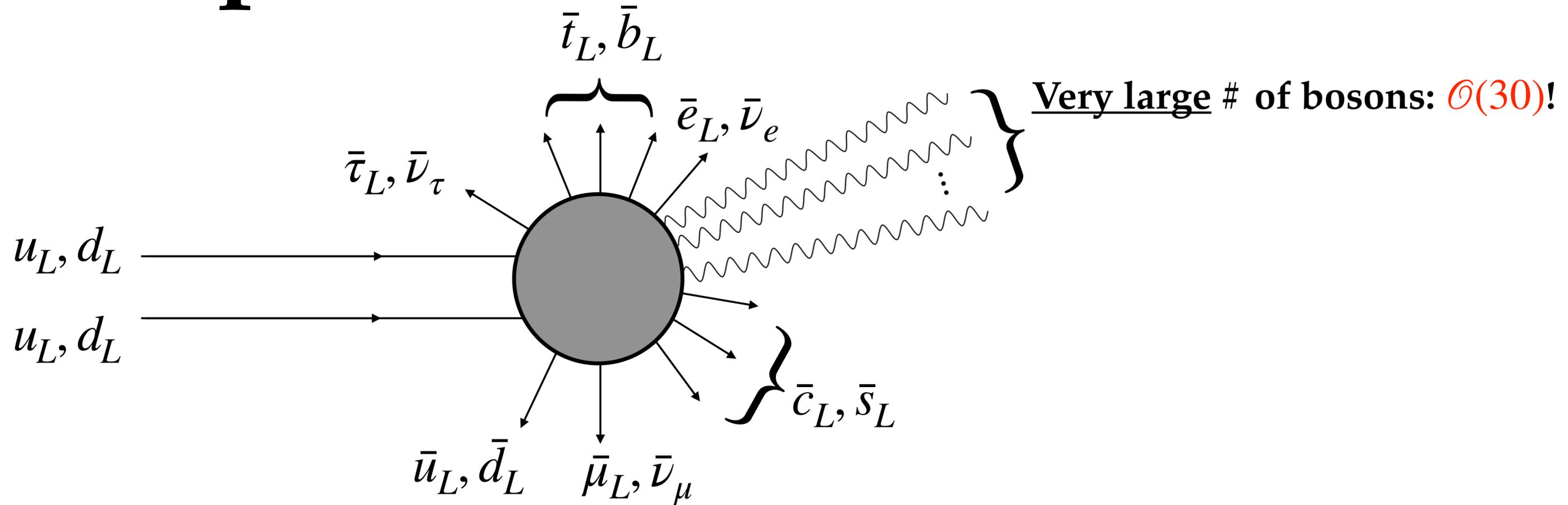


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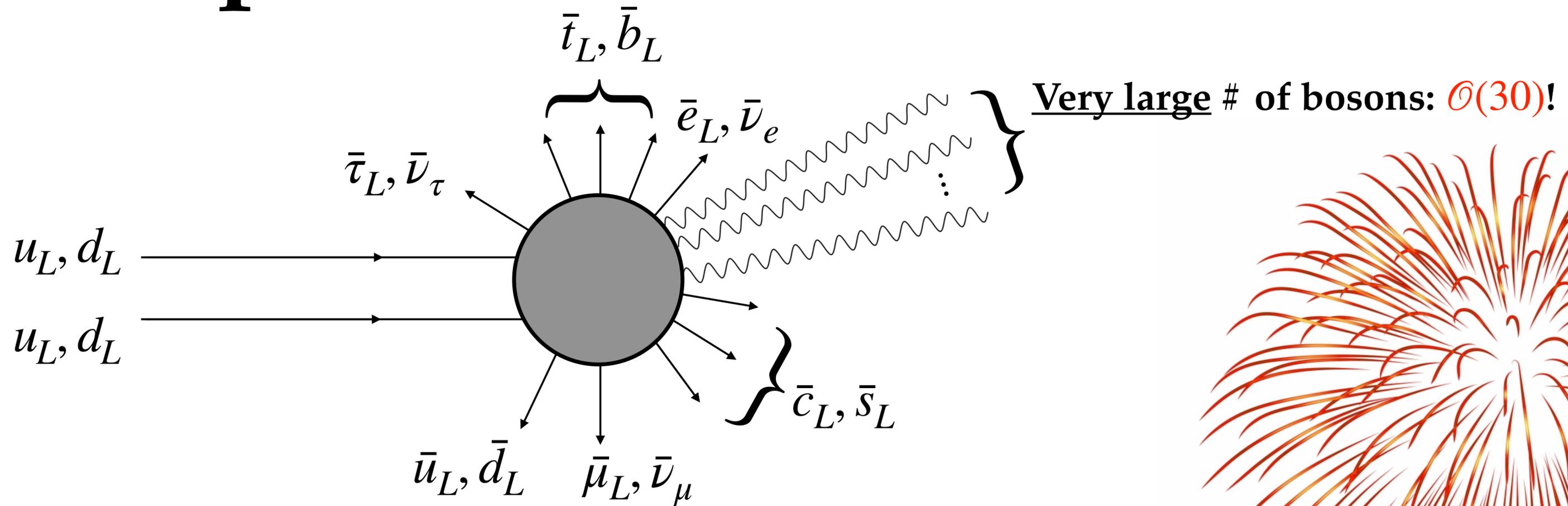
EW Sphalerons at Colliders?



