#### October 4 Math 2306 sec. 54 Fall 2021

#### **Section 11: Linear Mechanical Equations**

Simple Harmonic Motion

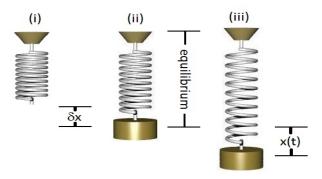


Figure: We track the displacement x from the equilibrium position. Recall  $\delta x$  is the displacement **in** equilibrium. We assume there are no damping nor external forces (for now).

# Simple Harmonic Motion

We derive the equation for displacement. Given initial displacement  $x_0$  and initial velocity  $x_1$ ,

$$x'' + \omega^2 x = 0$$
,  $x(0) = x_0$ ,  $x'(0) = x_1$ , (1)

where  $\omega^2 = \frac{k}{m} = \frac{g}{\delta x}$ . The solution is

$$x(t) = x_0 \cos(\omega t) + \frac{x_1}{\omega} \sin(\omega t)$$
 (2)

called the equation of motion<sup>1</sup>.

*k*-spring constant, *m*-object's mass, *g*-acceleration due to gravity.

Some use this phrase to refer the IVP (1). Others use it to refer to the **solution** to the IVP such as (2).

<sup>&</sup>lt;sup>1</sup>Caution: The phrase equation of motion is used differently by different authors.

# Simple Harmonic Motion

$$x(t) = x_0 \cos(\omega t) + \frac{x_1}{\omega} \sin(\omega t)$$

Characteristics of the system include

- the period  $T = \frac{2\pi}{\omega}$ ,
- the frequency  $f = \frac{1}{7} = \frac{\omega}{2\pi}^2$
- the circular (or angular) frequency  $\omega$ , and
- the amplitude or maximum displacement  $A = \sqrt{x_0^2 + (x_1/\omega)^2}$

<sup>&</sup>lt;sup>2</sup>Various authors call f the natural frequency and others use this term for  $\omega$ .

### Amplitude and Phase Shift

We can formulate the solution in terms of a single sine (or cosine) function. Letting

$$x(t) = x_0 \cos(\omega t) + \frac{x_1}{\omega} \sin(\omega t) = A \sin(\omega t + \phi)$$

requires

$$A = \sqrt{x_0^2 + (x_1/\omega)^2},$$

and the **phase shift**  $\phi$  must be defined by

$$\sin \phi = \frac{x_0}{A}$$
, with  $\cos \phi = \frac{x_1}{\omega A}$ .



# Amplitude and Phase Shift (alternative definition)

We can formulate the solution in terms of a single sine (or cosine) function. Letting

$$x(t) = x_0 \cos(\omega t) + \frac{x_1}{\omega} \sin(\omega t) = A \cos(\omega t - \hat{\phi})$$

requires

$$A = \sqrt{x_0^2 + (x_1/\omega)^2},$$

and this **phase shift**  $\hat{\phi}$  must be defined by

$$\cos \hat{\phi} = \frac{x_0}{A}$$
, with  $\sin \hat{\phi} = \frac{x_1}{\omega A}$ .



# Example

An object stretches a spring 6 inches in equilibrium. Assuming no driving force and no damping, set up the differential equation describing this system.

No driving and no domping => simple harmonic motion. The ODE is Mx"+ kx = 0, in Standard form  $X'' + \omega^2 X = 0$ we need to identify w2. We don't have enough in to to find

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on and le, but we can use

$$\omega^2 = \frac{9}{8 \times 10^{-3}}$$

The displacement in equilibrium 5x=6in.

$$\delta x = 6 in \left(\frac{14t}{12in}\right) = \frac{1}{2} ft$$

The ODE is

$$X'' + G4X = 0$$

### Example

A 4 pound weight stretches a spring 6 inches. The mass is released from a position 4 feet above equilibrium with an initial downward velocity of 24 ft/sec. Find the equation of motion, the period, amplitude, phase shift, and frequency of the motion. (Take q = 32 ft/sec<sup>2</sup>.)