- Possible enhancement if large number of bosons,

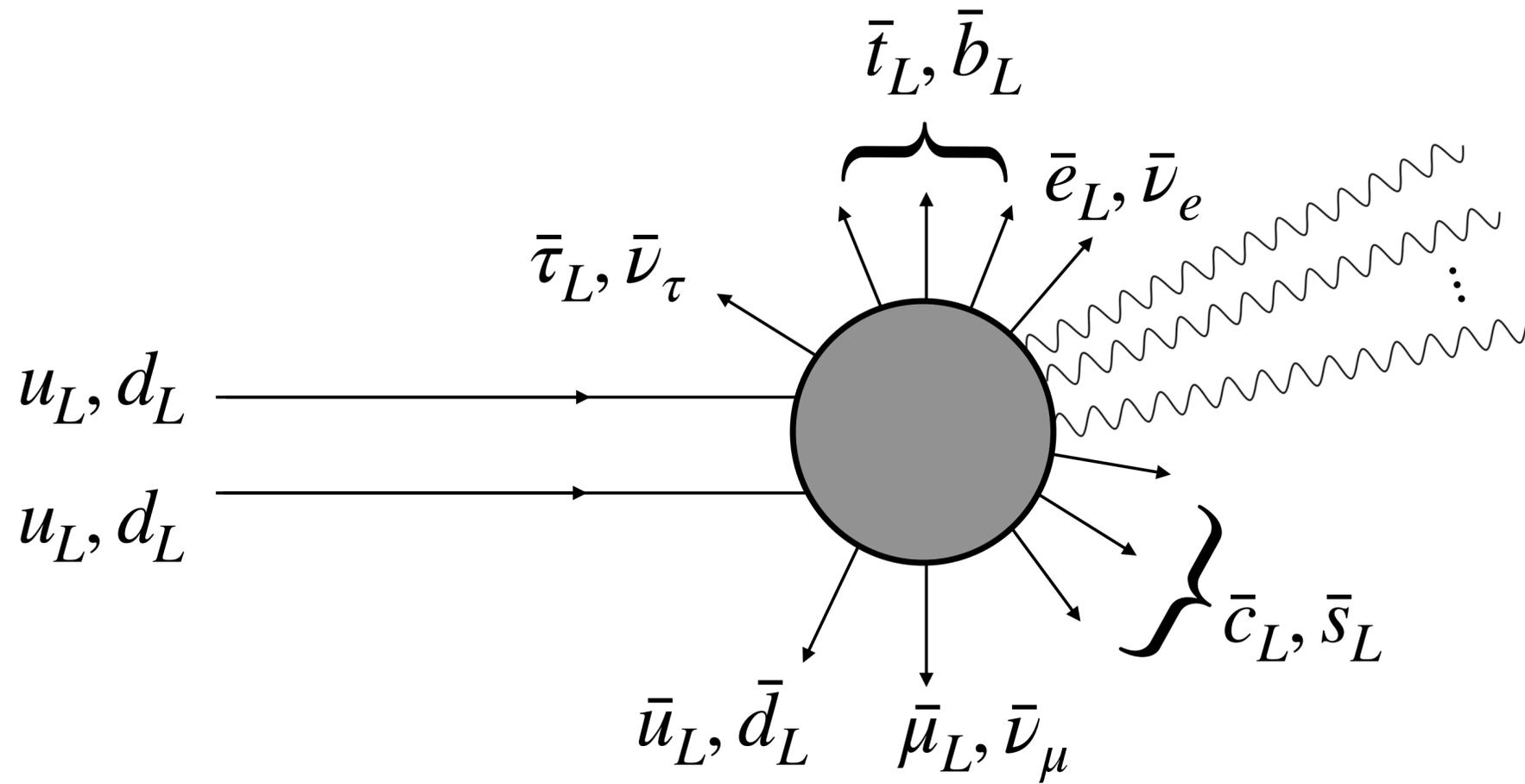
⇒ Events would **spectacularly light up detectors** at experiments!

EW Sphalerons at Colliders?