The oDE is 
$$mx'' + kx = 0$$
, or  $m$  standard form  $x'' + \omega^2 x = 0$ .  
Since  $8x = 6in$ , we know  $\omega^2 = 64$ .  
Let's find  $m$  and  $k$ .  
The mass  $m = \frac{W}{2} = \frac{416}{32} \frac{1}{12} \frac{1}{$ 

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The spring constant  $k = \frac{W}{\delta x} = \frac{41b}{2ft} = 8 \frac{16}{ft}$ 

$$\Rightarrow x'' + 64x = 0$$

We have initial conditions

4ft above

24 Alsec vard

The characteristic equation (in r) is

$$(^{2} + 64 = 0)$$

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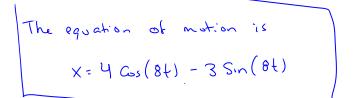
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$$\times'(t) = -8C_1 \sin(8t) + 8C_2 \cos(8t)$$
  
 $\times(6) = C_1 \cos(6) + C_2 \sin(6) = 4$ 

$$X'(0) = -845 \cdot x(0) + 86 \cdot x \cdot C_{0}(0) = -24$$
  
 $86 \cdot x - 24 \Rightarrow C_{0} = -3$ 



The period 
$$T = \frac{2\pi}{\omega} = \frac{2\pi}{8} = \frac{\pi}{4}$$

The amplitude
$$A = \sqrt{4^2 + (-3)^2} = 5$$

$$Sin \varphi = \frac{4}{5}$$
 and  $Cos \varphi = \frac{-3}{5}$ 

$$\phi = \cos^{1}\left(\frac{-3}{5}\right) \approx 2.21$$

### Free Damped Motion

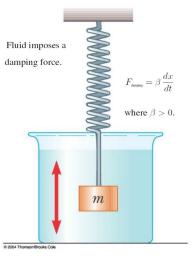


Figure: If a damping force is added, we'll assume that this force is proportional to the instantaneous velocity.

# Free Damped Motion

Now we wish to consider an added force corresponding to damping—friction, a dashpot, air resistance.

Total Force = Force of spring + Force of damping

$$m\frac{d^2x}{dt^2} = -\beta\frac{dx}{dt} - kx \implies \frac{d^2x}{dt^2} + 2\lambda\frac{dx}{dt} + \omega^2x = 0$$

where

$$2\lambda = \frac{\beta}{m}$$
 and  $\omega = \sqrt{\frac{k}{m}}$ .

Three qualitatively different solutions can occur depending on the nature of the roots of the characteristic equation

$$r^2 + 2\lambda r + \omega^2 = 0$$
 with roots  $r_{1,2} = -\lambda \pm \sqrt{\lambda^2 - \omega^2}$ .



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# Case 1: $\lambda^2 > \omega^2$ Overdamped

$$x(t) = e^{-\lambda t} \left( c_1 e^{t\sqrt{\lambda^2 - \omega^2}} + c_2 e^{-t\sqrt{\lambda^2 - \omega^2}} \right)$$

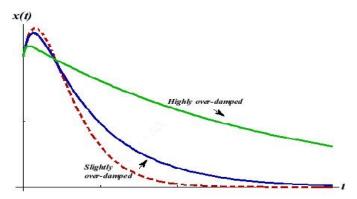


Figure: Two distinct real roots. No oscillations. Approach to equilibrium may be slow.

# Case 2: $\lambda^2 = \omega^2$ Critically Damped

$$x(t) = e^{-\lambda t} \left( c_1 + c_2 t \right)$$

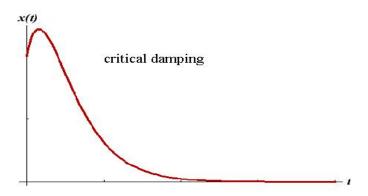


Figure: One real root. No oscillations. Fastest approach to equilibrium.

# Case 3: $\lambda^2 < \omega^2$ Underdamped

$$\mathbf{X}(t) = \mathbf{e}^{-\lambda t} \left( \mathbf{C}_1 \cos(\omega_1 t) + \mathbf{C}_2 \sin(\omega_1 t) \right), \quad \omega_1 = \sqrt{\omega^2 - \lambda^2}$$

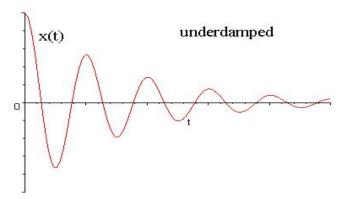


Figure: Complex conjugate roots. Oscillations occur as the system approaches (resting) equilibrium.



# Comparison of Damping

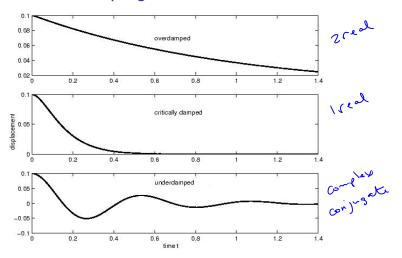


Figure: Comparison of motion for the three damping types.