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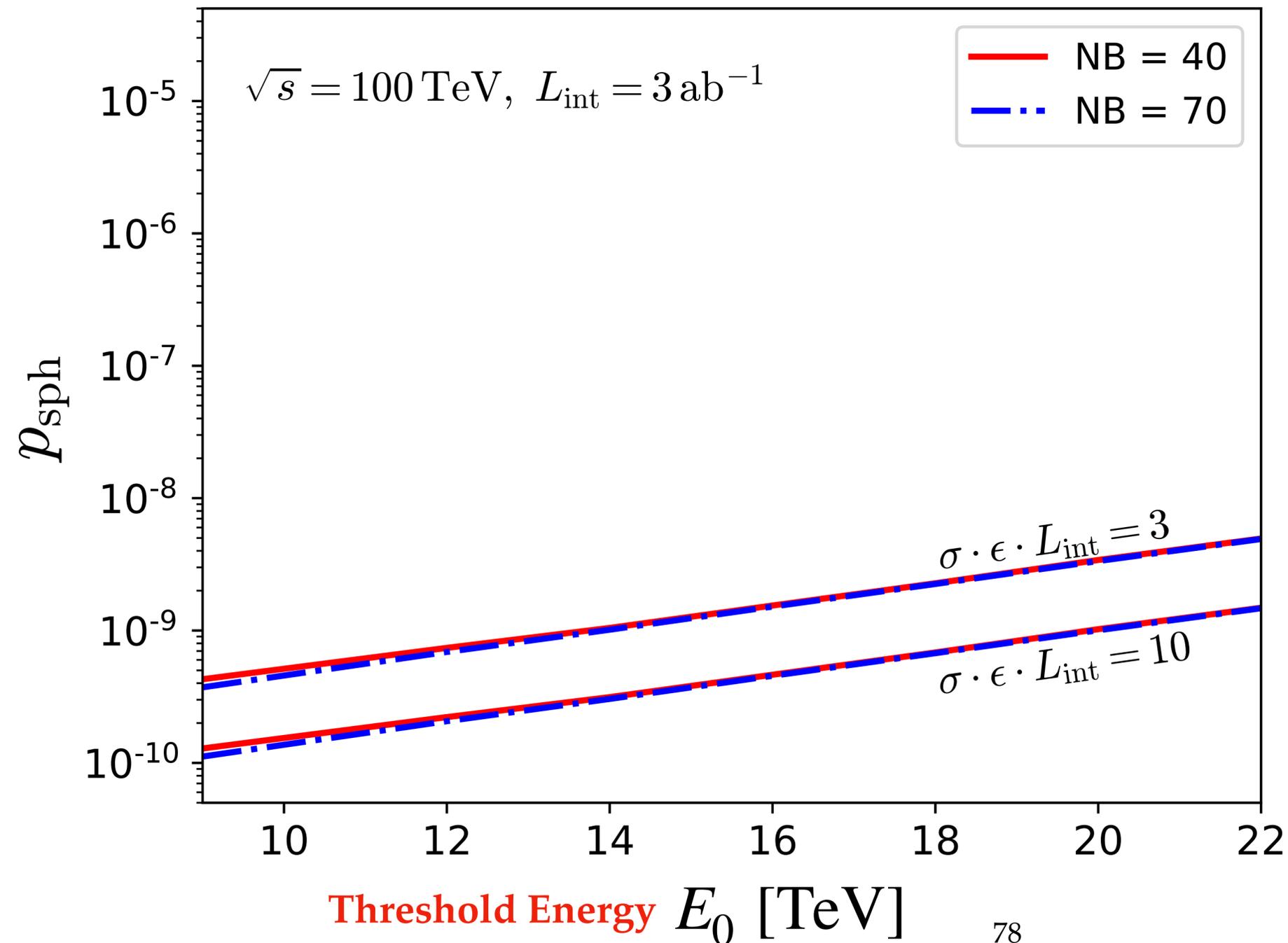
⇒ Events would **spectacularly light up detectors** at experiments!



Sphalerons at the FCC

- Parametrise parton-parton cross section by P_{sph} :

$$\hat{\sigma}(E) = \frac{P_{\text{sph}}}{m_W^2} \Theta(E - E_0)$$



→ Event Generator within HERWIG 7.

[[AP](#), Sakurai, Plätzer, [arXiv:1910.4761](#)]

EW Sphalerons at Colliders?

[[AP](#), Sakurai, Plätzer, [arXiv:1910.4761](#)]

EW Sphalerons at Colliders?

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Homework:

- (i) What can we learn about the Higgs sector and EWBG?
- (ii) New theoretical features in Sphaleron MC.
- (iii) Model discrimination, e.g. VS micro-black holes.
- (iv) Collaboration with experimentalists for measurements.

